

Effective Field Theory in LUX Run04

Shaun Alsum

1

Yes, I re-used the fall theme...
It's a good theme...

What Is EFT?



Wick Haxton
Wednesday

Effective Theory Treatments of
Dark Matter Direct Detection

- Basic Scales
- Nonrelativistic Effective Theory Description
- Implications for experiment: number and type

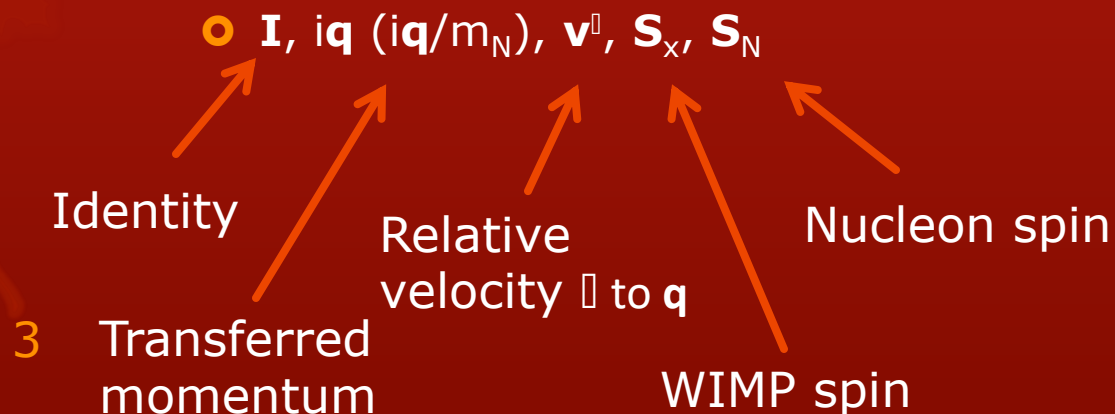
Wick Haxton October 25, 2017

The slide features a background image of a city at night with a bridge. At the bottom, there are three logos: CU Boulder, Office of Science (U.S. Department of Energy), and University of Rochester.

http://teacher.pas.rochester.edu:8080/wiki/pub/Lux/AWGworkshop171023_Agenda/DM_LUX.pdf

For those without time machine access...

- Standard Spin Independent and Spin Dependent neglect other possible interactions
- In general, we want to account for any possible interaction that is allowed.
- There are 5 galilean and hermitian invariant quantities from which these operators can be built.



Still livin' in the past

- These invariant quantities can be combined into 16 operators up to 2nd order in momentum transfer and spin-exchange of 1 or less
- These are:

$\mathcal{O}_1 = 1_X 1_N$	$\mathcal{O}_{12} = \vec{S}_X \cdot (\vec{S}_N \times \vec{v}^\perp)$
$\mathcal{O}_2 = (v^\perp)^2$	$\mathcal{O}_{13} = i(\vec{S}_X \cdot \vec{v}^\perp)(\vec{S}_N \cdot \frac{\vec{q}}{m_N})$
$\mathcal{O}_3 = i\vec{S}_N \cdot (\frac{\vec{q}}{m_N} \times \vec{v}^\perp)$	$\mathcal{O}_{14} = i(\vec{S}_X \cdot \frac{\vec{q}}{m_N})(\vec{S}_N \cdot \vec{v}^\perp)$
$\mathcal{O}_4 = \vec{S}_X \cdot \vec{S}_N$	$\mathcal{O}_{15} = -(\vec{S}_X \cdot \frac{\vec{q}}{m_N})((\vec{S}_N \times \vec{v}^\perp) \cdot \frac{\vec{q}}{m_N})$
$\mathcal{O}_5 = i\vec{S}_X \cdot (\frac{\vec{q}}{m_N} \times \vec{v}^\perp)$	$\mathcal{O}_{16} = -((\vec{S}_X \times \vec{v}^\perp) \cdot \frac{\vec{q}}{m_N})(\vec{S}_N \cdot \frac{\vec{q}}{m_N})$
$\mathcal{O}_6 = (\vec{S}_X \cdot \frac{\vec{q}}{m_N})(\vec{S}_N \cdot \frac{\vec{q}}{m_N})$	
$\mathcal{O}_7 = \vec{S}_N \cdot \vec{v}^\perp$	
$\mathcal{O}_8 = \vec{S}_X \cdot \vec{v}^\perp$	
$\mathcal{O}_9 = i\vec{S}_X \cdot (\vec{S}_N \times \frac{\vec{q}}{m_N})$	
$\mathcal{O}_{10} = i\vec{S}_N \cdot \frac{\vec{q}}{m_N}$	
$\mathcal{O}_{11} = i\vec{S}_X \cdot \frac{\vec{q}}{m_N}$	

More recap...

- These operators can be expressed as combinations of 6 nuclear responses:

