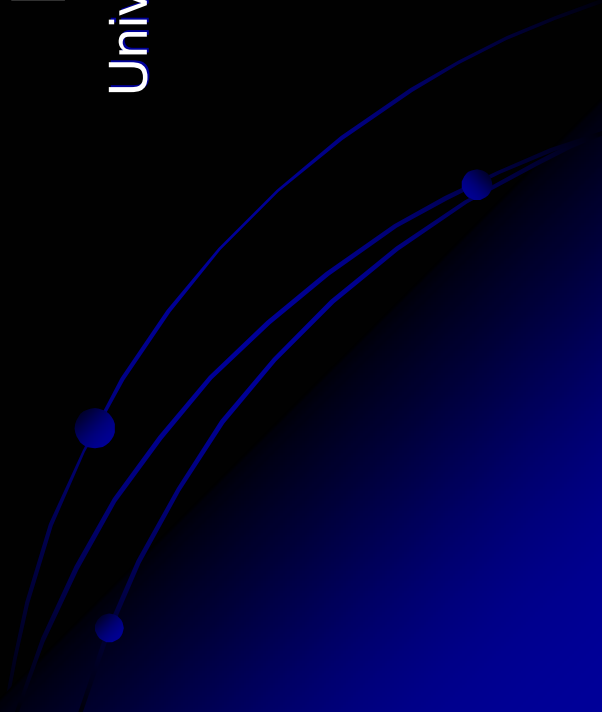
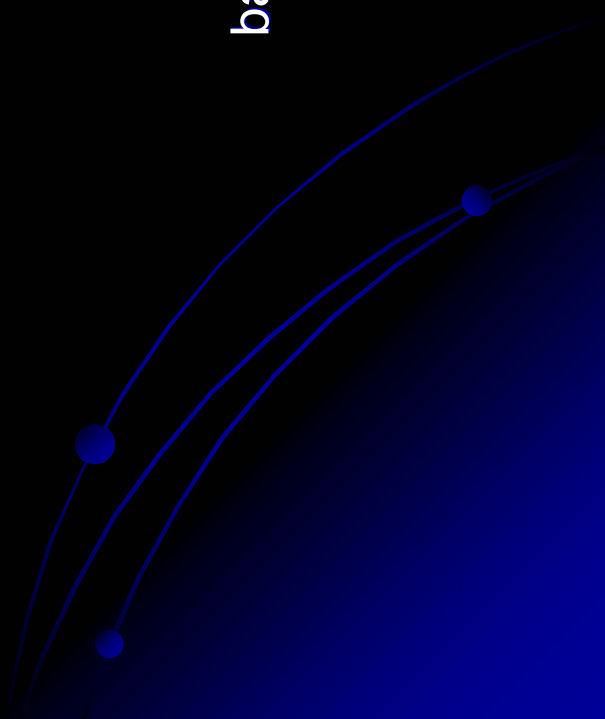


QCD corrections to the electroproduction of hadrons with high p_T

Rodolfo Sassot
Universidad de Buenos Aires

DIS05, Madison April 27, 2005





In collaboration with A. Daleo and D. de Florian,

based on Phys. Rev. D71 034013 (2005)
and references therein.

Outline:

- theory & phenomenological motivation & background

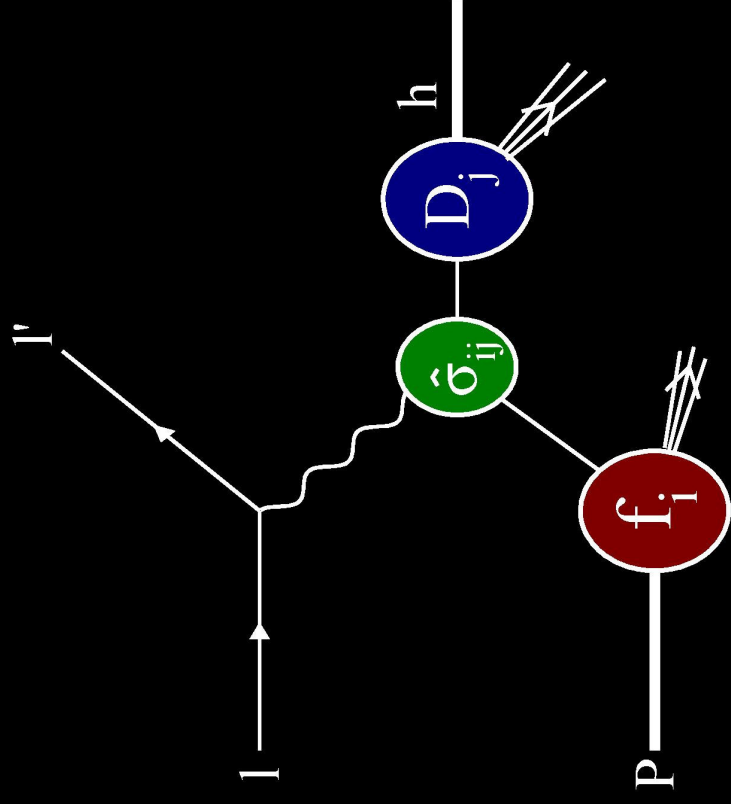
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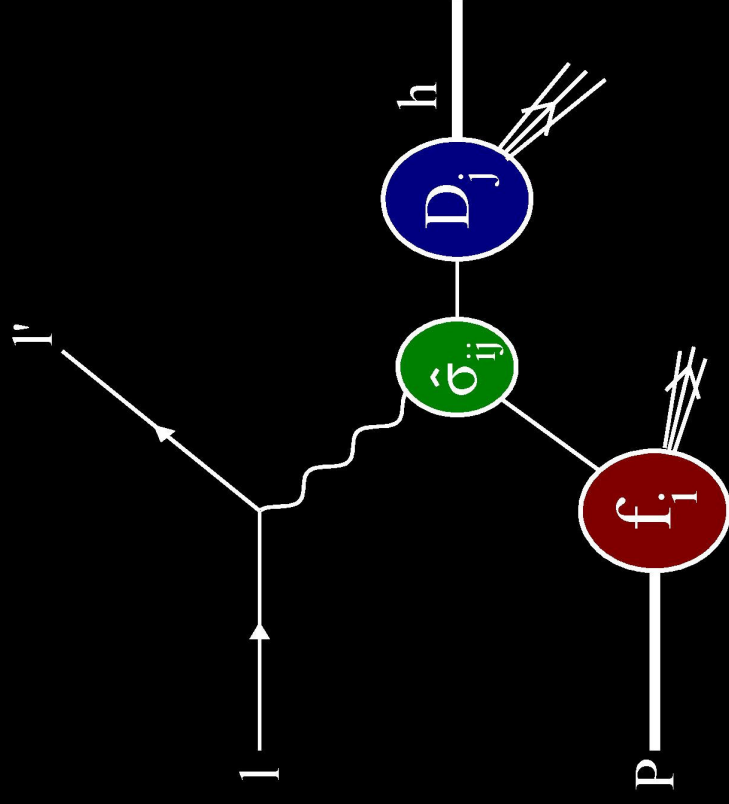
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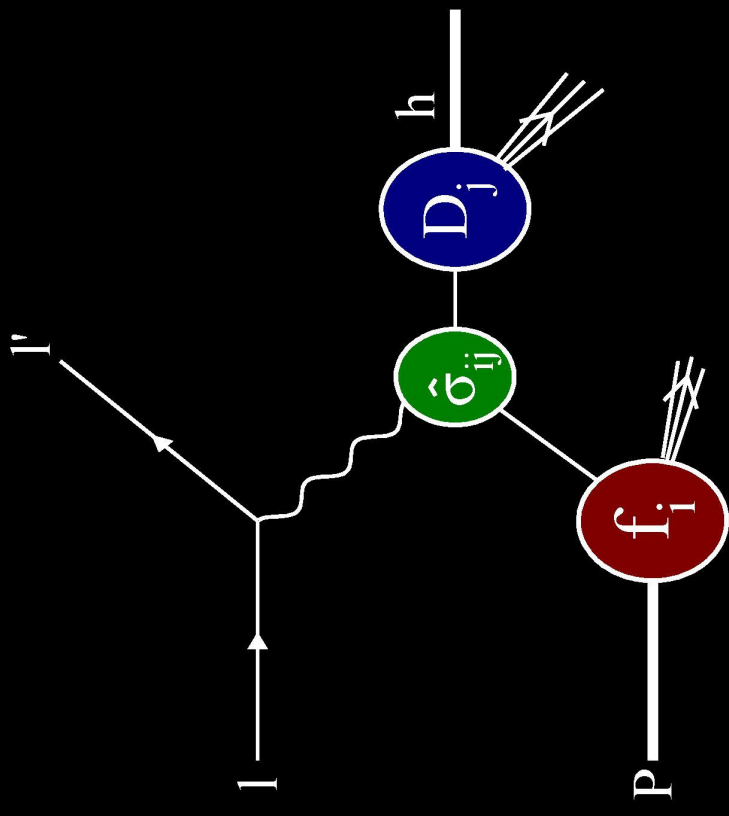
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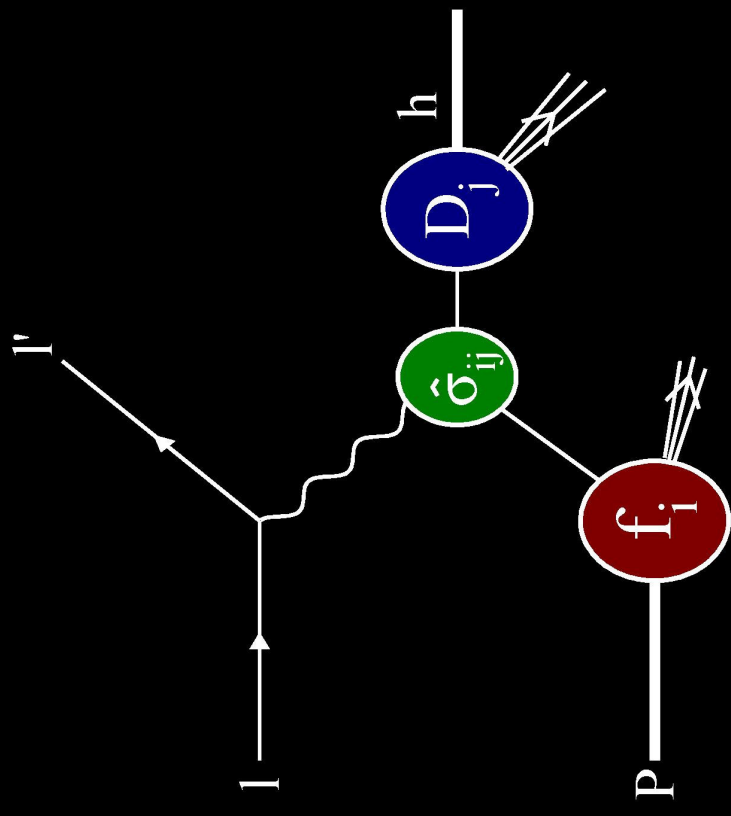
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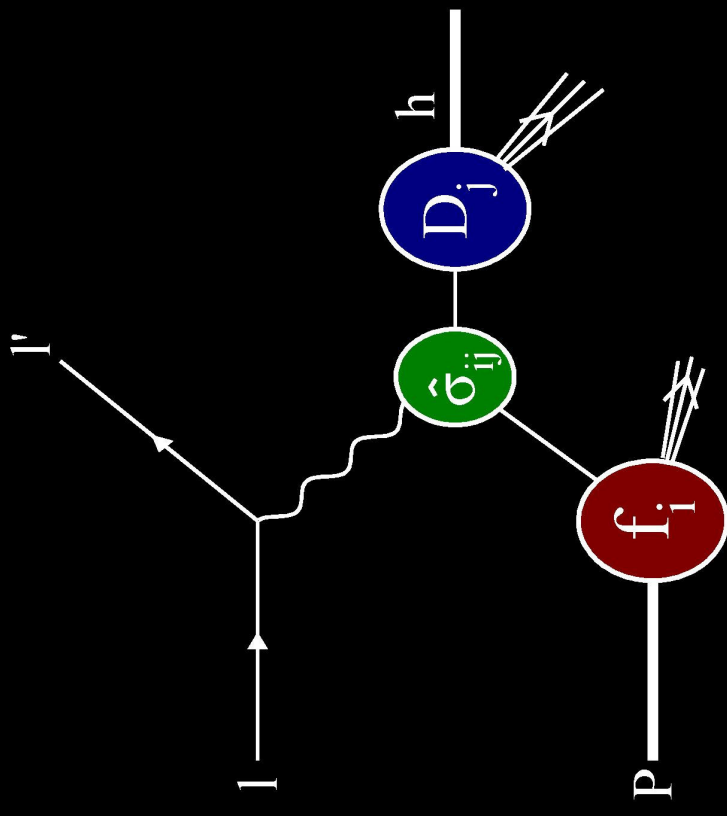
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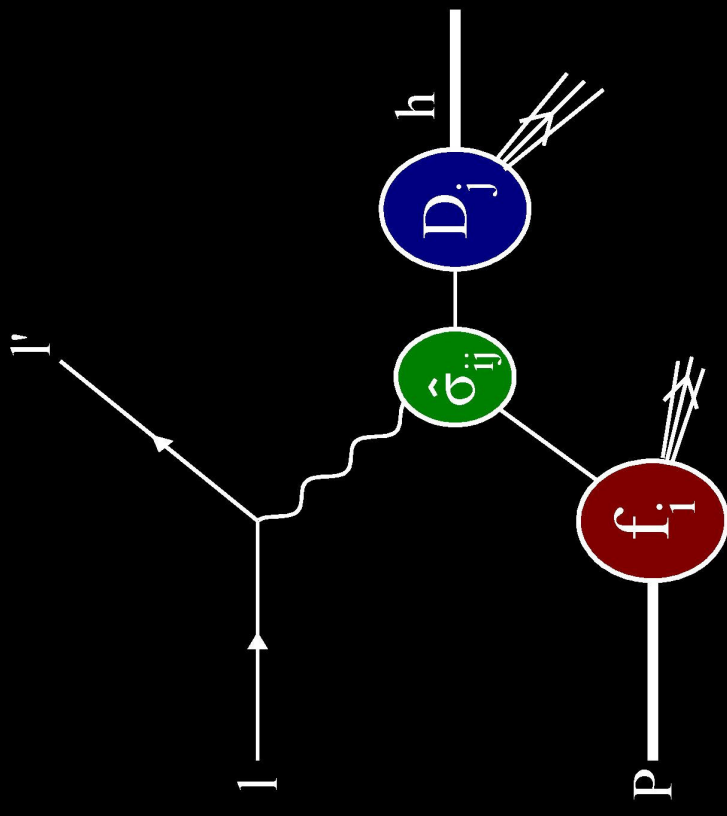
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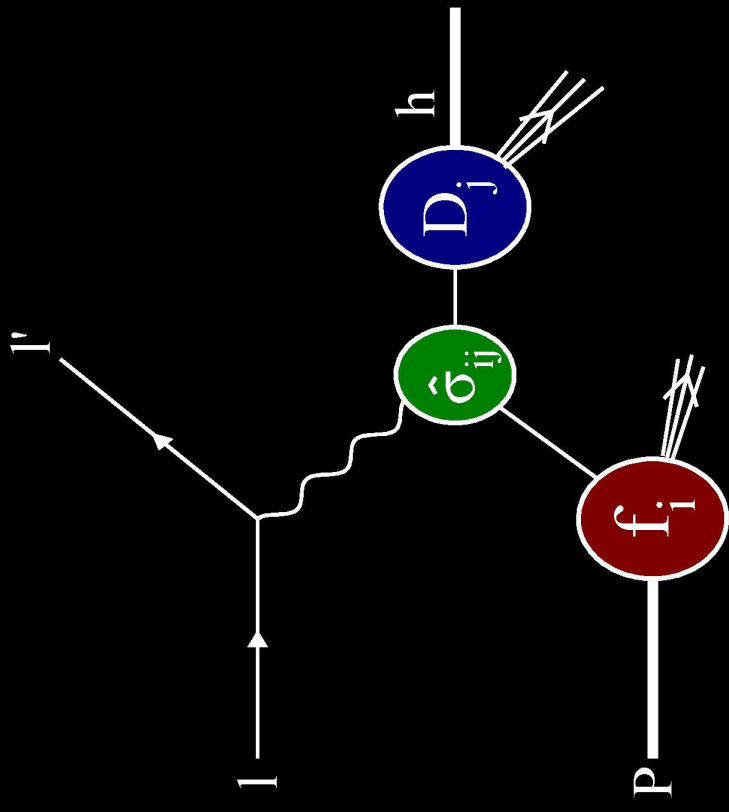
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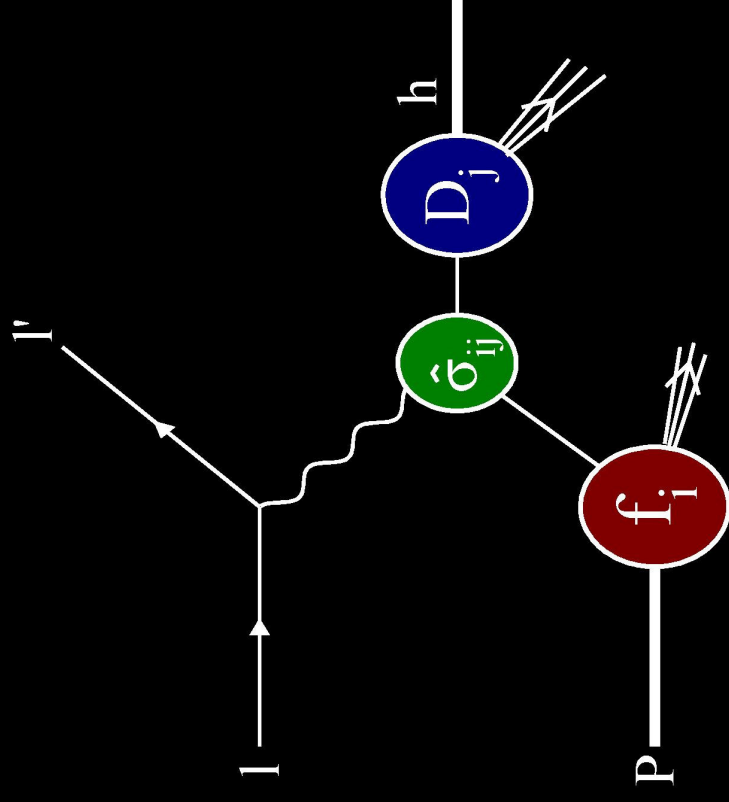
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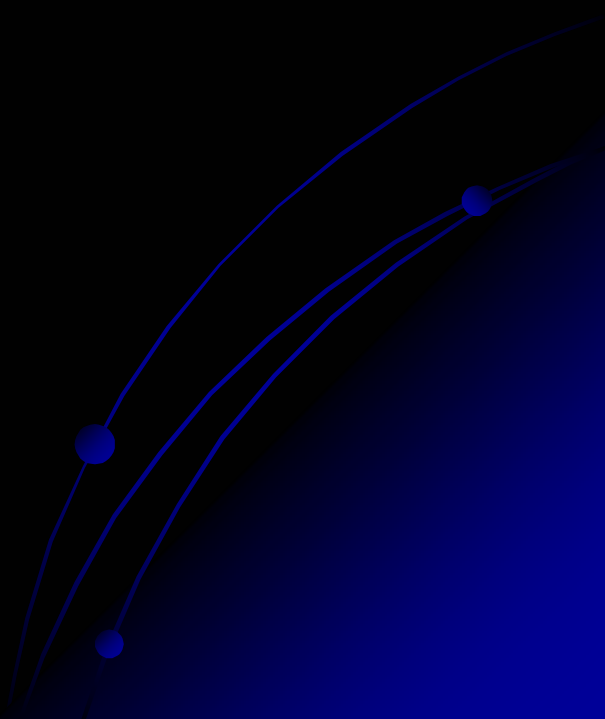
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- Analytic computation of higher order QCD corrections to OPI process have always lagged behind more inclusive observables: **singularity structure**.

