

FerMINI - Fermilab Search for Millicharged Particle & Strongly Interacting Dark Matter

Yu-Dai Tsai, Fermilab/U.Chicago (WH674)

with Magill, Plestid, Pospelov (<u>1806.03310</u>, *PRL '19*),

with Kelly (<u>1812.03998</u>, **PRD '19**)

Email: <u>ytsai@fnal.gov</u>; arXiv: <u>https://arxiv.org/a/tsai_y_1.html</u>

FerMINI Collaboration



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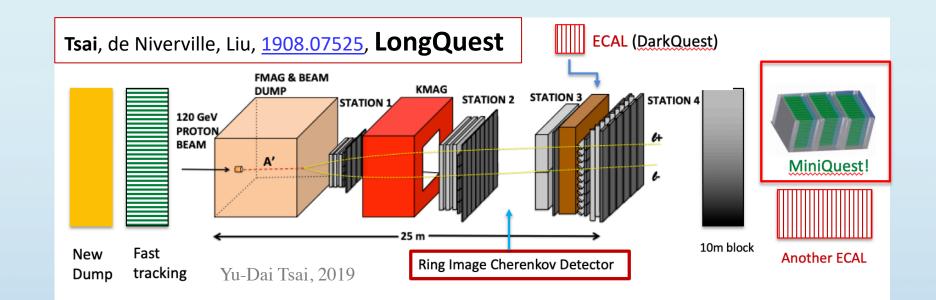


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Joe Bramante Queen's U



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Long-Lived Particles in Proton Fixed-Target Experiments

- Light Scalar & Dark Photon at BoreXino & LSND, <u>1706.00424</u> (proton-charge radius anomaly)
- Dipole Portal Heavy Neutral Lepton, <u>1803.03262</u> (LSND/MiniBooNE anomalies)
- Dark Neutrino at Scattering Exp: CHARM-II & MINERvA! <u>1812.08768</u> (MiniBooNE Anomaly)
- Closing dark photon, inelastic dark matter, and muon g-2 windows; & the LongQuest Proposal! <u>1908.07525</u> (muon g-2 Anomaly)

Outline

- Motivations & Intro to Millicharged Particle (MCP)
- The FerMINI Experiment:
 Proton Fixed-Target Scintillation Experiment to
 Search for Minicharged Particles
- Link to Strongly Interacting Dark Matter

Millicharged Particles

Is electric charge quantized? Other Implications

Yu-Dai Tsai, Fermilab, 2019

Finding Minicharge

- Is electric charge quantized and why? A long-standing question!
- SM U(1) allows arbitrarily small (any real number) charges. Why don't we see them? Motivates
 Dirac quantization, Grand Unified Theory (GUT), etc, to explain such quantization (anomaly cancellations fix some SM U(1)_Y charge assignments)
- Testing if **e/3 is the minimal charge**
- MCP could have natural link to **dark sector** (dark photon, etc)
- Could account for **dark matter (DM) abundance**
- Used for the cooling of gas temperature to explain the EDGES anomaly [EDGES collab., Nature, (2018); Barkana, Nature, (2018)].
 A small fraction of the DM as MCP can potentially explain EDGES observation

Millicharged Particle: Models

Yu-Dai Tsai, Fermilab, 2019

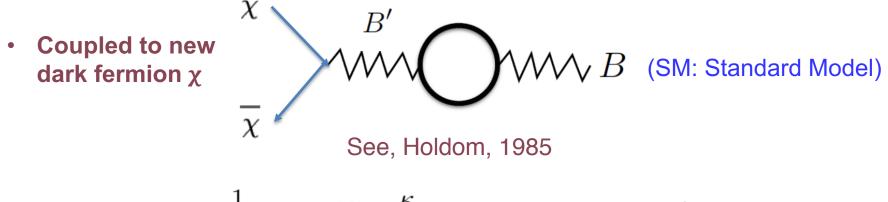
MCP Model

• A particle fractionally charged under a U(1) hypercharge

$$\mathcal{L}_{\rm MCP} = i\bar{\chi}(\partial \!\!\!/ - i\epsilon' e B \!\!\!/ + M_{\rm MCP})\chi$$

- Can just consider these Lagrangian terms by themselves (no extra mediator, i.e., dark photon), one can call this a "pure" MCP
- Or this could be from **Kinetic Mixing**
 - give a nice origin to the above term
 - an example that gives rise to dark sectors
 - easily compatible with Grand Unification Theory
 - I will not spend too much time on the model

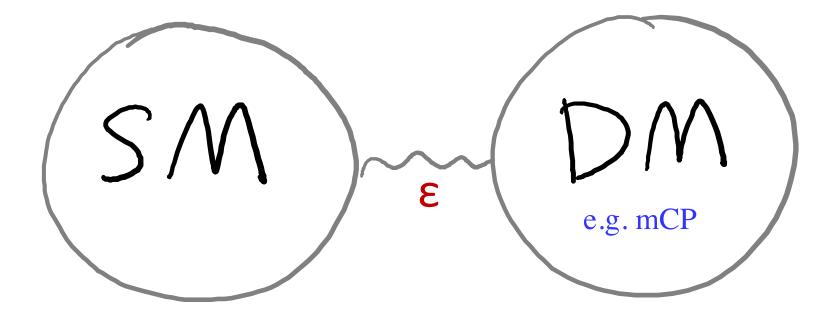
Kinetic Mixing and MCP Phase



$$\mathcal{L} = \mathcal{L}_{\rm SM} - \frac{1}{4} B'_{\mu\nu} B'^{\mu\nu} - \frac{\kappa}{2} B'_{\mu\nu} B^{\mu\nu} + i\bar{\chi}(\partial \!\!\!/ + ie' B' + iM_{\rm MCP})\chi$$

- New fermion χ charged under new gauge boson B'.
- Millicharged particle (MCP) can be a low-energy consequence of massless dark photon (a new U(1) gauge boson) coupled to a new fermion (become MCP in a convenient basis.)
- See Holdom, 1985; or <u>arXiv:1806.03310</u>

The Rise of Dark Sector



Yu-Dai Tsai, Fermilab, 2019

Important Notes!

- Our search is simply a search for particles (fermion χ) with {mass, electric charge} = { $m_{\chi}, \epsilon e$ }
- Minimal theoretical inputs/parameters

(harder to probe in MeV – GeV+ mass regime)

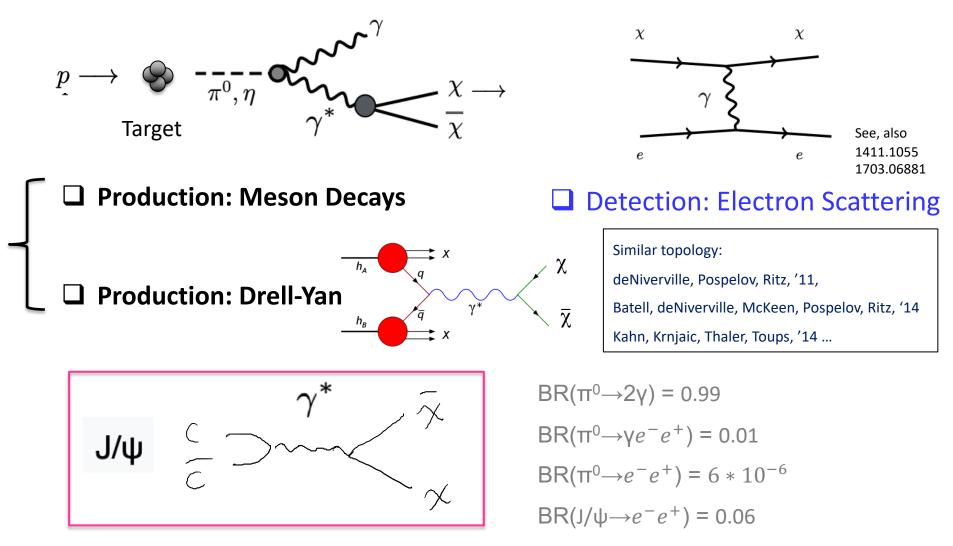
- MCPs do not have to be DM in our searches
- The bounds we derive still put constraints on DM (SIDM) as well as dark sector scenarios.
- Not considering bounds on dark photon (not necessary for MCP particles)
- Similar bound/sensitivity applies to scalar MCPs