$$\begin{aligned} M_{JM}(q\vec{x}) & \\ \Delta_{JM}(q\vec{x}) & \equiv \vec{M}_{JJ}^M(q\vec{x}) \cdot \frac{1}{q} \vec{\nabla} \\ \Sigma'_{JM}(q\vec{x}) & \equiv -i \left\{ \frac{1}{q} \vec{\nabla} \times \vec{M}_{JJ}^M(q\vec{x}) \right\} \cdot \vec{\sigma} = [J]^{-1} \left\{ -\sqrt{J} \vec{M}_{JJ+1}^M(q\vec{x}) + \sqrt{J+1} \vec{M}_{JJ-1}^M(q\vec{x}) \right\} \cdot \vec{\sigma} \\ \Sigma''_{JM}(q\vec{x}) & \equiv \left\{ \frac{1}{q} \vec{\nabla} M_{JM}(q\vec{x}) \right\} \cdot \vec{\sigma} = [J]^{-1} \left\{ \sqrt{J+1} \vec{M}_{JJ+1}^M(q\vec{x}) + \sqrt{J} \vec{M}_{JJ-1}^M(q\vec{x}) \right\} \cdot \vec{\sigma} \\ \tilde{\Phi}'_{JM}(q\vec{x}) & \equiv \left(\frac{1}{q} \vec{\nabla} \times \vec{M}_{JJ}^M(q\vec{x}) \right) \cdot \left(\vec{\sigma} \times \frac{1}{q} \vec{\nabla} \right) + \frac{1}{2} \vec{M}_{JJ}^M(q\vec{x}) \cdot \vec{\sigma} \\ \Phi''_{JM}(q\vec{x}) & \equiv i \left(\frac{1}{q} \vec{\nabla} M_{JM}(q\vec{x}) \right) \cdot \left(\vec{\sigma} \times \frac{1}{q} \vec{\nabla} \right) \end{aligned}$$

These are very complicated... don't ask me to explain exactly what each is ☺

- 5
- These nuclear responses can be evaluated for a given isotope.

What does it all matter?

- WIMP recoil spectrum is determined by:

$$\frac{dR}{dE_R} = \left\langle \frac{\rho_\chi m_T}{\mu_T^2 m_\chi v} \frac{d\sigma}{d\cos\theta} \right\rangle$$

- The lagrangian includes all possible interactions:

$$\mathcal{L} = \sum_{i=1}^{12} c_i^{(n)} \mathcal{O}_i^{(n)} + c_i^{(p)} \mathcal{O}_i^{(p)}$$

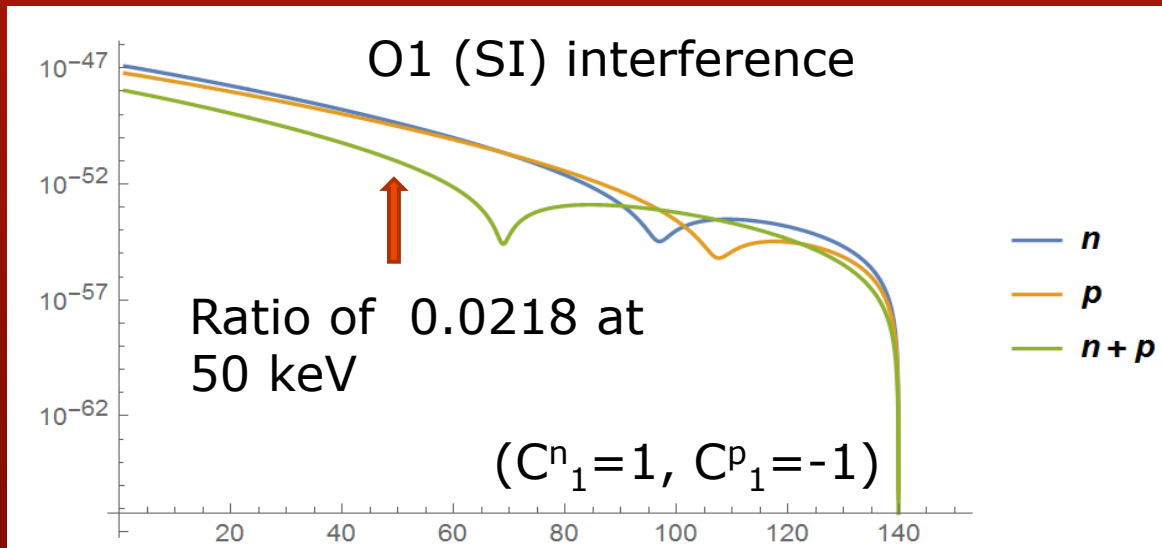
- The cross section is proportional to the matrix element squared

$$\frac{dR_D}{dE_R} = N_T \frac{\rho_\chi m_T}{32\pi m_\chi^3 m_N^2} \left\langle \frac{1}{v} \sum_{ij} \sum_{N,N'=p,n} c_i^{(N)} c_j^{(N')} F_{ij}^{(N,N')}(v^2, q^2) \right\rangle$$

- The F_{ij} is a form factor corresponding to operators \mathcal{O}_i and \mathcal{O}_j , which is a known combination of the known nuclear responses referred to on the previous slide.

That's a lot of factors! What are we to do?

- Interference can conspire to nullify almost any limits for a single parameter
 - Only certain combinations of operators can interfere because they must satisfy symmetries (parity, etc)
 - The proton and neutron couplings can interfere for any operator, however.



Still not fully sure...

