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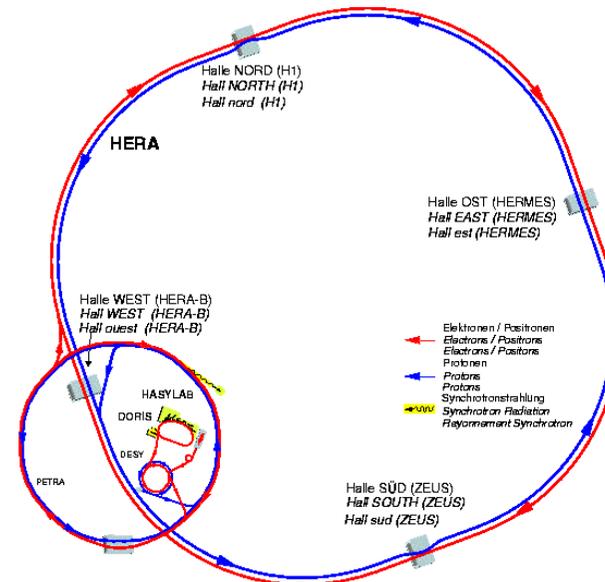
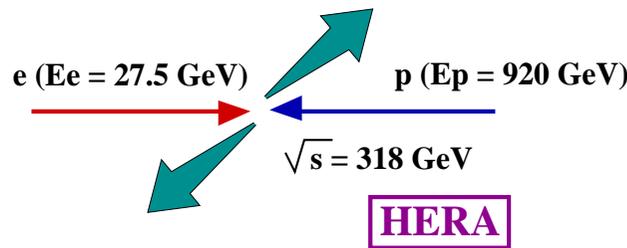
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COLOR DYNAMICS IN PHOTOPRODUCTION OF JETS AT HERA

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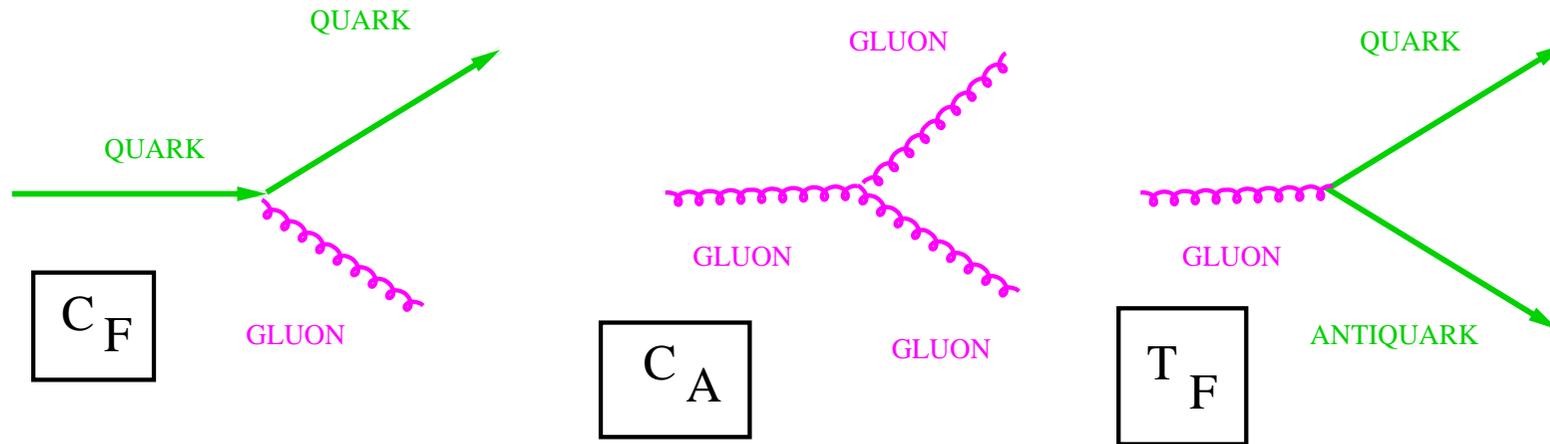


ZEUS Collaboration



Color Factors

- The color factors (C_F , C_A , T_F) represent the relative strength of the processes



- Their values are predicted by the underlying gauge-group structure
 - for $SU(N)$: $C_F = (N^2 - 1)/2N$, $C_A = N$, $T_F = 1/2$
- The color factors determine the relative contributions of the various subprocesses
 - e.g. in $e^+e^- \rightarrow 4$ jets the contributions from $q\bar{q}gg$ and $q\bar{q}q'\bar{q}'$
- Since the couplings qqg and ggg have different spin structures, the color factors give rise to a specific pattern of angular correlations between the final-state jets

Color Factors

- The effects of color factors have been extensively studied at LEP by measuring the angular correlations between the final state jets in