Millicharged Particle: Signature

Yu-Dai Tsai, Fermilab, 2019

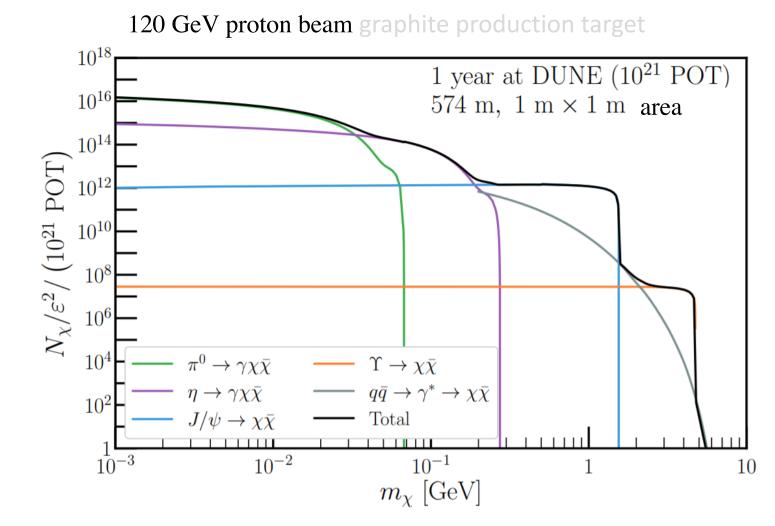
Production & Detection:

MCP (or light DM with massless mediator):



Heavy mesons are important for high-mass mCP's in high-energy beams

MCP Production/Flux

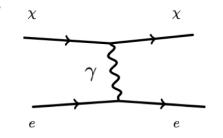


MCP Detection: Electron Scattering & Ionization

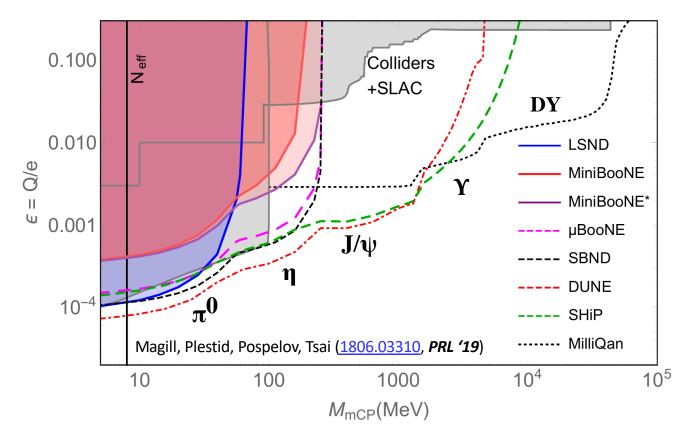
- MCP scattering with electron prefers **low-momentum transfer**, Q^2
- Expressed in **recoil energy threshold**, $E_e^{(min)}$, we have

$$\sigma_{e\chi} \simeq 2.6 \times 10^{-25} \text{cm}^2 \times \epsilon^2 \times \frac{1 \text{ MeV}}{E_e^{(\text{min})} - m_e}$$

- Sensitivity greatly enhanced by accurately measuring low energy electron recoils for mCP's & electron scattering w/ light mediator
- See Magill, Plestid, Pospelov, YT, <u>1806.03310</u> (MCP in neutrino Experiments) & deNiverville, Frugiuele, <u>1807.06501</u> (for sub-GeV DM)
- Very low-energy scattering: Ionization (eV-level)!



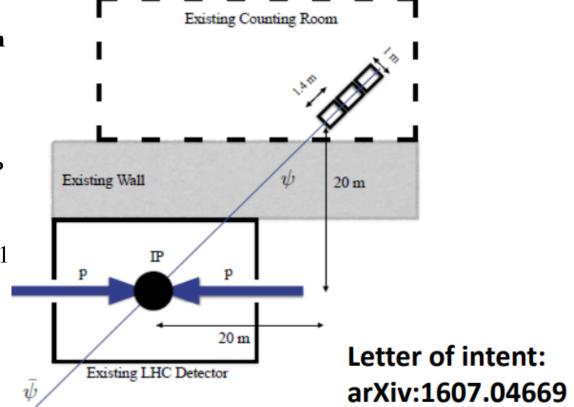
Sensitivity at Neutrino Detectors



- Electron recoil-energy threshold: MeV to 100 MeV
- Can use timing information to improve sensitivity (see ANGELICO's talk)
- Harnik, Liu, Palamara: double-hit to reduce background + Ivan Lepetic (ArgoNeuT+DUNE) '19

MilliQan @ LHC: General Idea

- Require triple coincidence in small time window
 (15 nanoseconds)
- Q down to 10⁻³ e, each MCP produce averagely ~ 1 photoelectron (PE) observed per ~ 1 meter long scintillator
- Long axis points at the CMS Interaction Point (P5).



Andrew Haas, Fermilab (2017)

Andy Haas, Christopher S. Hill, Eder Izaguirre, Itay Yavin, 1410.6816, PRD '15

FerMINI:

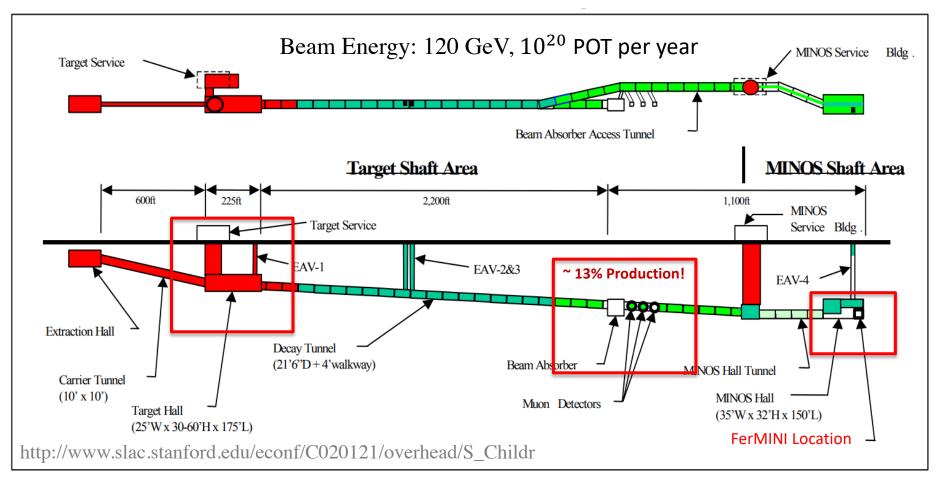
A Fermilab Search for **MINIcharged Particle** Kelly, **Tsai**, arXiv:1812.03998 (PRD`19)

visually "a detector made of stacks of light sabers,"

can also potentially probe new physics scenarios like **small-electric-dipole dark fermions**, or **quirks**, etc

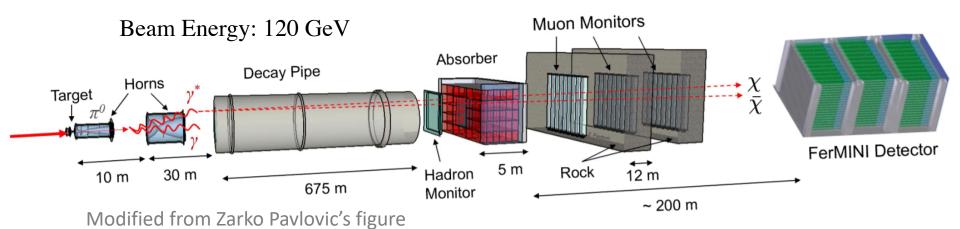
Yu-Dai Tsai, Fermilab, 2019

Site 1: NuMI Beam & MINOS ND Hall



NuMI: Neutrinos at the Main Injector MINOS: Main Injector Neutrino Oscillation Search, ND: Near Detector See YONEHARA's talk

FerMINI @ NuMI-MINOS Hall



An illustration of the FerMINI experiments utilizing the NuMI facility.