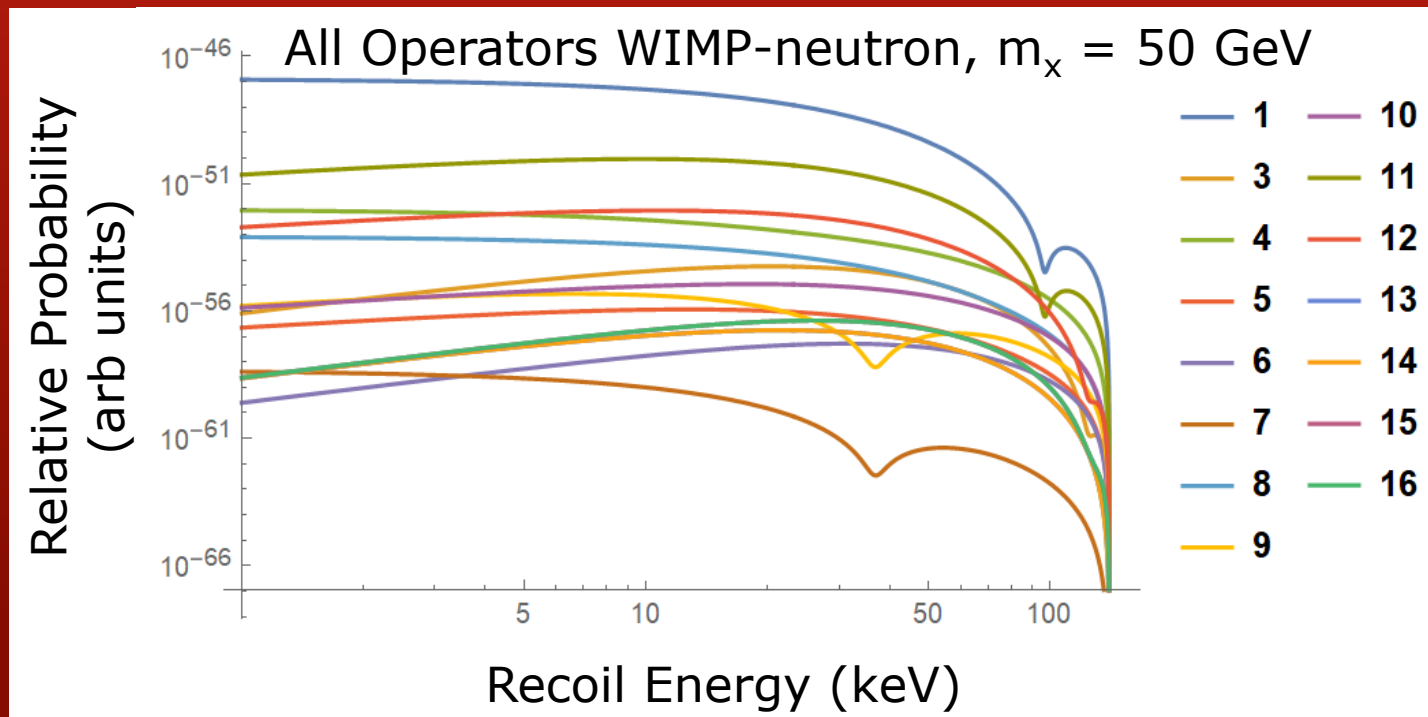
- To get a proper limit, would want to re-write

$\sum_i c_i^n O_i + c_i^p O_i^p$ as $O_a^n + \frac{c_a^p}{c_a^n} O_a^p + \sum_{i \neq a} \frac{c_i^n}{c_a^n} O_i^n + \frac{c_i^p}{c_a^n} O_i^p$ for an arbitrary choice of c_a^n for each interference class. Then minimize the above to get the worst-case value for each $\frac{c_i^{(N)}}{c_a^n}$. Then use the PLR to determine a limit on c_a^n which would determine a worst-case limit on each other parameter.

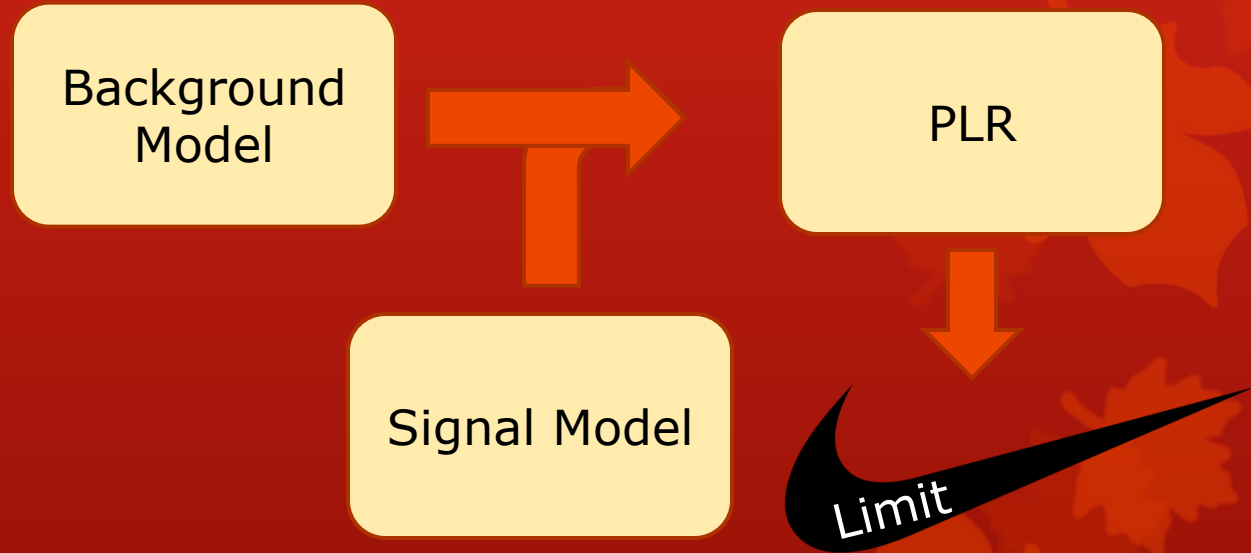
- This is hard, so for now we just work on 1 parameter at a time while setting all others to 0

