

# Charmonium production at NLO in two-photon collisions

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# Contents

1. Introduction
2. LO results
3. NLO calculation
  - (a) Virtual corrections
  - (b) Real corrections
4. Phenomenology
  - (a)  $\gamma\gamma \rightarrow J/\psi + j + X^1$
  - (b)  $\gamma\gamma \rightarrow J/\psi + \gamma + X^2$
5. Summary and outlook

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<sup>1</sup>M. Klasen, B.A.K., L.N. Mihaila, M. Steinhauser, Nucl. Phys. **B713** (2005) 487.

<sup>2</sup>M. Klasen, B.A.K., L.N. Mihaila, M. Steinhauser, Phys. Rev. **D71** (2005) 014016.

## 1. Introduction

Heavy quarkonium  $H$ :

- useful testing ground for QCD
- simplest laboratory to study the binding of quarks to hadrons

Color Singlet Model: E.L. Berger, D. Jones, Phys. Rev. **D23** (1981) 1521; R. Baier, R. Rückl, Phys. Lett. **102B** (1981) 364.

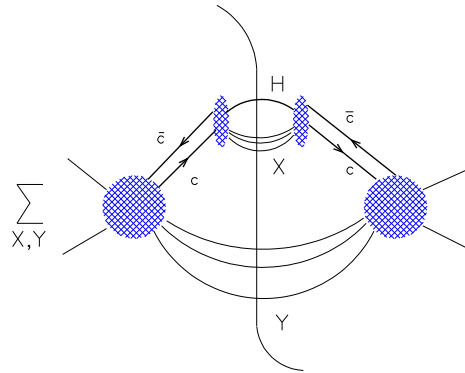
- Factorization in perturbative short-distance part  $[\alpha_s(m_Q), m_Q]$  and nonperturbative long-distance part  $[R_S(0), R'_P(0)]$ .
- $Q\bar{Q}$  pair must be in CS state and have same  $^{2S+1}L_J$  quantum numbers as  $H$ .
- No general argument for validity in higher orders; logarithmic IR divergences in  $P$ -wave decays to light hadrons and relativistic corrections to  $S$ -wave annihilation.
- Explains direct  $J/\psi$  photoproduction at HERA in NLO, but falls far short of Tevatron  $p\bar{p}$ , LEP2  $\gamma\gamma$ , and HERA  $ep$  DIS data.

Color Evaporation Model: H. Fritzsche, Phys. Lett. **B67** (1977) 217; F. Halzen, Phys. Lett. **B69** (1977) 105; M. Glück, J.F. Owens, E. Reya, Phys. Rev. **D17** (1978) 2324.

Hard comover scattering: P. Hoyer, S. Peigné, Phys. Rev. **D59** (1999) 034011; N. Marchal, S. Peigné, P. Hoyer, Phys. Rev. **D62** (2000) 114001.

## NRQCD factorization formalism

W. E. Caswell, G. P. Lepage, Phys. Lett. **B167** (1986) 437; G.T. Bodwin, E. Braaten, G.P. Lepage, Phys. Rev. **D51** (1995) 1125; **D55** (1997) 5853 (E).



- Rigorous theoretical framework, renormalizable, predictive.
- Factorization in perturbative short-distance coefficients and long-distance matrix elements:

$$d\sigma(a + b \rightarrow H + d) = \sum_n \langle \mathcal{O}^H[n] \rangle d\sigma(a + b \rightarrow Q\bar{Q}[n] + d),$$

where  $n = {}^{2S+1}L_J^{(c)}$ .

- Relative importance of  $\langle \mathcal{O}^H[n] \rangle$  governed by velocity scaling rules.  $\leadsto$  Double expansion in  $\alpha_s(m_H)$  and  $v \approx \alpha_s(m_H)$ .  
 $v^2 \approx 0.3$  (0.1) for  $J/\psi$  ( $\Upsilon$ ).
- CSM recovered for  $v \rightarrow 0$ .
- Predicts CO processes, to be identified in other kinds of experiments.

