

# Averaging Robertson-Walker Cosmologies

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Motivation

Standard Cosmology  
Averaging in  
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# Motivation

# Standard Cosmology

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- Copernican Principle + CMB observations  $\Rightarrow$  Universe homogeneous and isotropic.

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- Copernican Principle + CMB observations  $\Rightarrow$  Universe homogeneous and isotropic.
- Robertson-Walker cosmology: foliate spacetime with maximally-symmetric three-spaces
  - Line element:  $ds^2 = -dt^2 + a^2(t)\delta_{ij}dx^i dx^j$
  - Friedmann equation:  $(\dot{a}/a)^2 = (8\pi G/3)\bar{\rho} + \Lambda/3$
  - Raychaudhuri equation:  $\ddot{a}/a = -(4\pi G/3)(\bar{\rho} + \bar{p}) + \Lambda/3$
  - Perturb metric with  $\mathcal{O}(\epsilon) \approx 10^{-5}$

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  - Perturb metric with  $\mathcal{O}(\epsilon) \approx 10^{-5}$
- We have assumed the existence of an average and added perturbations

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- An implicit averaging in cosmology transfers local equations to global cosmology; should be made explicit
- $\langle \partial_t \rho \rangle \neq \partial_t \langle \rho \rangle \Rightarrow$  Naïve EFE for assumed averages does not reflect a true average of small-scale physics.

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- An implicit averaging in cosmology transfers local equations to global cosmology; should be made explicit
- $\langle \partial_t \rho \rangle \neq \partial_t \langle \rho \rangle \Rightarrow$  Naïve EFE for assumed averages does not reflect a true average of small-scale physics.
- We should be using

$$\langle G_{\mu\nu}(g_{\mu\nu}) \rangle = 8\pi G \langle T_{\mu\nu} \rangle + \Lambda \langle g_{\mu\nu} \rangle$$

instead of

$$G_{\mu\nu}(\langle g_{\mu\nu} \rangle) = 8\pi G \langle T_{\mu\nu} \rangle + \Lambda \langle g_{\mu\nu} \rangle .$$

- “Backreaction” may not be dark energy, but *all* cosmological models should be properly averaged
- Aim: Express Buchert equations in general form, apply to range of perturbed Robertson-Walker models from radiation domination to present day.

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# Formalism: 3+1 Split

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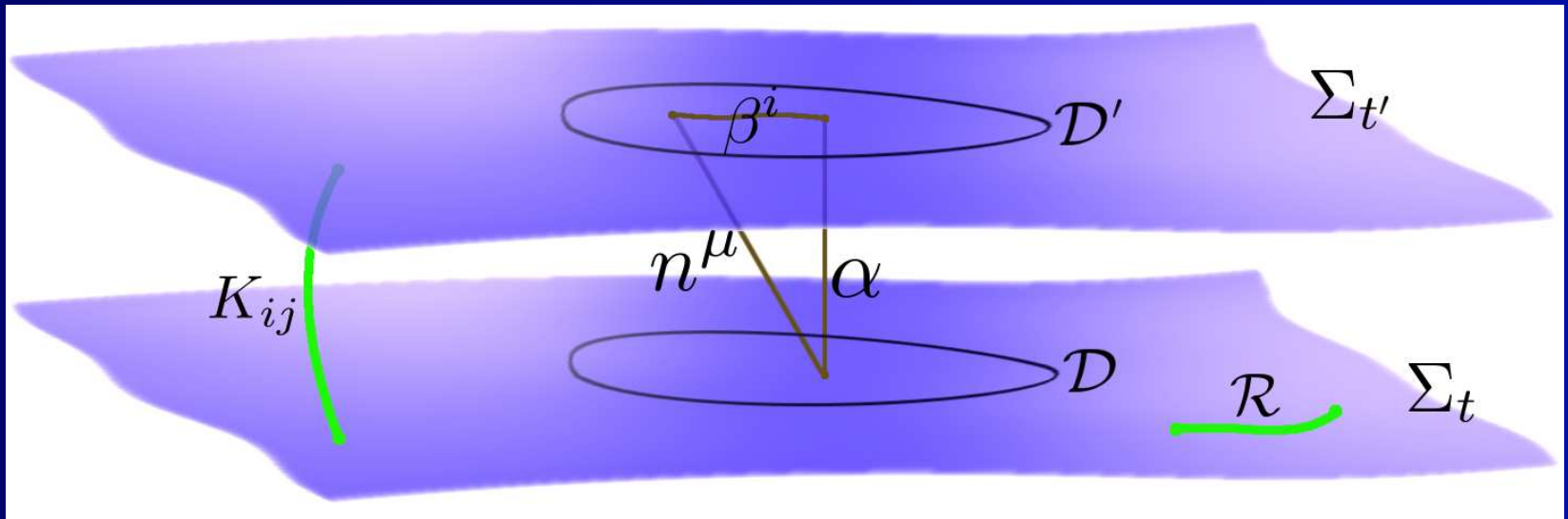
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- Line-element:  $ds^2 = -\alpha^2 dt^2 + h_{ij} dx^i dx^j$
- Fluids ( $\Lambda, \phi, b, \text{CDM}, \gamma, \nu$ ):  $\rho = n^\mu n^\nu T_{\mu\nu}, j_i = -n^\mu T_{i\mu}, S_{ij} = T_{ij}$
- Perfect fluids,  $T_{\mu\nu} = (\rho + p)u_\mu u_\nu + pg_{\mu\nu}$ :  

