# Gravity and the cosmological constant as superconducting phenomena

Gianluca Calcagni



August 28th, 2008

(日) (日) (日) (日) (日) (日) (日)

S. Alexander, G.C., Superconducting loop quantum gravity and the cosmological constant [0806.4382].

◆□▶ ◆□▶ ◆□▶ ◆□▶ ● ● ● ●

S. Alexander, G.C., Quantum gravity as a Fermi liquid [0807.0225].

# Aims of the talk

◆□▶ ◆□▶ ◆ □▶ ◆ □▶ ● □ ● ● ● ●

• To throw a rock.



- To throw a rock.
- Not to hide the hand.

# The rock

◆□▶ ◆□▶ ◆ □▶ ◆ □▶ ● □ ● ● ● ●

• Setup: Loop quantum gravity with  $\Lambda$ , no matter, and the Chern–Simons state as ground state.

▲□▶ ▲□▶ ▲□▶ ▲□▶ = 三 のへで

• Setup: Loop quantum gravity with  $\Lambda$ , no matter, and the Chern–Simons state as ground state.

◆□▶ ◆□▶ ▲□▶ ▲□▶ □ のQ@

 Assumption: Deform the topological sector (then Λ becomes dynamical).

- Setup: Loop quantum gravity with  $\Lambda$ , no matter, and the Chern–Simons state as ground state.
- Assumption: Deform the topological sector (then Λ becomes dynamical).
- Result 1: spacetime degenerate (1 + 1 dimensions), Hamiltonian modified by a quantum counterterm.

- Setup: Loop quantum gravity with  $\Lambda$ , no matter, and the Chern–Simons state as ground state.
- Assumption: Deform the topological sector (then Λ becomes dynamical).
- Result 1: spacetime degenerate (1 + 1 dimensions), Hamiltonian modified by a quantum counterterm.
- Result 2: Gravity behaves as a Fermi liquid, in particular BCS.

(日) (日) (日) (日) (日) (日) (日)

◆□▶ ◆□▶ ◆ □▶ ◆ □▶ ● □ ● ● ● ●

Bardeen, Cooper, and Schrieffer, Phys. Rev. 108, 1175 (1957).

Bardeen, Cooper, and Schrieffer, Phys. Rev. **108**, 1175 (1957). Nobel Prize in 1972 (Bardeen's second!).

Bardeen, Cooper, and Schrieffer, Phys. Rev. **108**, 1175 (1957). Nobel Prize in 1972 (Bardeen's second!).



▲□ > ▲圖 > ▲目 > ▲目 > ▲目 > ● ④ < @

Bardeen, Cooper, and Schrieffer, Phys. Rev. **108**, 1175 (1957). Nobel Prize in 1972 (Bardeen's second!).



1958-1971: 1 Nobel Prize for studies on condensed matter (Landau).

Bardeen, Cooper, and Schrieffer, Phys. Rev. **108**, 1175 (1957). Nobel Prize in 1972 (Bardeen's second!).



1958-1971: 1 Nobel Prize for studies on condensed matter (Landau). 1973-2007: 10 Prizes awarded in this field.

## The ripples (to be explored)

▲□▶▲圖▶★≧▶★≧▶ 差 の�?

 Λ is exponentially suppressed, nonperturbative phenomenon; at small scales gravity is perturbative, like in QCD confinement.

# The ripples (to be explored)

- A is exponentially suppressed, nonperturbative phenomenon; at small scales gravity is perturbative, like in QCD confinement.
- Geometrical measurements amount to counting Cooper pairs.

# The ripples (to be explored)

- Λ is exponentially suppressed, nonperturbative phenomenon; at small scales gravity is perturbative, like in QCD confinement.
- Geometrical measurements amount to counting Cooper pairs.
- Classically, Cooper pairs are microscopic nonlocal d.o.f. living on the dS boundary (wormholes).

- A is exponentially suppressed, nonperturbative phenomenon; at small scales gravity is perturbative, like in QCD confinement.
- Geometrical measurements amount to counting Cooper pairs.
- Classically, Cooper pairs are microscopic nonlocal d.o.f. living on the dS boundary (wormholes).
- Four dimensions recovered by the spin network defined by the superfluid theory.