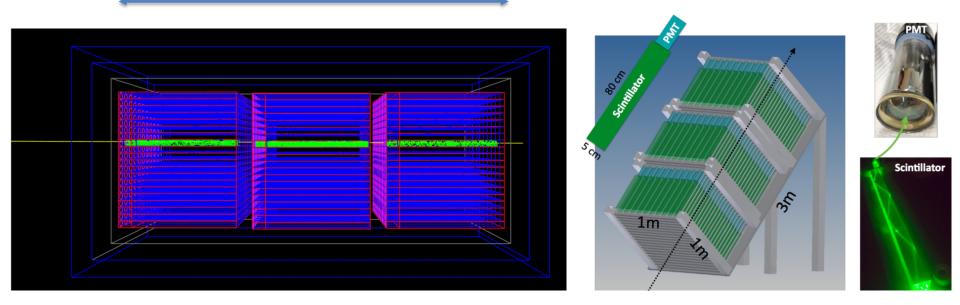
Yu-Dai Tsai Fermilab

MINOS hall downstream of NuMI beam

MCP Production/Flux 120 GeV proton beam hitting graphite production target 10^{18} 1 year at NuMI (6×10^{20} POT) 10^{16} 1040 m, $1 \text{ m} \times 1 \text{ m}$ det. area 10^{14} $N_{\chi}/\varepsilon^2/(6 \times 10^{20} \text{ POT})$ 10^{12} 10^{10} 10^{8} 10^{6} 10^{4} $\pi^0 \to \gamma \chi \bar{\chi}$ $\Upsilon \to \chi \bar{\chi}$ $q\bar{q} \to \gamma^* \to \chi\bar{\chi}$ $\eta \to \gamma \chi \bar{\chi}$ 10^{2} $J/\psi \to \chi \bar{\chi}$ Total 10^{-2} 10^{-1} 10 m_{χ} [GeV]

Detector Concept

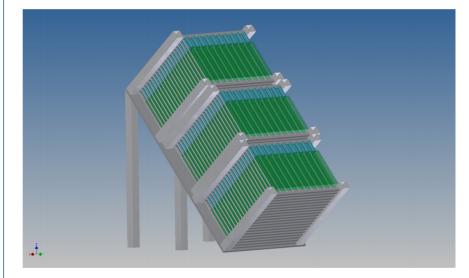
$$(\Delta t)_{\text{offline}} = 15$$



See arXiv:1607.04669; arXiv:1810.06733

Detector: Details of the Nominal Design

- Total: 1 m × 1 m (transverse plane) × 3 m (longitudinal) plastic scintillator array.
- 3 sections each containing 400 5 cm × 5 cm
 × 80 cm scintillator bars optically coupled
 to high-gain photomultiplier (PMT).
- Signature: triple-coincidence within a 15
 ns time window along longitudinally
 contiguous bars in each of the 3 sections as
- Major Background: dark-current noise reduced by requiring triple coincidence



- Scintillator: Saint-Gobain BC-408 plastic scintillator
- PMT: Hamamatsu R329-02 PMT

Photoelectrons (PE) from Scintillation

• The averaged number of photoelectron (PE) seen by the

detector from single MCP is:

$$N_{PE} \propto \left\langle -\frac{dE}{dx} \right\rangle \times l_{scint}, \ \left\langle -\frac{dE}{dx} \right\rangle \propto \epsilon^2.$$

 $\langle dE/dx\rangle$ is the "mass stopping power" (PDG 2018)

One can use Bethe-Bloch Formula to get a good approximation

- $N_{PE} \sim \epsilon^2 \times 10^6$ for 1 meter plastic scintillation bar
- $\epsilon \sim 10^{-3}$ roughly gives one PE



Signature: Triple Coincidence

• Based on Poisson distribution, zero event in each bar correspond to $P_0 = e^{-N_{PE}}$, so the probability of seeing triple

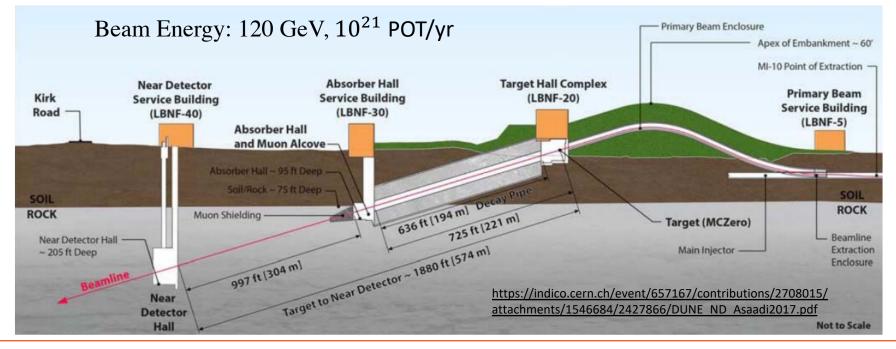
 $I_0 - c_{i}$, so the probability of seeing triple

incident of one or more photoelectron is:

$$P = \left(1 - e^{-N_{PE}}\right)^3$$

• $N_{x,detector} = N_x \times P$.

Site 2: LBNF Beam & DUNE ND Hall



Jonathan Asaadi – University of Texas Arlington

LBNF: Long-Baseline Neutrino Facility

There are many other **new physics opportunities** in the **near detector hall**! Combine with **DUNE PRISM**?

Detector Background

- We will discuss two major detector
 backgrounds and the reduction technique
- SM charged particles from background radiation (e.g., cosmic muons):
 - Offline veto of events with > 10 PEs
 - Offset middle detector
- Dark current: triple coincidence

Dark Current Background @ PMT

• Major Background (BG) Source!

- dark-current frequency to be $v_B = 500 \text{ Hz}$ for estimation (1607.04669)
- For each tri-PMT set, the background rate for triple incidence is

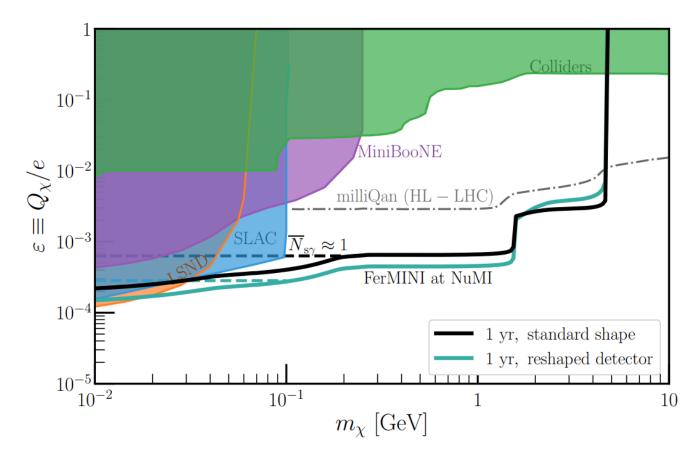
$$v_B^3 \Delta t^2 = 2.8 \times 10^{-8}$$
 Hz, for $\Delta t = 15$ ns.

• Consider 400 such PMT sets:

the total background rate is 400 x 2.8 x $10^{-8} \sim 10^{-5}$ Hz

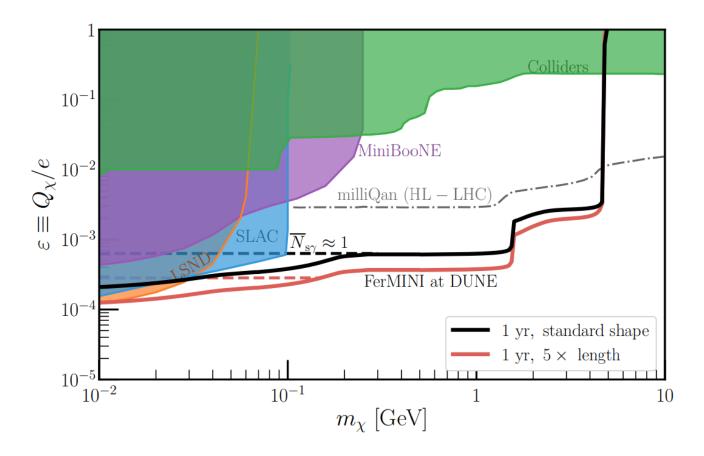
- ~ **300 events** in one year of trigger-live time
- Quadruple coincidence can reduce this BG to essentially zero!