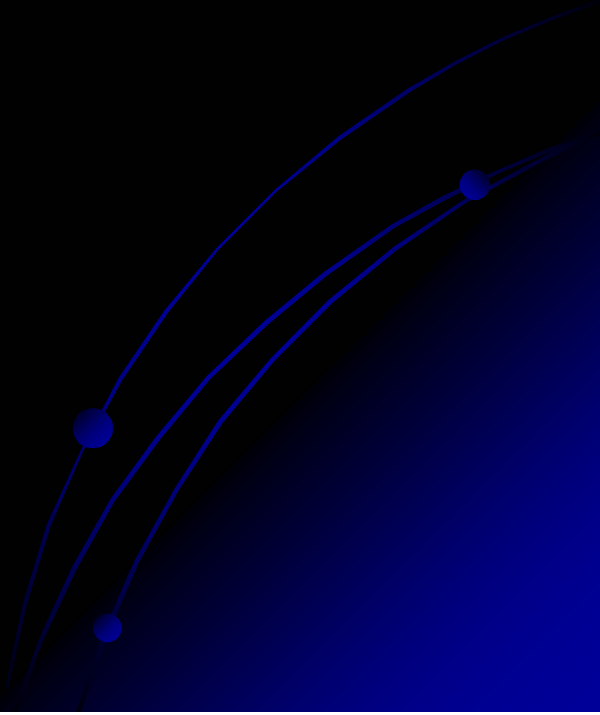


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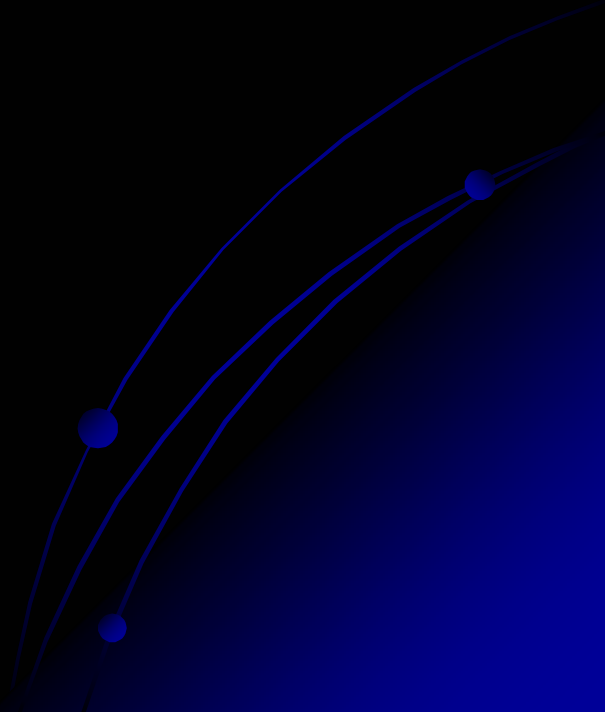


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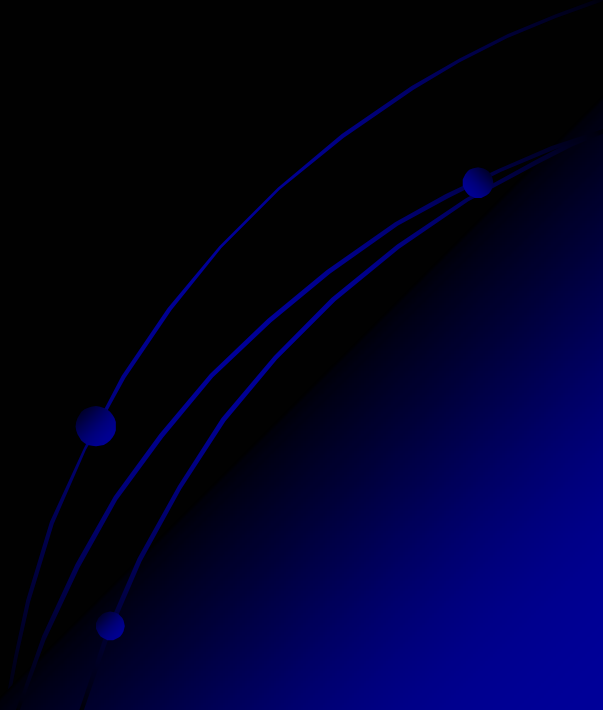
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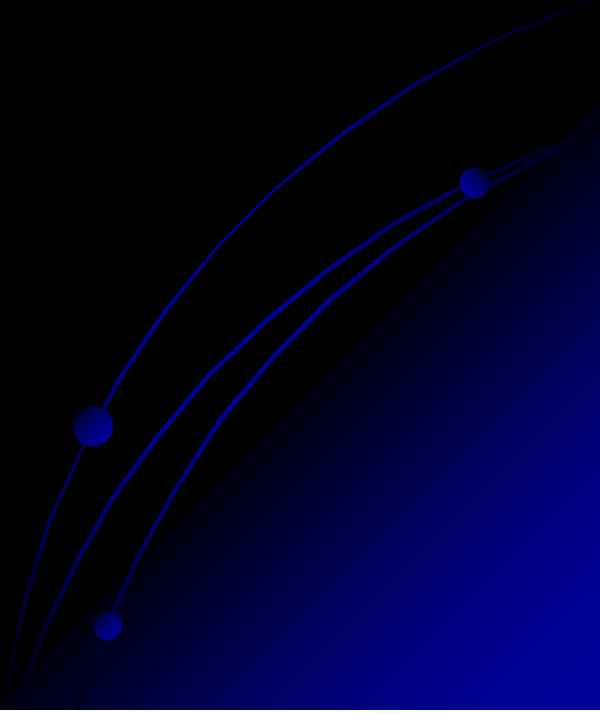
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- $\mathcal{O}(\alpha_s^1)$ approximation may not be good enough!