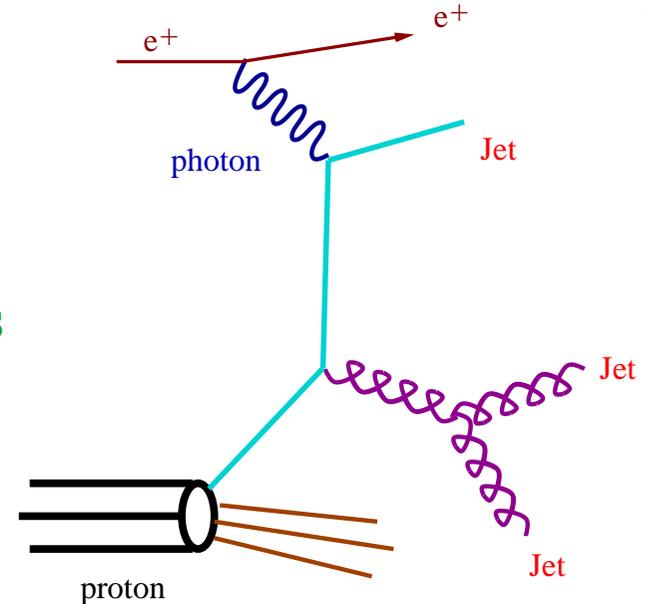
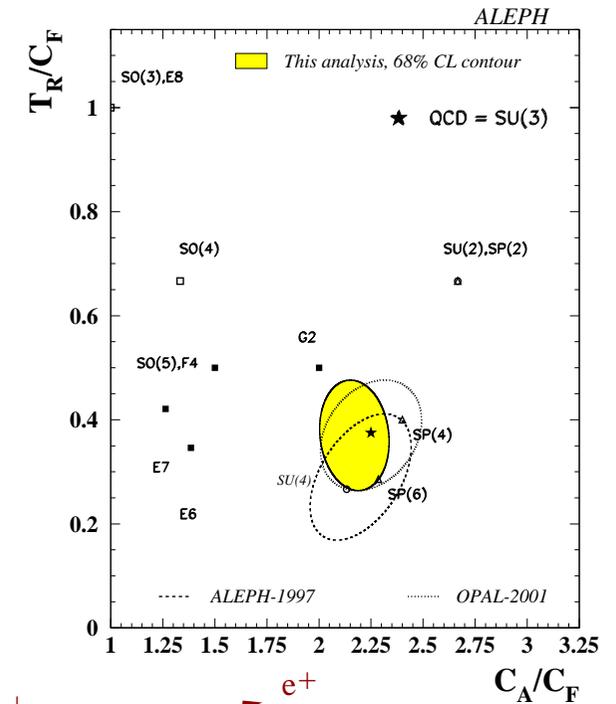
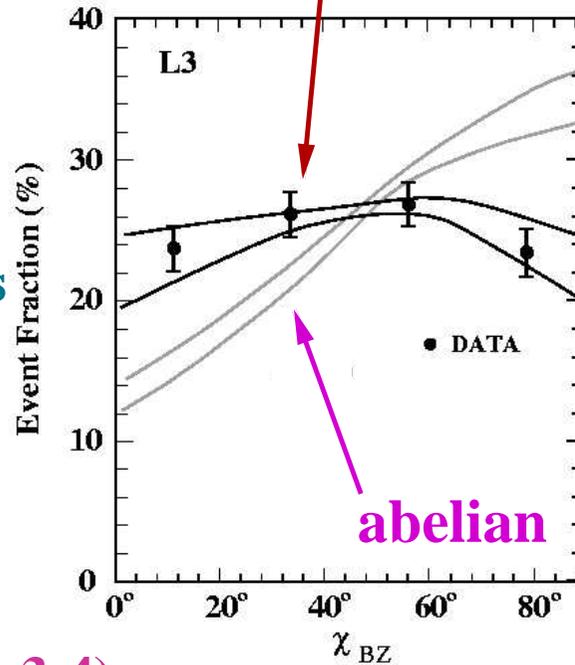
$$e^+ e^- \rightarrow 4 \text{ jets}$$

- (e.g. using the Bengtsson-Zerwas variable χ_{BZ} : angle between the plane of jets 1-2 and that of jets 3-4)

⇒ Determination of the ratios C_A/C_F and T_F/C_F

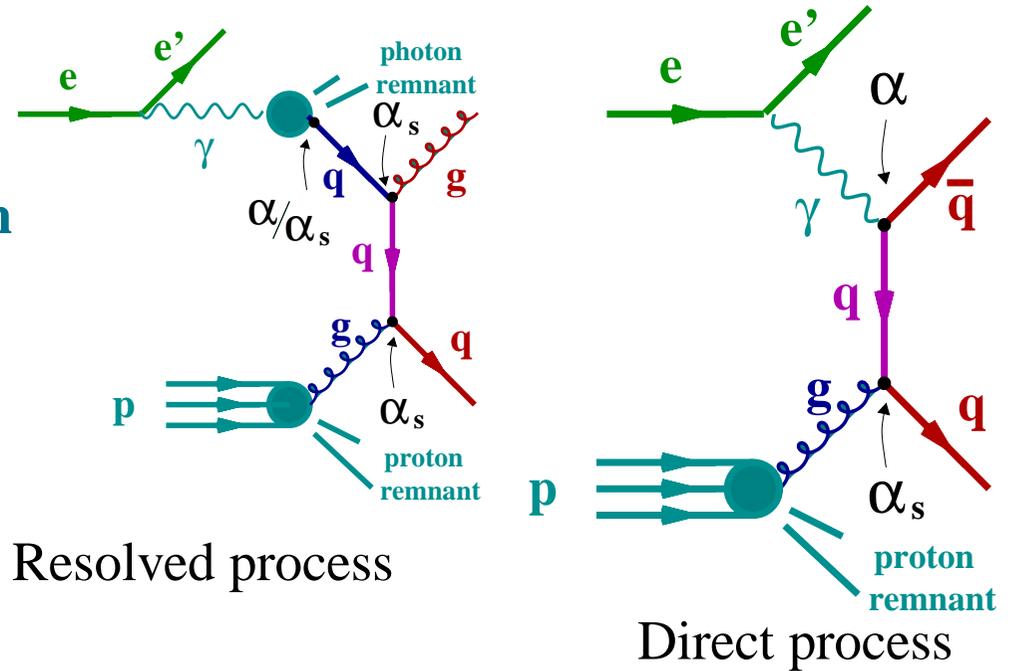
- In ep collisions at HERA, the effects of the color factors are expected in the photoproduction of three-jet events

QCD (non-abelian)



Photoproduction of Jets

- Production of jets in γp collisions has been measured via ep scattering at $Q^2 \approx 0$
- At lowest order QCD, two hard scattering processes contribute to jet production \Rightarrow
- pQCD calculations of jet cross sections



$$\sigma_{jet} = \sum_{a,b} \int_0^1 dy f_{\gamma/e}(y) \int_0^1 dx_\gamma f_{a/\gamma}(x_\gamma, \mu_{F\gamma}^2) \int_0^1 dx_p f_{b/p}(x_p, \mu_{Fp}^2) \hat{\sigma}_{ab \rightarrow jj}$$

longitudinal momentum fraction of γ/e^+ (y), **parton a/γ** (x_γ), **parton b /proton** (x_p)

