

# Small-X effects in forward-jet production at HERA

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DIS 2005, Madison

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# Introduction

# BFKL evolution and saturation

- example:  $\gamma^* - \gamma^*$  total cross-section  
 $\gamma^* \gamma^* \rightarrow X$  suitable to test fixed-scale evolution **BFKL prediction:**

$$\sigma_{BFKL}(Q_1, Q_2, \Delta\eta) \propto \exp(\lambda \Delta\eta)$$

Saturation: restores unitarity

- High-energy behavior:  
determined by dipole-dipole scattering
- Timneanu, Kwiecinski and Motyka (2002)  
Kozlov and Levin (2003)

# Forward-jet production

- $p + \gamma^* \rightarrow \text{jet} + X$   
 $Q, k_T \gg \Lambda_{\text{QCD}}$     $\exp(\Delta\eta) \gg 1$   
 $\sigma_{BFKL}(k_T, Q, \Delta\eta) \propto \exp(\lambda \Delta\eta)$
- Same kind of process than  $\gamma^* - \gamma^*$   
but more statistics  
**data from H1 and ZEUS**
- Is the BFKL resummation seen at  
HERA?
- Are saturation effects sizable at  
HERA?

The high-energy behavior is also determined by dipole-dipole scattering

# Gluon emission by a hadron

C. M., Nucl. Phys. B 705 (2005) 319

It can be formulated in terms of the gluon dipole cross-section  $\sigma_{(gg)t}$ :

$$\frac{d\sigma^{h+t \rightarrow J+X}}{d^2 k d x_J} = \frac{1}{2k^2} g_h(x_J, \mu^2) \int dr J_0(kr) \frac{\partial}{\partial r} r \frac{\partial}{\partial r} \sigma_{(gg)t}(r)$$

- $x_J$  is the longitudinal momentum fraction of the gluon-jet
- $k$  is the transverse momentum of the gluon-jet
- $g_h$  is the gluon density inside the incident hadron
- $\mu$  is the factorization scale
- This formula is valid in the leading-logarithmic approximation
- One recovers the kT-factorization formula by restricting  $\sigma_{(gg)t}$  to a two-gluon exchange (BFKL)

# The forward-jet cross-section

- $x, y, Q_2$ : usual kinematic variables of DIS
- $x_J, k$ : longitudinal and transverse momentum of the jet
- $\Delta\eta = \log(x_J/x)$  : rapidity interval

$$\frac{d\sigma}{dx dQ^2 dx_J dk^2} = \frac{\alpha g_p(x_J, k^2)}{2x k^2 Q^2} \left\{ \frac{y^2 d\sigma_T}{2 dk^2} \right. \\ \left. + (1-y) \left( \frac{d\sigma_T}{dk^2} + \frac{d\sigma_L}{dk^2} \right) \right\}$$

with the hard cross-section given by

$$\frac{d\sigma_{TL}}{dk^2} = \int d^2 r_1 d^2 r_2 |\psi_{TL}(r_1, Q^2)|^2 \frac{J_0(kr_2)}{2\pi r_2^2} \frac{\partial}{\partial r_2} r_2^2 \frac{\partial}{\partial r_2} \sigma_{(gg)d}(\Delta\eta, r_1, r_2)$$

# Parametrizations of $\sigma_{(gg)d}$

# BFKL parametrization

- The BFKL dipole-dipole cross-section:

$$\sigma_{(gg)d}(\Delta\eta, r_1, r_2) = 2\pi\alpha_S^2 r_1^2 \int \frac{dy}{2i\pi} \left| \frac{r_1/r_2}{1-y} \right|^{2\gamma} \frac{\exp(\bar{\alpha}\chi(y)\Delta\eta)}{\gamma^2(1-y)^2}$$

with the BFKL kernel  $\chi(y) = 2\psi(1) - \psi(y) - \psi(1-y)$

- We fit the parameter  $\lambda=4\bar{\alpha}\ln(2)$  and a normalization

# Saturation parametrization

- An extension of the GBW model

Tămneanu, Kwieciński and Motyka (2002)

$$\sigma_{(gg)d}(\Delta\eta, r_1, r_2) / \sigma_0 = 1 - \exp \left[ -r_{eff}^2 / 4R_0^2(\Delta\eta) \right]$$