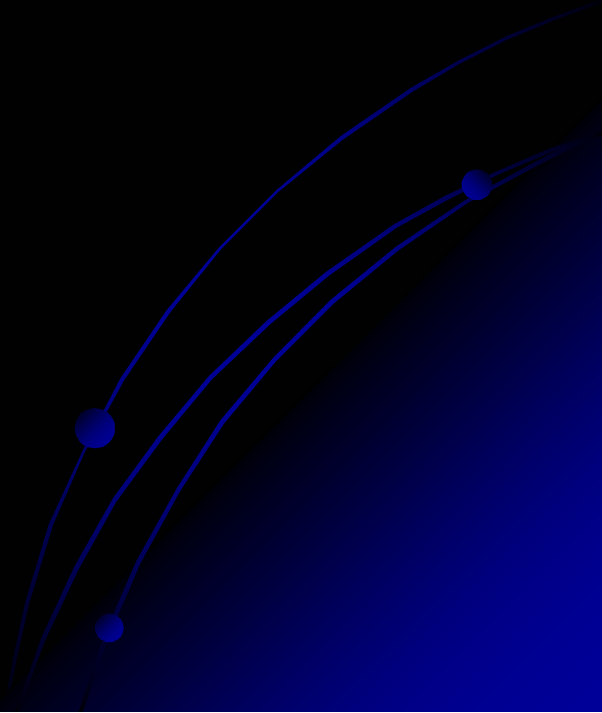
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$$s = \xi S - Q^2 (1 - \xi)$$

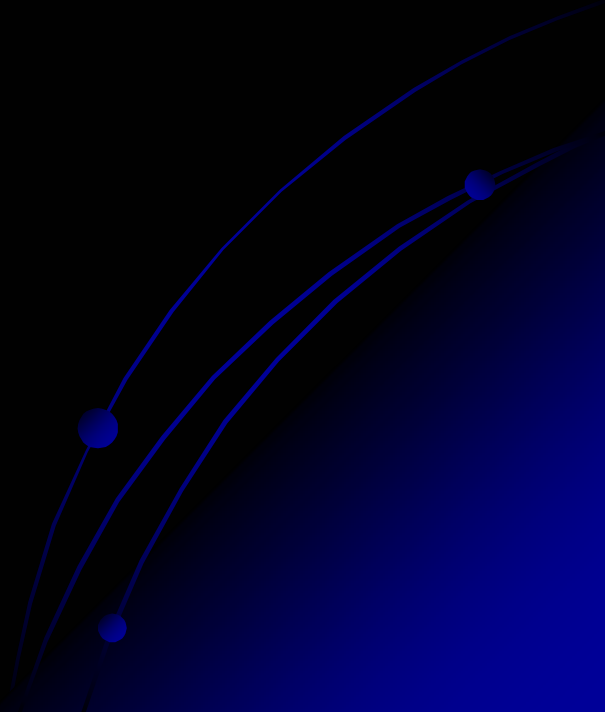
$$t = \frac{T}{\zeta} \qquad u = \frac{\xi}{\zeta} U$$

Kinematics II:

$$y \equiv -\frac{u}{Q^2+s}$$

$$z \equiv \frac{(Q^2+s)(s+t+u)}{s(Q^2+s+u)}$$

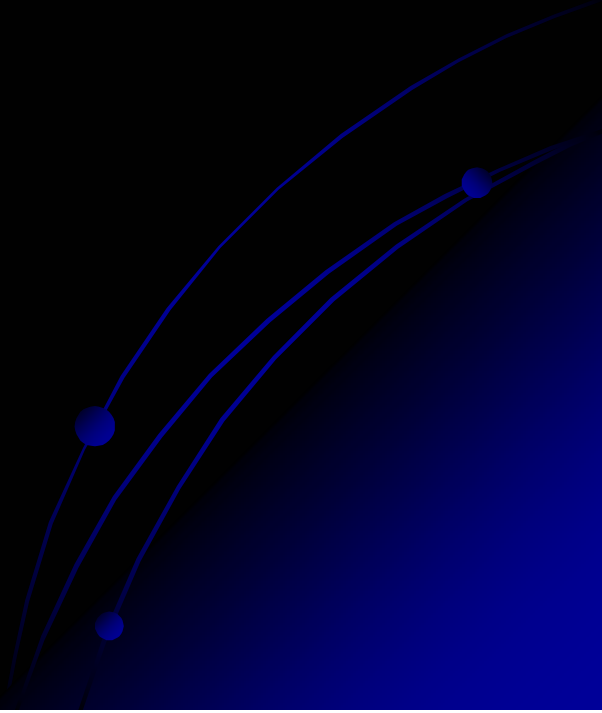
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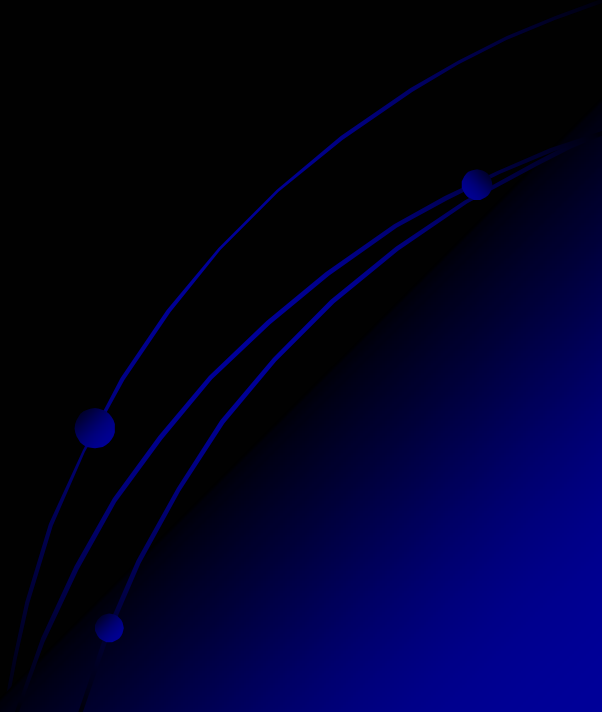
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$$\text{IR singularities} \begin{cases} y = 0, 1 \\ z = 0, 1 \end{cases} \quad (1)$$

if $|p_T| > 0$ only $z = 0$

Kinematics IV:

$$\frac{d\sigma_{ij}^{(2)}}{dx_B dQ^2 dy dz} = \left\{ \frac{1}{\epsilon} \mathcal{P}_{1ij}^{(2)}(e, y, z) + C_{ij}^{(2)}(e, y, z) + \mathcal{O}(\epsilon) \right\}^{(1)}$$

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$C_{ij}^{(2)}(e, y, z)$ coefficients are finite but rather long because of $\mathcal{O}(\alpha_s^2)$

NLO: $\mathcal{O}(\alpha_s^2)$

Real contributions

$$\left\{ \begin{array}{ll} \gamma + q(\bar{q}) & \rightarrow g + g + q(\bar{q}) \\ \gamma + q_i(\bar{q}_i) & \rightarrow q_i(\bar{q}_i) + q_j + \bar{q}_j \quad (i \neq j) \\ \gamma + q_i(\bar{q}_i) & \rightarrow q_i(\bar{q}_i) + q_i + \bar{q}_i \\ \gamma + g & \rightarrow g + q + \bar{q} \end{array} \right.$$