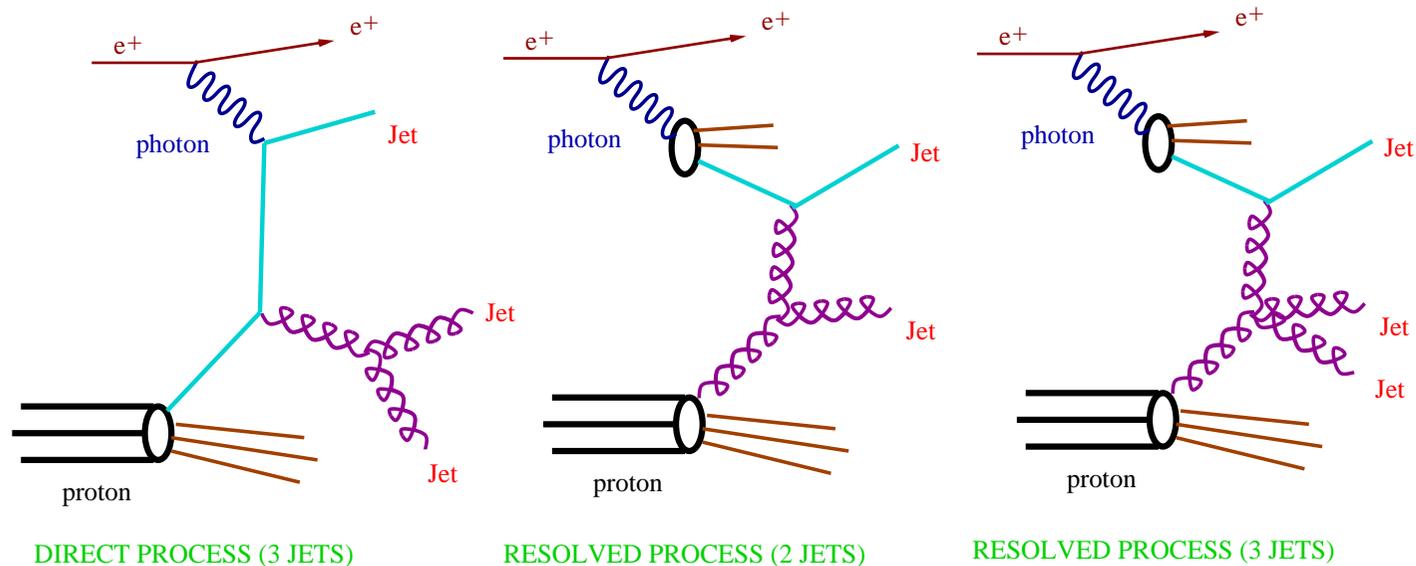
$\rightarrow f_{\gamma/e}(y)$ = flux of photons in the positron (WW approximation)

$\rightarrow f_{a/\gamma}(x_\gamma, \mu_{F\gamma}^2)$ = **parton densities in the photon** (for direct processes $\delta(1 - x_\gamma)$)

$\rightarrow f_{b/p}(x_p, \mu_{Fp}^2)$ = **parton densities in the proton**

$\rightarrow \sigma_{ab \rightarrow jj}$ **subprocess cross section; short-distance structure of the interaction**

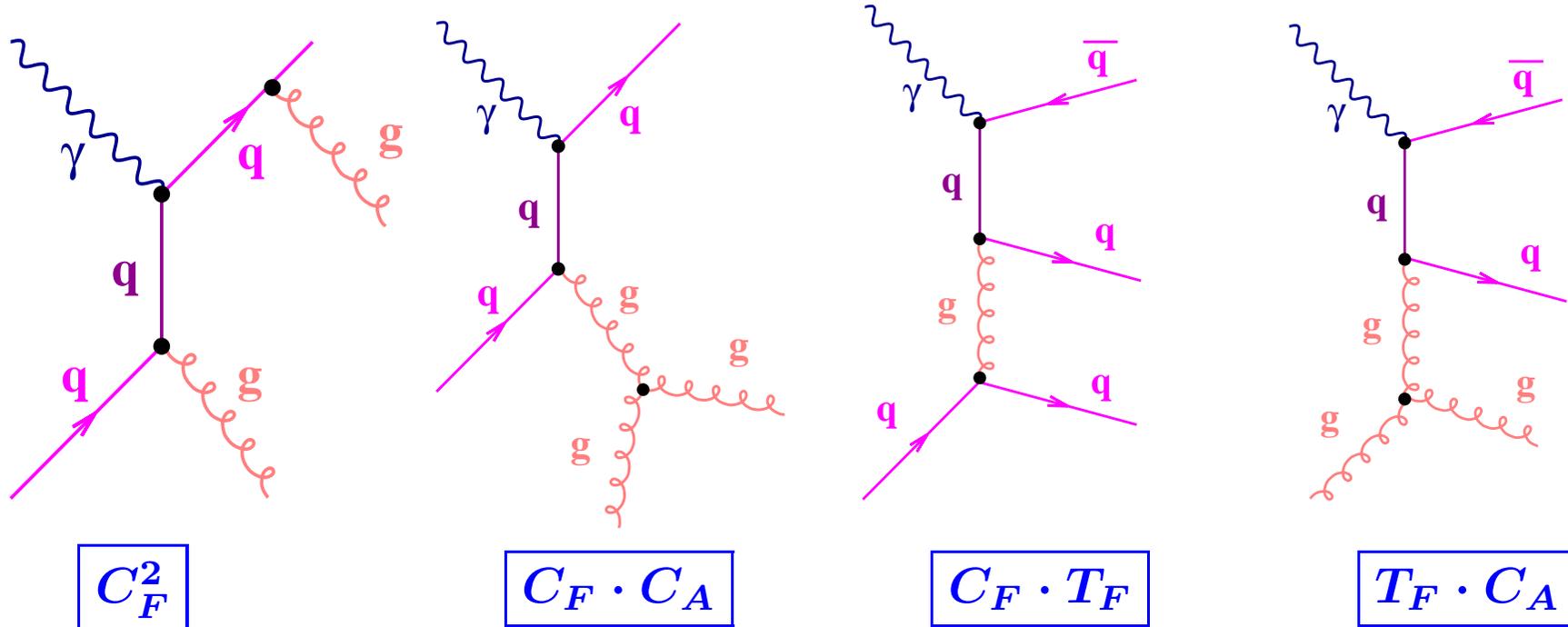
Color Factors: Direct vs Resolved Processes