FerMINI @ MINOS



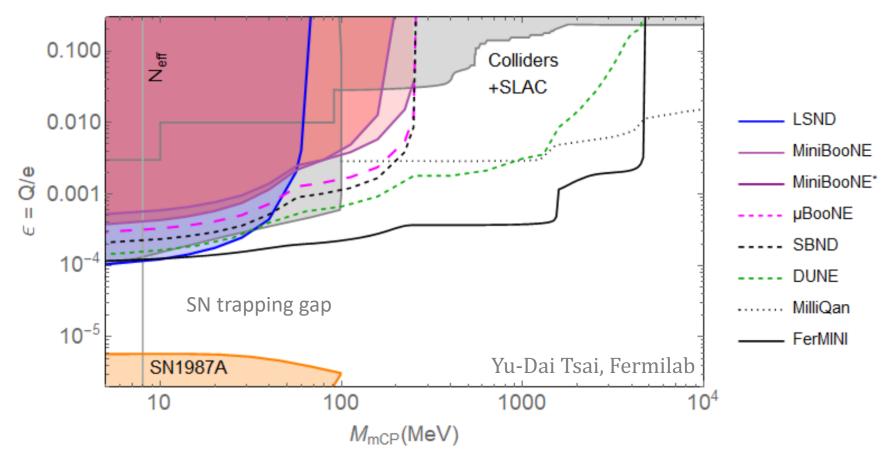
Yu-Dai Tsai, Fermilab

FerMINI @ DUNE



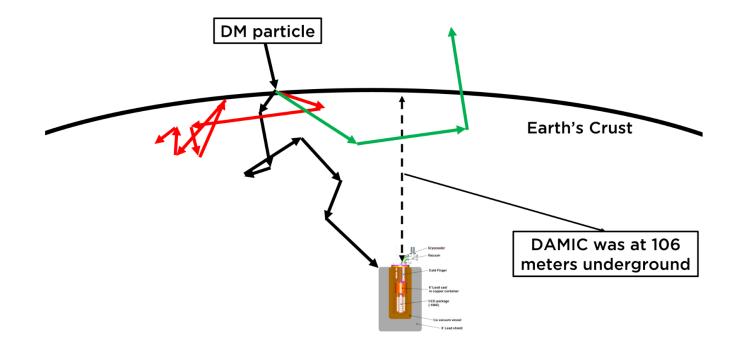
Yu-Dai Tsai,
Fermilab
+DUNE PRISM? Combine with DUNE to get timing?

Compilation of MCP Probes



- One can combine the MCP detector with neutrino detector to improve sensitivity or reduce background
- Filling up the MCP "cavity"

Extended discussion: Strongly Interacting Dark Matter DM-SM Interaction too strong that attenuation stop the particles from reach the direct detection detector



DMATIS (Dark Matter ATtenuation Importance Sampling), Mahdawi & Farrar '17

Strongly Interacting Dark Matter

See, e.g., arXiv:1905.06348 (Emken, Essig, Kouvaris, Sholapurkar '19)

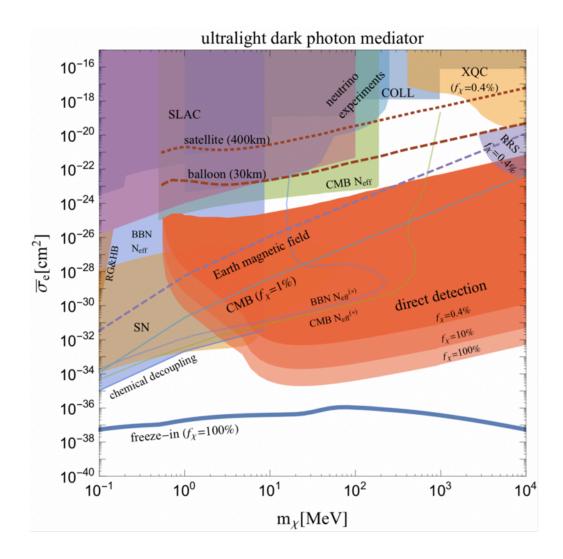
Scatterings both on electrons and nuclei in the Earth's crust, atmosphere, and shielding material attenuate the expected local dark matter flux at a terrestrial detector, so that such experiments lose sensitivity to dark matter above some critical cross section.

Limits of the underground Direct Detection (DD) Experiments, including SENSEI, CDMS-HVeV, XENON10, XENON100, and DarkSide-50

One can call the DM that could escape the DD bound this way as **Strongly Interacting Dark Matter (SIDM)**

Not to confuse with Self Interacting Dark Matter (also SIDM)

Millicharged (with ultralight A') SIDM Window



From arXiv:1905.06348, they defined **reference cross section:**

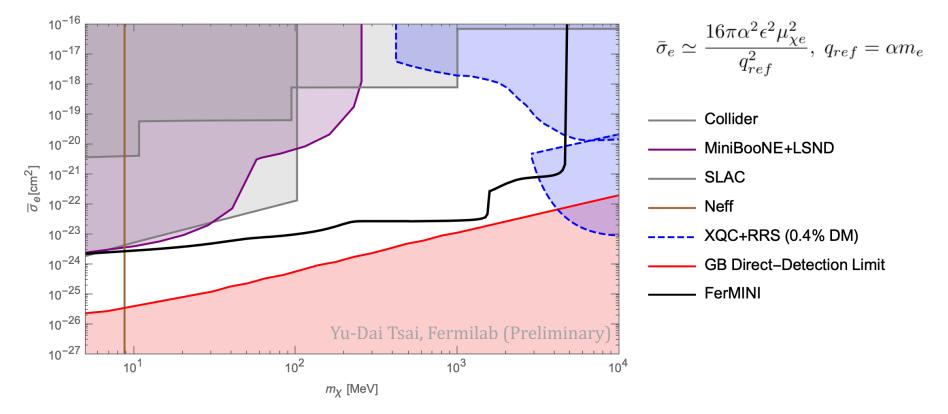
$$\bar{\sigma}_e \equiv \frac{16\pi\alpha\alpha_D\kappa^2\mu_{\chi e}^2}{(q_{ref}^2 + m_{A'}^2)^2},$$
$$m_{A'} \to 0, \ q_{ref} = \alpha m_e$$

 q_{ref} is chosen as the typical momentum transfer in DM-electron collisions for noble-liquid / semiconductor targets.

Agonistic to the abundance setting mechanism for the SIDM window.

FerMINI Probe of Millicharged SIDM

MCP / LDM with ultralight dark photon mediators, all curves except FerMINI are from arXiv:1905.06348



- Here we plot the **electron-scattering Millicharged SIDM** from 1905.06348 (Emken, Essig, Kouvaris, Sholapurkar)
- FerMINI can help close the Millicharged SIDM window!

Advantages of FerMINI: Timeliness, Low-cost, Movable, Tested, Easy to Implement, ...

- 1. LHC entering long shutdown
- 2. NuMI operating, shutting down in 5 years (DO IT NOW! Fermilab! USA!)
- 3. Broadening the physics case for fixed-target facilities
- 4. **DUNE near detector design** still underway
- 5. Can develop at NuMI/MINOS and then move to DUNE
- 6. Sensitivity better than milliQan for MCP up to 5 GeV and don't have to wait for HL-LHC
- Synergy between dark matter, neutrino, and collider community.
 Join us on the proposal! (ytsai@fnal.gov)

More on MCP/DM & 21-cm Cosmology

Some more reference of Millicharged DM (mDM) and constraints.

See, e.g.,

McDermott, Yu, Zurek, 1011.2907;

Muñoz, Dvorkin, Loeb, 1802.10094, 1804.01092;

Berlin, Hooper, Krnjaic, McDermott, 1803.02804;

Kovetz, Poulin, Gluscevic, Boddy, Barkana, Kamionkowski, 1807.11482;

Liu, Outmezguine, Redigolo, Volansky, 1908.06986:

"Reviving Millicharged Dark Matter for 21-cm Cosmology,"

Introduces a long-range force between a subdominant mDM and the dominant cold dark matter (CDM) components. Leads to efficient cooling of baryons in the early universe. Extend the range of viable mDM masses for EDGES explanation to ~ 100 GeV.

Thank You!

Additional Motivations

• Won't get into details, but it's interesting to find

"pure" MCP, that is WITHOUT a massless or ultralight dark photon (finding MCP in the regime where ultralight/massless A' is strongly constrained by cosmology!)

- More violent violation of the charge quantization
 (if not generating millicharge through kinetic mixing)
- Test of GUT models, and String Compactifications see Shiu, Soler, Ye, arXiv:1302.5471, PRL '13 for more detail.

FerMINI: Alternative Designs & New Ideas

New Ideas ...