Virtual contributions

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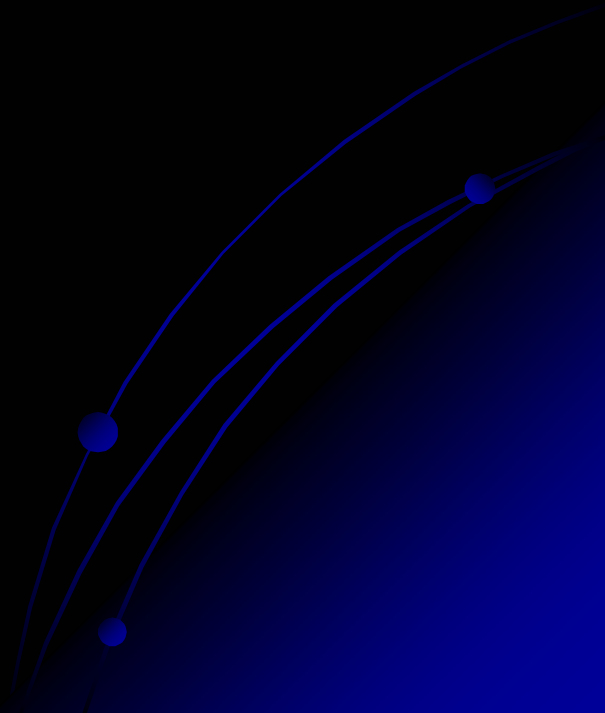
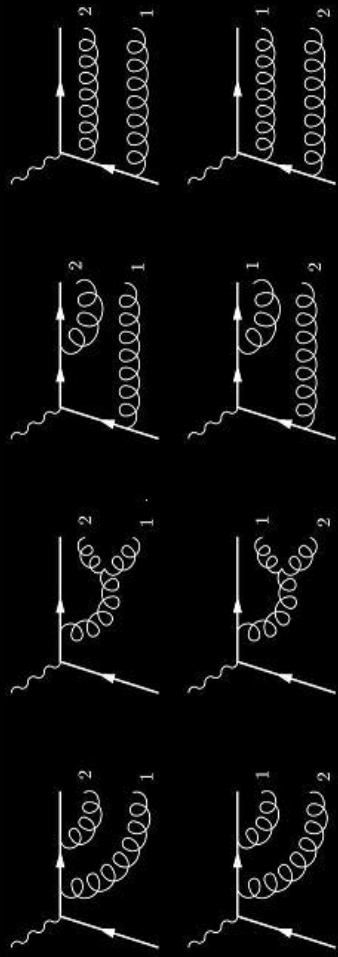
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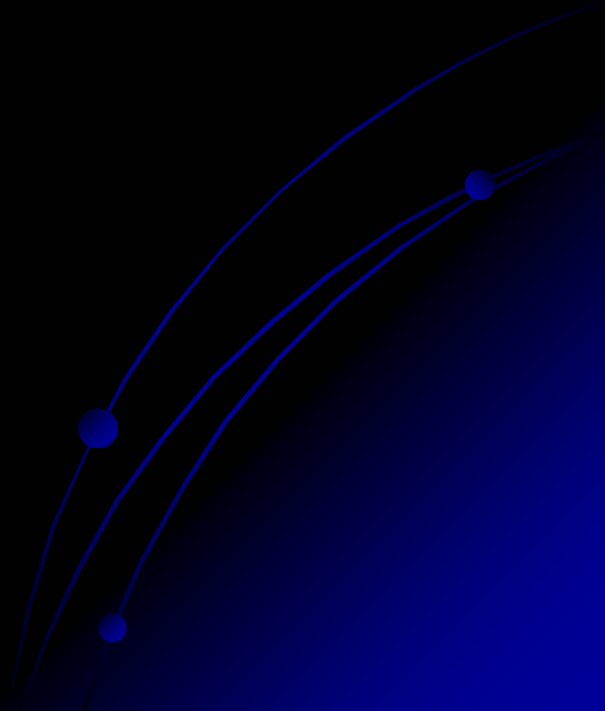
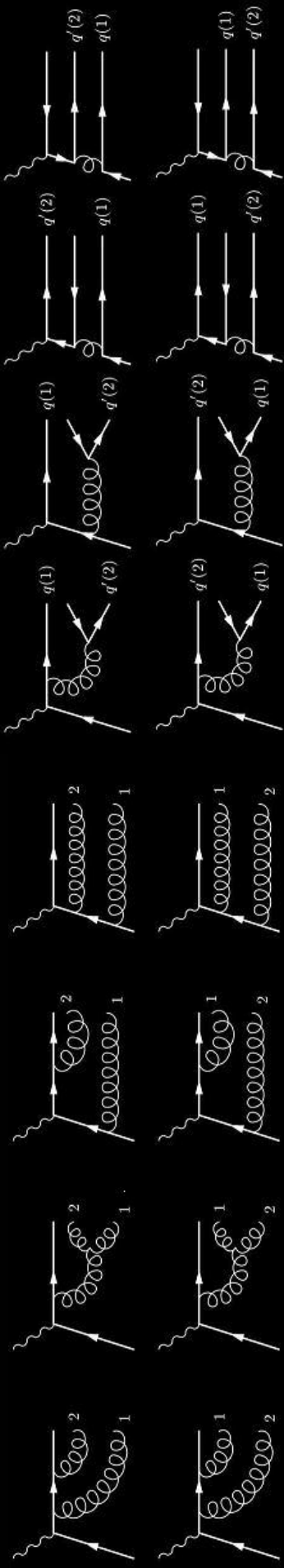
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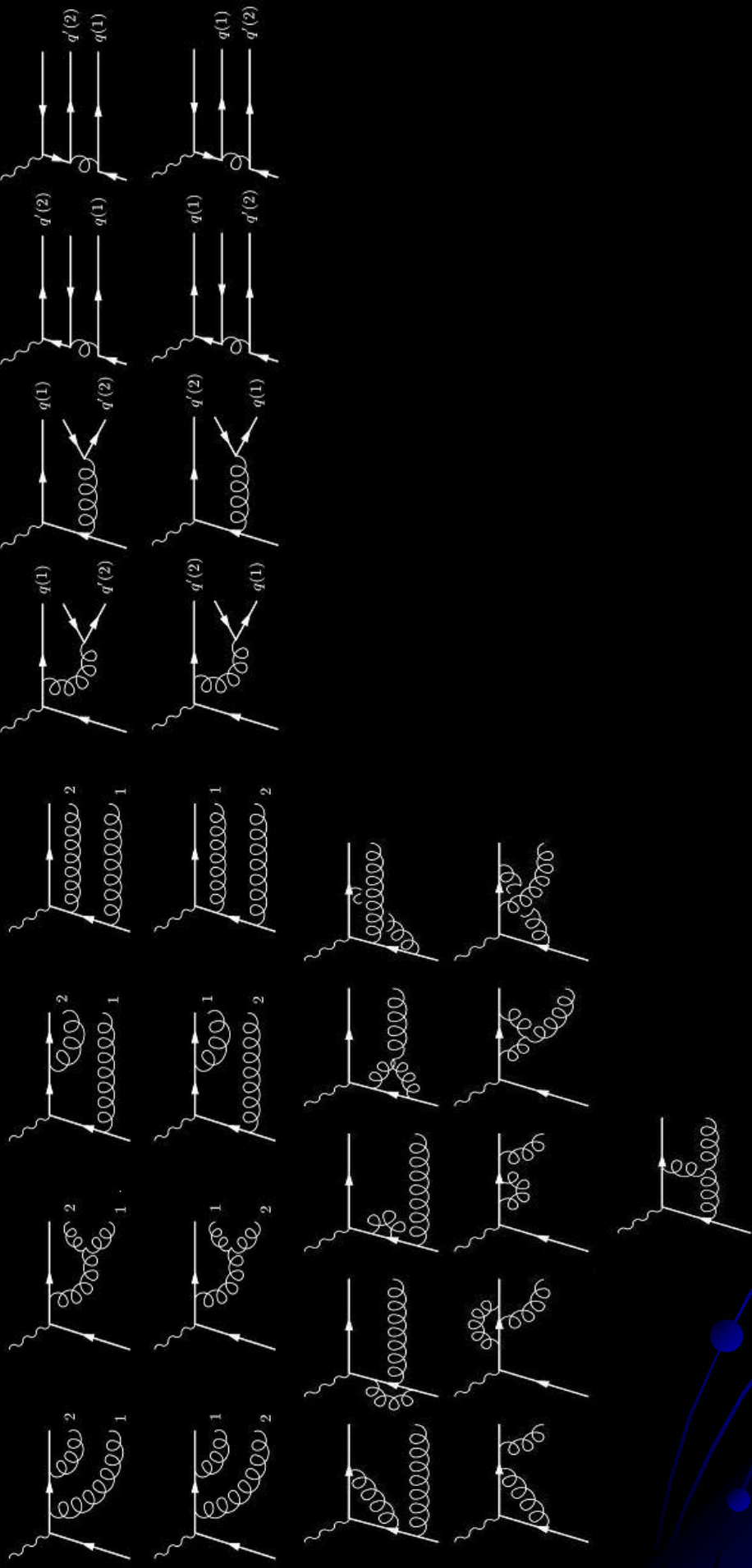
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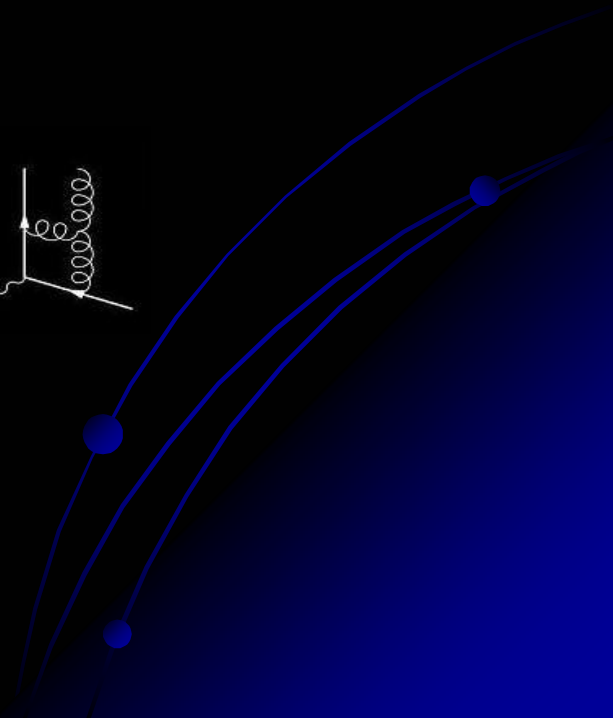
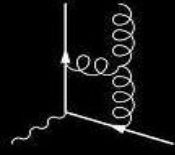
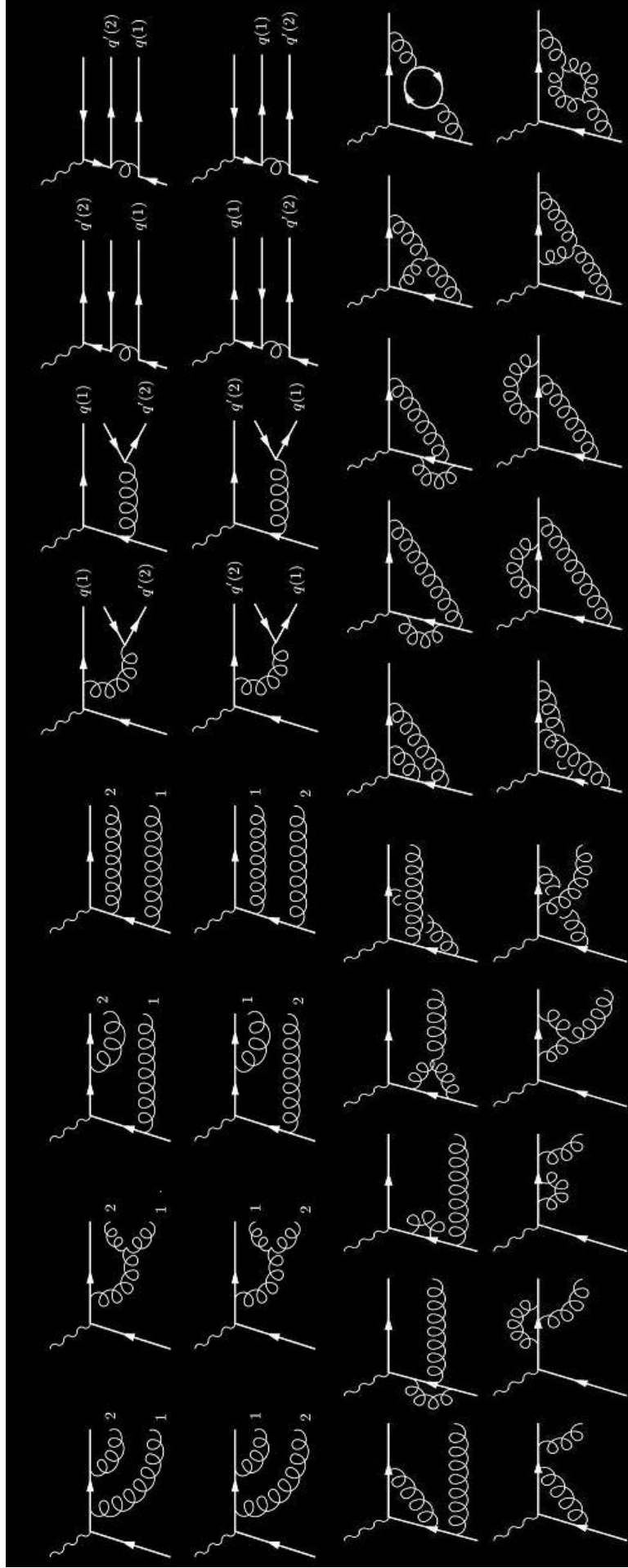
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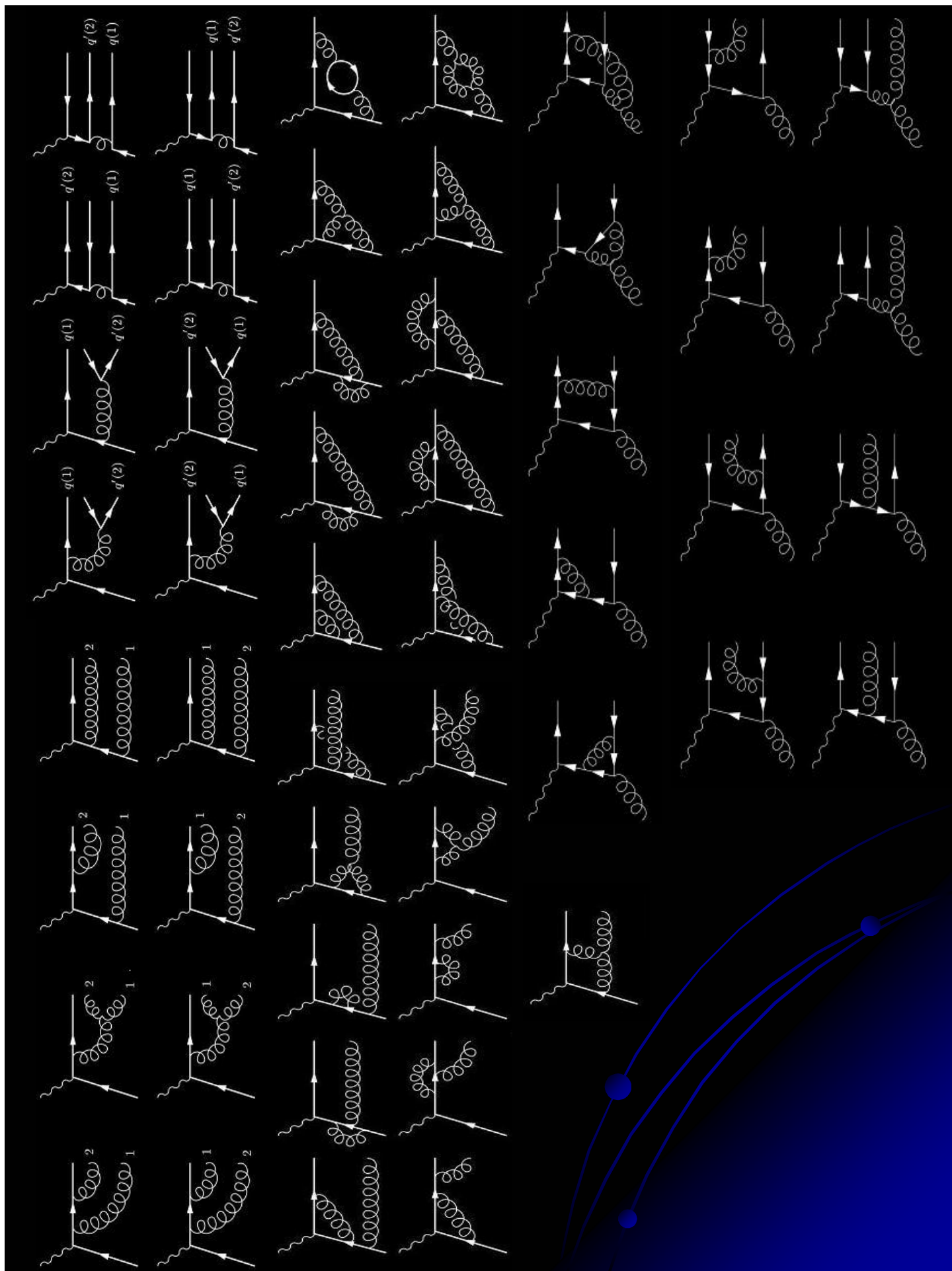
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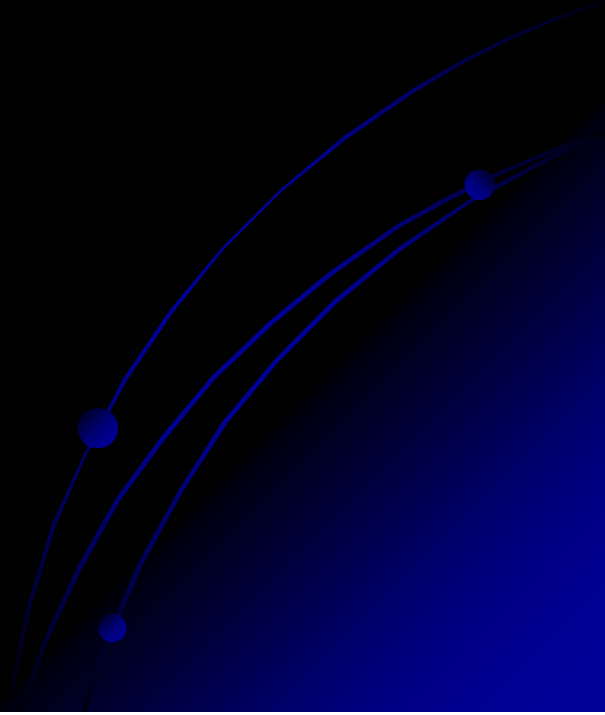


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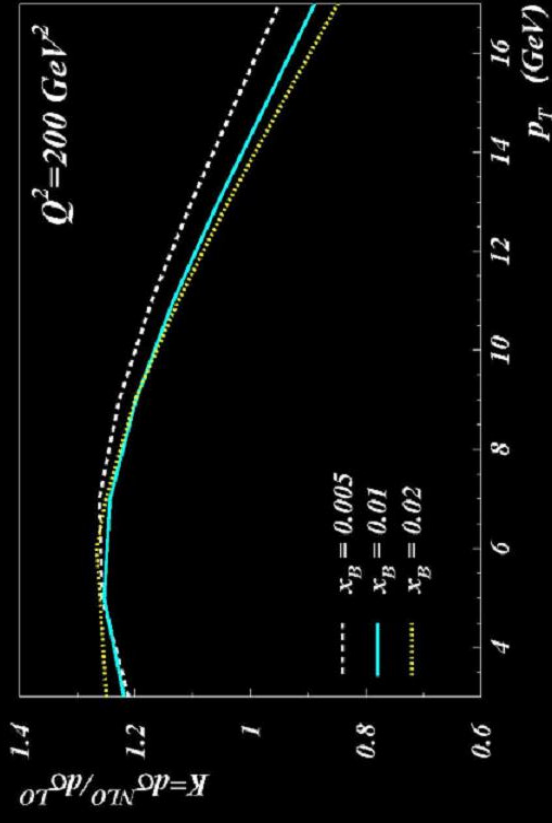
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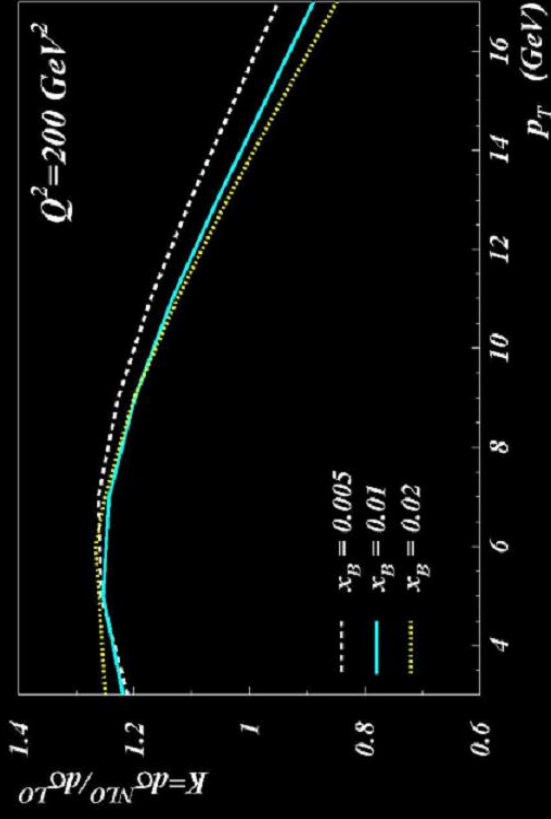