- Resolved processes (dijet): the triple-gluon-vertex contribution already appears in the production of dijet events → **not a distinct feature in the final state jets**
- Resolved processes (trijet): new color factors contribute to the production of three-jet events → a rather complex situation to disentangle the effects of C_F , C_A , ...
- Direct processes (trijet): sensitive to the colour factors C_F , C_A and T_F through the angular correlations between the final-state jets (clean!)
- Observable to separate resolved/direct ⇒
$$x_\gamma^{obs}(3jets) = \frac{1}{2E_\gamma} \sum_{i=1}^3 E_T^{jet_i} e^{-\eta^{jet_i}}$$

Photoproduction of Three-Jet Events

- **Direct processes** provide a clean way to study the effects of the different color configurations



- The predicted cross section at $\mathcal{O}(\alpha\alpha_s^2)$ can be written as

$$\sigma_{ep \rightarrow 3\text{jets}} = C_F^2 \cdot \sigma_A + C_F C_A \cdot \sigma_B + C_F T_F \cdot \sigma_C + T_F C_A \cdot \sigma_D$$

Photoproduction of Three-Jet Events

- **Variables** to highlight the contributions from the different color configurations

⇒ **angular correlations between the jets**

→ θ_H , the angle between the plane determined by the highest E_T^{jet} and the beam and the plane determined by the two lowest E_T^{jet} jets (Muñoz-Tapia, Stirling);

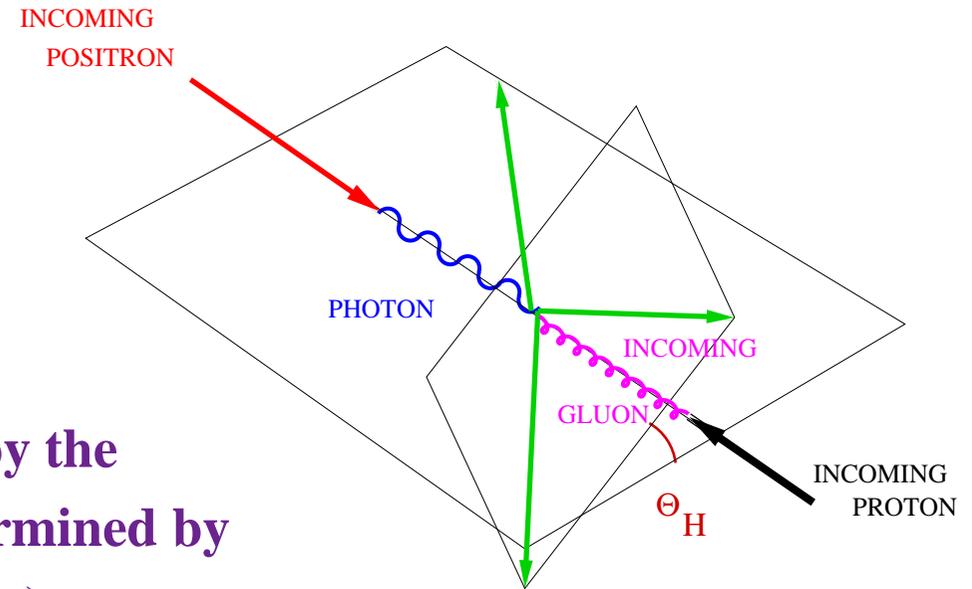
→ α_{23} , the angle between the two lowest E_T^{jet} jets;
(inspired by the variable $\alpha_{34}^{e^+e^-}$ for $e^+e^- \rightarrow 4\text{jets}$)

→ β_{KSW} , defined by

$$\cos \beta_{KSW} = \cos \frac{1}{2} [\angle[(\vec{p}_1 \times \vec{p}_3), (\vec{p}_2 \times \vec{p}_B)] + \angle[(\vec{p}_1 \times \vec{p}_B), (\vec{p}_2 \times \vec{p}_3)]],$$

where \vec{p}_i is the momentum of jet i (ordered according to decreasing E_T^{jet}) and \vec{p}_B is a unit vector in the direction of the proton beam;