$$\rho = n^\mu n^\nu T_{\mu\nu} = \rho (n^\mu u_\mu)^2 + p \left( (n^\mu u_\mu)^2 - 1 \right),$$

$$S = T_i^i = 3p + (\rho + p) u^i u_i .$$

# Formalism: Buchert Averaging

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## ■ Select average

$$\langle A \rangle = \frac{1}{V} \int_{\mathcal{D}} A \sqrt{h} d^3 \mathbf{x},$$

Define averaged “scale factor” and Hubble rate by

$$3H_{\mathcal{D}} = 3 \frac{\dot{a}_{\mathcal{D}}}{a_{\mathcal{D}}} = \frac{\dot{V}}{V} = -\frac{1}{V} \int_{\mathcal{D}} \alpha K \sqrt{h} d^3 \mathbf{x} = -\langle \alpha K \rangle = \langle \mathcal{H} \rangle,$$

## ■ Buchert equations:

$$\left( \frac{\dot{a}_{\mathcal{D}}}{a_{\mathcal{D}}} \right)^2 = \frac{8\pi G}{3} \langle \alpha^2 \varrho \rangle + \frac{\Lambda}{3} \langle \alpha^2 \rangle - \frac{1}{6} (\mathcal{Q}_{\mathcal{D}} + \mathcal{R}_{\mathcal{D}})$$

$$\frac{\ddot{a}_{\mathcal{D}}}{a_{\mathcal{D}}} = -\frac{4\pi G}{3} \langle \alpha^2 (\varrho + S) \rangle + \frac{\Lambda}{3} \langle \alpha^2 \rangle + \frac{1}{3} (\mathcal{Q}_{\mathcal{D}} + \mathcal{P}_{\mathcal{D}})$$

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- Kinematical “backreaction”:

$$Q_{\mathcal{D}} = \left\langle \alpha^2 \left( K^2 - K_j^i K_i^j \right) \right\rangle - \frac{2}{3} \langle \alpha K \rangle^2$$

- Dynamical “backreaction”:  $\mathcal{P}_{\mathcal{D}} = \langle \dot{\alpha} K \rangle + \langle \alpha D^i D_i \alpha \rangle$
- Curvature contribution:  $\mathcal{R}_{\mathcal{D}} = \langle \alpha^2 \mathcal{R} \rangle$
- Deviation from average density and pressure:

$$\frac{3\mathcal{T}_{\mathcal{D}}^{(a)}}{8\pi G} = \langle \alpha^2 \varrho_{(a)} \rangle - \bar{\rho}_{(a)}, \quad \frac{3\mathcal{S}_{\mathcal{D}}^{(a)}}{4\pi G} = \langle \alpha^2 S_{(a)} \rangle - \bar{S}_{(a)}$$

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- The Buchert equations can then be written as

$$\left(\frac{\dot{a}_{\mathcal{D}}}{a_{\mathcal{D}}}\right)^2 = \frac{8\pi G}{3} \sum_a \bar{\rho}_{(a)} + \frac{\Lambda}{3} + \frac{8\pi G}{3} \bar{\rho}_{\text{eff}},$$

$$\frac{\ddot{a}_{\mathcal{D}}}{a_{\mathcal{D}}} = -\frac{4\pi G}{3} \sum_a \left(\bar{\rho}_{(a)} + \bar{S}_{(a)}\right) + \frac{\Lambda}{3} - \frac{4\pi G}{3} (\bar{\rho}_{\text{eff}} + \bar{S}_{\text{eff}})$$

with effective correction fluid

$$\frac{8\pi G}{3} \bar{\rho}_{\text{eff}} = \sum_a \mathcal{T}_{\mathcal{D}}^{(a)} + \langle \alpha^2 - 1 \rangle \frac{\Lambda}{3} - \frac{1}{6} (\mathcal{Q}_{\mathcal{D}} + \mathcal{R}_{\mathcal{D}}),$$