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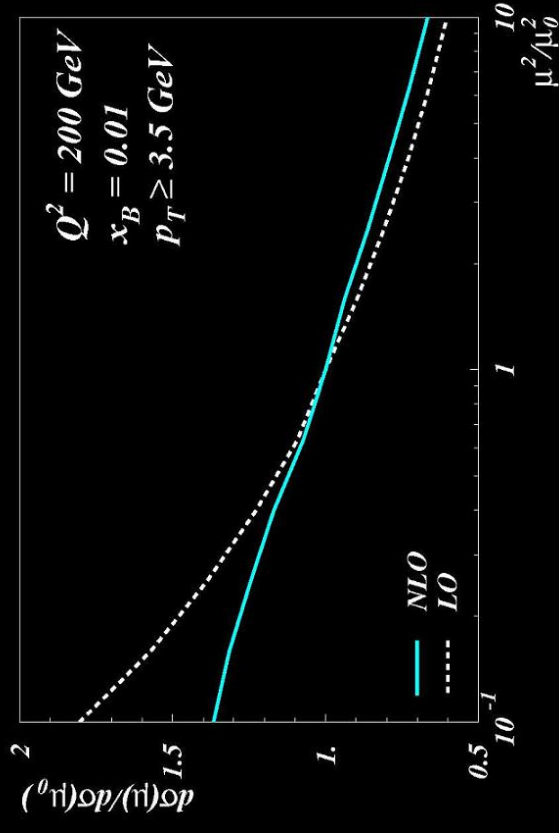
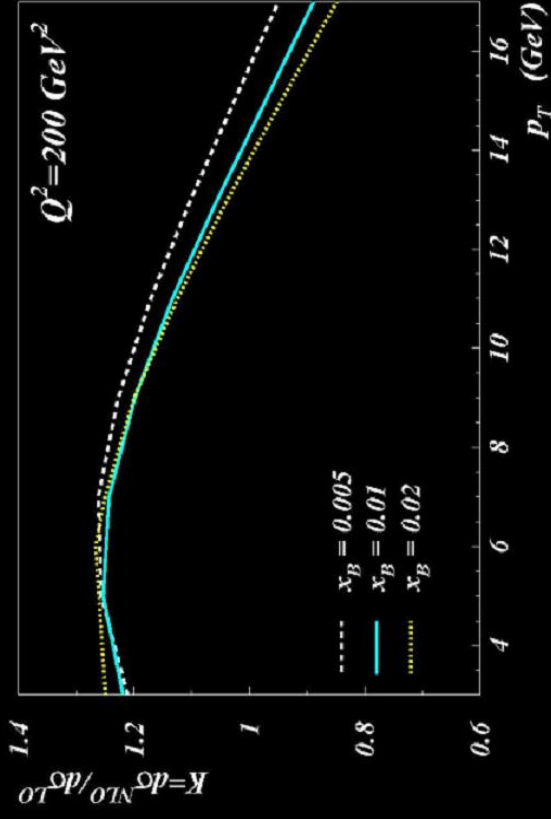


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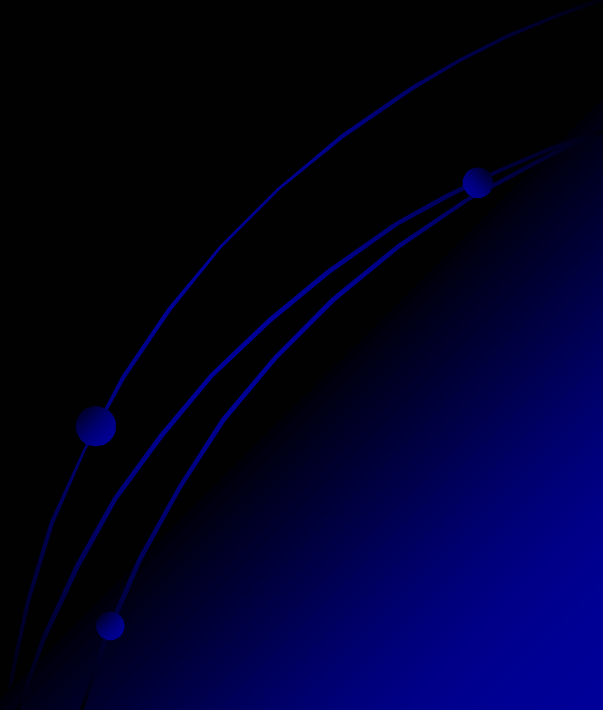
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$2.5 < p_T < 15 \text{ GeV}$

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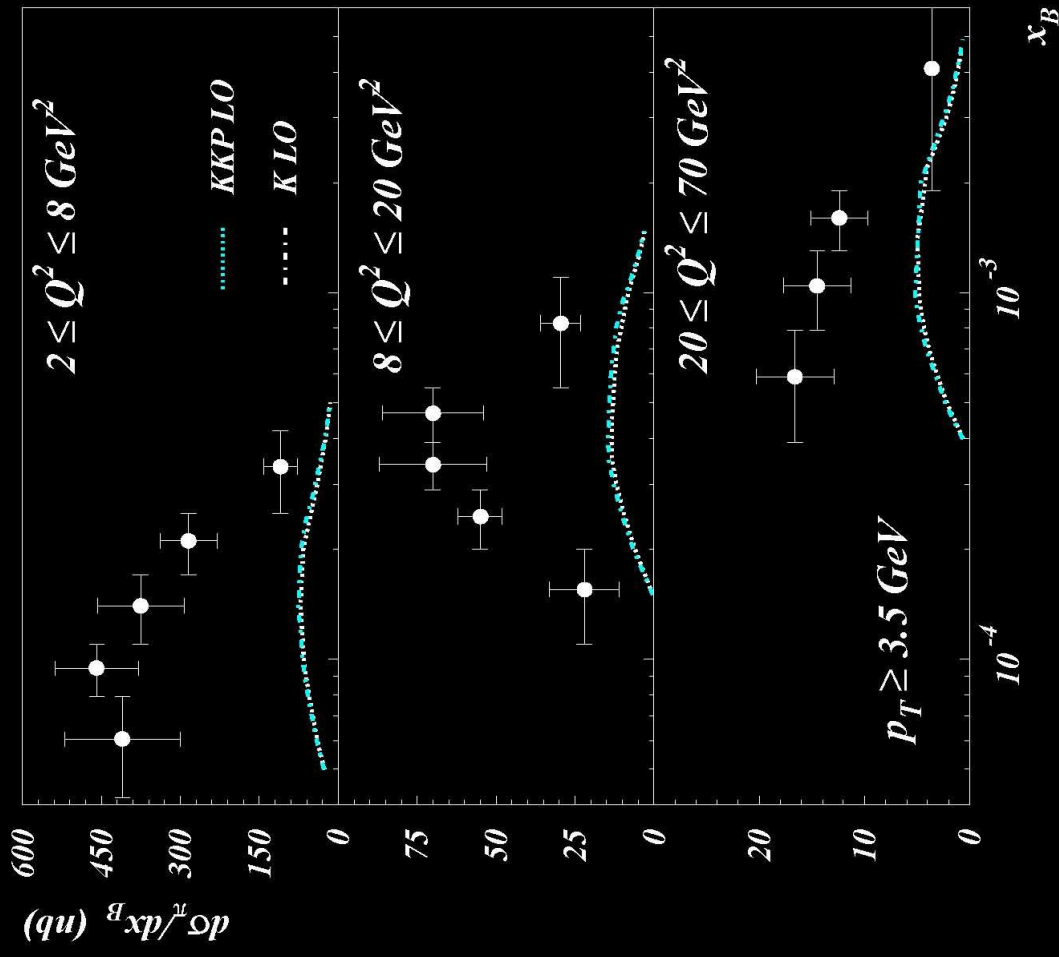
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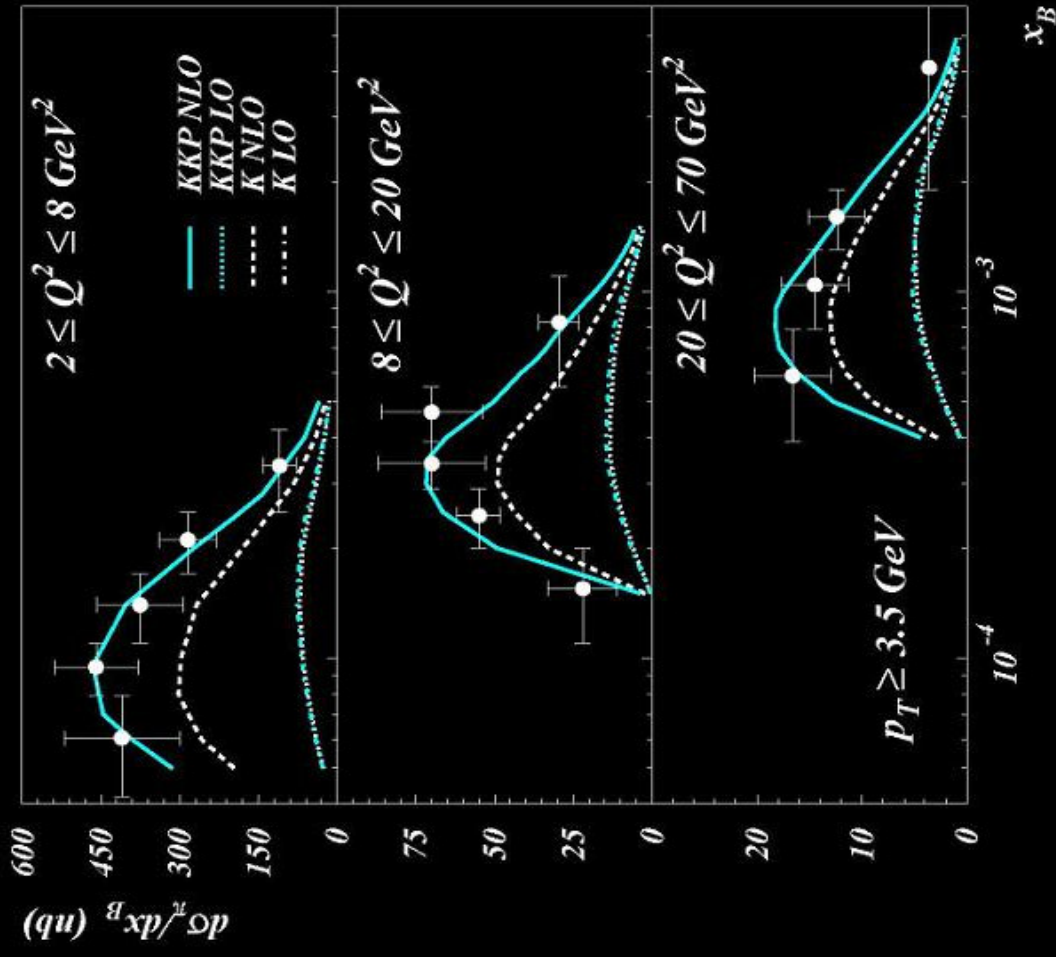
$e^+ @ 27.6 \text{ GeV}, p @ 820 \text{ GeV}$

$2.5 < p_T < 15 \text{ GeV}$

$\theta_\pi \in [5^\circ, 25^\circ]$

$x_\pi = E_\pi / E_p > 0.01$

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'KKP' B. A. Kniehl et al. Nucl.Phys.B582 (2000) 514.