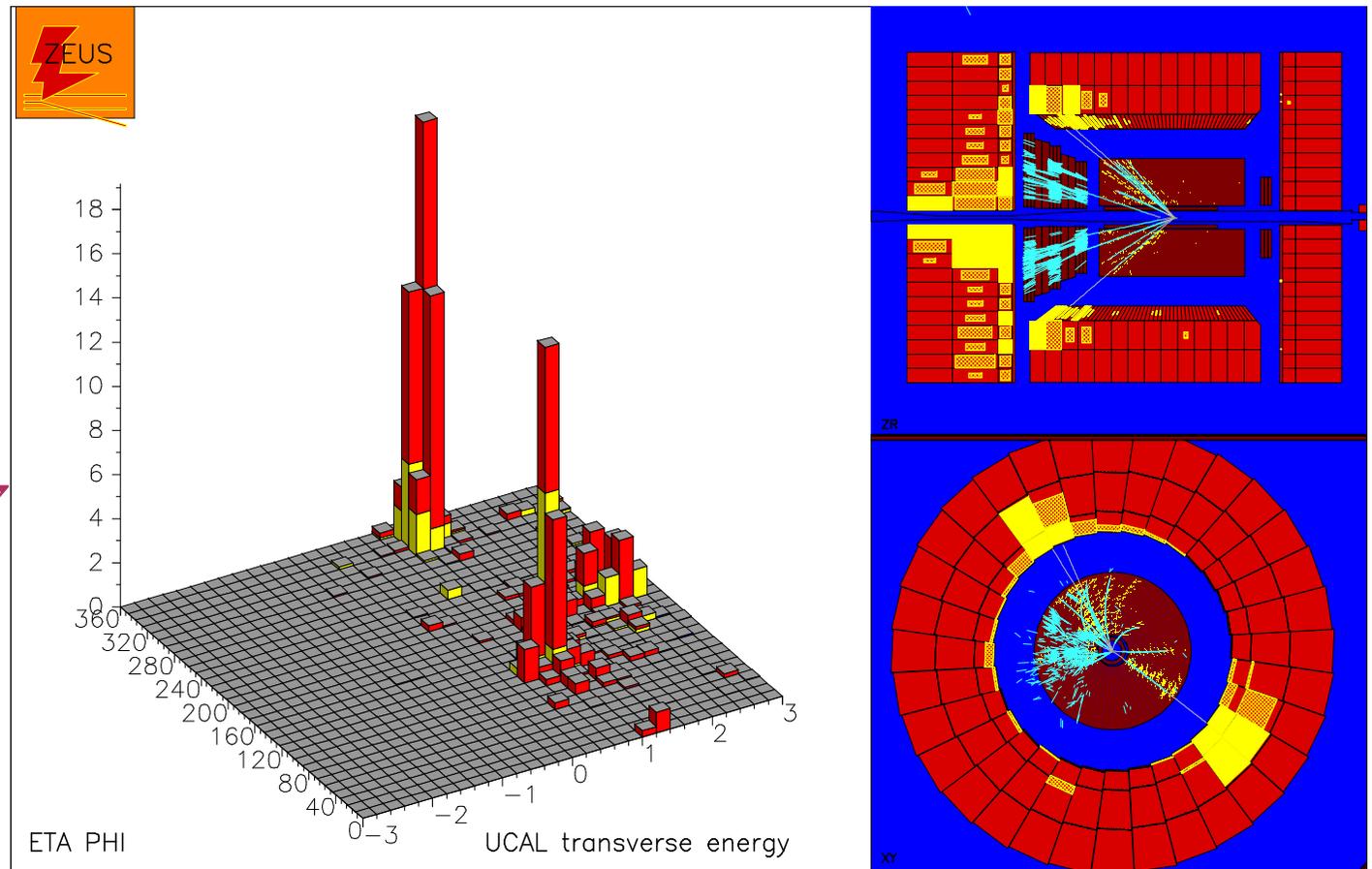
(inspired by the Körner-Schierholz-Willrodt angle $\Phi_{KSW}^{e^+e^-}$ for $e^+e^- \rightarrow 4\text{jets}$)



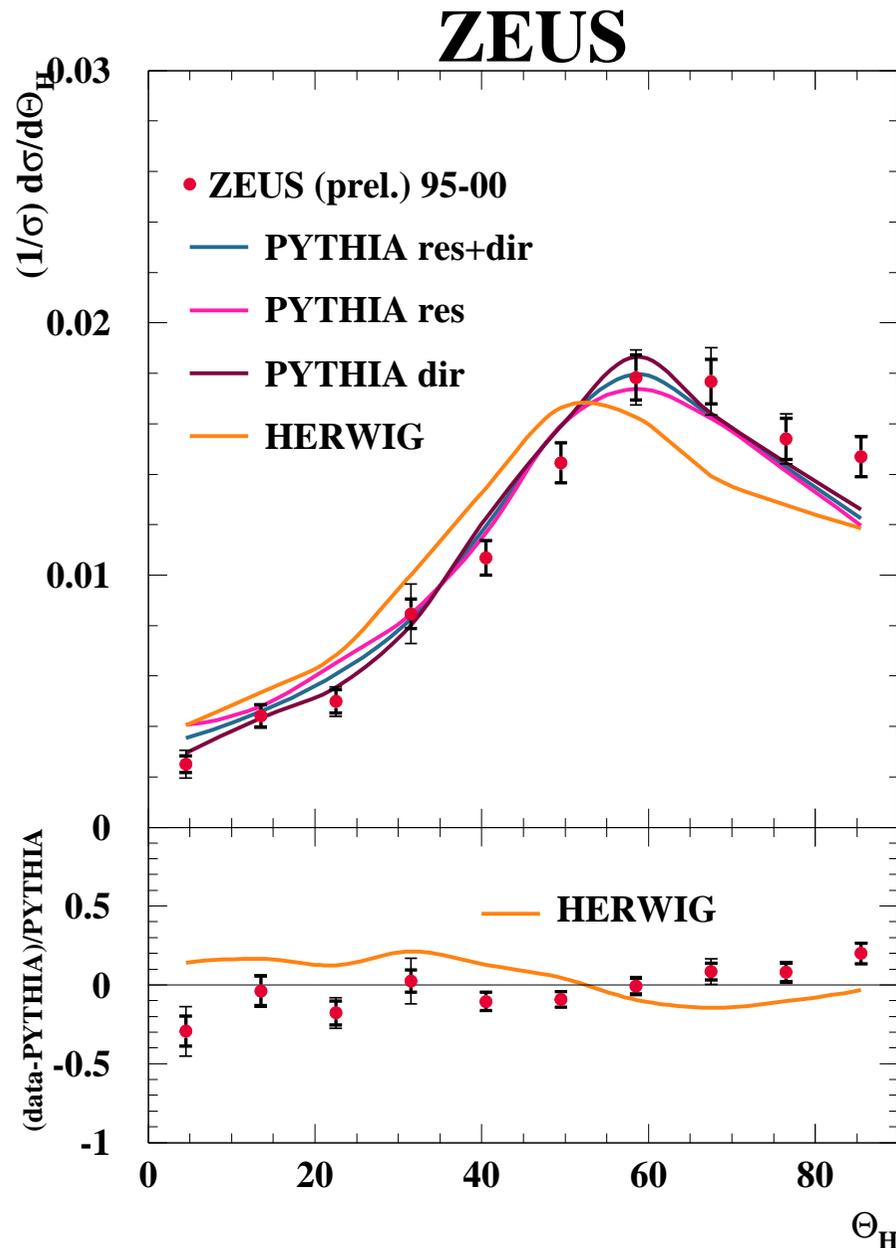
Measurements of Three-Jet Events in Photoproduction