$$16\pi G \bar{p}_{\text{eff}} = 4 \sum_a \mathcal{S}_{\mathcal{D}}^{(a)} - 2\Lambda \langle \alpha^2 - 1 \rangle + \frac{1}{3} (\mathcal{R}_{\mathcal{D}} - 3\mathcal{Q}_{\mathcal{D}} - 4\mathcal{P}_{\mathcal{D}}),$$

$$w_{\text{eff}} = \frac{1}{3} \frac{\mathcal{R}_{\mathcal{D}} - 3\mathcal{Q}_{\mathcal{D}} - 4\mathcal{P}_{\mathcal{D}} + 12 \sum_a \mathcal{S}_{\mathcal{D}}^{(a)} - 6\Lambda \langle \alpha^2 - 1 \rangle}{\mathcal{R}_{\mathcal{D}} + \mathcal{Q}_{\mathcal{D}} - 6 \sum_a \mathcal{T}_{\mathcal{D}}^{(a)} - 2\Lambda \langle \alpha^2 - 1 \rangle}.$$

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- Identify ADM and Newtonian co-ordinates (c.f. Mukhanov et. al.)

$$ds^2 = -(1+2\Psi)dt^2 + a^2(t)(1-2\Phi)\delta_{ij}dx^i dx^j = -\alpha^2 dt^2 + h_{ij}dx^i dx^j$$

- $a_{\mathcal{D}}(t)$  is “observational”,  $a(t)$  is “physical” – drawback of re-averaging assumed average (Kolb, Marra, Matarrese 08; IB, Behrend, Robbers 08)

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- $a_{\mathcal{D}}(t)$  is “observational”,  $a(t)$  is “physical” – drawback of re-averaging assumed average (Kolb, Marra, Matarrese 08; IB, Behrend, Robbers 08)

- Quickly find

$$\frac{\dot{a}_{\mathcal{D}}}{a_{\mathcal{D}}} = \frac{\dot{a}}{a} - \left\langle \dot{\Phi} (1 + 2\Phi) \right\rangle$$

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## ■ Kinematical and dynamical backreactions:

$$Q_{\mathcal{D}} = 6 \left( \langle \dot{\Phi}^2 \rangle - \langle \dot{\Phi} \rangle^2 \right),$$

$$\begin{aligned} \mathcal{P}_{\mathcal{D}} = & \frac{1}{a^2} \langle \nabla^2 \Psi - (\nabla \Psi)^2 + 2\Phi \nabla^2 \Psi - (\nabla \Phi) \cdot (\nabla \Psi) \rangle \\ & + 3 \frac{\dot{a}}{a} \langle \dot{\Psi} - 2\Psi \dot{\Psi} \rangle - 3 \langle \dot{\Psi} \dot{\Phi} \rangle \end{aligned}$$

## ■ Curvature correction:

$$\mathcal{R}_{\mathcal{D}} = \frac{2}{a^2} \langle 2\nabla^2 \Phi + 3(\nabla \Phi)^2 + 4(2\Phi + \Psi)\nabla^2 \Phi \rangle.$$

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## ■ Fluid corrections:

$$\mathcal{T}_{\mathcal{D}} = \frac{8\pi G}{3} \bar{\rho} \langle \delta + 2\Psi + (1 + \bar{w})a^2 v^2 + 2\Psi\delta \rangle,$$

$$\mathcal{S}_{\mathcal{D}} = \frac{4\pi G}{3} \bar{\rho} \langle 3c_s^2 \delta + 6\bar{w}\Psi + (1 + \bar{w})a^2 v^2 + 6c_s^2 \Psi\delta \rangle$$



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### ■ Fluid corrections:

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- Note: alternative gauges – uniform density to simplify  $\mathcal{T}_{\mathcal{D}}$  and  $\mathcal{S}_{\mathcal{D}}$ , uniform curvature to remove  $\mathcal{R}_{\mathcal{D}}$ , synchronous gauge to remove  $\mathcal{P}_{\mathcal{D}}$ .  $\mathcal{Q}_{\mathcal{D}}$  cannot be entirely removed except in EdS matter domination.