$k$	$\langle \mathcal{O}^H[n] \rangle \propto v^k$			
	$\eta_c$	$\psi(nS)$	$h_c$	$\chi_{cJ}$
3	$^1S_0^{(1)}$	$^3S_1^{(1)}$	—	—
5	—	—	$^1P_1^{(1)}, ^1S_0^{(8)}$	$^3P_J^{(1)}, ^3S_1^{(8)}$
7	$^1S_0^{(8)}, ^3S_1^{(8)}, ^1P_1^{(8)}$	$^1S_0^{(8)}, ^3S_1^{(8)}, ^3P_J^{(8)}$	—	—

### Motivation for NLO:

- Reduction of renormalization and factorization scale dependence.
- Sizeable effects, e.g. due to opening of new partonic production channels.
- Ultimate test of NRQCD factorization by global NLO fit.
- High-statistics data from HERA II, Tevatron Run II, LHC, ILC.

### Previous NLO calculations:

- $\gamma p \rightarrow J/\psi + X$  w/ direct  $\gamma$  and  $J/\psi$  for  $p_T > 0$  in CSM M. Krämer, J. Zunft, J. Steegborn, P.M. Zerwas, Phys. Lett. **B348** (1995) 657; M. Krämer, Nucl. Phys. **B459** (1996) 3.
- $\gamma p \rightarrow J/\psi + X$  w/ direct  $\gamma$  and  $J/\psi$  for  $p_T = 0$  in NRQCD F. Maltoni, M.L. Mangano, A. Petrelli, Nucl. Phys. **B519** (1998) 361.
- $p\bar{p} \rightarrow J/\psi + X$  for  $p_T = 0$  in NRQCD A. Petrelli, M. Cacciari, M. Greco, F. Maltoni, M.L. Mangano, Nucl. Phys. **B514** (1998) 245.

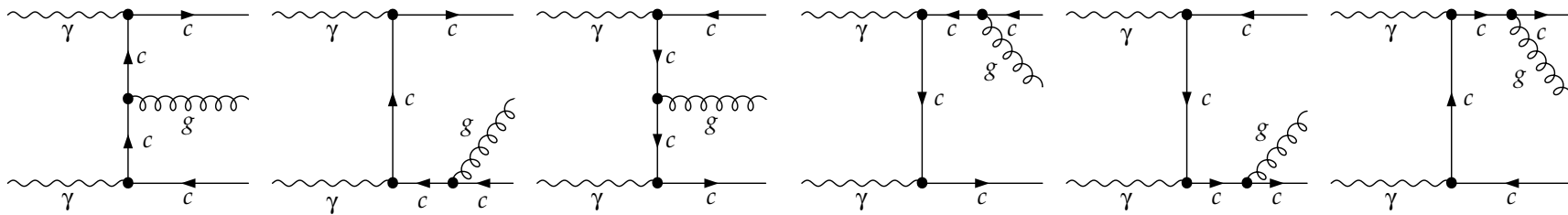
### Here:

- $\gamma\gamma \rightarrow J/\psi + X$  w/ direct  $\gamma$ 's and prompt  $J/\psi$  for  $p_T > 0$  in NRQCD
- $X$  purely hadronic: compensate  $\mu_R$  dependence of LO single-resolved contribution
- $X$  w/ prompt  $\gamma$ : direct photoproduction dominant