- **Combine with neutrino detector**: behind, in front, or sandwich them
- Combine with **DUNE PRISM**: moving up and down
- FerMINI + DUNE 3-D scintillation detector (3DST)
- Combine with **SPS/SHiP facilities**
- Can potentially probe (electric) dipole portal dark fermion, quirks, etc.
- Join the Proposal: <u>ytsai@fnal.gov</u>

Looking Ahead

- Exploring Energy Frontier of the Intensity Frontier (complementary to and before HL-LHC upgrade)
- Cosmology-driven models / more motivated models.
- Near-future (and almost free) opportunity (NuMI Facility, SBN program, DUNE Near Detector, etc.)
- Other new low-cost alternatives/proposals (~ \$1M) to probe hidden particles and new forces (FerMINI is just a beginning!)
- Dark sectors in neutrino telescopes

Some anomalies involving MeV-GeV+ Explanations



- Muon g-2
- Proton charge radius anomaly
- LSND & MiniBooNE anomaly
- EDGES result

Below ~ MeV there are also strong astrophysical/cosmological bounds

Kinetic Mixing and MCP Phase

 Coupled to new dark fermion (scalar) χ

$$B' \bigvee O \bigvee B$$
 (SM: Standard Model)

See, Holdom, 1985

$$\mathcal{L} = \mathcal{L}_{\rm SM} - \frac{1}{4} B'_{\mu\nu} B'^{\mu\nu} - \frac{\kappa}{2} B'_{\mu\nu} B^{\mu\nu} + i\bar{\chi}(\partial \!\!\!/ + ie' B' + iM_{\rm MCP})\chi$$

- New Fermion χ charged under dark U(1)'
- Field redefinition into a more convenient basis for massless B', $B' \rightarrow B' + \kappa B$
- new fermion acquires an small EM charge Q (the charge of mCP χ): $Q = \kappa e' \cos \theta_W \quad \epsilon \equiv \kappa e' \cos \theta_W / e.$

Not all bounds are created with equal assumptions

Accelerator-based: Collider, Fixed-Target Experiments Some other ground based experiments

Astrophysical productions (not from ambient DM): energy loss/cooling, etc: Rely on modeling/observations of (extreme/complicated/rare) systems (SN1987A)

Dark matter direct/indirect detection: abundance, velocity distribution, etc

Zdifferent

techinica

Cosmology: assume cosmological history, species, etc

Or, how likely is it that theorists would be able to argue our ways around them

Yu-Dai Tsai, FNAL, 2019 • Astrophysical/cosmological observations are important to reveal the actual story of dark matter (DM).

Backup Slides

Potential Detection Limitation: $N_{photon} \leq 1$

- Define: ε_{low} as $N_{sintilator photon} = 1$
- Roughly around or below this, one really have to worry about scintillator performance
- One can elongate the scintillator or consider alternative materials to help.

Material	Photons/keV	Density (g/cm ³)	* Length needed (cm)	Speed (ns)	Cost for 5x5 cm (\$)	Notes	
Plastic BC408	10	1.03	145	~2	~200	Current choice	
Nal	38	3.67	11	~230	~800	Slow, fragile	
LaBr3(Ce)	63	5.08	5	~16	~3000	Radioactive	
Liquid Xe	62	2.95	8	~2 / ~34	~1000?	Cryogenic, ultraviolet	

• Andy Haas, Fermilab, 2017

* Length needed to get 3 photons for charge 1/1000 e

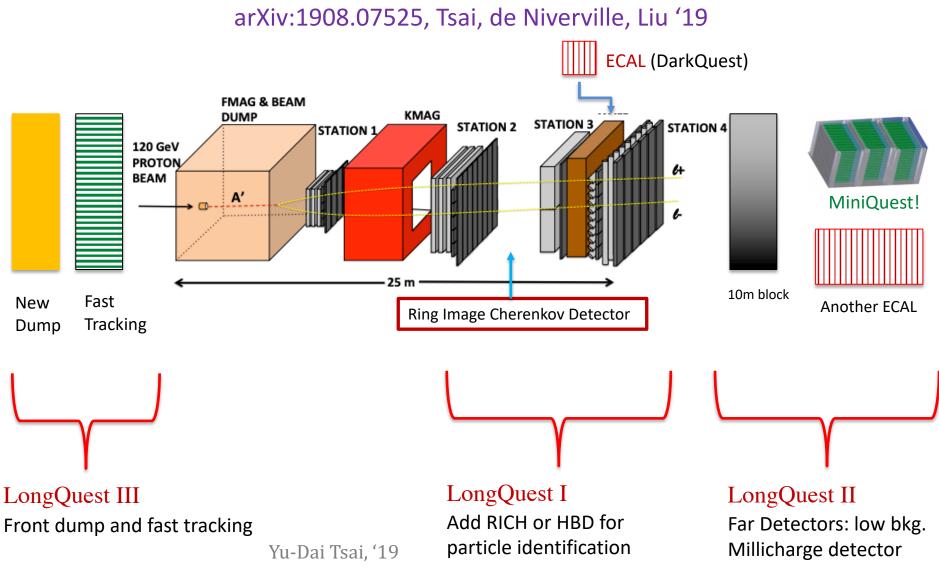
NuMI (MINOS) / LBNF (DUNE) Now and the future bests in POTs

- LSND: total of 10^{23} POT (beam: 800 MeV)
- Fermilab (FT):
- NuMI beam: 1 4 x 10^{20} POT/yr (120 GeV)
- LBNF beam: $1 2 \times 10^{21}$ POT/yr (120 GeV)
- CERN SPS (FT):
- NA62: up to 3×10^{18} POT/yr (400 GeV)
- SHiP: up to 10^{19} POT/yr (400 GeV)
- FASER (collider, forward): 10^{16} - 10^{17} POT/yr

Yu-Dai Tsai Fermilab

much higher energy

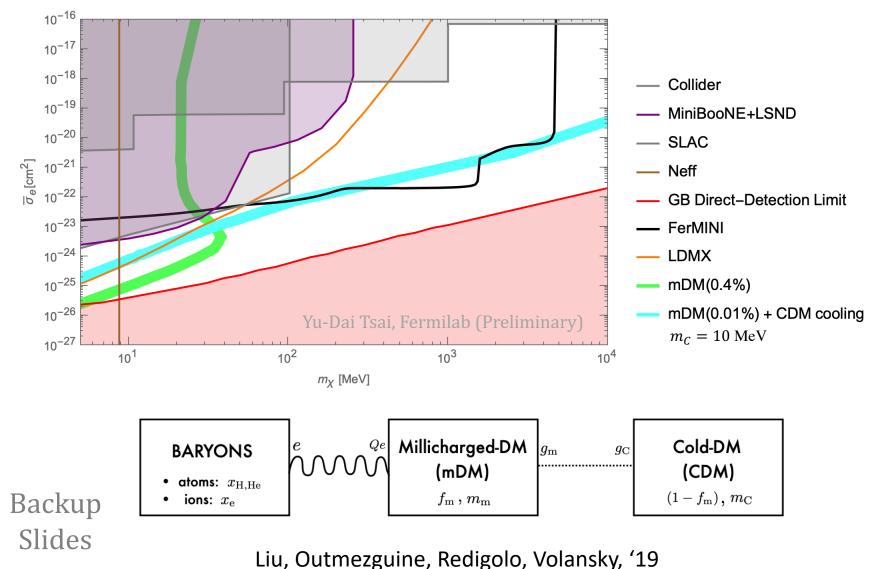
LongQuest: Three Stage Retool of SpinQuest, as Dedicated Long-Lived Particle Experiment



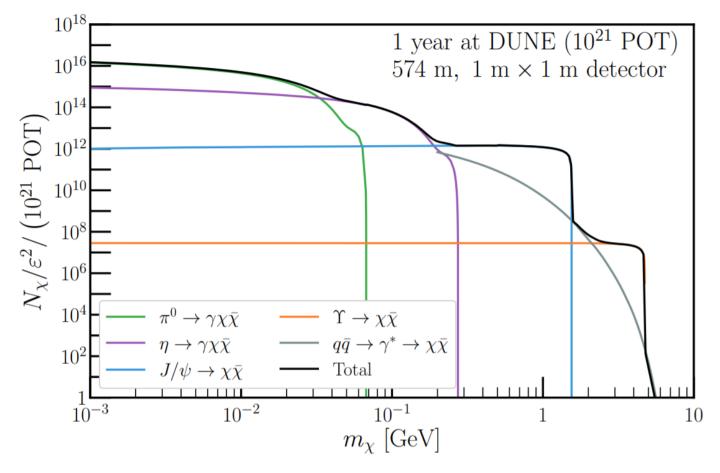
XQC & RRS

- 1. X-ray Quantum Calorimeter: X-ray detector aboard a sounding rocket
- 2. RRS (RICH, ROCCHIA, SPIRO), Ahlen et al., Harvard-Smithsonian Observatory pre- print 2292 (1986).
- 3. RRS is on balloon

