Dark Energy's Dual Personality



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Dark energy evidence I - Geometric distance measures



Dark Energy evidence II – Growth of structure

So far lots focus on background evolution, little on perturbations... Yet the growth of structure provides lots of information !

Weak lensing







How do we interpret dark energy?



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How do we interpret dark energy?

Focus of this talk:

- 1. Can modified gravity be a viable dark energy solution?
- 2. Can we distinguish modified gravity from L and matter?
- 3. Can we constrain dark sector interactions?





New matter?



Why modify gravity?

- Many theories predict deviations from GR
- Conceivably, we don't yet have the correct large (and small?) scale theory of gravity?
- The starting point... the action describing motion in GR



Why modify gravity?

- Simplistically modifications -> scale/ time dependent G
- Scalar tensor gravity gravity mediated by scalar field ϕ ?
 - e.g. Brans-Dicke theory



- Higher order terms?
 - e.g. f(R) theories



f(R) = A/R+... Low curvature (late time) modifications - dark energy $f(R) = AR^2+...$ High curvature (early time) modifications - inflation

Why modify gravity?

Gravity is a higher dimensional theory?
e.g. Dvali, Gabadaze, Porrati PLB 485:208(2000)

$$\int d^4x \sqrt{-g} \frac{m_p^2}{2} R \quad \Longrightarrow \quad \int d^5x \sqrt{-g^{(5)}} \frac{m_p^{(5)3}}{2} R^{(5)}$$

- Induced intrinsic curvature on brane
- $I >> I_c 5D$ gravity (1/r²), $I >> I_c 4D$ gravity (1/r)





Constraining modifications to gravity



Cosmological evolution in modified gravity

Modified Friedmann and acceleration equations \rightarrow acceleration with normal matter (P>0)

Example: f(R) gravity

Standard GR

modifications

$$\frac{H^2 + \frac{f}{6} + H\dot{f}_R = \frac{\kappa^2}{3}\rho}{\frac{\ddot{a}}{a} - H^2 f_R + \frac{a^2}{6}f + \frac{3}{2}H\dot{f}_R + \frac{1}{2}\ddot{f}_R = -\frac{\kappa^2}{6}\left(\rho + 3P\right)}$$

An alternative perspective

<u>"conformal transformation"</u> Coordinate change (redefine the metric) from Jordan frame to Einstein frame

$$g^{(J)}_{\mu\nu} = e^{\sqrt{\frac{2}{3}}\kappa C\phi}g^{(E)}_{\mu\nu}$$

Jordan frame

- Modified gravity $F(\Phi)R$, or F(R)
- Minimally coupled matter
- Matter follows geodesics



Einstein frame

- Gravity is GR
- Matter coupled to scalar $\boldsymbol{\varphi}$ with strength C
- Matter doesn't follow geodesics
- Scalar potential $V(\phi)$ derived from F

Evolution viewed from the Einstein frame

• GR + extra matter (scalar field)

$$\frac{\ddot{a}}{a} = -\frac{\kappa^2}{6} \left(\rho + 3P + 2\dot{\phi}^2 - 2V(\phi)\right)$$

• Matter and scalar coupled together as fifth force with strength C

$$\ddot{\phi} + 3H\dot{\phi} + V'(\phi) = \sqrt{\frac{2}{3}}C\kappa(\rho - 3P)$$
$$\rho \propto e^{-\sqrt{\frac{2}{3}}\kappa C\phi}a^{-3(1+w)}$$

Strength of coupling C
determined by theoryGRC = 0F(R)F(R)C = 1/2 $F(\Phi)R$ C can take range of values

Acceleration in the Einstein frame perspective

• Acceleration when scalar sits in minimum of effective potential

$$\ddot{\phi} + 3H\phi + V'(\phi) = \sqrt{\frac{2}{3}}C\kappa(\rho - 3P)$$



EF shows possible instabilities in structure growth

• Jean's instability in a general fluid undergoing gravitational collapse

If coupling too strong (minimum too steep) get catastrophic instabilities
 Jean's instability in matter growth with negative speed of sound

$$\ddot{\delta} + 2H\dot{\delta} - 4\pi G_{eff}
ho\delta = 0$$

$$\frac{G_{eff}}{G_N} = e^{-2\kappa\beta\phi} \left[1 + \frac{k^2\beta^2}{4\pi \left(k^2 + V''e^{-\kappa\beta\phi}\right)} \right]$$

N. Afshordi et al. PRD 2005 (MaVaNs) RB, E. Flanagan and M. Trodden PRD 2007

Which frame to use, Jordan or Einstein?

- <u>Benefits of Einstein frame</u>
 - Equations more intuitive (at least to me) in this frame
 - Evolution also easier to understand
- <u>Benefits of Jordan frame</u>
 - Observations e.g. redshifts typically interpreted in Jordan frame
 - Assumes atomic physics in distant galaxy/ supernovae or the CMB is the same as on Earth (I.e. minimally coupled)
- No frame more `physical' than the other
 - all a matter of convenience

Dynamical attractors powerful in classifying viable models

- Finds stable,saddle and unstable fixed points
- Evolution tends to fixed points largely independent of initial conditions
- Gives conditions on theory for viable evolution
 - Matter dominated era Saddle point
 - late time acceleration- Stable or saddle point
- Behavior often generic irrespective of the specific form of the action



Attractors and observational implications



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Dynamical attractors in scalar-tensor theories $F(\Phi)R$

• Perform dynamical analysis in Jordan or Einstein frame

$$x_1 = -\frac{d\ln F}{d\ln a}, \quad x_2 = -\frac{\int F d\Phi}{6FH^2}, \quad x_3 = \frac{\Phi}{6H^2}$$

- Require evolution to satisfy
 - 1) CDM dominated era with $w^{(J)} \sim 0$ saddle point
 - 2) Late time accelerated era saddle point or stable
- Dynamical attractor solutions place restrictions on
 - Size of coupling C In Einstein frame
 - Allowed trajectories in Jordan frame of two variables, m and r

$$m = rac{\Phi}{F} rac{dF}{d\Phi}$$
 $r = -rac{\Phi F}{\int F d\Phi}$

Amendola et al PRD 2006 (f(R) theories) Nishant Agarwal and RB CQG 2008

A viable trajectory



Can observations distinguish between modified gravity & Λ ?