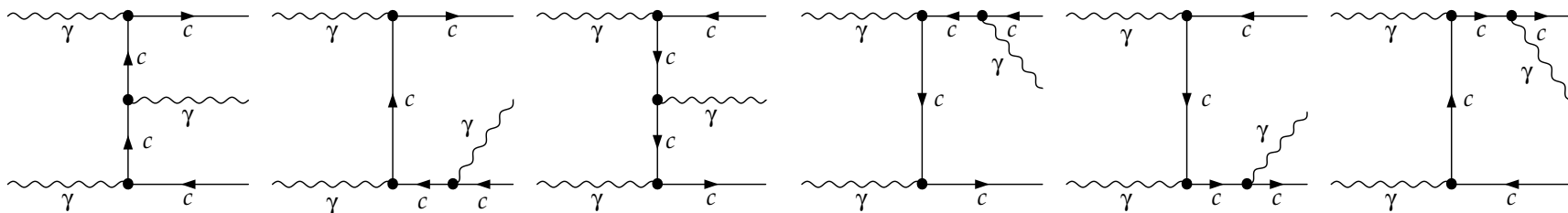
## 2. LO results

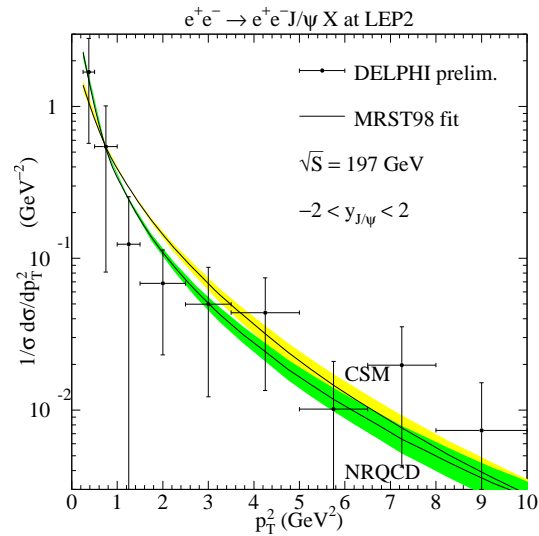
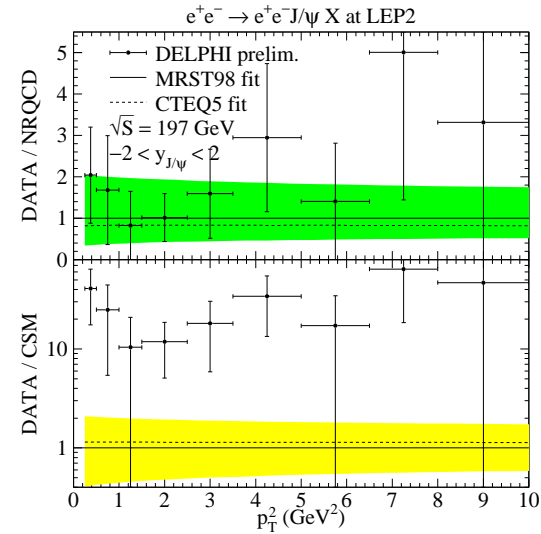
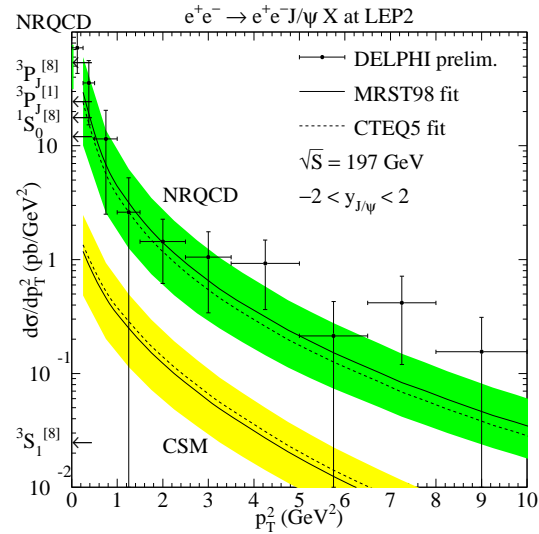
M. Klasen, B.A.K., L.N. Mihaila, M. Steinhauser, Phys. Rev. Lett. **89** (2002) 032001; R.M. Godbole, D. Indumathi, M. Krämer, Phys. Rev. **D65** (2002) 074003.

Subprocess:  $\gamma + \gamma \rightarrow c\bar{c} \left[ {}^3S_1^{(8)} \right] + g$



Subprocess:  $\gamma + \gamma \rightarrow c\bar{c} \left[ {}^3S_1^{(1)} \right] + \gamma$