$$\text{with } r_{eff}^2 = r_<^2 \left[ 1 + \ln(r_>/r_<) \right]$$

$$r_< = \min(r_1, r_2) \quad r_> = \max(r_1, r_2)$$

- The saturation radius is

$$R_0(\Delta\eta) = \frac{1}{Q_0} \exp \left[ -\frac{\lambda}{2}(\Delta\eta - \Delta\eta_0) \right]$$

$$Q_0 \equiv 1 \text{ GeV}$$

we fit the parameters  $\lambda$ ,  $\Delta\eta_0$  and a normalization

# Fitting the HERA data

J. Contreras, R. Peschanski and C. Royon, Phys. Rev. D 62 (2000) 034006

C.M., R. Peschanski and C. Royon, Phys. Lett. B 599 (2004) 236

# Results of the fits

fit	$\lambda$	$\Delta\eta_0$	$\chi^2$ (/dof)
BFKL	0.430	*****	12 (/13)
sat.	0.402	-0.82	6.8 (/11)
weak sat.	0.370	8.23	8.3 (/11)

- The BFKL fit is a measure of the hard Pomeron intercept
- We obtain two  $\chi^2$  minima for the saturation fit:

The « sat. » solution corresponds to significant saturation effects

The « weak sat. » solution corresponds to weak saturation effects

The intercept  $\lambda$  is in both cases higher than what was found for  $F_2$  ( $\lambda_{GBW} = 0.288$ )

# The BFKL fit

# The Saturation fit

# The saturation scales

- The saturation scale is

$$Q_S \equiv 1/R_0(\Delta\eta)$$

- The plot represents

$$\log(Q_S^2/Q_0^2) = \lambda(\Delta\eta - \Delta\eta_0)$$

- The weak saturation solution is compatible with the  $F_2$  parametrization

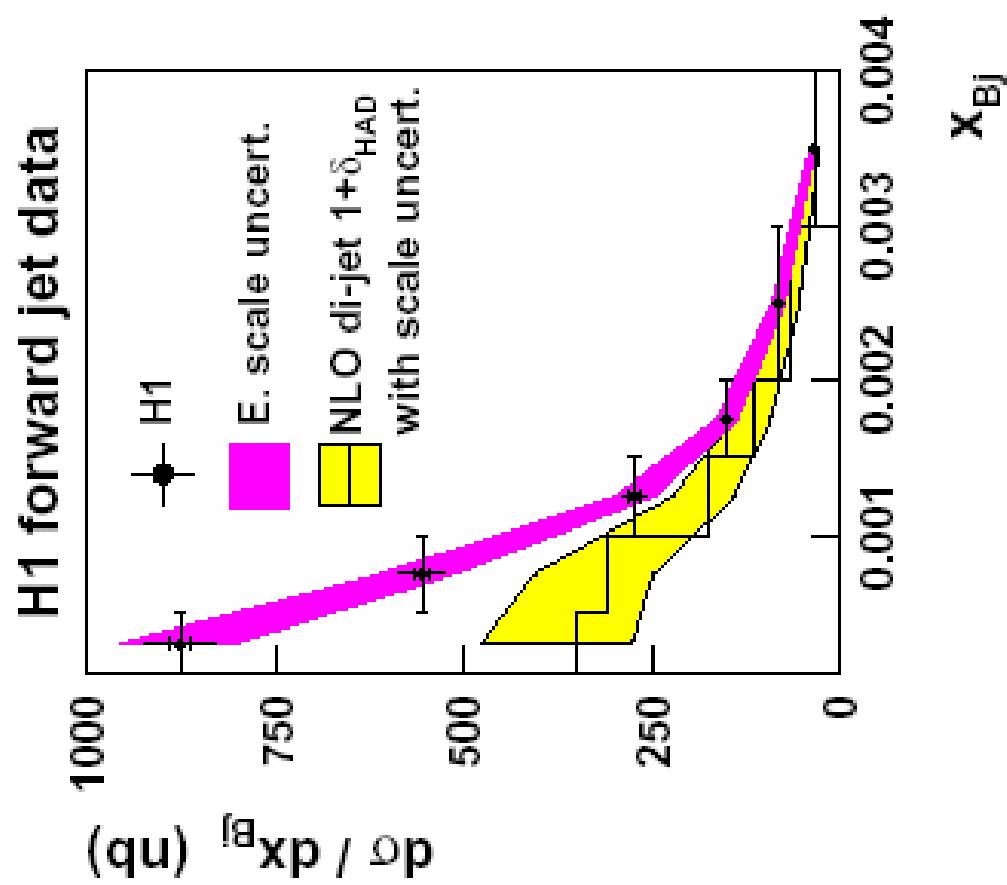
- The other solution shows a harder saturation scale