- A is exponentially suppressed, nonperturbative phenomenon; at small scales gravity is perturbative, like in QCD confinement.
- Geometrical measurements amount to counting Cooper pairs.
- Classically, Cooper pairs are microscopic nonlocal d.o.f. living on the dS boundary (wormholes).
- Four dimensions recovered by the spin network defined by the superfluid theory.

(日) (日) (日) (日) (日) (日) (日)

Matter is 'hidden' in gravity?

## Nonlocal degrees of freedom on dS horizons



◆□▶ ◆□▶ ◆豆▶ ◆豆▶ □ ● の々で

• Ashtekar variables: connection  $\mathbb{C}$ -field  $A \equiv A^i_{\alpha} \tau_i dx^{\alpha}$  and real triad  $E^i_{\alpha}$ .

▲□▶ ▲□▶ ▲ 三▶ ▲ 三▶ - 三 - のへぐ

- Ashtekar variables: connection  $\mathbb{C}$ -field  $A \equiv A^i_{\alpha} \tau_i dx^{\alpha}$  and real triad  $E^i_{\alpha}$ .
- Scalar, vector, and Gauss constraints:

$$\mathcal{H} = \epsilon_{ijk} E^i \cdot E^j \times \left( B^k + \frac{\Lambda}{3} E^k \right) = 0,$$
  
$$\mathcal{V}_{\alpha} = (E_i \times B^i)_{\alpha} = 0, \quad \mathcal{G}_i = D_{\alpha} E_i^{\alpha} = 0.$$

▲□▶ ▲□▶ ▲三▶ ▲三▶ - 三 - のへで

- Ashtekar variables: connection  $\mathbb{C}$ -field  $A \equiv A^i_{\alpha} \tau_i dx^{\alpha}$  and real triad  $E^i_{\alpha}$ .
- Scalar, vector, and Gauss constraints:

$$\mathcal{H} = \epsilon_{ijk} E^i \cdot E^j \times \left( B^k + \frac{\Lambda}{3} E^k \right) = 0,$$
  
 $\mathcal{V}_{\alpha} = (E_i \times B^i)_{\alpha} = 0, \mathcal{G}_i = D_{\alpha} E_i^{\alpha} = 0.$ 

(日) (日) (日) (日) (日) (日) (日)

• Quantum theory:  $E \rightarrow \hat{E}_i^{\alpha} = -\delta/\delta A_{\alpha}^i$ ,  $\hat{A}_{\alpha}^i$  multiplicative.

- Ashtekar variables: connection  $\mathbb{C}$ -field  $A \equiv A^i_{\alpha} \tau_i dx^{\alpha}$  and real triad  $E^i_{\alpha}$ .
- Scalar, vector, and Gauss constraints:

$$\mathcal{H} = \epsilon_{ijk}E^i \cdot E^j \times \left(B^k + \frac{\Lambda}{3}E^k\right) = 0,$$
  
 $\mathcal{V}_{\alpha} = (E_i \times B^i)_{\alpha} = 0, \mathcal{G}_i = D_{\alpha}E_i^{\alpha} = 0.$ 

• Quantum theory:  $E \rightarrow \hat{E}_i^{\alpha} = -\delta/\delta A_{\alpha}^i$ ,  $\hat{A}_{\alpha}^i$  multiplicative. Constraints annihilated by the Chern–Simons state

$$\Psi_{\rm CS} = \exp\left[rac{i heta}{8\pi^2}\int_{S^3} {
m tr}(A\wedge dA + rac{2}{3}A\wedge A\wedge A)
ight]\,,\qquad heta\equiv rac{6\pi^2}{i\Lambda}\,,$$

▲□▶▲圖▶▲≣▶▲≣▶ ≣ のへで

- Ashtekar variables: connection  $\mathbb{C}$ -field  $A \equiv A^i_{\alpha} \tau_i dx^{\alpha}$  and real triad  $E^i_{\alpha}$ .
- Scalar, vector, and Gauss constraints:

$$\mathcal{H} = \epsilon_{ijk}E^i \cdot E^j \times \left(B^k + \frac{\Lambda}{3}E^k\right) = 0,$$
  
 $\mathcal{V}_{\alpha} = (E_i \times B^i)_{\alpha} = 0, \mathcal{G}_i = D_{\alpha}E_i^{\alpha} = 0.$ 

• Quantum theory:  $E \rightarrow \hat{E}_i^{\alpha} = -\delta/\delta A_{\alpha}^i$ ,  $\hat{A}_{\alpha}^i$  multiplicative. Constraints annihilated by the Chern–Simons state

$$\Psi_{\rm CS} = \exp\left[rac{i heta}{8\pi^2}\int_{S^3}{
m tr}(A\wedge dA+rac{2}{3}A\wedge A\wedge A)
ight]\,,\qquad heta\equivrac{6\pi^2}{i\Lambda}\,,$$

• Different sectors of Euclidean gravity  $(\theta \rightarrow i\theta)$  connected by large gauge transformations.