M. Klasen, B.A.K., L.N. Mihaila, M. Steinhauser, Phys. Rev. Lett. **89** (2002) 032001.



### 3. NLO calculation: (a) Virtual corrections

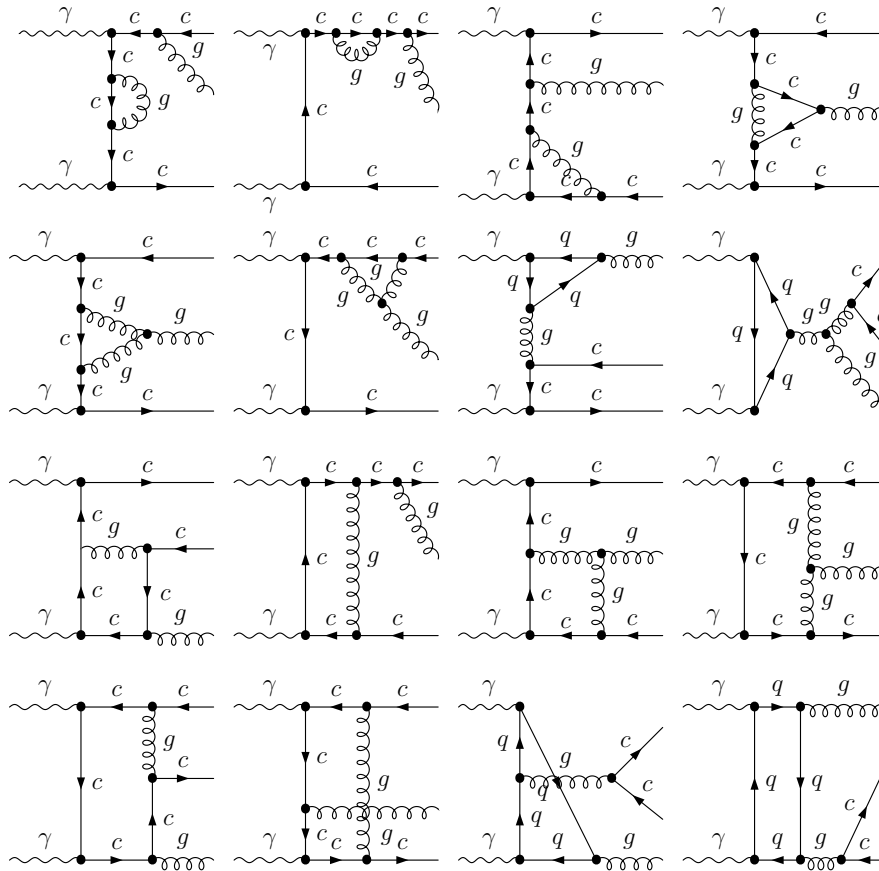
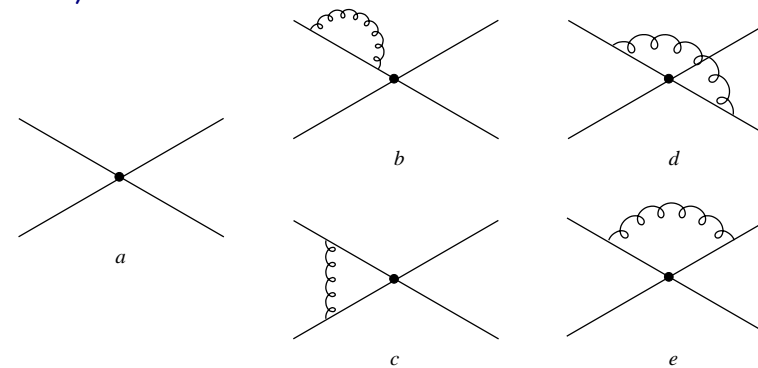


diagram	UV	IR	Coulomb
self energy	×		
triangle	×	×	
box		×	
pentagon		×	×

divergence	regularization	removal
UV	dim. reg. ( $1/\epsilon_{UV}$ )	renormalization of $g_s, m_Q, \psi_Q, G_\mu$ ( $\rightsquigarrow \mu$ ) and $\langle \mathcal{O}^H[n] \rangle$ ( $\rightsquigarrow \lambda$ )
IR	dim. reg. ( $1/\epsilon_{IR}$ )	real corrections $\langle \mathcal{O}^H[n] \rangle$
Coulomb	$1/v$	