• Yes, look at how structure (galaxies, cluster of galaxies) grows...

$$\rho(x,a) = \bar{\rho}(a) \left[1 + \delta(x,a)\right]$$

• Growth in δ related to gravitational metric perturbations (Φ and Ψ)

$$ds^{2} = -(1+2\Psi)dt^{2} + a^{2}(t)(1-2\Phi)dx^{2}$$

Interactions between matter and gravitational potential

 1) <u>Poisson equation</u>: How potential related to local density

$$abla^2 \Phi(x,a) = rac{ar
ho(a)}{2m_p^2} \delta(x,a)$$



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• 2) <u>Peculiar acceleration equation:</u> How matter responds to potential

$$g_i =
abla_i \Psi$$



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• 3) <u>Relationship between two potentials:</u> Presence of shear stresses

$$\Phi = \Psi$$



These equations alter in modified gravity:

1) Can modify how potential is related to density (Poisson's equation)

$$\nabla^2 \Phi = \frac{\bar{\rho}}{2m_p^2} Q(x,a) \delta$$

2) Can introduce intrinsic shear stress

$$\Psi = [1 + \eta(x, a)]\Phi$$

Shear stress alters how matter responds to the potential (peculiar accelⁿ)

$$g_i = \nabla_i (1+\eta) \Phi$$

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Example : DGP

• Scale independent modifications

$$egin{array}{l} Q(a) = 1 - \displaystylerac{1}{3eta(a)} \ \eta(a) = \displaystylerac{2}{3eta(a)-1} \end{array}$$

$$l_c = \frac{m_p^{(4)2}}{2m_p^{(5)3}}$$

(1)0

$$\beta(a) \equiv 1 - \frac{2H(a)^2 l_c^2}{2H(a)l_c - 1}$$

- Different observations measure different quantities
 - Weak lensing distortions $\Phi+\Psi$ -> Q(1+ $\eta/2$)
 - Galaxy number counts Φ -> Q
 - Peculiar velocities Ψ -> Q(1+ η)
- By comparing observations we can distinguish modifications from GR

$$\hat{E}_G \propto \frac{\text{lensing} - \text{galaxy correlation}}{\text{velocity} - \text{galaxy correlation}}$$

$$\hat{E}_G = \left. rac{
abla^2(\Psi+\Phi)}{3H_0^2rac{d\delta}{d\ln a}}
ight|_{k=rac{l}{\chi},ar{z}}$$

Zhang, Liguori, RB, Dodelson PRL 2007



Measuring η





Error bars for SKA.

ADEPT, HSHS 21cm, and LSST could provide more immediate constraints

Zhang, RB, Liguori, Dodelson in preparation

Accurate non-linear growth predictions

- Weak lensing and galaxy correlations measured when over-densities large
 - <u>"non-linear growth"</u>
- <u>N-body simulations</u> needed to model growth numerically
- Analytical fits done of GR simulations
 - Smith and Peacock
 - Peacock and Dodds
- <u>Need simulations of non-linear growth in</u> <u>modified gravity theories</u> to accurately predict effect on weak lensing and galaxy surveys



F. Stabenau and B. Jain PRD 2006 I. Laszlo and RB PRD 2007

A simple 5D gravity model



I. Laszlo and RB, PRD 2007

Constraining a coupling between dark energy and dark matter?

• What if purely dark sector interactions? Just couple scalar to CDM.

$$\ddot{\phi} + 3H\dot{\phi} + V'(\phi) = \sqrt{\frac{2}{3}}C\kappa(\rho - 3P)$$
$$\rho \propto e^{-\sqrt{\frac{2}{3}}\kappa C\phi}a^{-3(1+w)}$$

• Interaction acts as species dependent G

$$\frac{G_{cc}G_{bb}}{G_{bc}^2} = 1 + \frac{4}{3}C^2$$

RB, E. Flanagan, I. Laszlo, M. Trodden, arxiv: 0808.1105

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Constraining a coupling between dark energy and dark matter?

Attractor and tracker solutions again important



RB, E. Flanagan, I. Laszlo, M. Trodden, arxiv: 0808.1105

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CMB, LSS and SN observations place very tight constraints on any such interaction

Coupling < 0.07 at 95% confidence level



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Conclusions

- 1. We should establish empirically if dark energy caused by Λ , or not
 - Late time acceleration can arise in matter or gravity based theories
 - Can affect the growth of structure differently
- Dynamical attractor analysis powerful for establishing conditions for viability
 Constrain broad properties of modified gravity and coupled theories
- 3. We <u>can</u> hope to differentiate between modified gravity and Λ
 - By contrasting different observations of large scale structure
- 4. Theorists should also aim to minimize systematic errors
 - Shouldn't assume a priori that fits to GR simulations hold here
 - However, weakly non-linear growth well fit by current analytical fits