- Measurements using HERA I data: $\mathcal{L} = 127 \text{ pb}^{-1}$
- Jets identified using the longitudinally invariant k_T cluster algorithm in the laboratory frame: at least three jets with $E_T^{\text{jet}} > 14 \text{ GeV}$, $-1 < \eta^{\text{jet}} < 2.5$
- Kinematic region: $Q^2 < 1 \text{ GeV}^2$ and $0.2 < y < 0.85$

~ 2200 three-jet events
satisfying $x_\gamma^{\text{obs}}(3\text{jets}) > 0.7$



Measurement of the Distribution in θ_H



- Measurement of the normalised cross section $1/\sigma d\sigma/d\theta_H$ for the production of events with three jets satisfying

$$E_T^{\text{jet}} > 14 \text{ GeV}, -1 < \eta^{\text{jet}} < 2.5$$

$$\text{and } x_\gamma^{\text{obs}}(3\text{jets}) > 0.7$$

- in the kinematic region

$$Q^2 < 1 \text{ GeV}^2 \text{ and } 0.2 < y < 0.85$$

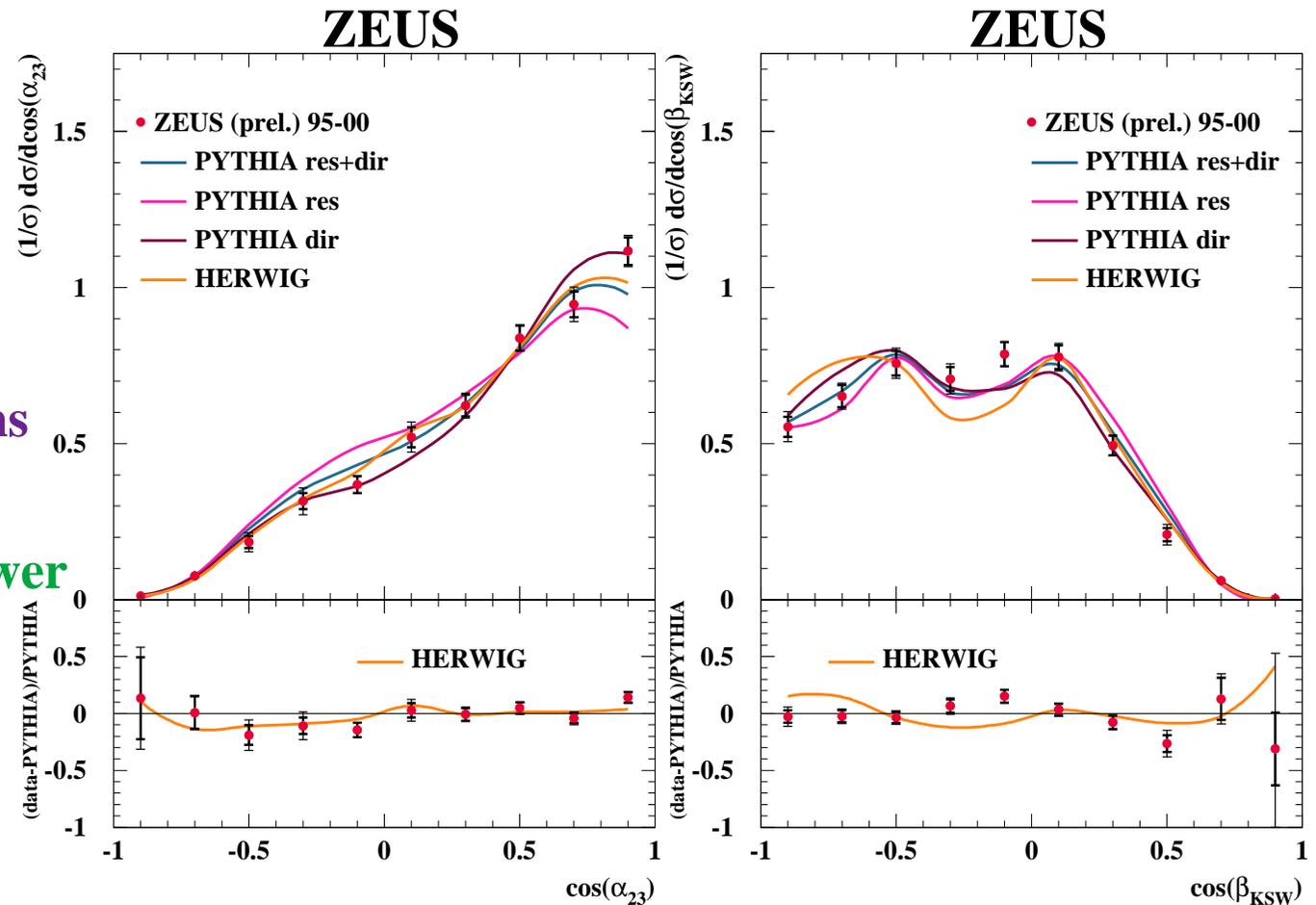
- Comparison to the predictions of PYTHIA and HERWIG (leading-logarithm parton-shower calculations based on SU(3))

→ PYTHIA reproduces reasonably well the measured distribution

→ The distributions for direct and resolved processes ($\sim 34\%$) are very similar