### Operator renormalization



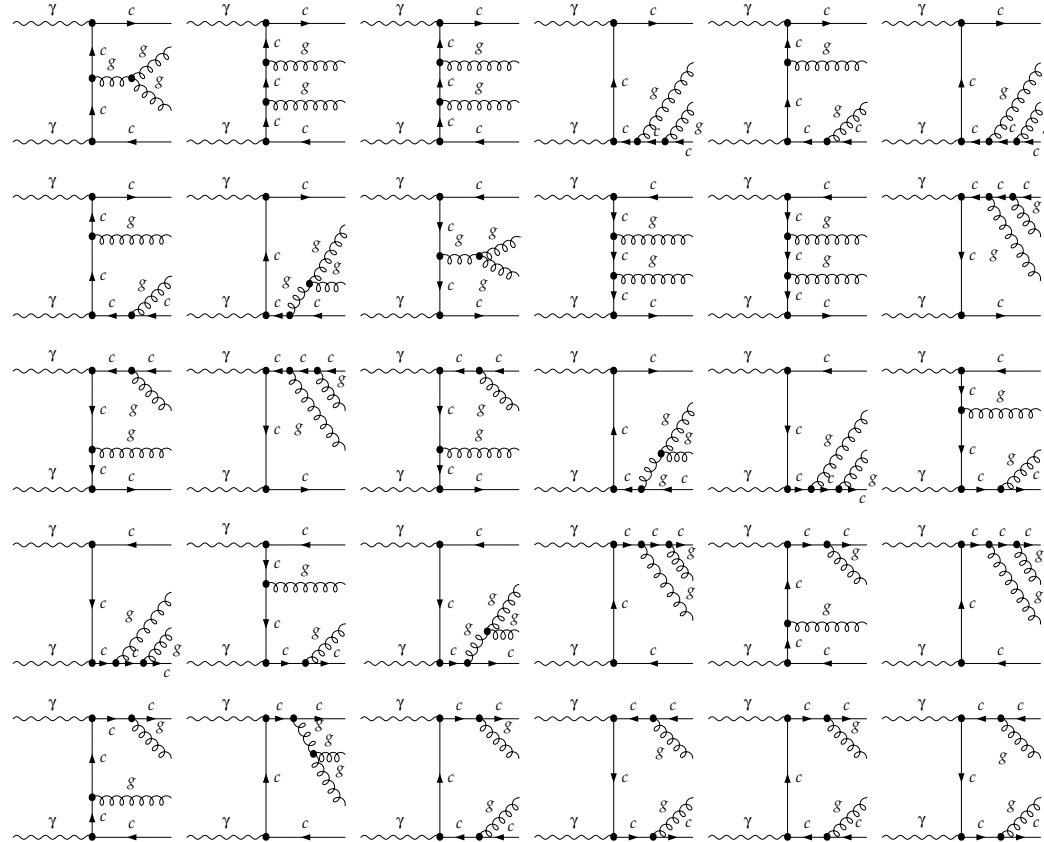
$$\begin{aligned}
 \langle \mathcal{O}^H [{}^3S_1^{(8)}] \rangle_1 &= \langle \mathcal{O}^H [{}^3S_1^{(8)}] \rangle_0 \left[ 1 + \left( C_F - \frac{C_A}{2} \right) \frac{\pi \alpha_s}{2v} \right] + \frac{4\alpha_s}{3\pi m^2} \left( \frac{4\pi \mu^2}{\lambda^2} \right)^\epsilon \\
 &\times \exp(-\epsilon \gamma_E) \left( \frac{1}{\epsilon_{UV}} - \frac{1}{\epsilon_{IR}} \right) \sum_{J=0}^2 \left( C_F \langle \mathcal{O}^H [{}^3P_J^{(1)}] \rangle + B_F \langle \mathcal{O}^H [{}^3P_J^{(8)}] \rangle \right)
 \end{aligned}$$

**(b) Real corrections**

Subprocesses:

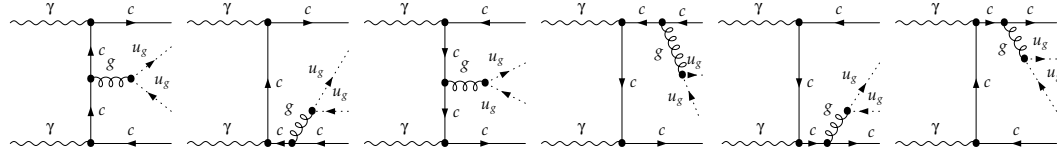
$$\gamma + \gamma \rightarrow c\bar{c}[n] + g + g$$

$$n = {}^3P_J^{(1)}, {}^1S_0^{(8)}, {}^3S_1^{(8)}, {}^3P_J^{(8)}$$



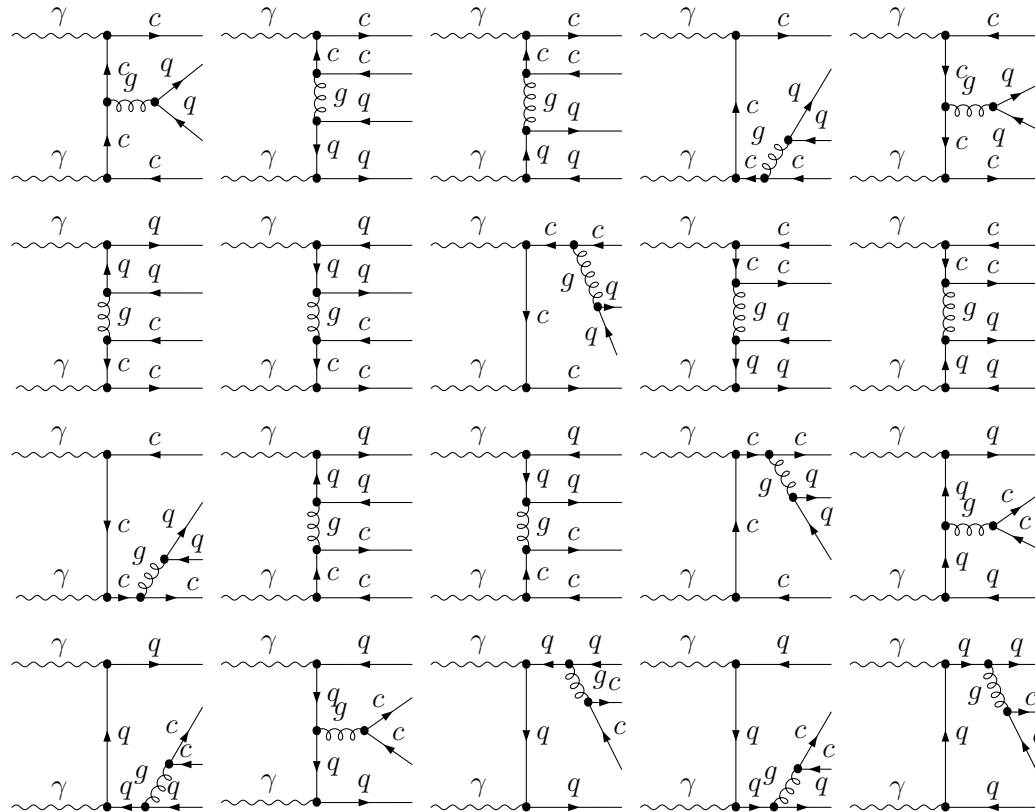
$$\gamma + \gamma \rightarrow c\bar{c}[n] + u_g + \bar{u}_g$$

$$n = {}^3S_1^{(8)}$$



$$\gamma + \gamma \rightarrow c\bar{c}[n] + q + \bar{q}$$

$$n = {}^1S_0^{(8)}, {}^3S_1^{(8)}, {}^3P_J^{(8)}$$



Soft and/or collinear singularities:

- Introduce IS and FS phase-space slicing parameters  $\delta_i, \delta_f$ .
- Integrate over soft and/or collinear regions analytically in  $d$  dimensions.
- Factorize IS collinear singularities into  $\gamma$  PDFs of single-resolved process ( $\leadsto M$ ).
- Integrate over hard region numerically in 4 dimensions.