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- Boltzmann codes are 1-d, averages are 3-d, so take  $\mathcal{D}$  large enough to employ ergodic principle

# Ergodic Averaging

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- Boltzmann codes are 1-d, averages are 3-d, so take  $\mathcal{D}$  large enough to employ ergodic principle

- Corrections to standard case to be evaluated with cmbeasy:

$$Q_{\mathcal{D}} = 6 \int \mathcal{P}_{\psi}(k) \left| \dot{\Phi} \right|^2 (dk/k) \text{ etc.}$$

- Integration domain  $k \in (1/\eta, 100\text{Mpc}^{-1})$

# $\Lambda$ CDM and EdS: Friedmann and Raychaudhuri Equations

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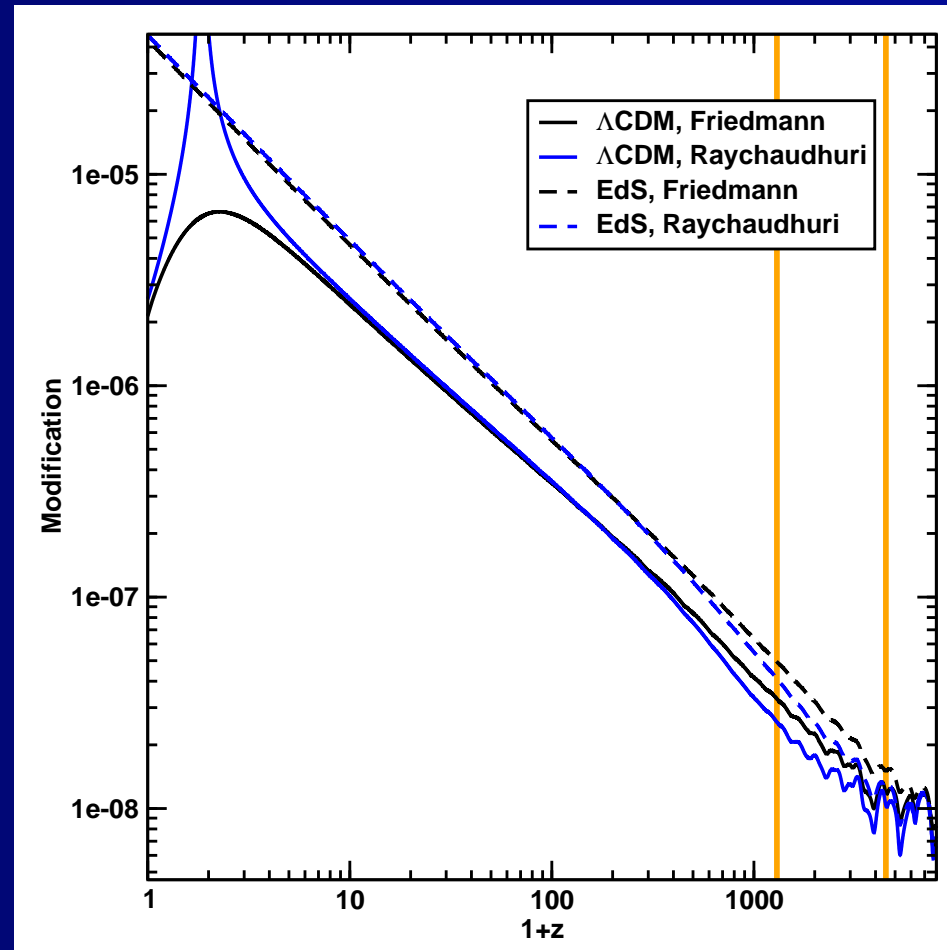
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- Impact at recombination  $\sim 10^{-8}$  – potentially observable with Planck?
- Boosts with Halofit not significant

# $\Lambda$ CDM: $w$

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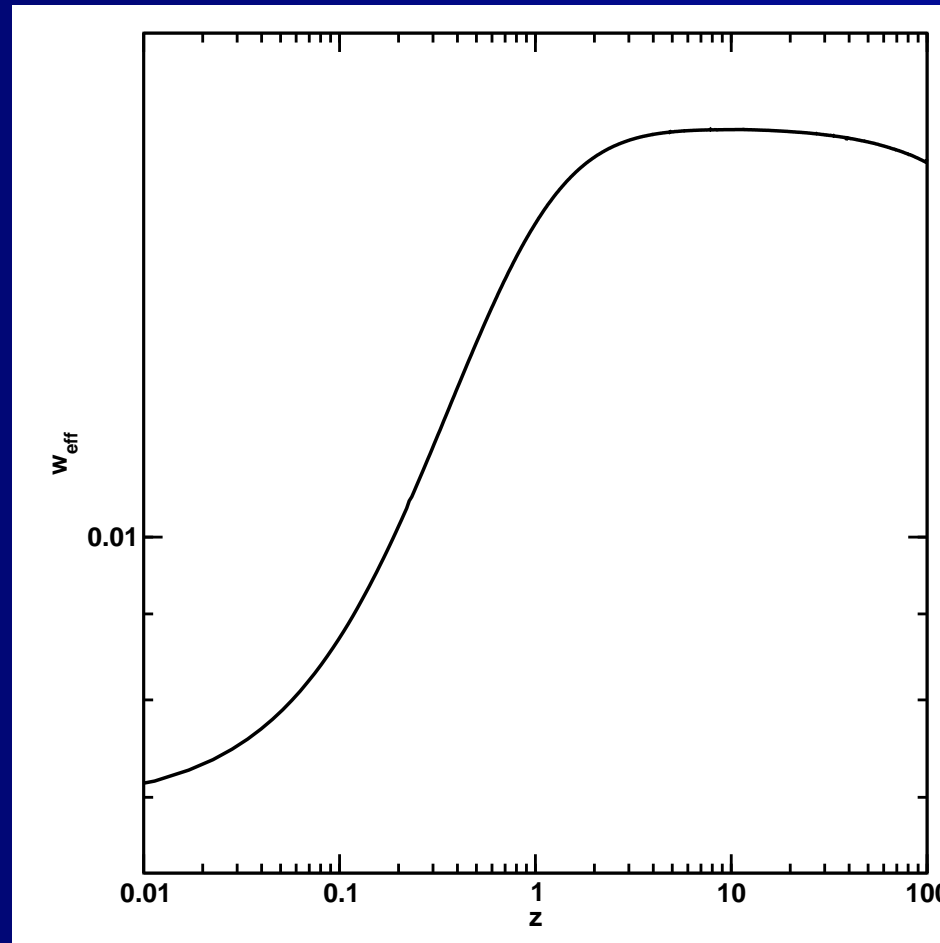
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- For  $z > \sim 100$  find increase up to  $w_{\text{eff}} \approx 0.2$  at  $z \approx 8000$