Measurement of the Distributions in α_{23} and β_{KSW}

- The distribution in α_{23} peaks at $\alpha_{23} \sim 0$
- The distribution in $\cos \beta_{KSW}$ has a broad peak between -0.5 and 0
- Comparison to the predictions of PYTHIA and HERWIG (leading-logarithm parton-shower calculations based on SU(3))
- PYTHIA reproduces reasonably well the measured distribution



→ The distributions for direct and resolved processes are similar

- Experimental uncertainties: the statistical uncertainties are dominant

Fixed-order Calculations for the Photoproduction of Three-Jet Events

- Fixed-order ($\mathcal{O}(\alpha\alpha_s^2)$) calculations using the program by Klasen, Kleinwort and Kramer

→ direct processes only

→ choice of scales: $\mu_R = \mu_F = E_T^{\max}$ (highest E_T^{jet})

→ proton PDFs: MRST99 parametrisations

- Parton-to-hadron corrections estimated with PYTHIA

⇒ small effect on the angular distributions ($< 5\%$)

- Small theoretical uncertainties

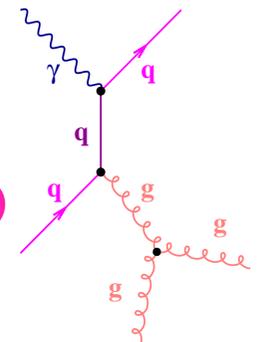
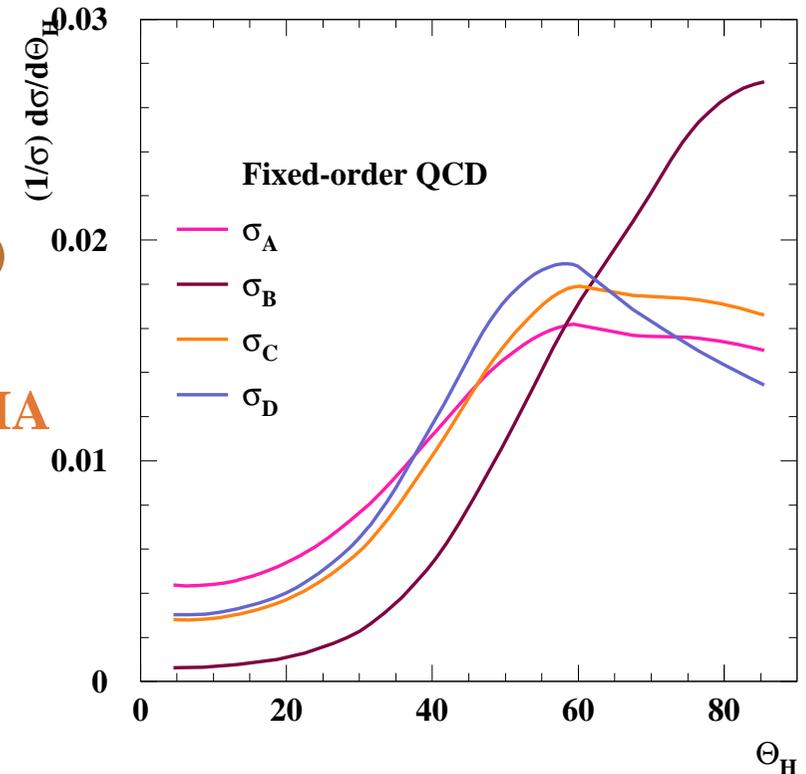
→ higher-order terms (by varying μ_R)

→ uncertainties on the proton PDFs

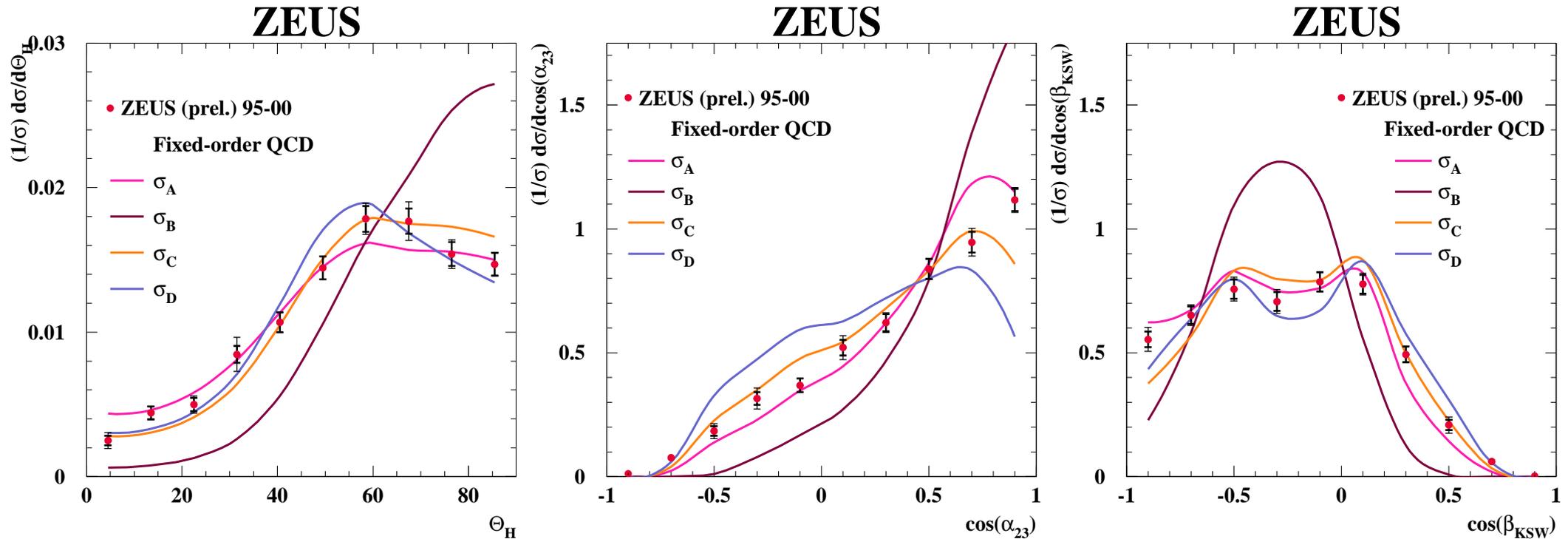
- The program has been modified to allow the calculation for given input values of the color factors (C_F, C_A, T_F):

$$\sigma_{ep \rightarrow 3\text{jets}} = C_F^2 \cdot \sigma_A + C_F C_A \cdot \sigma_B + C_F T_F \cdot \sigma_C + T_F C_A \cdot \sigma_D$$