Process	IS coll. ( $\delta_i$ )	FS coll. ( $\delta_f$ )	soft ( $\delta_f$ )
$\gamma\gamma \rightarrow c\bar{c} \left[ \begin{array}{l} {}^3P_J^{(1)}, {}^3P_J^{(8)} \end{array} \right] gg$			×
$\gamma\gamma \rightarrow c\bar{c} \left[ \begin{array}{l} {}^1S_0^{(8)} \end{array} \right] gg$			
$\gamma\gamma \rightarrow c\bar{c} \left[ \begin{array}{l} {}^3S_1^{(8)} \end{array} \right] gg$		×	×
$\gamma\gamma \rightarrow c\bar{c} \left[ \begin{array}{l} {}^3S_1^{(8)} \end{array} \right] q\bar{q}$	×	×	
$\gamma\gamma \rightarrow c\bar{c} \left[ \begin{array}{l} {}^1S_0^{(8)}, {}^3P_J^{(8)} \end{array} \right] q\bar{q}$	×		

Assembly of  $\gamma\gamma \rightarrow H + X$  cross section:

$$d\sigma(\mu, \lambda, M) = d\sigma_0(\mu, \lambda) [1 + \delta_{\text{vi}}(\mu; \epsilon_{\text{UV}}, \epsilon_{\text{IR}}, v) + \delta_{\text{ct}}(\mu; \epsilon_{\text{UV}}, \epsilon_{\text{IR}}) + \delta_{\text{op}}(\mu, \lambda; \epsilon_{\text{IR}}, v) + \delta_{\text{fs}}(\mu; \epsilon_{\text{IR}}, \delta_f)] + d\sigma_{\text{is}}(\mu, \lambda, M; \delta_i) + d\sigma_{\text{so}}(\mu, \lambda; \epsilon_{\text{IR}}, \delta_f) + d\sigma_{\text{ha}}(\mu, \lambda; \delta_i, \delta_f)$$

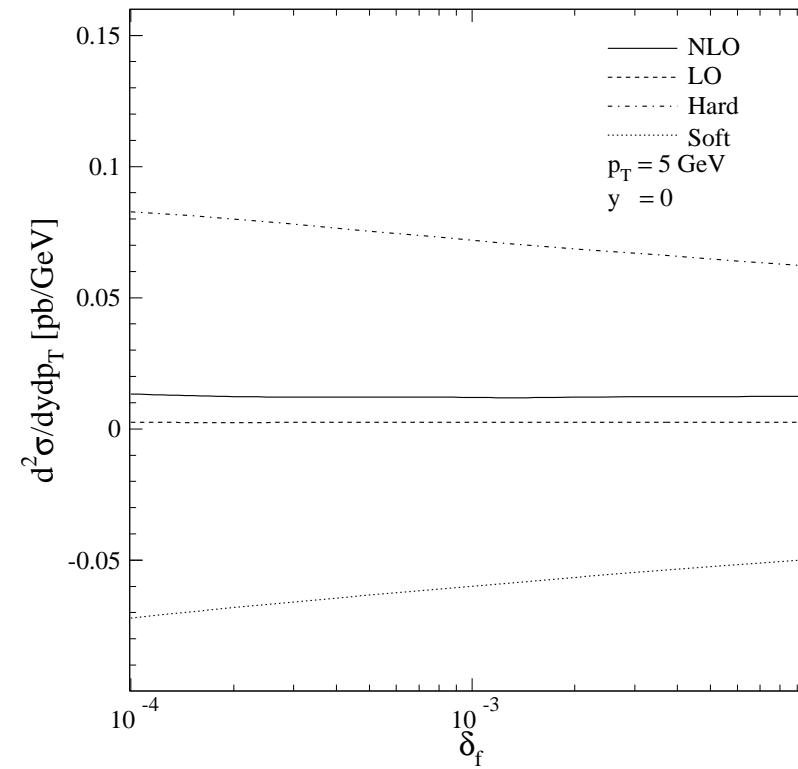
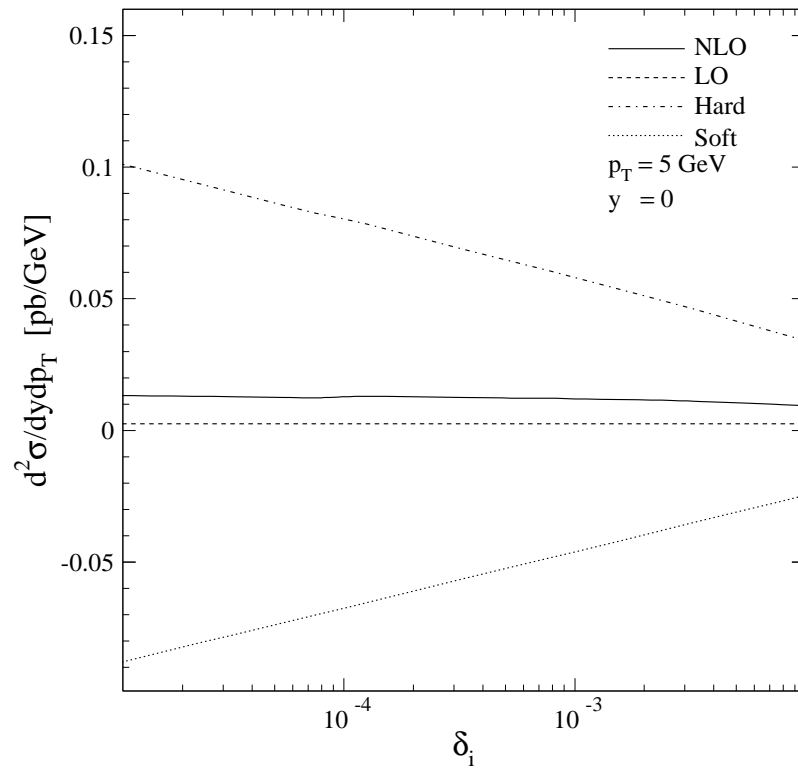
## 4. Phenomenology