# Quintessence Cosmology

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- Models tested
  - Early dark energy parameterisation
  - Exponential potential
  - Inverse power-law potential
- Still linear analysis  $\Rightarrow$  still expect small impacts on the observed evolution
- Expect  $w_{\text{eff}}$  to increase with dark matter perturbations – so  $w_{\text{eff}}$  clearest discriminant

# Early Dark Energy

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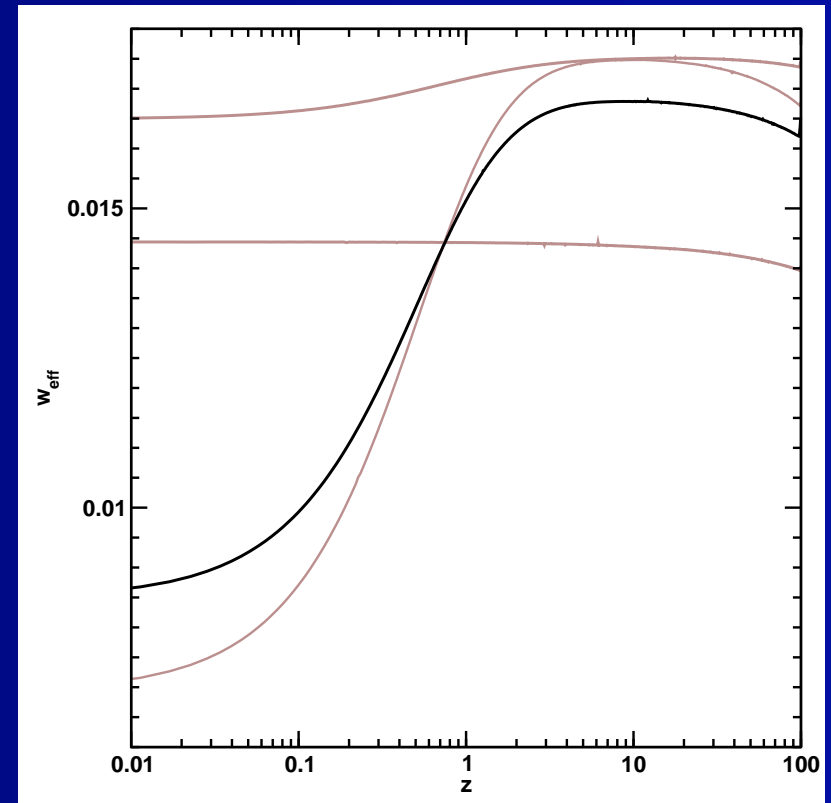
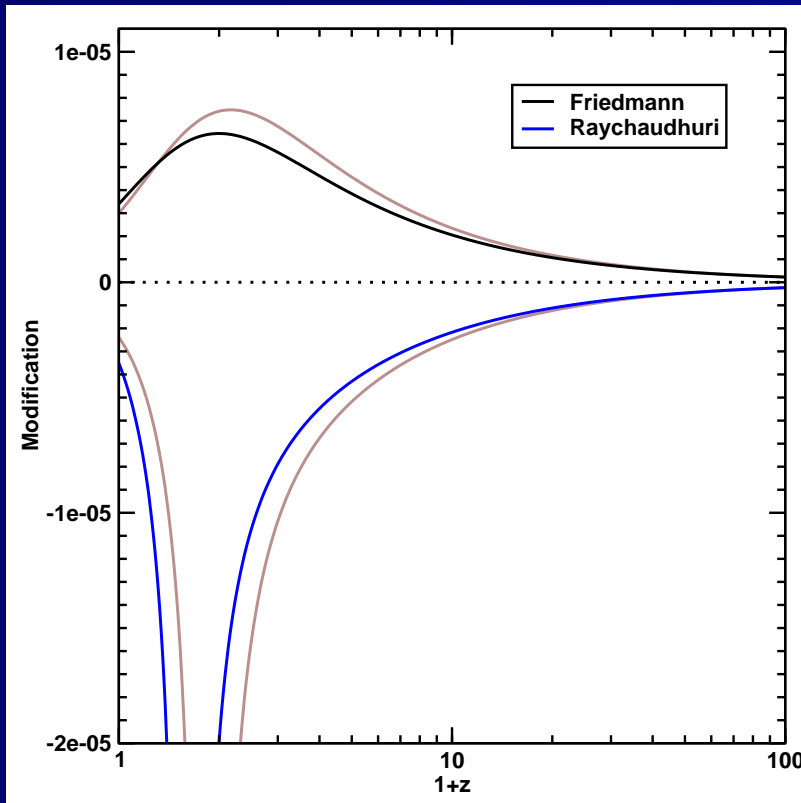
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- Rough model of early dark energy;  $\Omega_{\phi}^{z=0} = 0.7$ ,  $\Omega_{\phi}^{z=\infty} = 0.05$ ,  $w_0 = -0.95$
- Very similar to  $\Lambda$ CDM; larger at present day, smaller at peak