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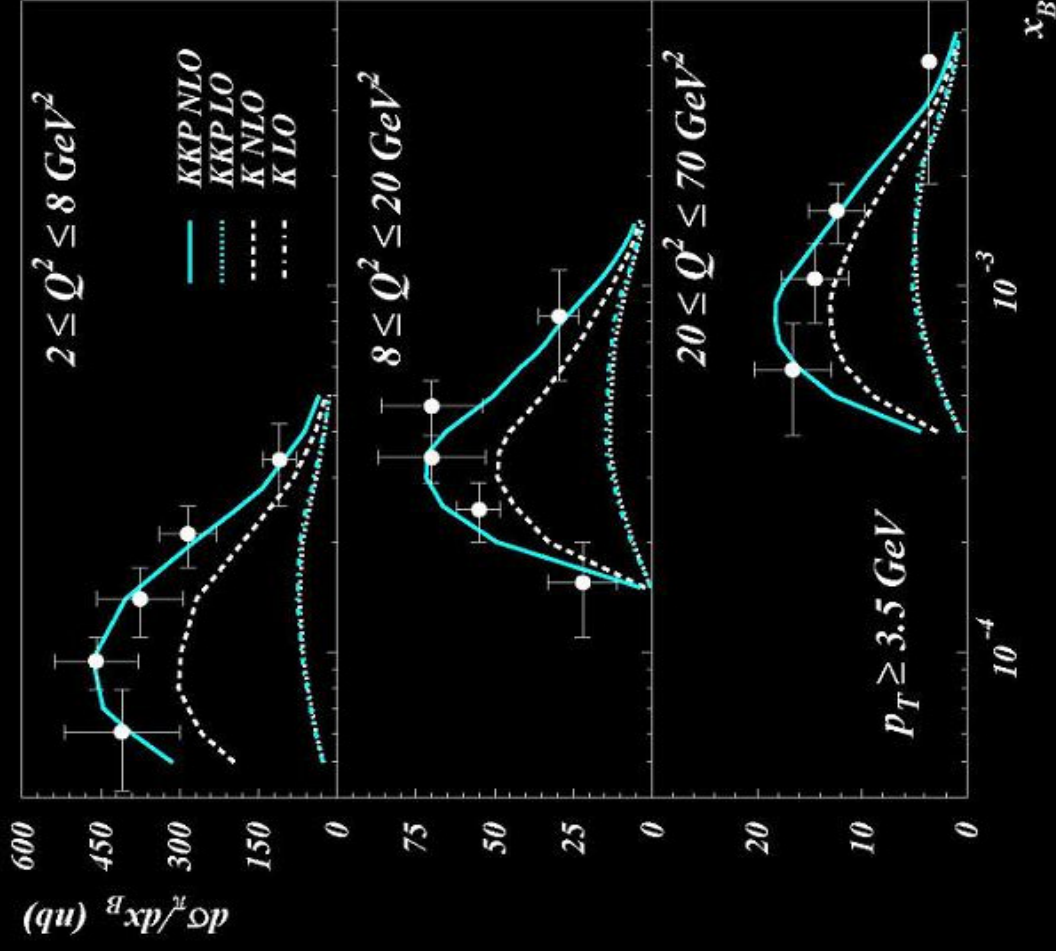
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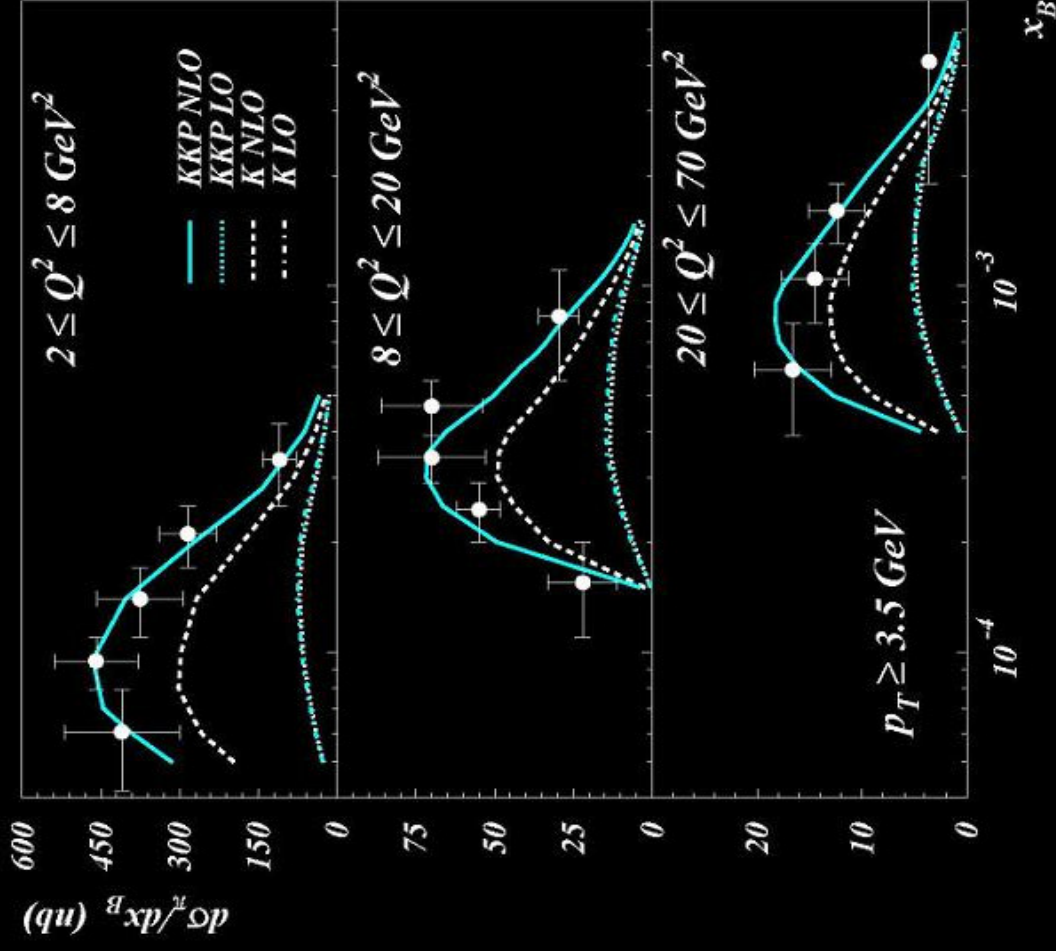
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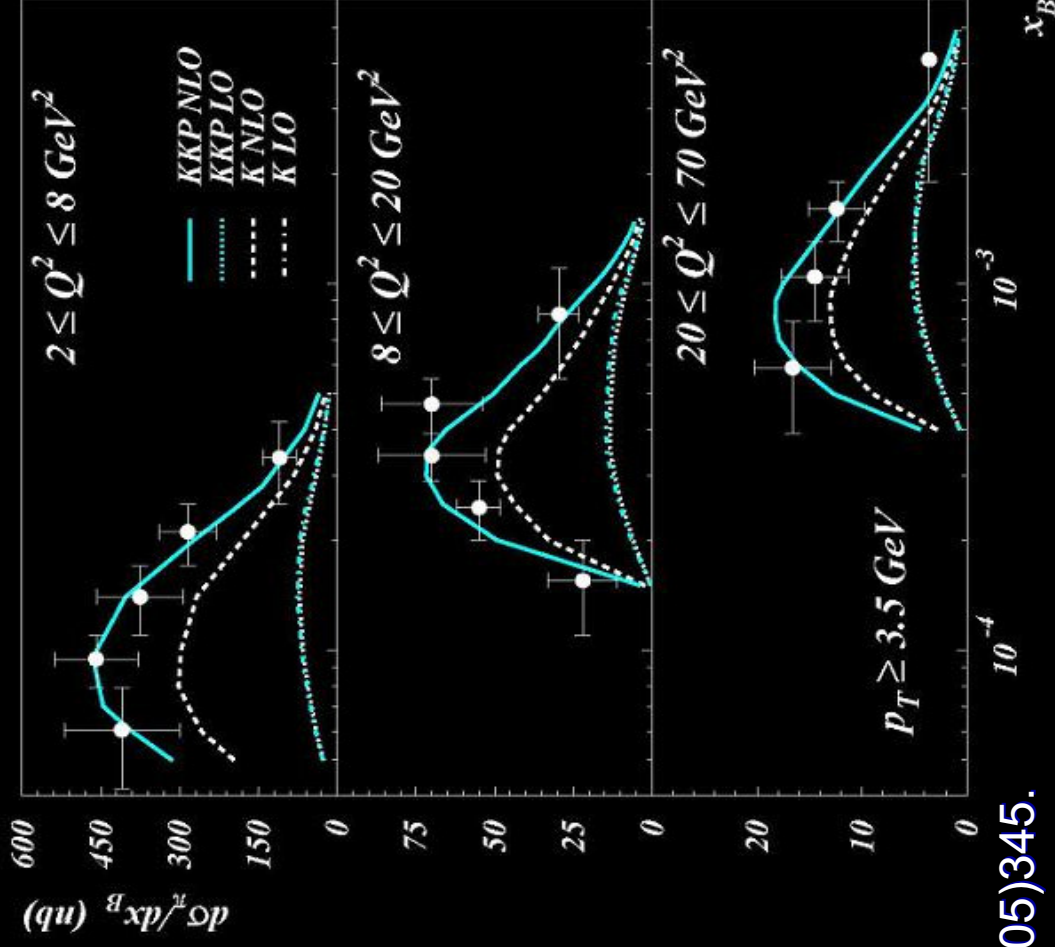
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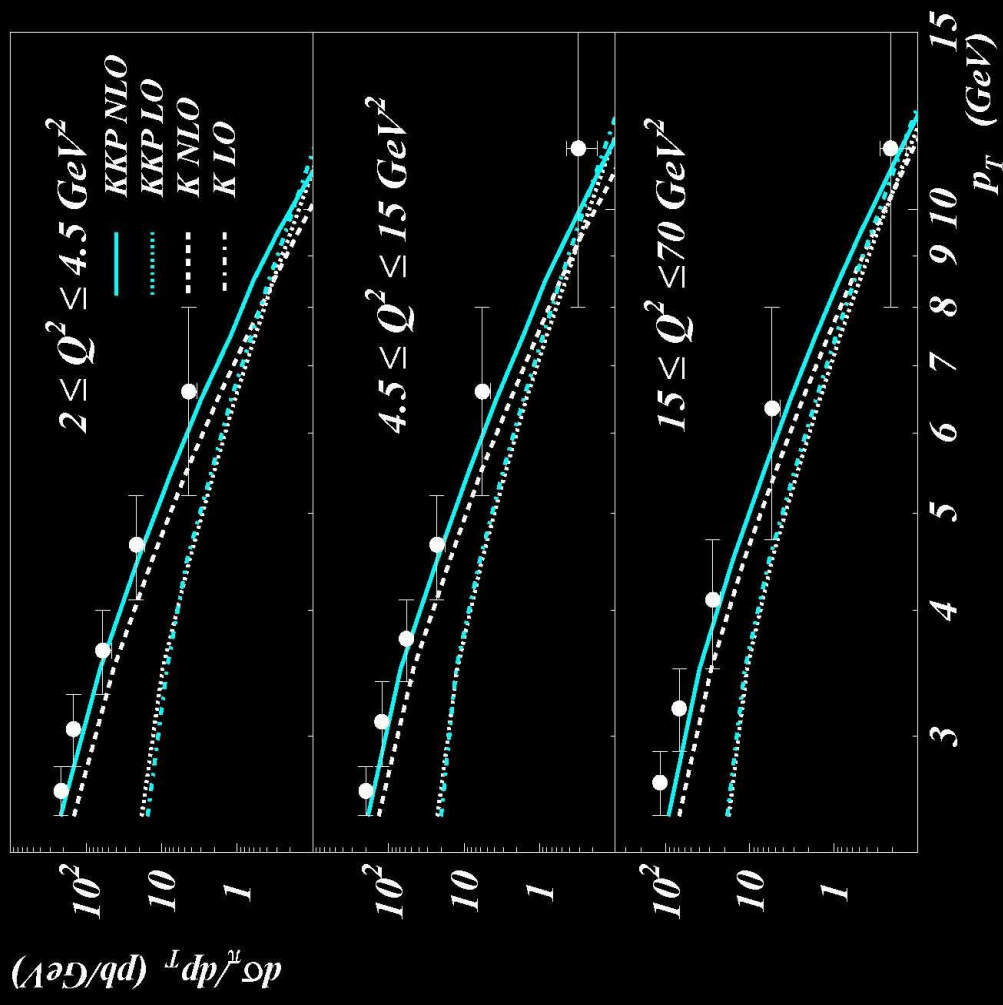
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Similar conclusions by:

B. Kniehl et al. Nucl.Phys.B711 (2005)345.

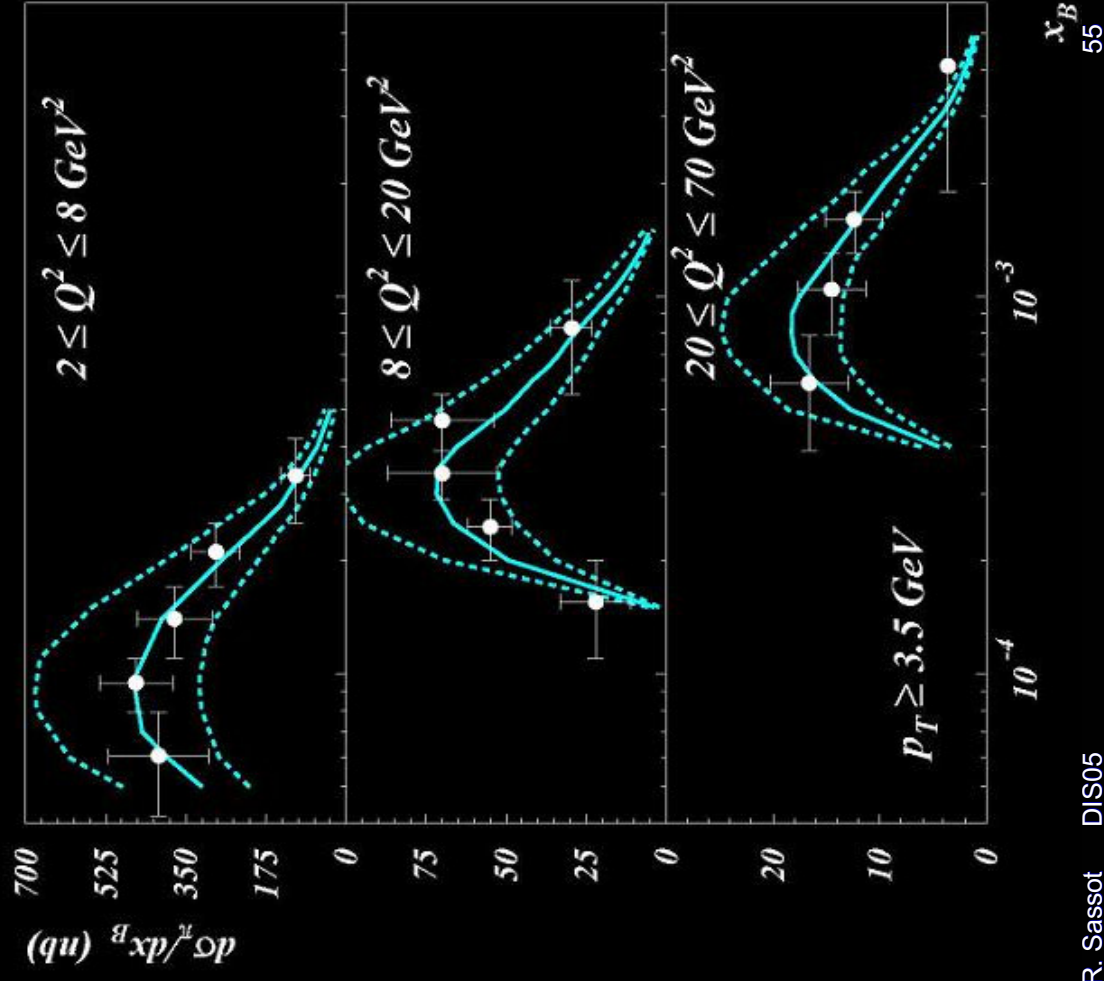


Phenomenology III-



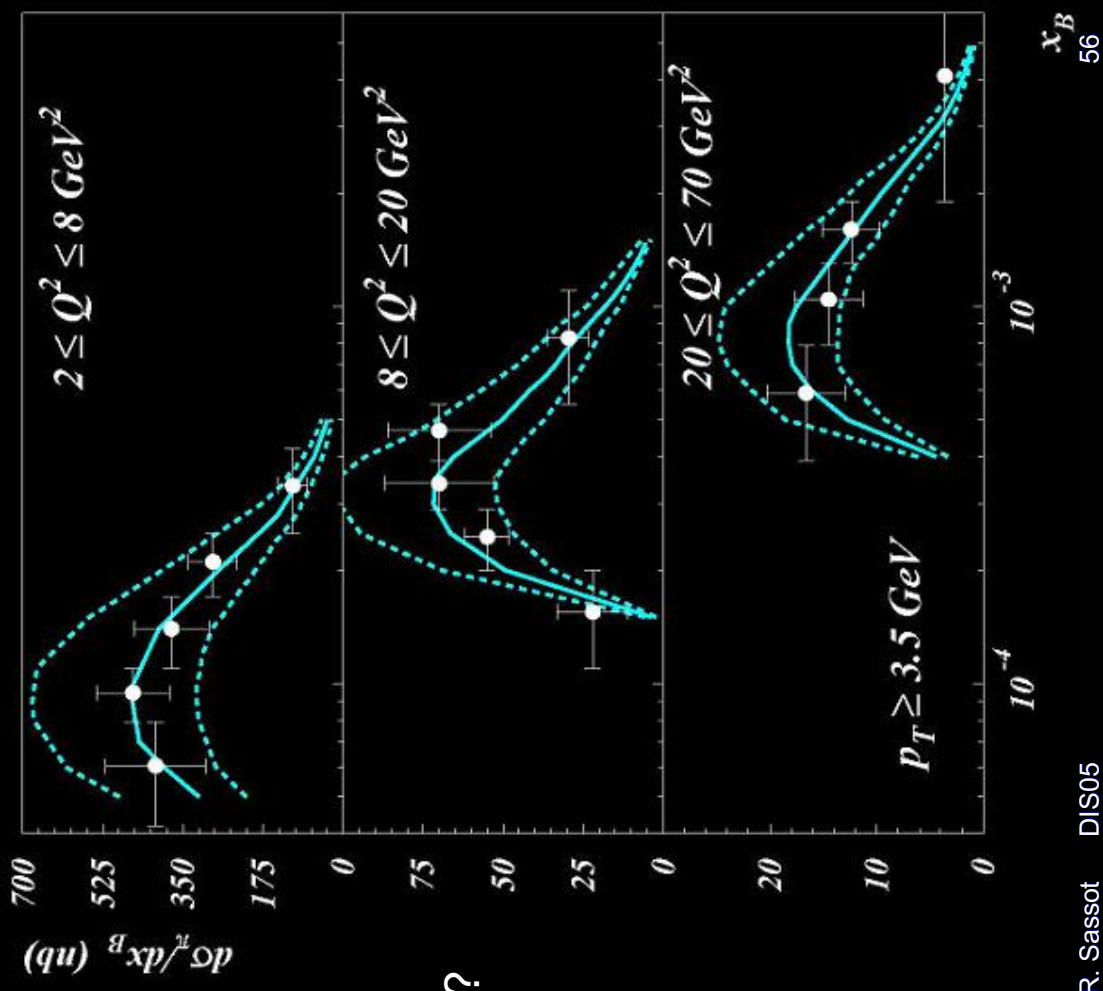
Phenomenology IV:

- factorization scale dependence



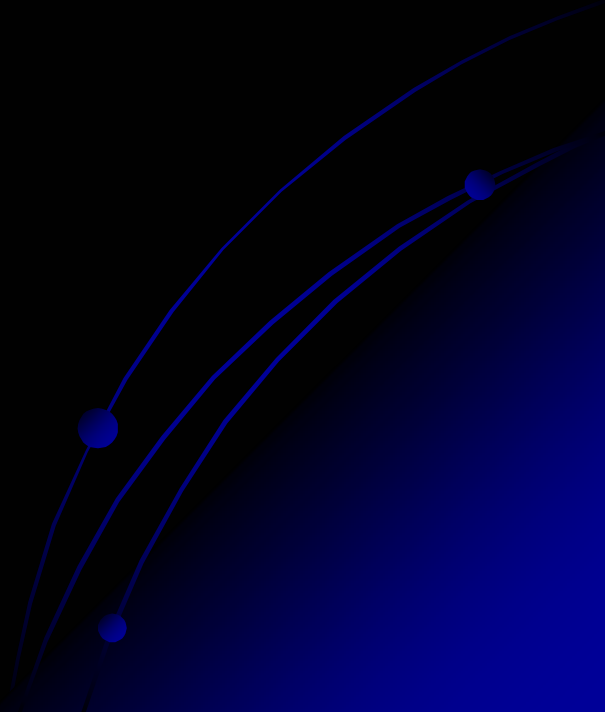
Phenomenology IV:

- factorization scale dependence
- significant corrections at next order?
NNLO



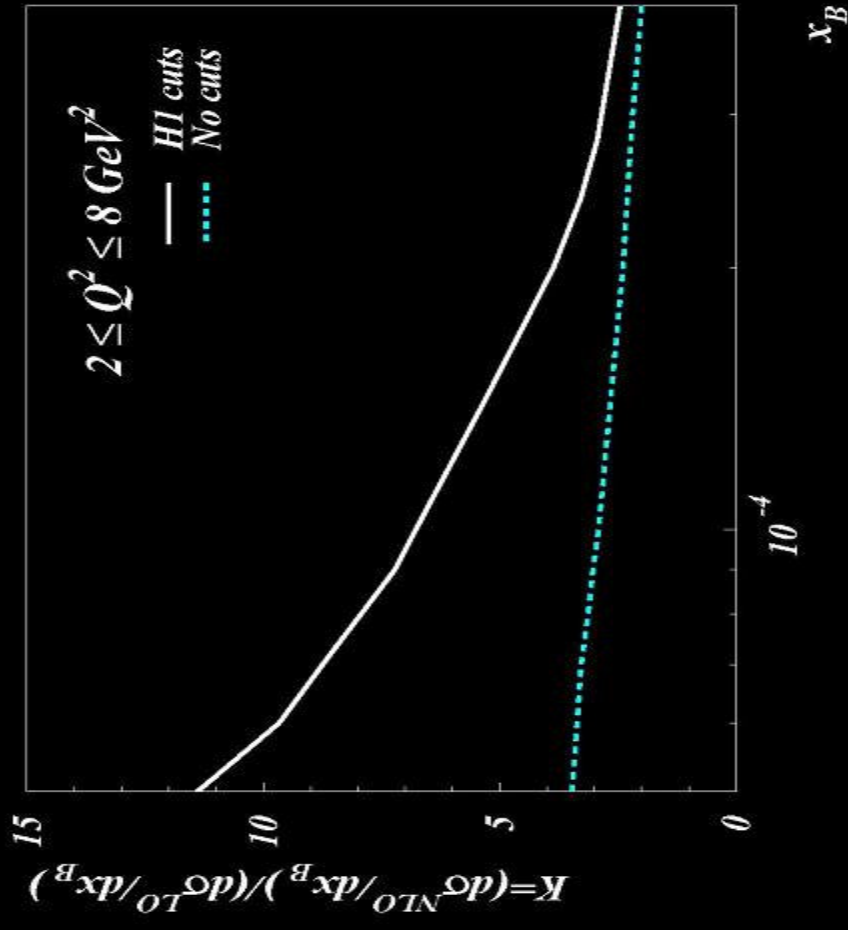
Phenomenology V:

- why so large??



Phenomenology V:

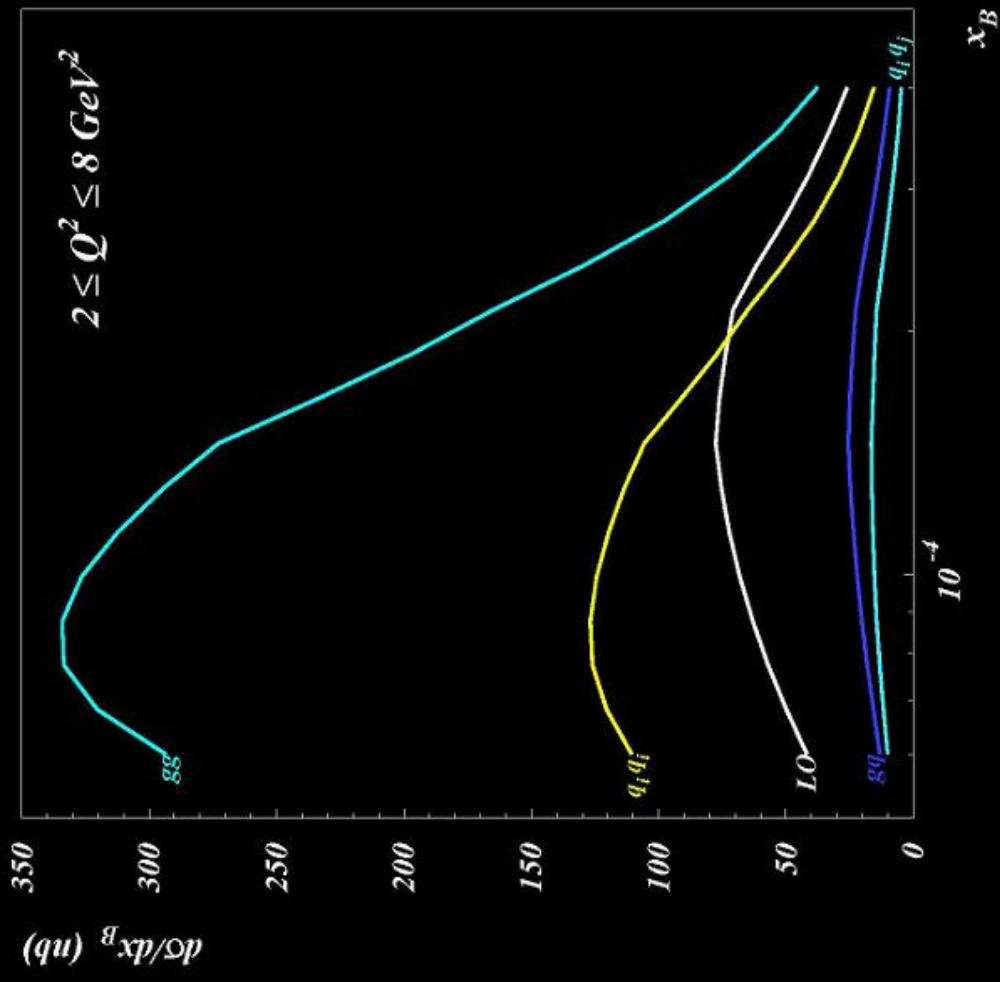
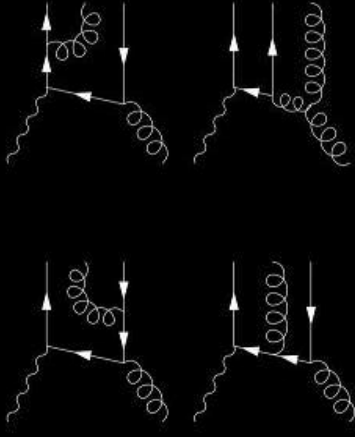
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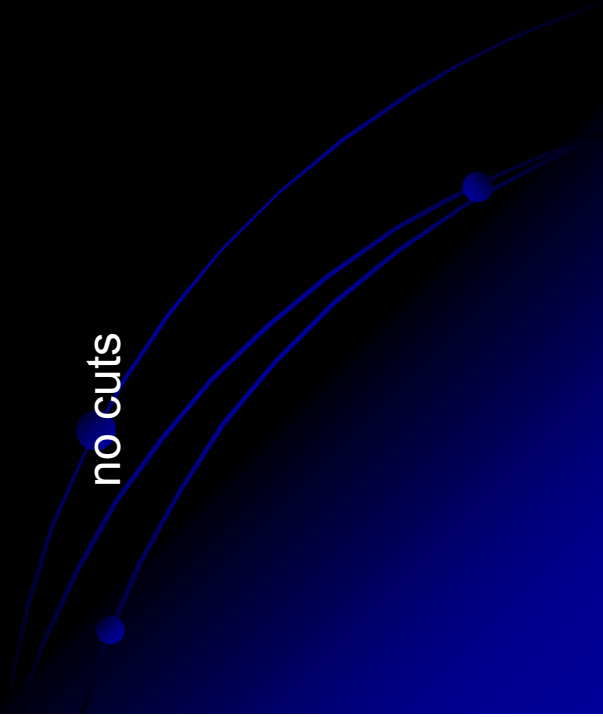
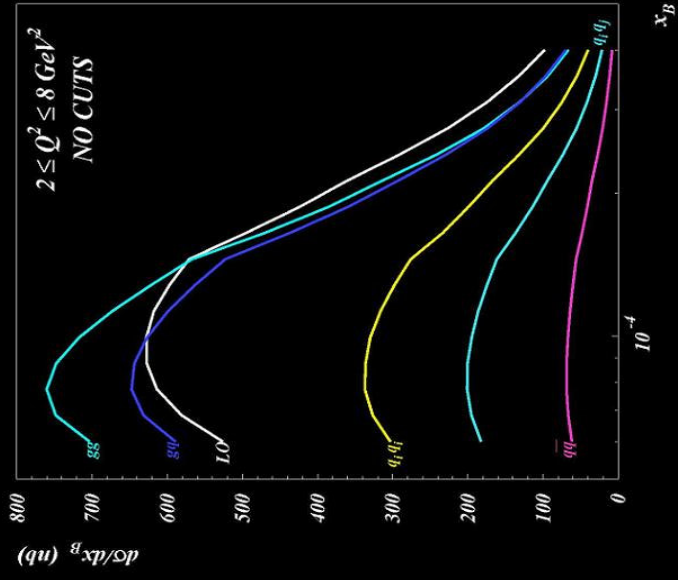