# The hand: 2. Deformation of $\theta$

▲□▶▲圖▶▲≣▶▲≣▶ ≣ のQ@

$$\theta \to \theta(A),$$

▲□▶ ▲□▶ ▲□▶ ▲□▶ = 三 のへで

 $\theta \to \theta(A),$ 

thus **breaking large-gauge** U(1) **invariance** (analogy with Peccei–Quinn invariance in QCD).

▲□▶ ▲□▶ ▲□▶ ▲□▶ = 三 のへで

$$\theta \to \theta(A),$$

thus **breaking large-gauge** U(1) **invariance** (analogy with Peccei–Quinn invariance in QCD).  $\Lambda$  is promoted to an evolving functional  $\Lambda(A)$ .

$$\theta \to \theta(A),$$

thus **breaking large-gauge** U(1) **invariance** (analogy with Peccei–Quinn invariance in QCD).  $\Lambda$  is promoted to an evolving functional  $\Lambda(A)$ .

◆□▶ ◆□▶ ◆□▶ ◆□▶ ● ● ● ●

No matter introduced by hand!

$$\theta \to \theta(A),$$

thus **breaking large-gauge** U(1) **invariance** (analogy with Peccei–Quinn invariance in QCD).  $\Lambda$  is promoted to an evolving functional  $\Lambda(A)$ .

- No matter introduced by hand!
- The only sectors compatible with this step and the Gauss constraint are **degenerate**: det E = 0, no metric!

(日) (日) (日) (日) (日) (日) (日)

E.o.m. for A can be written as the (1 + 1)-dimensional Dirac equation  $\gamma^0 \dot{\psi} + \gamma^z \partial_z \psi = 0$ 

▲□▶ ▲□▶ ▲ 三▶ ▲ 三▶ - 三 - のへぐ

E.o.m. for *A* can be written as the (1 + 1)-dimensional Dirac equation  $\gamma^0 \dot{\psi} + \gamma^z \partial_z \psi = 0$ , where

$$\psi \equiv \begin{pmatrix} iA_1^1 \\ A_2^1 \\ A_1^2 \\ A_1^2 \\ iA_2^2 \end{pmatrix}$$

٠

▲□▶ ▲□▶ ▲ 三▶ ▲ 三▶ - 三 - のへぐ

E.o.m. for *A* can be written as the (1 + 1)-dimensional Dirac equation  $\gamma^0 \dot{\psi} + \gamma^z \partial_z \psi = 0$ , where



E.o.m. for *A* can be written as the (1 + 1)-dimensional Dirac equation  $\gamma^0 \dot{\psi} + \gamma^z \partial_z \psi = 0$ , where



A model for *V* interactions and physical interpretation naturally emerge at quantum level.

# The hand: 4. Suppression of $\Lambda$

 $\gamma^0 \dot{\psi} + \gamma^z \partial_z \psi + im\psi = 0.$ 

▲□▶ ▲□▶ ▲ 三▶ ▲ 三▶ - 三 - のへぐ

$$\gamma^0 \dot{\psi} + \gamma^z \partial_z \psi + im\psi = 0.$$

▲□▶ ▲□▶ ▲ 三▶ ▲ 三▶ - 三 - のへぐ

• Mass term  $m = -2i\bar{\psi}\gamma^5\partial_z\psi$ .

$$\gamma^0 \dot{\psi} + \gamma^z \partial_z \psi + im\psi = 0.$$

- Mass term  $m = -2i\bar{\psi}\gamma^5\partial_z\psi$ .
- The simplest nonperturbative solution requires