# Inverse Power Law Potential

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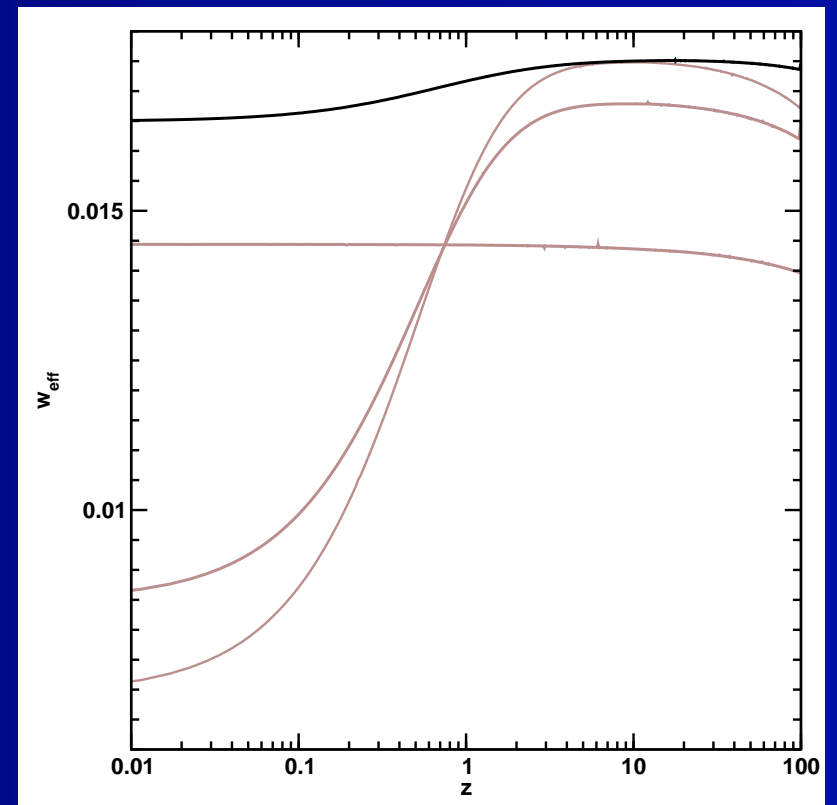
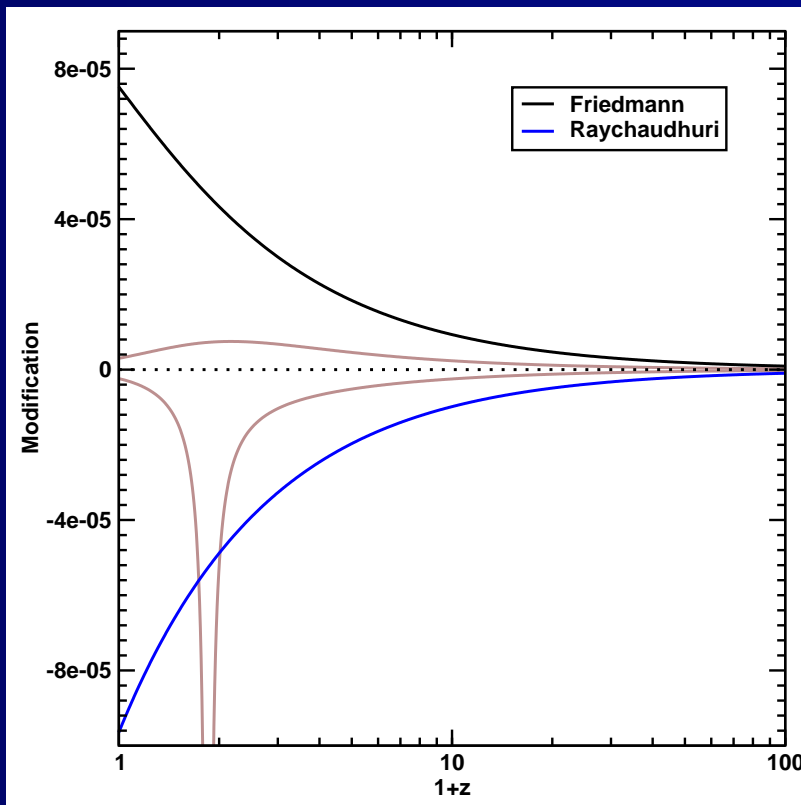
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- Inverse-Power Law Potential (Ratra-Peebles);  $\Omega_\phi = 0.118$ ,  $\Omega_b = 0.046$ ,  $\Omega_c = 0.837$ ,  $n = -2$
- Similar to but smaller than EdS for these parameters



