

Conservation and evolution of the curvature perturbation

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Based on

- D. J. H. Chung and [JG](#), in preparation
- K.-Y. Choi, [JG](#) and D. Jeong, in preparation

Outline

- 1 Introduction
 - Curvature perturbation?
- 2 Single field inflation
 - Equation of motion
 - Subtleties
- 3 Multi field inflation
 - Single and multi field inflation
 - Evolution after multi field inflation
- 4 Conclusions

Curvature perturbation... Err?

Gauge invariant curvature perturbation ζ

$$-\zeta = \psi + H \frac{\delta\rho}{\dot{\rho}}$$

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Generic hypersurface \rightarrow **the one with uniform energy density**

- ① Wiggles in the spatial curvature
- ② Strongly constrained by observations

$$\mathcal{P}_\zeta \sim 10^{-5}, \quad n \sim 0.96, \quad \left| \frac{dn}{d\log k} \right| \lesssim 0.01, \dots$$

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Q: What do we know about ζ ?

Equation of motion of ζ

From **perturbed Klein-Gordon equation** / **Einstein equation** / **variation principle** we can obtain

$$\ddot{\zeta} + \left(\frac{2\ddot{\phi}}{\dot{\phi}} - \frac{2\dot{H}}{H} + 3H \right) \dot{\zeta} - \frac{\nabla^2}{a^2} \zeta = 0$$

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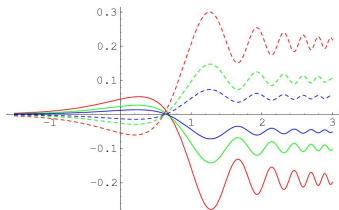
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Conservation of ζ : Slow-roll is **necessary** even for single field case

$$\ddot{\zeta} + \mathcal{O}(H)\dot{\zeta} = 0$$

Case of particle production (1/2)

Conservation of ζ \longleftrightarrow conservation of energy

$$\Delta\mathcal{L} = -\frac{1}{2}g^2\phi^2\chi^2$$

At $\phi_\star = m_\chi/g$, χ particles are resonantly produced

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$$\chi = \cancel{\chi_0} + \delta\chi$$

χ field is **quantum**: How to treat it?

Case of particle production (2/2)

- With $Q = \delta\phi + (\dot{\phi}/H)\psi = -(\dot{\phi}/H)\zeta$,

$$\ddot{Q}_k = -3H\dot{Q}_k - \left\{ \frac{k^2}{a^2} + V'' + g^2 \mathcal{N} \langle \chi^2 \rangle + \frac{1}{m_{\text{Pl}}^2 H} \left[\left(3H + \frac{\dot{H}}{H} \right) \dot{\phi} + 2V' - g \mathcal{N} (m_\chi - g\phi) \langle \chi^2 \rangle \right] \right\} Q_k$$

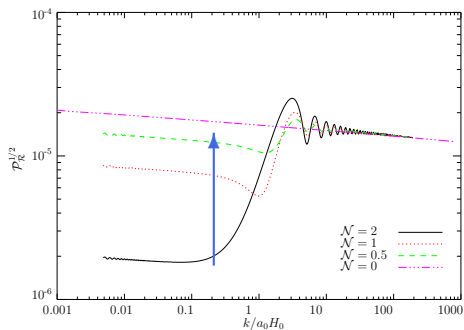
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- $\rho_\chi \lll \rho \rightarrow$ **second order effect**

From conservation equation: $\dot{\mathcal{P}}_\zeta = -\frac{\rho_\chi^2 H}{6(\rho + p)^2} \mathcal{P}_\chi$

Many subtleties

- How to “**turn on**” χ ?
- What is the **range** of wavenumber and time?

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We are trying to find a new perspective on ζ regarding its conservation

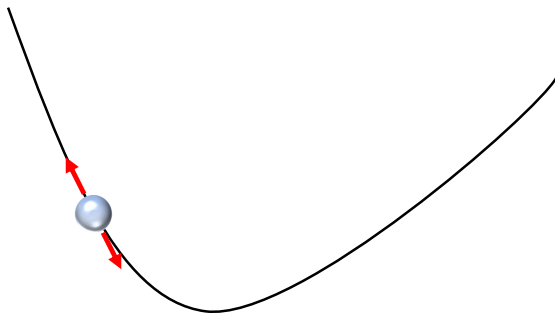


What is different from single field inflation?