 $\Lambda = \Lambda_0 \exp(-\bar{\psi}\gamma^5 \gamma^z \psi)$ 

▲□▶ ▲□▶ ▲三▶ ▲三▶ - 三 - のへで

$$\gamma^0 \dot{\psi} + \gamma^z \partial_z \psi + im\psi = 0.$$

- Mass term  $m = -2i\bar{\psi}\gamma^5\partial_z\psi$ .
- The simplest nonperturbative solution requires

$$\Lambda = \Lambda_0 \exp(-\bar{\psi}\gamma^5 \gamma^z \psi)$$

◆□▶ ◆□▶ ◆□▶ ◆□▶ ● ● ● ●

•  $j^{5\alpha}$  associated with a chiral transformation of the fermion  $\psi$  and not conserved in the presence of *m*.

$$\gamma^0 \dot{\psi} + \gamma^z \partial_z \psi + im\psi = 0.$$

- Mass term  $m = -2i\bar{\psi}\gamma^5\partial_z\psi$ .
- The simplest nonperturbative solution requires

$$\Lambda = \Lambda_0 \exp(-\bar{\psi}\gamma^5 \gamma^z \psi)$$

(日) (日) (日) (日) (日) (日) (日)

•  $j^{5\alpha}$  associated with a chiral transformation of the fermion  $\psi$  and not conserved in the presence of *m*. *P* symmetry is broken.

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 のへぐ

• Perturbative regime



• Perturbative regime (small values of the connection)

• Perturbative regime (small values of the connection):  $|\langle j^{5z} \rangle| \ll 1, \langle \Lambda \rangle \approx \Lambda_0 (1 - \langle j^{5z} \rangle) = O(1).$ 

< □ > < 同 > < Ξ > < Ξ > < Ξ > < Ξ < </p>

• Perturbative regime (small values of the connection):  $|\langle j^{5z} \rangle| \ll 1, \langle \Lambda \rangle \approx \Lambda_0 (1 - \langle j^{5z} \rangle) = O(1).$ 

(ロ) (同) (三) (三) (三) (○) (○)

Nonperturbative regime

• Perturbative regime (small values of the connection):  $|\langle j^{5z} \rangle| \ll 1$ ,  $\langle \Lambda \rangle \approx \Lambda_0 (1 - \langle j^{5z} \rangle) = O(1)$ .

(日) (日) (日) (日) (日) (日) (日)

Nonperturbative regime (large connection values)

- Perturbative regime (small values of the connection):  $|\langle j^{5z} \rangle| \ll 1$ ,  $\langle \Lambda \rangle \approx \Lambda_0 (1 \langle j^{5z} \rangle) = O(1)$ .
- Nonperturbative regime (large connection values): Condensate with v.e.v.  $\langle j^{5z} \rangle \sim O(10^2)$ .

(日) (日) (日) (日) (日) (日) (日)

- Perturbative regime (small values of the connection):  $|\langle j^{5z} \rangle| \ll 1, \langle \Lambda \rangle \approx \Lambda_0 (1 - \langle j^{5z} \rangle) = O(1).$
- Nonperturbative regime (large connection values): Condensate with v.e.v.  $\langle j^{5z} \rangle \sim O(10^2)$ .
- Smallness of  $\Lambda$  regarded as a large-scale nonperturbative quantum mechanism similar to quark confinement.

(日) (日) (日) (日) (日) (日) (日)

- Perturbative regime (small values of the connection):  $|\langle j^{5z} \rangle| \ll 1, \langle \Lambda \rangle \approx \Lambda_0 (1 - \langle j^{5z} \rangle) = O(1).$
- Nonperturbative regime (large connection values): Condensate with v.e.v.  $\langle j^{5z} \rangle \sim O(10^2)$ .
- Smallness of  $\Lambda$  regarded as a large-scale nonperturbative quantum mechanism similar to quark confinement.
- Quantizing  $\psi$  as a Majorana fermion,

$$\mathcal{H} \propto \mathcal{H}_{ extbf{BCS}} = \sum_{k,\sigma} \mathcal{E}_k c^{\dagger}_{k\sigma} c_{k\sigma} - \sum_{k,k'} oldsymbol{V}_{kk'} c^{\dagger}_{k+} c^{\dagger}_{k-} c_{k'-} c_{k'+}.$$

 Correspondence made rigorous using a deformed CFT (WZW model at critical level).