It's all about the recoil spectrum

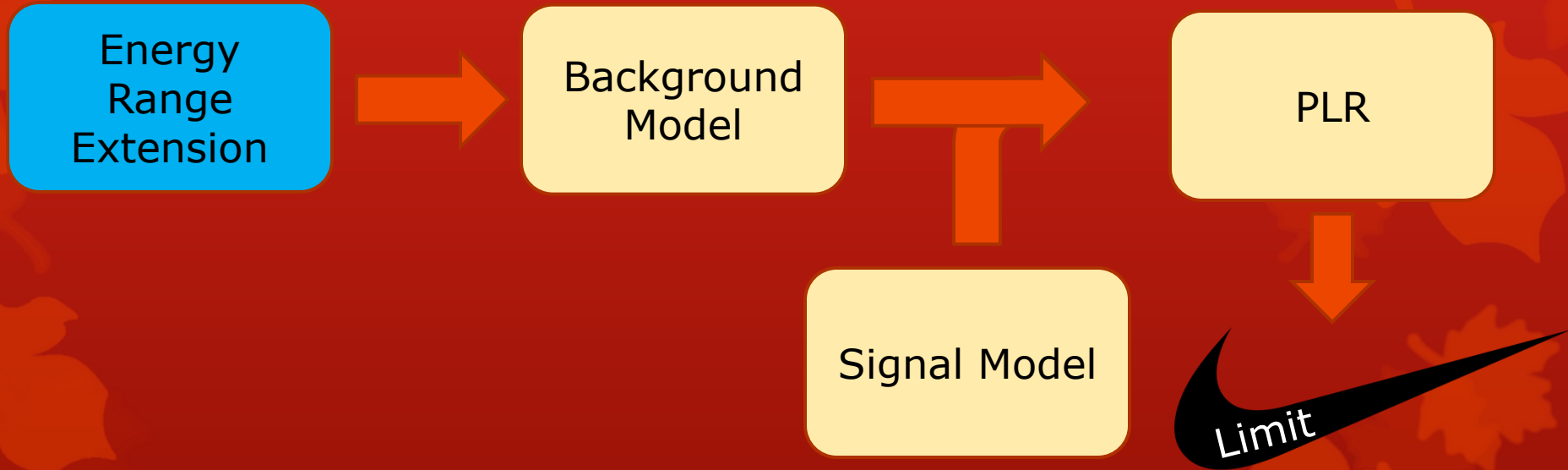
- Different combinations, then, have different recoil spectra as expected.
- Here are all of them for a 50 GeV WIMP coupling only to neutrons



What makes a limit?



Changes to the Background?



Background Model Modification

- Initial analysis only extended to 50 keV_{nr}.
- EFT signals can possibly extend much further than this for some heavier WIMP candidates.
- ^{83m}Kr , which is not a huge problem for traditional energy ranges, may be an issue.

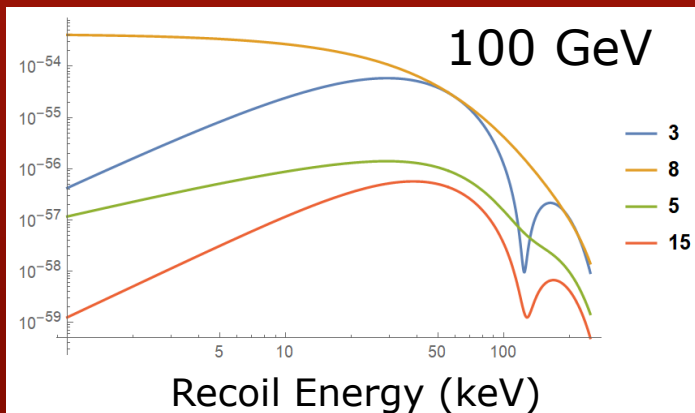
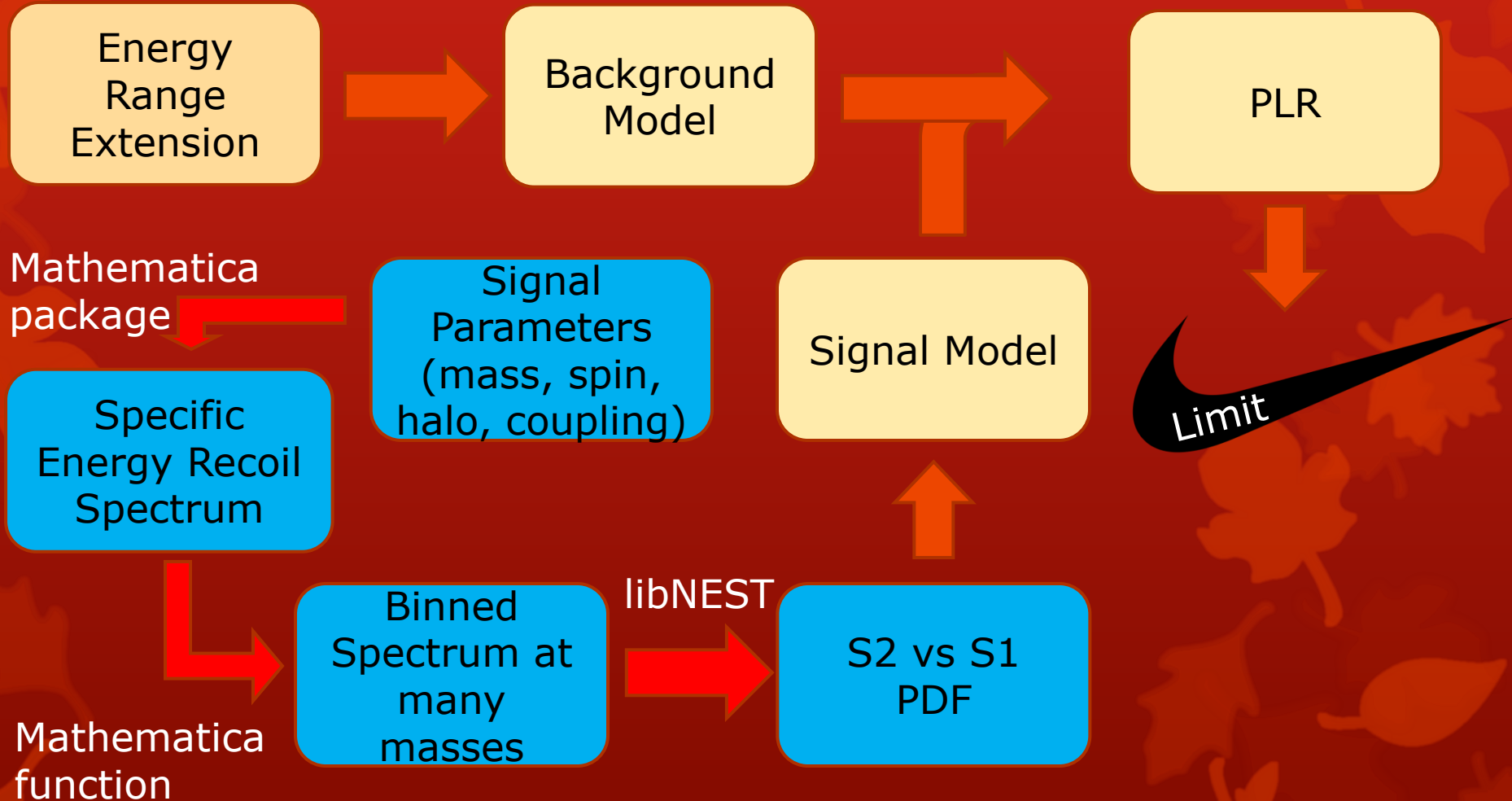