There are **more than** one orthogonal directions into which the field can be “*kicked*”

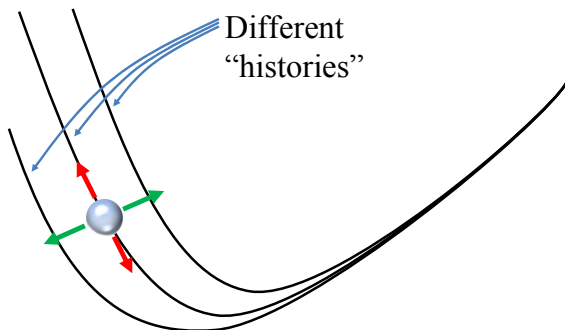
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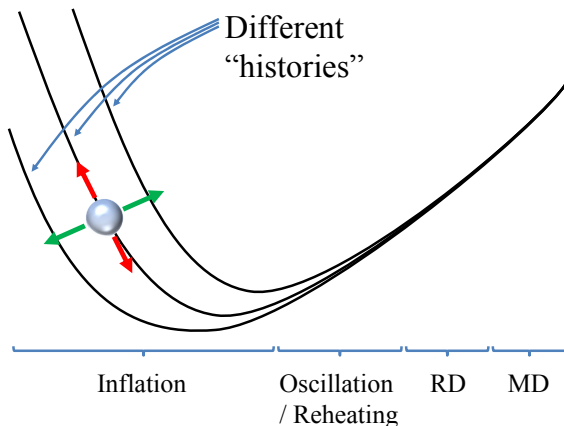
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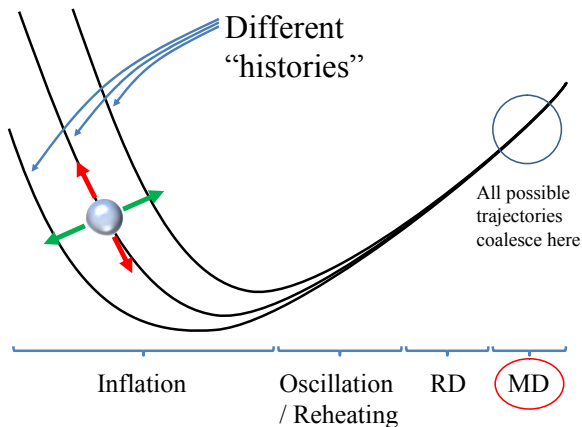
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Multi field inflation and afterwards

For multi field inflation, **inflationary estimates are not enough** and the **evolution after inflation** should be taken into account

e.g. **curvaton** σ should satisfy

- 1 flat potential: $m_\sigma \ll H$
- 2 non-zero amplitude: $\sigma \gtrsim 10^{-8} m_{\text{Pl}}$
- 3 small energy fraction: $V_\sigma \ll V_{\text{tot}}$

Easily satisfied by **individual inflaton field** after multi field inflation

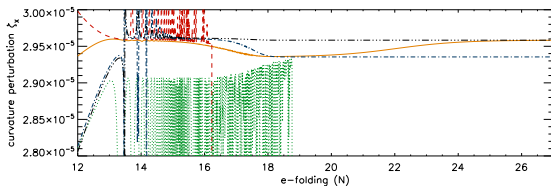
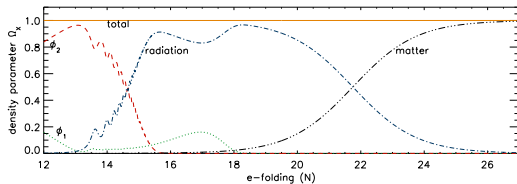
inflaton = curvaton

Evolution of ζ in multi field inflation

Multiple chaotic inflation: $V = \frac{1}{2} \sum_i m_i^2 \phi_i^2$ with $\phi_i \rightarrow \Gamma_\gamma^{(i)}, \Gamma_m^{(i)}$

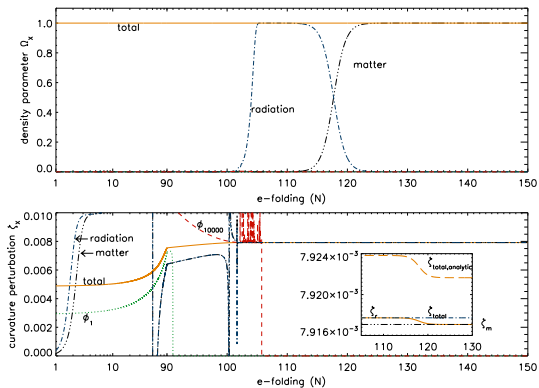
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Matter-radiation isocurvature perturbation

$$S_{\alpha\beta} = 3(\zeta_\alpha - \zeta_\beta) \Rightarrow S_{m\gamma} = \frac{\delta\rho_m}{\rho_m} - \frac{3}{4} \frac{\delta\rho_\gamma}{\rho_\gamma}$$

- 1 $S_{m\gamma} = 0$ in single field inflation
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We can find

- ① $S_{m\gamma} \rightarrow 0$ with large e -folds: $\zeta_i = \text{adiabatic} + \text{non-adiabatic}$
- ② ζ does change after inflation: non-zero δp_{nad} DO exist
- ③ \mathcal{P}_ζ and \mathcal{P}_S have slightly different scale dependence

Conclusions

Conservation of the curvature perturbation is not as simple as a piece of cake

- 1 Single field inflation
 - Slow-roll is *required* to ensure the conservation of ζ
 - But **many subtleties** regarding the conservation of ζ
- 2 Multi field inflation
 - ζ varies **throughout** the whole evolution of the universe
 - Inflationary estimates *may not* work
 - Possibly non-zero, detectable $S_{m\gamma}$