- Example: the distribution in θ_H for the contribution σ_B (term with $C_F \cdot C_A$) is particularly distinct from the other color configurations due to diagrams →



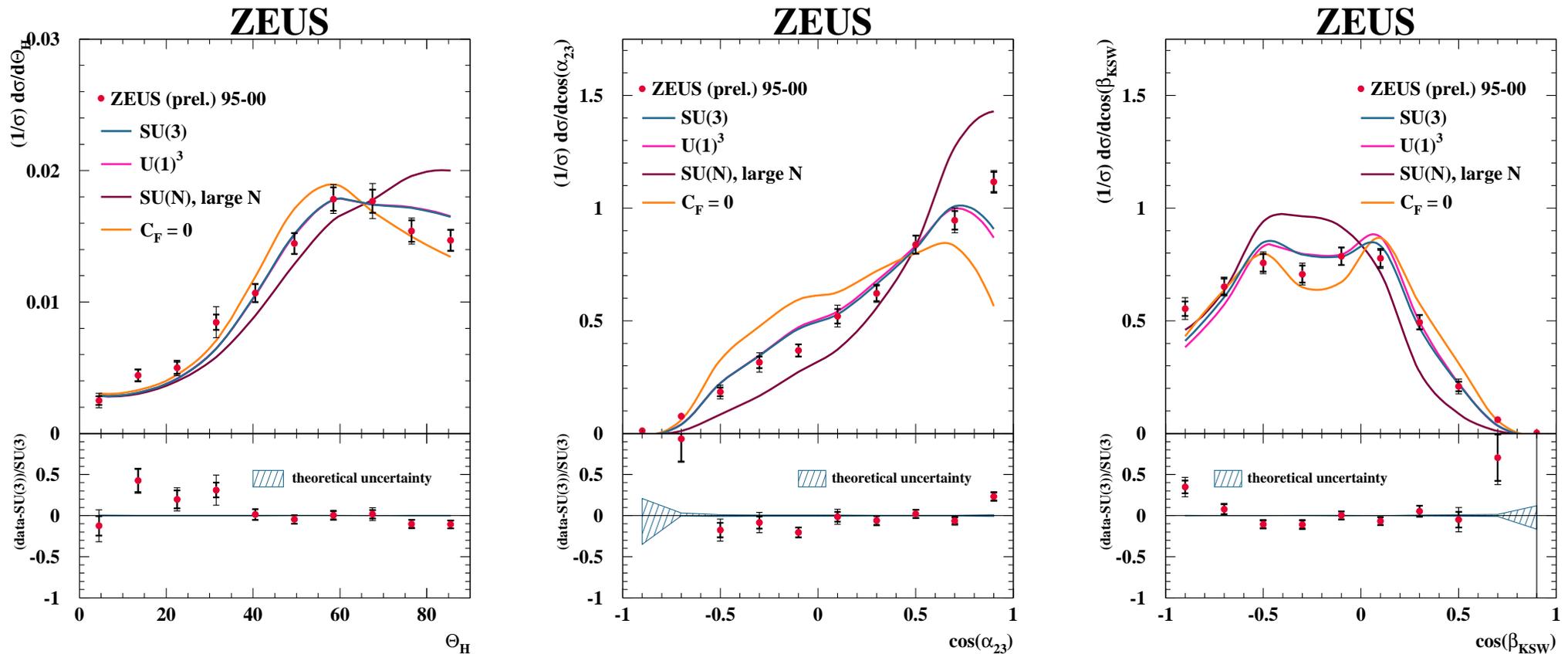
Measurements vs Fixed-order Calculations (Contributions)



- The contribution σ_B (term with $C_F \cdot C_A$) exhibits a very different shape than the other contributions in all three distributions ($\theta_H, \alpha_{23}, \beta_{KSW}$)
- The other contributions are better separated by the distribution of α_{23}
- The predicted relative contributions for $SU(3)$:

$$A (C_F^2) \rightarrow 13\%, B (C_F C_A) \rightarrow 10\%, C (C_F T_F) \rightarrow 45\%, D (T_F C_A) \rightarrow 32\%$$

Measurements vs Fixed-order Calculations (Complete)



- Comparison to fixed-order calculations based on different symmetry groups
 $SU(3)$, $SU(N)$ with $N \gg 1$, $U(1)^3$ and the extreme choice $C_F = 0$
- $U(1)^3$ vs $SU(3)$: similar shapes due to the smallness of σ_B
- The data disfavour $T_F/C_F \approx 0$ ($SU(N \gg 1)$) or $C_F = 0$
- The predictions of $SU(3)$ describe reasonably well the data

Summary

- Measurements of angular correlations between the final-state jets (and the beam) in the photoproduction of three-jet events → consistent with the admixture of color configurations predicted by $SU(3)$
- The measurements disfavour $T_F/C_F \approx 0$ as predicted by $SU(N)$ with large N or the extreme choice $C_F = 0$
- The observables θ_H , α_{23} and β_{KSW} → sensitivity to $C_F \cdot C_A$, but the predicted contribution is small (10%) → large contribution (32%) from $T_F \cdot C_A$, but not enough sensitivity ⇒ New variables are needed to enhance the contribution from the triple-gluon vertex in gluon induced processes