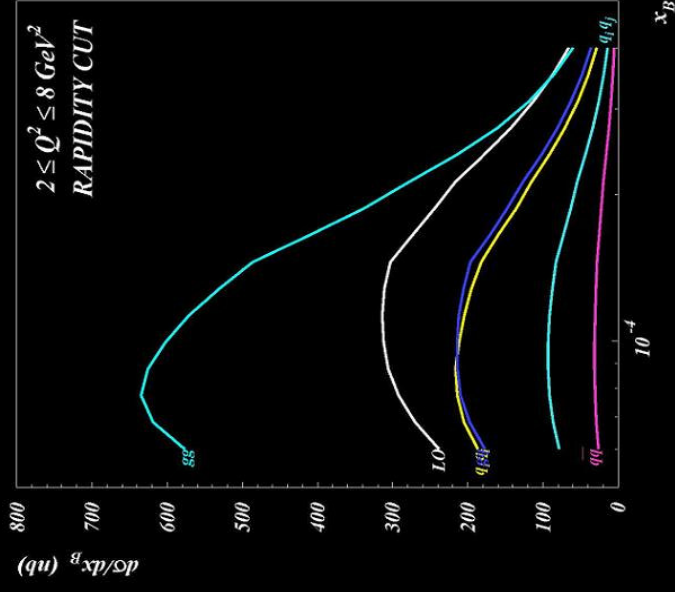
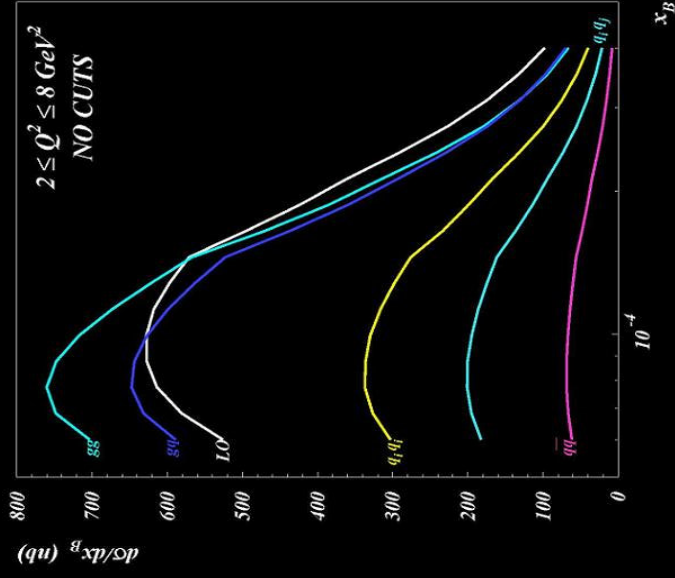
- effects from experimental cuts

Phenomenology VI:

- “new” channel

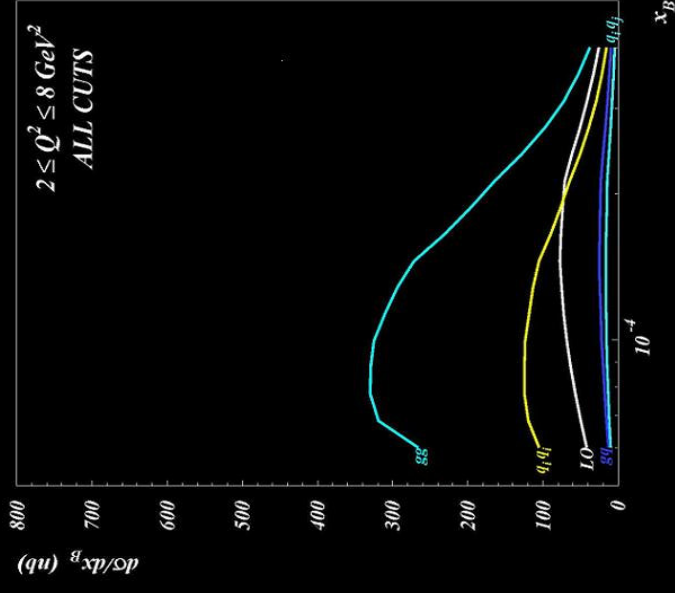
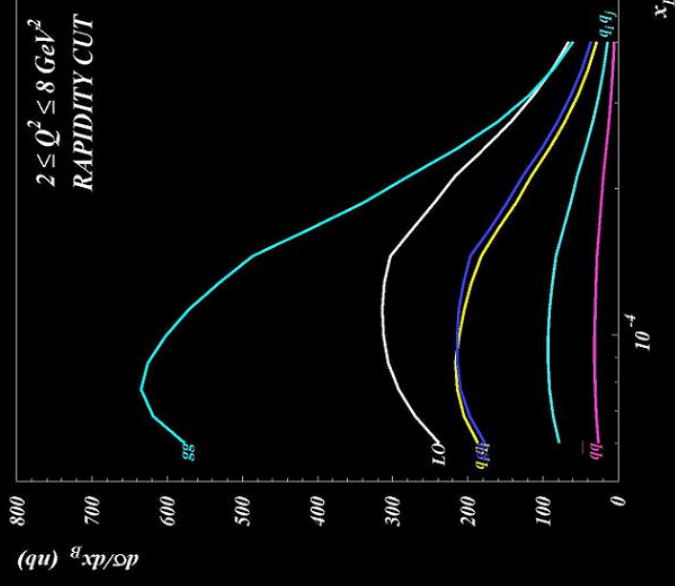
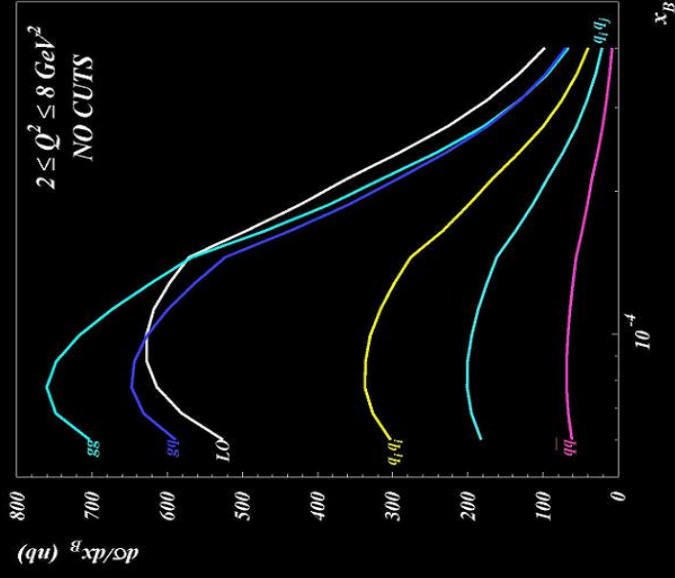






no cuts

rapidity cuts



no cuts

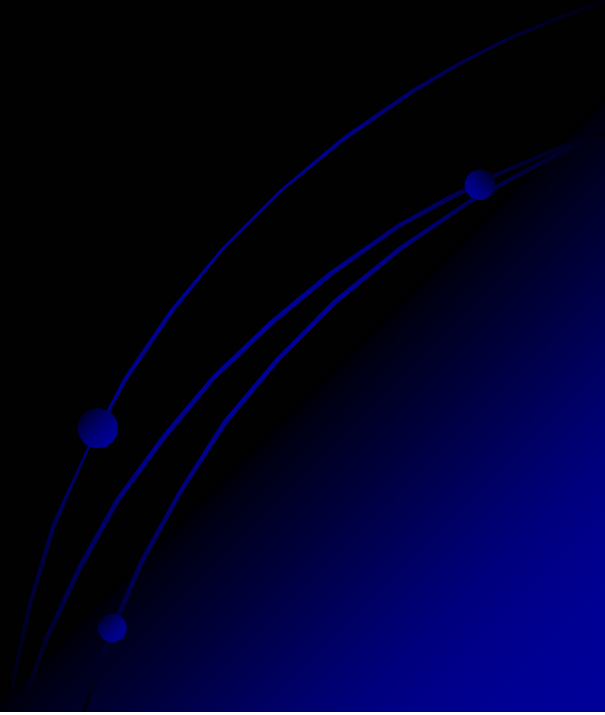
rapidity cuts

all cuts

Conclusions:

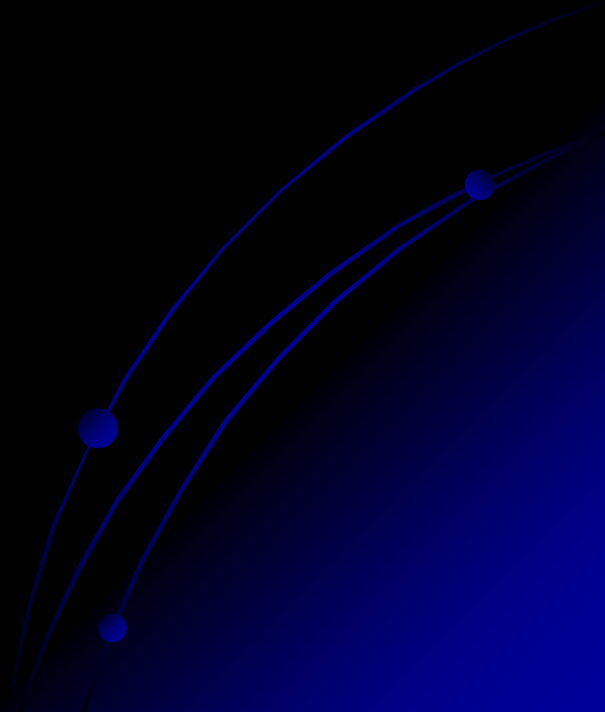
Conclusions:

- $\mathcal{O}(\alpha_s^2)$ QCD corrections to OPI DIS



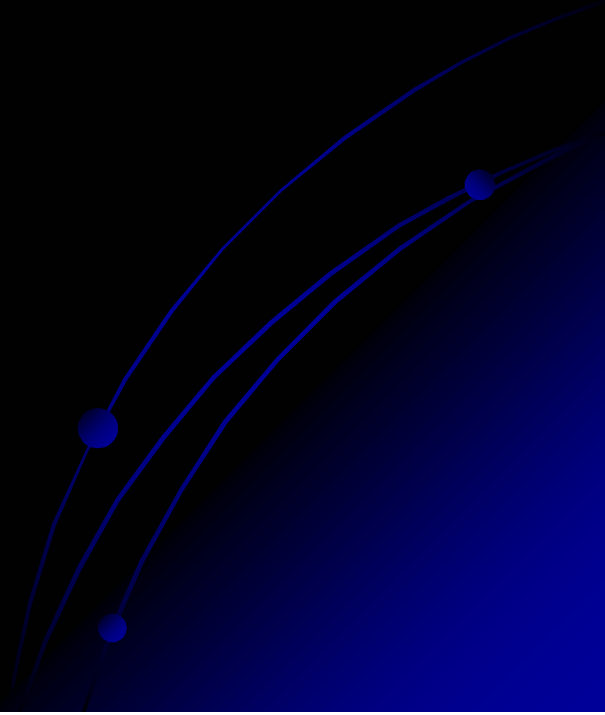
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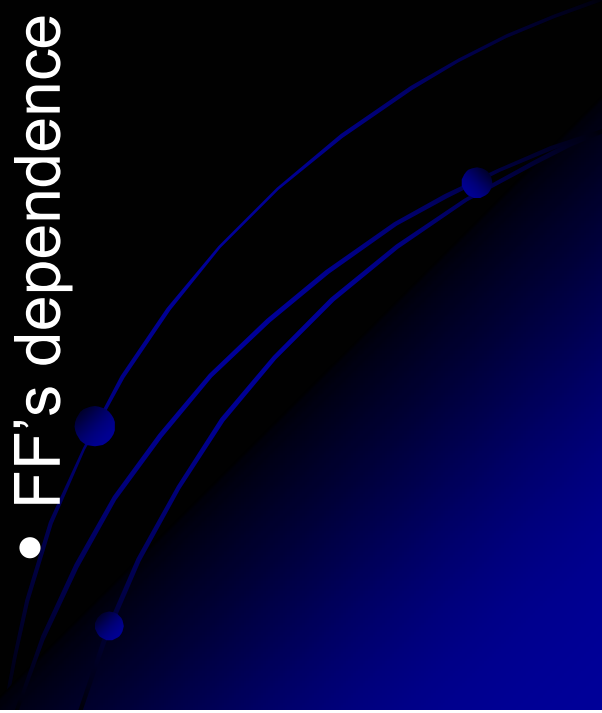
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- $\mathcal{O}(\alpha_s^2)$ QCD corrections to OPI DIS
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- FF's dependence
- DGLAP in good health (within present uncertainties)