Reviving mDM for EDGES

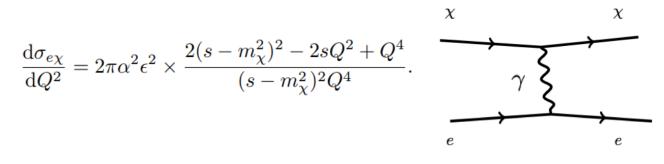


MCP Production/Flux



- Use PYTHIA to generate neutral meson Dalitz or direct decays from the pp collisions and rescale by considering, BR $(\mathcal{M} \to \chi \bar{\chi}) \approx \epsilon^2 \times BR(\mathcal{M} \to X e^+ e^-) \times f\left(\frac{m_{\chi}}{M}\right)$,
- M: mass of the parent meson, X:additional particles, $f(m_{\chi}/M)$: phase space factor
- We also include Drell-Yan production for the high mass MCPs (see <u>arXiv:1812.03998</u>)

Detection: MCP Elastic Scattering with Electrons



- Q^2 is the squared 4-momentum transfer.
- Integrate over Q^2 , total cross section dominated by the small Q^2 contribution, we have $\sigma_{e\chi} = 4\pi \alpha^2 \epsilon^2 / Q_{min}^2$.
- Light mediator: the total cross section is dominated by the small Q² contribution

DM Form Factor Defined in 1905.06348

$$F_{
m DM}(q) = rac{q_{
m ref}^2 + m_{A'}^2}{q^2 + m_{A'}^2}\,,$$

which parametrizes the q dependence.

Alternatives (Straightforward)

- Quadruple incidence: further background reduction, sacrifice event rate but potentially gain better control of background, reduce the background naively by 10⁻⁵ Basically zero dark-current background experiment?
- 2. Different lengths for each detectors
- 3. Different materials:

Material	Photons/keV	Density (g/cm³)	* Length needed (cm)	Speed (ns)	Cost for 5x5 cm (\$)	Notes	
Plastic BC408	10	1.03	145	~2	~200	Current choice	
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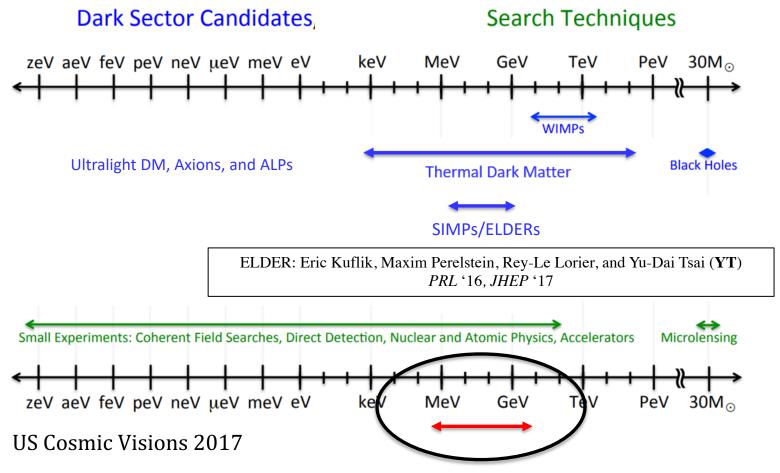
Backup Slides

• Andy Haas, Fermilab, <u>2017</u>

* Length needed to get 3 photons for charge 1/1000 e

Even More Backup Slides (Deleted Intro)

Exploration of Dark Matter & Dark Sector



- Astrophysical/cosmological observations are important to reveal the actual story of dark matter (DM).
- Why FT experiments? And why MeV GeV+?

Proton FT Experiments

- High statistics, e.g. LSND has 10²³ Protons on Target (POT)
- Shielded/underground: lower background
- Many of them existing and many to come:

strength in numbers

- Relatively high energy proton beams on targets exist
 O(100 400) GeV (I will compare Fermilab/CERN facilities)
- Produce hidden particles / involve less assumptions

Why study MeV – GeV+ dark sectors?

Backup Slides

Signals of discoveries grow from anomalies Maybe nature is telling us something so we don't have to search in the dark? (most likely systematics?)

Backup Slides

Some anomalies involving MeV-GeV+ Explanations



- Muon g-2
- Proton charge radius anomaly
- LSND & MiniBooNE anomaly
- EDGES result

Below ~ MeV there are also strong astrophysical/cosmological bounds

v Hopes for New Physics: Personal Trilogy

- Light Scalar & Dark Photon at Borexino & LSND
 Pospelov & YT, PLB '18, <u>1706.00424</u> (proton charge radius anomaly)
- Dipole Portal Heavy Neutral Lepton

Magill, Plestid, Pospelov & YT, PRD '18, <u>1803.03262</u> see also Coloma, Machado, Martinez-Soler, Shoemaker, <u>1707.08573</u> (LSND/MiniBooNE anomalies)

• Millicharged Particles in Neutrino Experiments

Magill, Plestid, Pospelov & **YT**, PRL '19, <u>1806.03310</u>

(EDGES 21-cm measurement anomaly)

New Physics in Proton FT Experiments

- Millicharged Particles in FerMINI Experiments
 - Kelly & **YT,** <u>1812.03998</u>
 - (EDGES Anomaly)
- Dark Neutrino at Scattering Experiments: CHARM-II & MINERvA! Argüelles, Hostert, YT, <u>1812.08768</u>, under *PRL* review (MiniBooNE Anomaly)
- Probing Dark Photon, Inelastic Dark Matter, and Muon g-2
 Windows with LongQuest Proposal! (Comin out Monday night!)

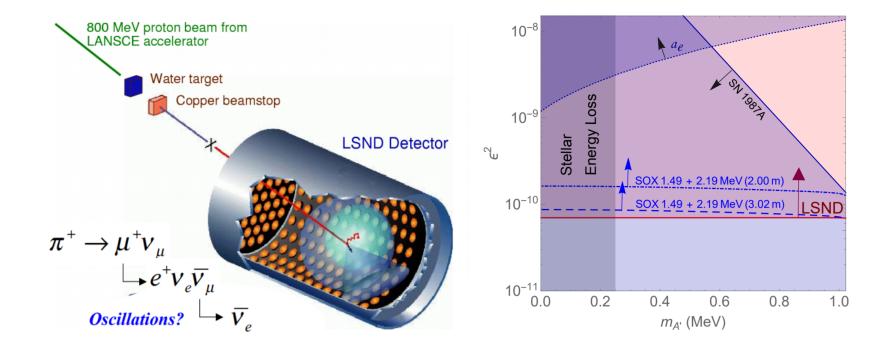
Other New Physics Probes

Backup Slides

Dark Photon @ LSND

Pospelov & YT, PLB '18, <u>1706.00424</u>

$$\mathcal{L}_{\rm d.ph.} = -\frac{1}{4} F'_{\mu\nu} F'^{\mu\nu} + \frac{1}{2} m_{A'}^2 (A'_{\mu})^2 + \epsilon A'^{\mu} J_{\mu}^{EM}$$