# Comparisons with new data and predictions

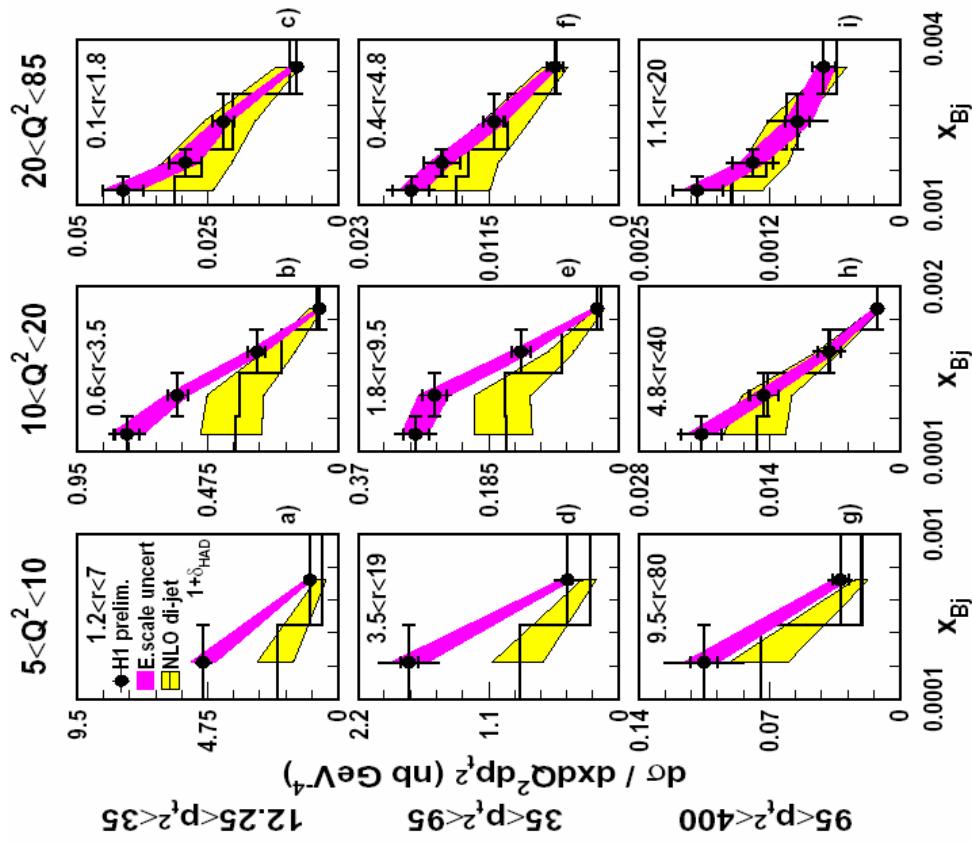
C. M. and C. Royon, in preparation

# Comparisons with ZEUS data

# Predictions for H1 preliminary data



# Predictions for $d\sigma/dx dQ^2 dk^2$



# Conclusion and outlook

- Studies of small- $x$  effects in forward jets at HERA using BFKL and GBW parametrizations
  - using the parameters fit to the old data for  $d\sigma/dx$ :
    - we describe well the new data for  $d\sigma/dx$
    - we make predictions for  $d\sigma/dx dQ^2 dk^2$  (H1)
- Study to be completed with  $d\sigma/dQ^2$ ,  $d\sigma/d\eta_J$ ,  $d\sigma/dk^2$  (ZEUS)
- At small- $x$ 
  - NLO QCD is below the data
  - BFKL LL agrees wellThe saturation model is compatible with the data, (as it is for  $F_2$ )
- But two different saturation scales work
  - we need other observables to distinguish them**Other forward-jet cross-sections? Mueller-Navelet at LHC?**