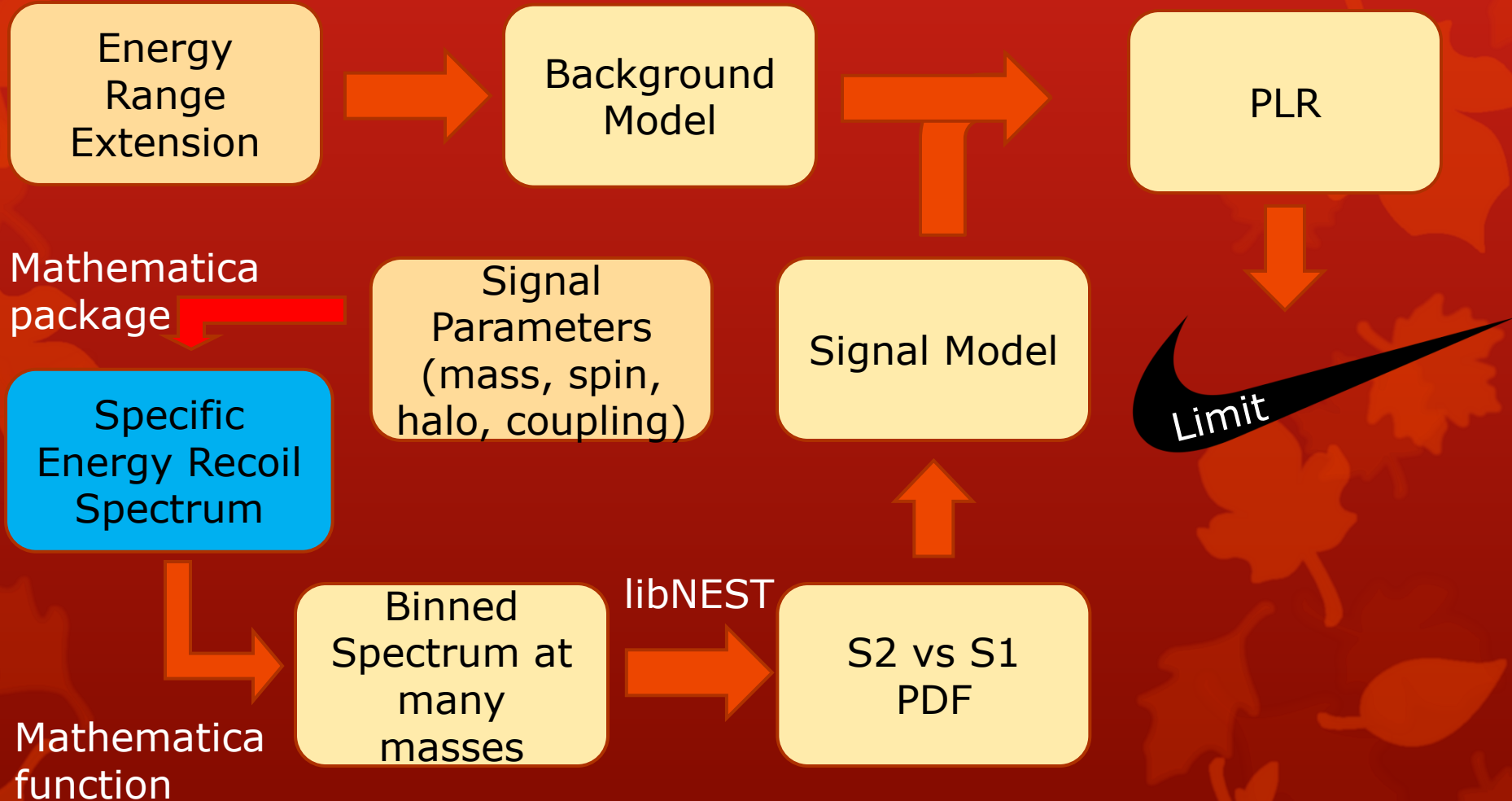
TABLE I. The upper energy threshold E_{max} (in keV_{nr}) for each of the effective field theory operators, such that an energy window from 0 to E_{max} captures either 50% or 90% of WIMP-neutron recoil events for the given operator and WIMP mass.