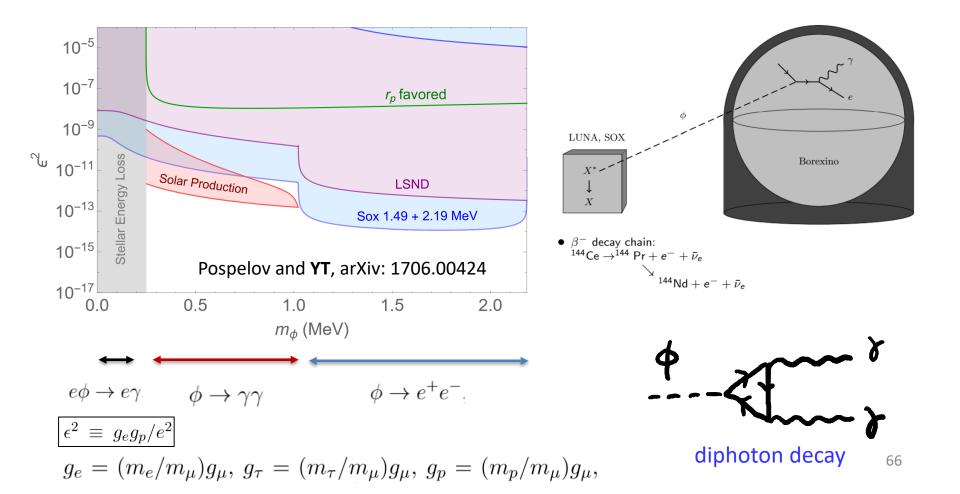


- Backup Slides - Major energy depositions: $e+A' \rightarrow e+\gamma$

Light Scalar @ LSND & Borexino

Pospelov & YT, PLB '18, 1706.00424

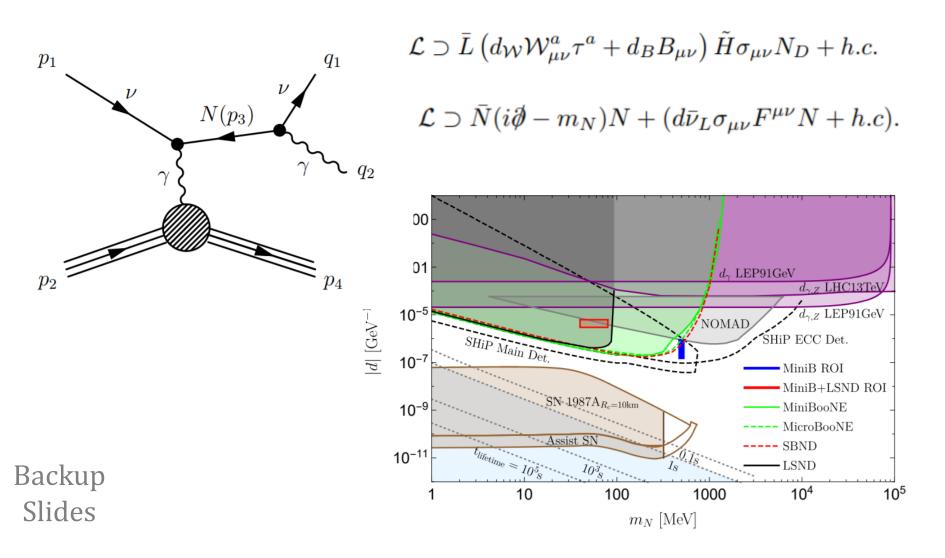
$$\mathcal{L}_{\phi} = \frac{1}{2} (\partial_{\mu}\phi)^2 - \frac{1}{2} m_{\phi}^2 \phi^2 + (g_p \bar{p}p + g_n \bar{n}n + g_e \bar{e}e + g_{\mu} \bar{\mu}\mu + g_{\tau} \bar{\tau}\tau)\phi$$



Dark Neutrino at CHARM & MINERvA Argüelles, Hostert, YT, <u>1812.08768</u>, under PRL review $\mathscr{L}_{\rm int} \supset g_D \overline{\nu}_D \gamma_\mu \nu_D Z'^\mu + e \varepsilon Z'^\mu J^{\rm EM}_\mu,$ $A \bigcirc^{Z'}$ $\nu_{\alpha} = \sum_{i=1}^{4} U_{\alpha i} \nu_i, \quad (\alpha = e, \mu, \tau, D).$ $m_{Z'} = 30 \text{ MeV}, \ \alpha \varepsilon^2 = 2 \times 10^{-10}, \ \alpha_D = 1/4$ 10^{-1} szemdeo 10^{-2} $\begin{array}{c|c} \text{MiniBooNE} \\ \hline & 1\sigma \\ \hline & 3\sigma \end{array}$ $O[1]{D_{\mu_4}}$ 10⁻³ 10^{-4} CHARM-II MINERvA 60% on axis 90% on axis Backup 10^{-1} 10^{0} Slides $m_4 \; (\text{GeV})$ 67

Dipole-Portal Heavy Neutral Lepton

Magill, Plestid, Pospelov & **YT** , PRD '18, <u>1803.03262</u>



(detail) Meson Production Details

- At LSND, the π 0 (135 MeV) spectrum is modeled using a Burman-Smith distribution
- Fermilab's Booster Neutrino Beam (BNB): π0 and η (548 MeV) mesons. π0's angular and energy spectra are modeled by the Sanford-Wang distribution. η mesons by the Feynman Scaling hypothesis.
- SHiP/DUNE: pseudoscalar meson production using the BMPT distribution, as before, but use a beam energy of 80 GeV
- J/ψ (3.1 GeV), we assume that their energy production spectra are described by the distribution from Gale, Jeon, Kapusta, PLB '99, nucl-th/9812056.
- Upsilon, Y (9.4 GeV): Same dist., normalized by data from HERA-B, I. Abt et al., PLB (2006), hep-ex/0603015.
- Calibrated with existing data [e.g. NA50, EPJ '06, nucl-ex/0612012, Herb et al., PRL '77]. and simulations from other groups [e.g. deNiverville, Chen, Pospelov, and Ritz, Phys. Rev. D95,
 Backup Slides

FerMINI: Beam Related Background

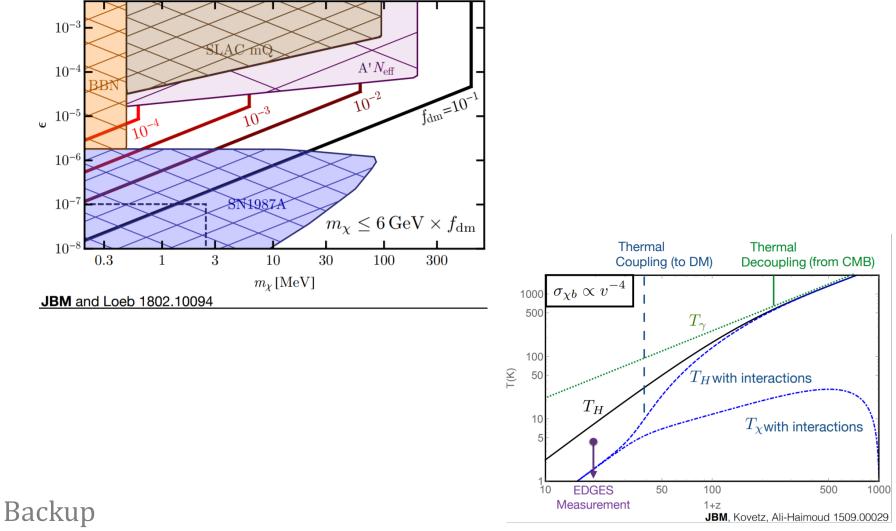
- Shielding: including absorber and rocks.
- Controlled: muon monitors.
- Can determine the SM charged particle rate on site
- Vetoed similar to the previous veto of cosmic muons.
- Neutrino produced hard-scattering background: **O(10⁻¹⁹)**, negligible.
- To be conservative, we assume the beam related background ≈ dark current background for our sensitivity determination.
- Based on **SENSEI experience**, beam produced charge background is weaker than cosmic, but of course energy dependence

Backup Assumed to be at the same level of detector background Slides

FerMINI: Increasing scintillation photons

- Elongating the scintillator bar does not affect the background from dark current (basically determined by the number of PMTs)
- So we estimate the sensitivity of FerMINI at DUNE for five times larger scintillation capability
- And estimate the sensitivity of FerMINI at NuMI for five time more scintillation capability but five times less scintillator
 bar-PMT sets (actually reduce dark current background!)