# Exponential Potential

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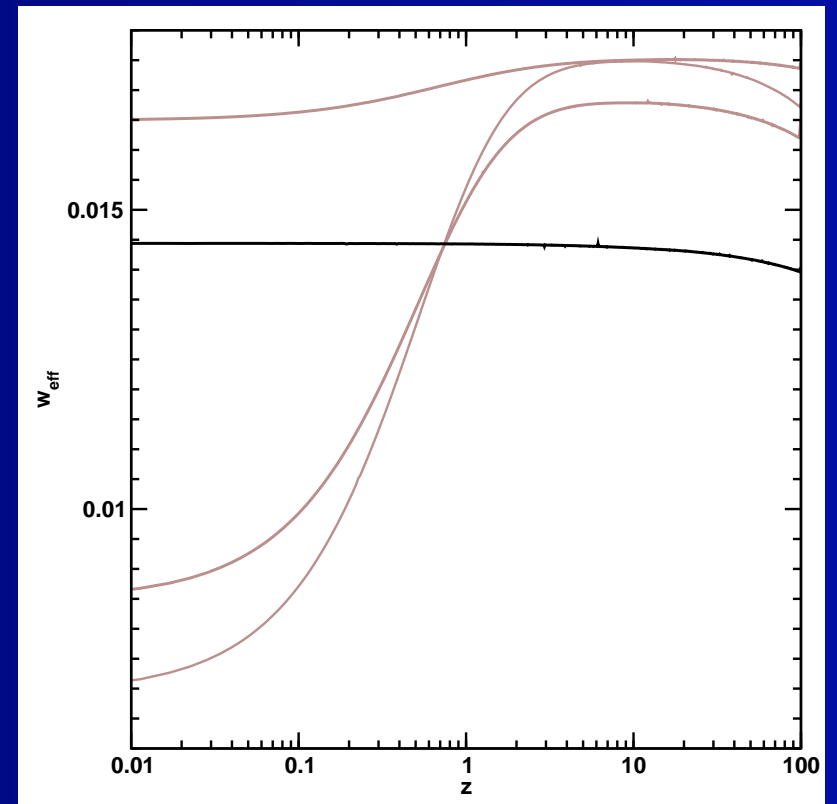
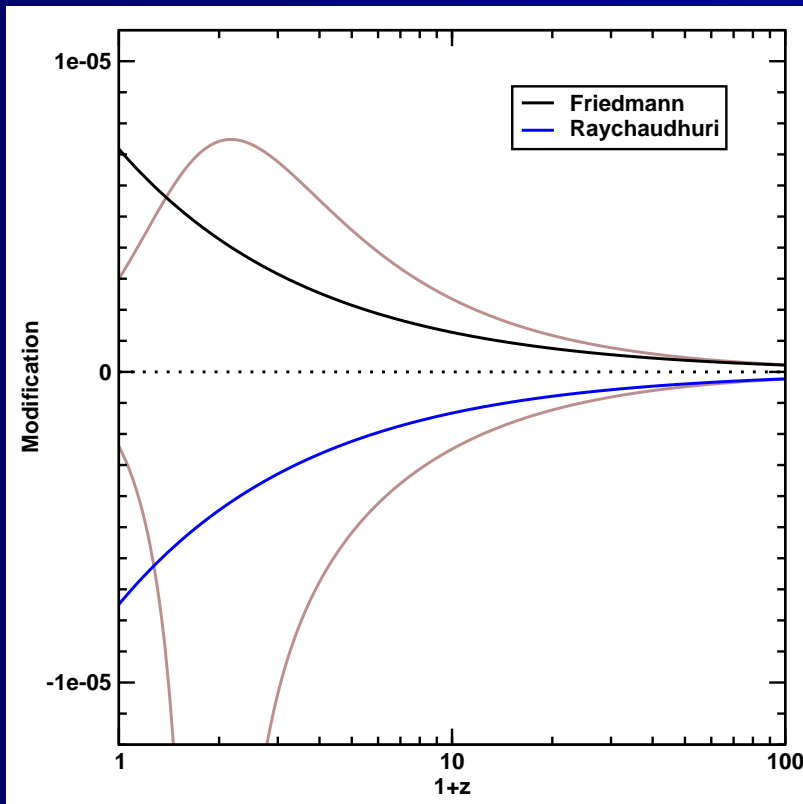
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- Exponential Potential;  $\Omega_\phi = 0.200$ ,  $\Omega_b = 0.041$ ,  $\Omega_m = 0.759$
- Similar to inverse power law