Operator	50-GeV WIMP		500-GeV WIMP	
	$E_{max}^{50\%}$ (keV _{nr})	$E_{max}^{90\%}$ (keV _{nr})	$E_{max}^{50\%}$ (keV _{nr})	$E_{max}^{90\%}$ (keV _{nr})
SI	10.8	27.3	16.6	44.7
\mathcal{O}_1	6.8	21.7	11.8	43.8
\mathcal{O}_3	26.4	49.1	148.1	344.4
SD	8.6	21.6	11.9	37.5
\mathcal{O}_4	7.0	24.0	32.8	299.6
\mathcal{O}_5	16.2	38.6	65.5	328.9
\mathcal{O}_6	33.6	64.0	267.3	433.7
\mathcal{O}_7	5.0	16.2	25.2	279.9
\mathcal{O}_8	6.8	22.2	14.5	64.8
\mathcal{O}_9	13.7	37.2	276.7	464.7
\mathcal{O}_{10}	21.7	48.6	112.6	340.4
\mathcal{O}_{11}	15.5	34.4	39.0	279.9
\mathcal{O}_{12}	17.4	38.1	34.8	176.5
\mathcal{O}_{13}	28.2	53.2	54.5	219.7
\mathcal{O}_{14}	11.9	27.9	240.9	400.0
\mathcal{O}_{15}	34.3	59.1	261.2	433.7

Changes to the Background?



Changes to the Background?



DMFormFactor - Mathematica

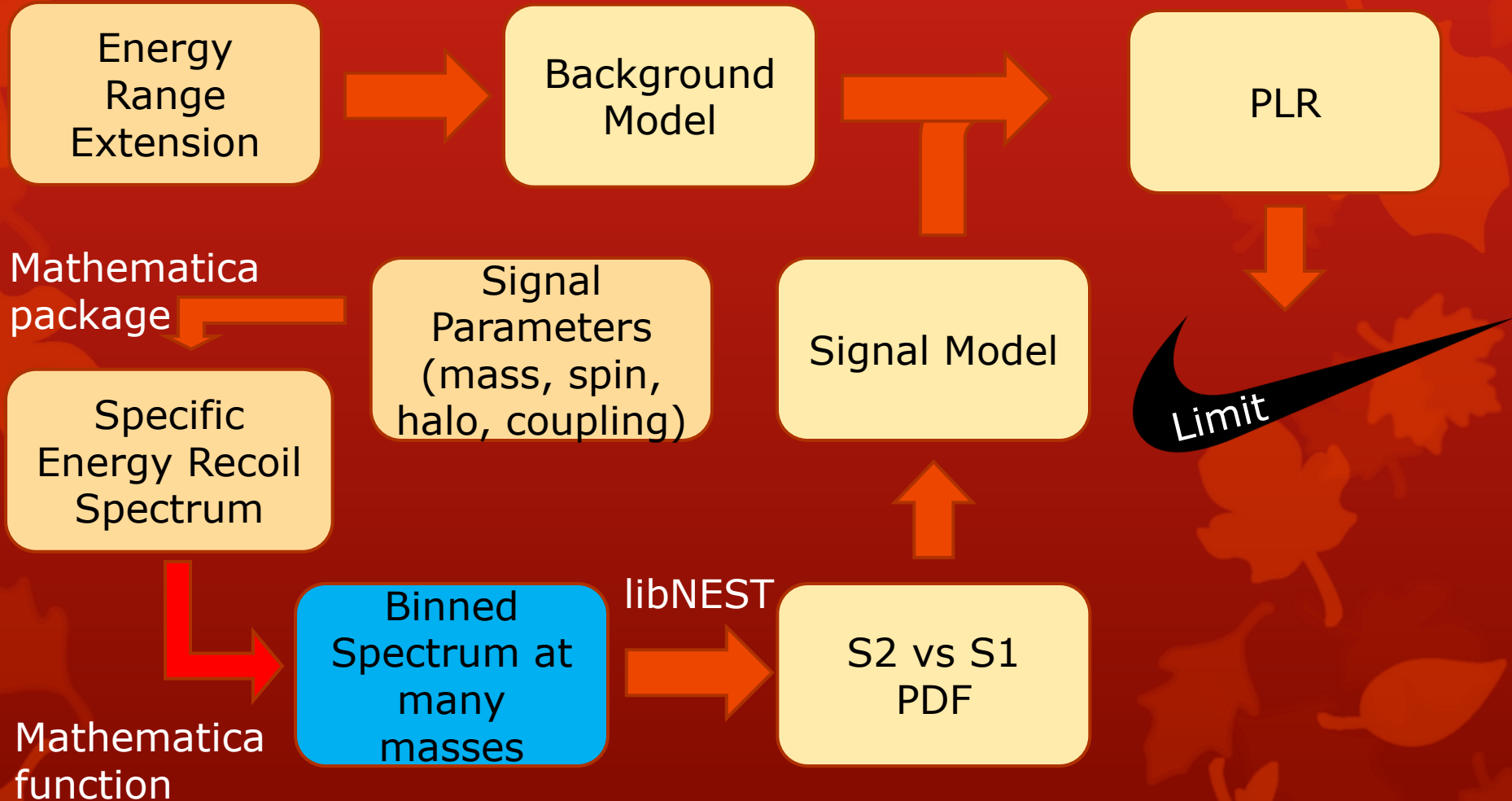
Welcome to

DMFormFactor version 1.1.

Functions are SetCoeffsNonrel, SetCoeffsRel, SetCoeffsNucl, ZeroCoeffs, SetJChi, SetMchi, SetIsotope, SetHALO, SetHelm, TransitionProbability, ResponseNucl, DiffCrossSection, ApproxTotalCrossSection, and EventRate.