EDGES ANOMALY and MCP Solution



Slides

(Detail) dE/dx formula

 For moderately small epsilon and heavy enough MCP (>> electron mass), one can use Bethe equation to estimate average energy loss.

$$\left\langle -\frac{dE}{dx}\right\rangle = Kz^2 \frac{Z}{A} \frac{1}{\beta^2} \left[\frac{1}{2}\ln\frac{2m_e c^2 \beta^2 \gamma^2 W_{\text{max}}}{I^2} - \beta^2 - \frac{\delta(\beta\gamma)}{2}\right]$$

- z charge number of incident particle
- Z atomic number of absorber
- A atomic mass of absorber $g \mod^{-1}$
- $K = 4\pi N_A r_e^2 m_e c^2$ (Coefficient for dE/dx)

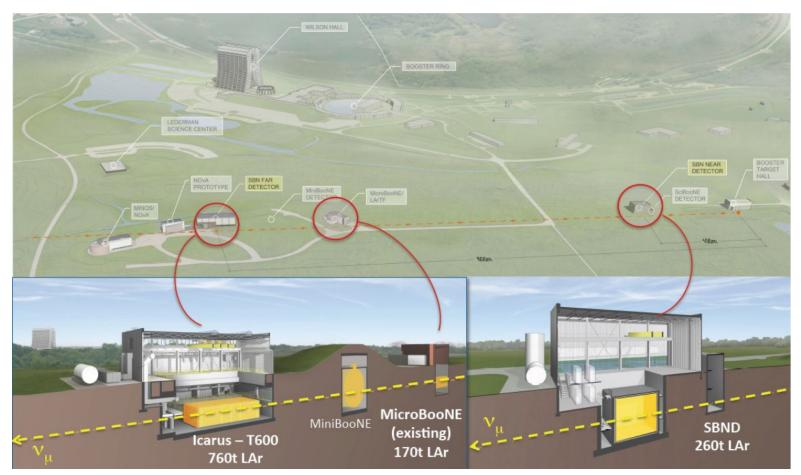
Ι

- $0.307\,075~{\rm MeV}~{
 m mol}^{-1}~{
 m cm}^2$
- (Coefficient for dE/dx) mean excitation energy eV (Nota bene!) $W_{\text{max}} = \frac{2m_e c^2 \beta^2 \gamma^2}{1 + 2\gamma m_e/M + (m_e/M)^2}$.
- $\delta(\beta\gamma)$ density effect correction to ionization energy loss
 - M: charged particle mass
- Backup
 For very small epsilon (related to the finite length effect), one have to consider most probable energy deposition & consider landau distribution for the energy transfer, see <u>arXiv:1812.03998</u>

MCP @ Neutrino Detectors

Backup Slides

Neutrino Experiments



<u>https://web.fnal.gov/collaboration/sbn_sharepoint/SitePages/Civil_Construction.aspx</u> SBND: Short Baseline Near Detector of Booster Beam MiniBooNE: Mini-Booster Neutrino Experiment <u>ICARUS (</u>Imaging Cosmic And Rare Underground Signals<u>):</u> Now a Far Detector of Booster Beam

MCP Signals

• signal events sevent

$$s_{\text{event}} \simeq \sum_{\text{Energies}} N_{\chi}(E_i) \times \frac{N_e}{\text{Area}} \times \sigma_{e\chi}(E_i; \ m_{\chi}) \times \mathcal{E}.$$

detection efficiency

- $N_{\chi}(E_i)$: number of mCPs with energy E_i arriving **at the detector**.
- N_e: total number of electrons inside the active volume of the detector
- Area: active volume divided by the average length traversed by particles inside the detector.
- $\sigma_{e\chi}(E_i)$: detection cross section consistent with the angular and recoil cuts in the experiment
- Here, $s_{event} \propto \varepsilon^4$. ε^2 from N_x and ε^2 from $\boldsymbol{\sigma}_{ex}$
- Throughout this paper, we choose a credibility interval of $1 \alpha = 95\%$ (~ 2 sigma)
- Roughly, $\varepsilon_{sensitivity} \propto E_{e, R, min}^{1/4} Bg^{1/8}$

MCP Bound/Sensitivity

• signal events sevent

$$s_{\text{event}} \simeq \sum_{\text{Energies}} N_{\chi}(E_i) \times \frac{N_e}{\text{Area}} \times \sigma_{e\chi}(E_i; m_{\chi}) \times \mathcal{E}.$$

- Our sensitivity curves are obtained by performing a standard sensitivity analysis [PDG, PLB 2010]:
- Given a number of background events b and data n, the number of signal events s_{event} . The (1α) credibility level is found by solving the equation $\alpha = \Gamma(1 + n, b + s_{event})/\Gamma(1 + n, b)$, where $\Gamma(x, y)$ is the upper incomplete gamma function.
- Throughout this paper, we choose a credibility interval of $1 \alpha = 95\%$ (~ 2 sigma)

Summary Table

		$N [\times 10^{20}]$		$A_{\rm geo}(m_{\chi})[\times 10^{-3}]$		Cuts [MeV]		
	Exp. (Beam Energy, POT)	π^0	η	$1 { m MeV}$	$100 {\rm ~MeV}$	E_e^{\min}	E_e^{\max}	Bkg
Existing	LSND (0.8 GeV, 1.7×10^{23})	130		20		18	52	300
	mBooNE (8.9 GeV, 2.4×10^{21})	17	0.56	1.2	0.68	130	530	2k
	mBooNE* (8.9 GeV, 1.9×10^{20})	1.3	0.04	1.2	0.68	75	850	0.4
	$\mu \text{BooNE} (8.9 \text{ GeV}, \ 1.3 \times 10^{21})$	9.2	0.31	0.09	0.05	2	40	16
	SBND (8.9 GeV, 6.6×10^{20})	4.6	0.15	4.6	2.6	2	40	230
	DUNE (80 GeV, 3.0×10^{22})	830	16	3.3	5.1	2	40	19k
	SHiP (400 GeV, 2.0×10^{20})	4.7	0.11	130	220	100	300	140

- $\varepsilon \propto E_{e,R,min}^{1/4} Bg^{1/8}$
- cos θ > 0 is imposed (*except for at MiniBooNE's DM run where a cut of cos θ > 0.99 effectively reduces backgrounds to zero [Dharmapalan, MiniBooNE, (2012)]).

• Efficiency of 0.2 for Cherenkov detectors, 0.5 for nuclear emulsion detectors, and 0.8 for liquid argon time projection chambers.

Recasting Existing Analysis:

LSND, MiniBooNE, and MiniBooNE* (DM Run)

- LSND: hep-ex/0101039. Measurement of electron-neutrino electron elastic scattering
- MiniBooNE: arXiv:1805.12028.

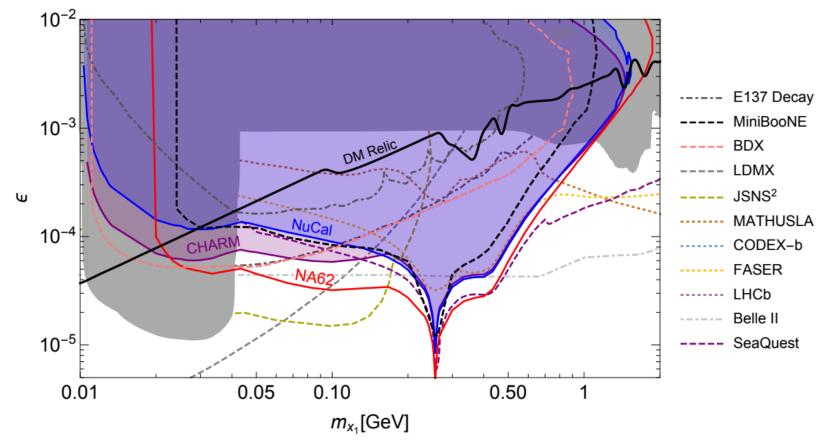
Electron-Like Events in the MiniBooNE Short-Baseline Neutrino Experiment, combines data from both neutrino and antineutrino runs and consider a sample of 2.4×10^{21} POT for which we take the single electron background to be 2.0×10^3 events and the measured rate to be 2.4×10^3

• MiniBooNE* (DM run): arXiv:1807.06137 (came out after our v1). Electron recoil analysis.

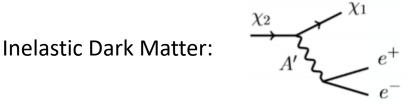
Thick target + no horn focusing +

A cut of $\cos \theta > 0.99$ effectively reduces backgrounds to basically zero [Dharmapalan, MiniBooNE, (2012)]).

Tsai, de Niverville, Liu, 1908.07525



(e) Compilation of relevant constraints and sensitivity projections for iDM with $\alpha_D = 0.1$ and $\Delta = 0.1$.



FerMINI

Low-cost fixed-target probes of dark sector/long-lived Particles