# Equations of State

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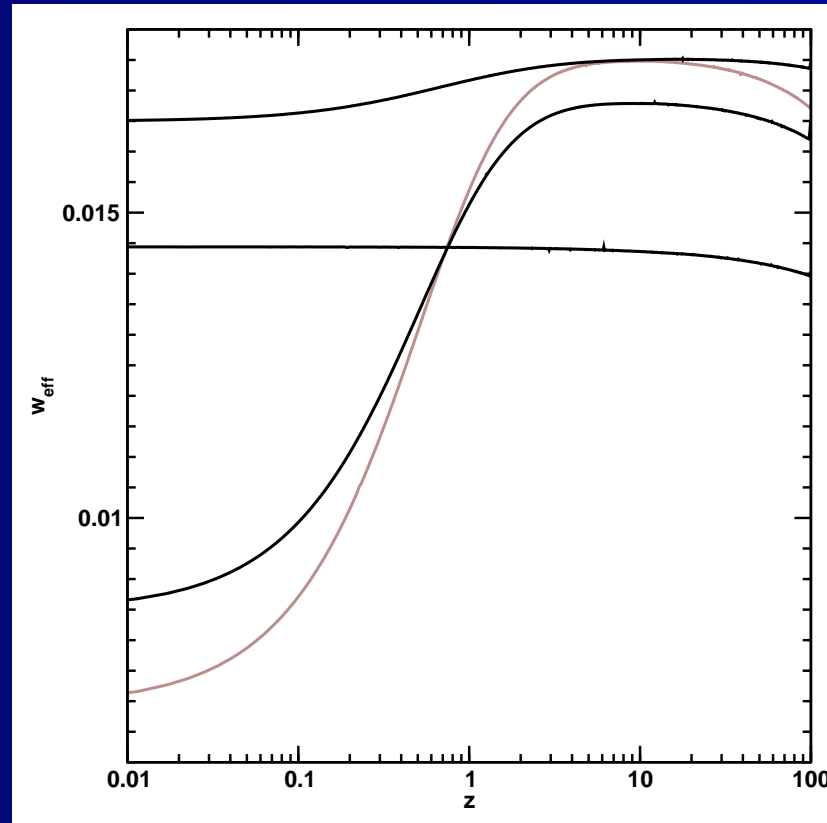
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- $w_{\text{eff}} > 0$  – as before acts against acceleration
- But: this includes quintessence perturbations!
- These differences far too small to observe, but smaller-scale study looks vital

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- Have expressed Buchert equations in multifluid form easily incorporated into general Boltzmann codes for wide variety of models
- Differing *linear* models barely change impact on Friedmann equation; on Raychaudhuri equation it's similar and remains  $\sim 10^{-5}$
- Impact at recombination is close to observable anisotropies  $\Rightarrow$  possible chance of detection?
  - $\Lambda$ CDM:  $w_{\text{eff}} \approx 0.007$
  - Early dark energy:  $w_{\text{eff}} \approx 0.009$
  - Exponential:  $w_{\text{eff}} \approx 0.014$
  - Inverse power law:  $w_{\text{eff}} \approx 0.016$
- Equation of state from quintessence perturbations  $> -1$ : is there a problem with clustering quintessence models? Small scale study is needed (c.f. Wetterich '02).
- CMB observables?
- Non-Linear models
- Modified averaging procedure (Behrend/Nachtmann?)