- Mathematica package written by Nikhil Anand, A. Liam Fitzpatrick, and W. C. Haxton.
- Returns the recoil energy (really momentum) spectrum in the form of a mathematica function given model inputs.
- Notable inputs:
 - Halo type, earth velocity, local DM mean velocity, escape vel
 - (set to typical values)
 - WIMP spin
 - Set to 1/2
 - WIMP mass
 - Target isotope
 - Operator coefficients (for each op and proton vs neutron coupling)

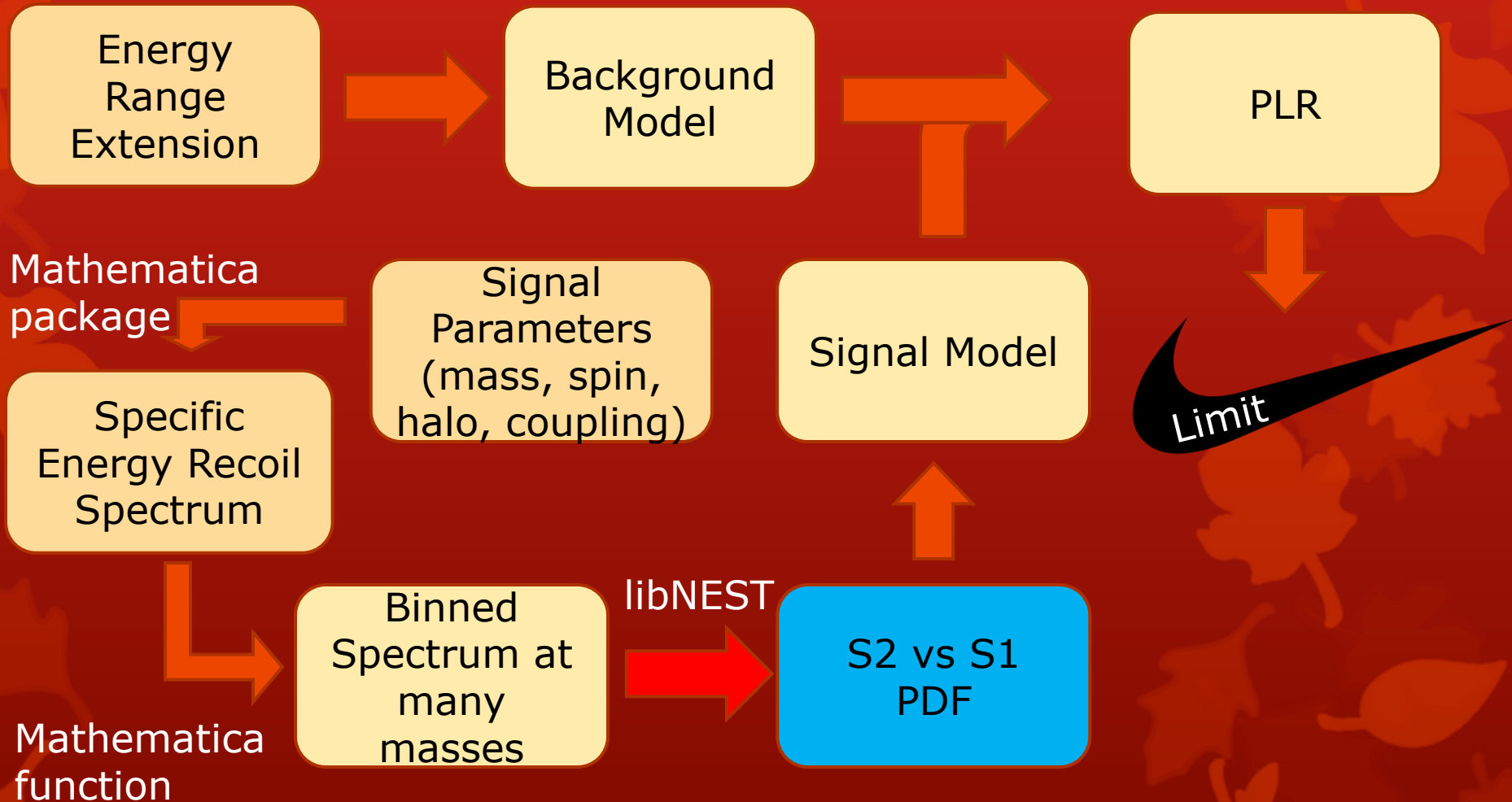
Changes to the Background?



Spectra generation

- Mathematica function calls DMFormFactor for different values of mass and different isotopes in an array.
- Resulting analytical spectra are integrated into bins used by the PLR analysis.
- Isotope Spectra are weighted by their abundance and then added together.
- The resulting array of varying masses and bins is stored in a text file whose name identifies the non-zero coupling constant used.

Changes to the Background?

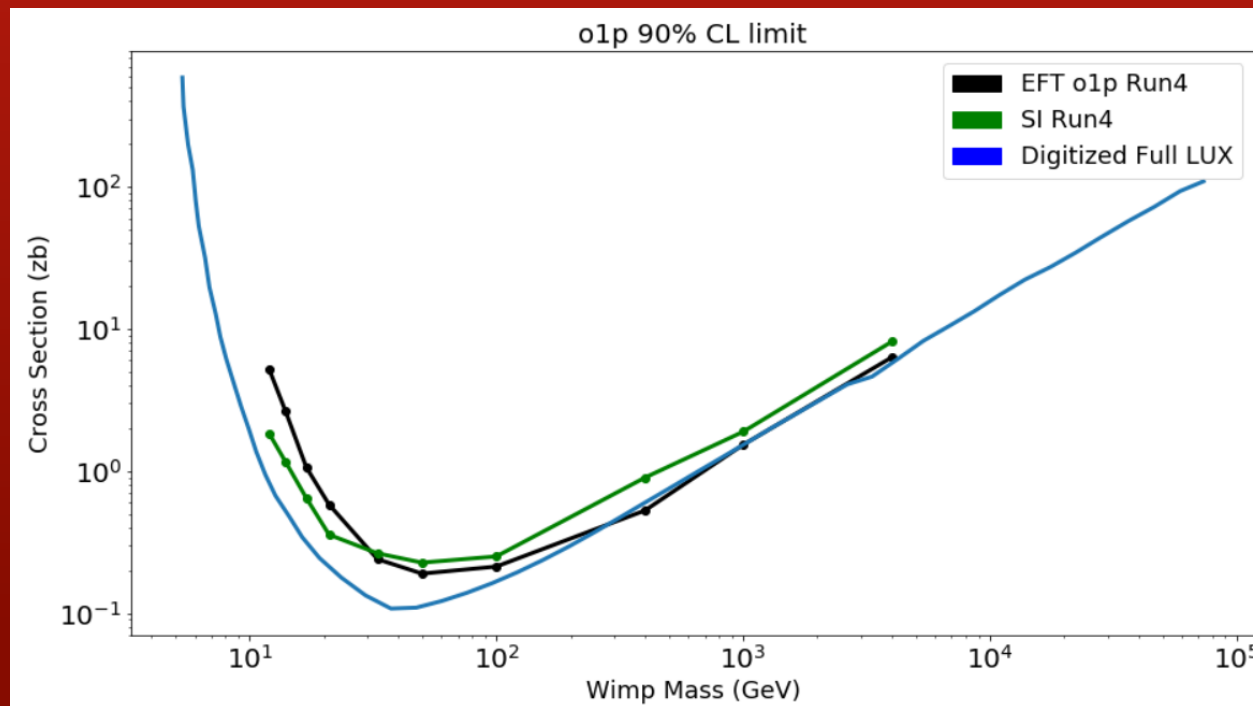


PDF creation

- libNEST is called for each mass in question.
- libNEST imports the recoil spectrum then simulates the S2 and S1 response given energies sampled from the imported spectrum.
- The result is stored as a root histogram to be used by the limit code.

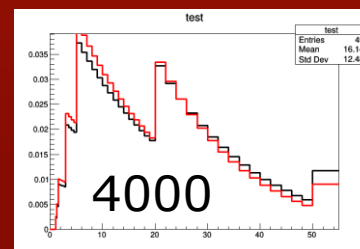
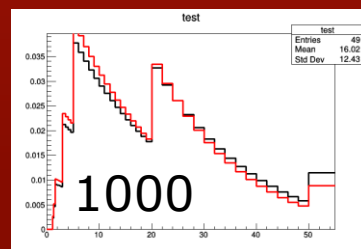
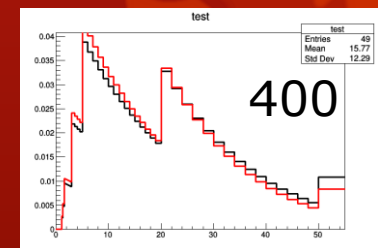
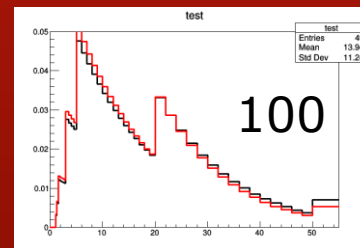
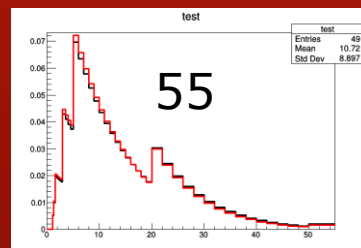
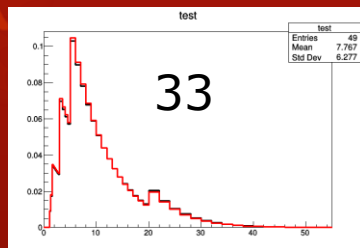
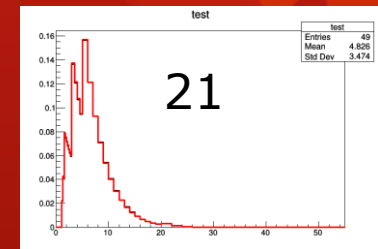
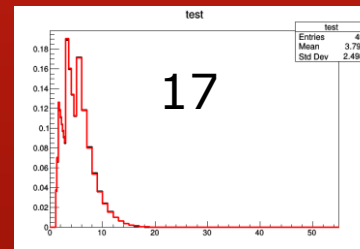
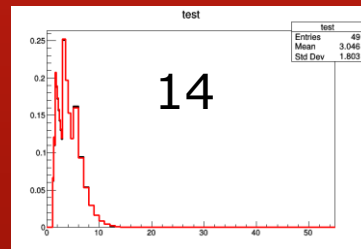
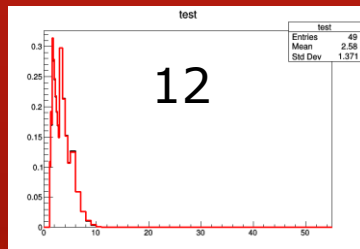
How do we do?

- Operator 1 should closely match the SI result.
- We attempted to validate our process by comparing the limit on the WIMP-proton coupled operator 1 to our SI result.



Not great, but why?

- Low statistics (1000 per nSig, mass point)?
- Model mismatch?



Conclusion

- Getting there...
- Process is in place
- Background model must be validated and incorporated
- Difference between WIMP proton operator 1 limit and SI results must be understood
- After single-parameter limits, other interesting things could be done.
 - Proton-neutron interference looked into?
 - Joint PLR with experiments with other detector mediums?