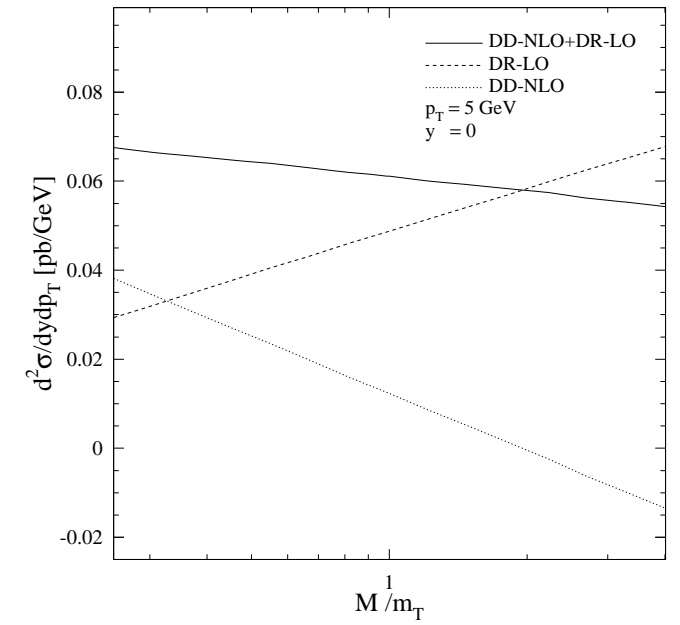
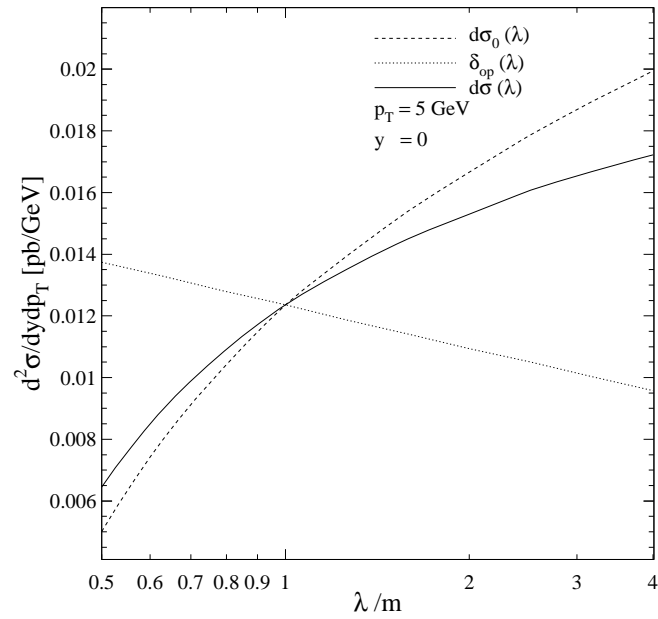
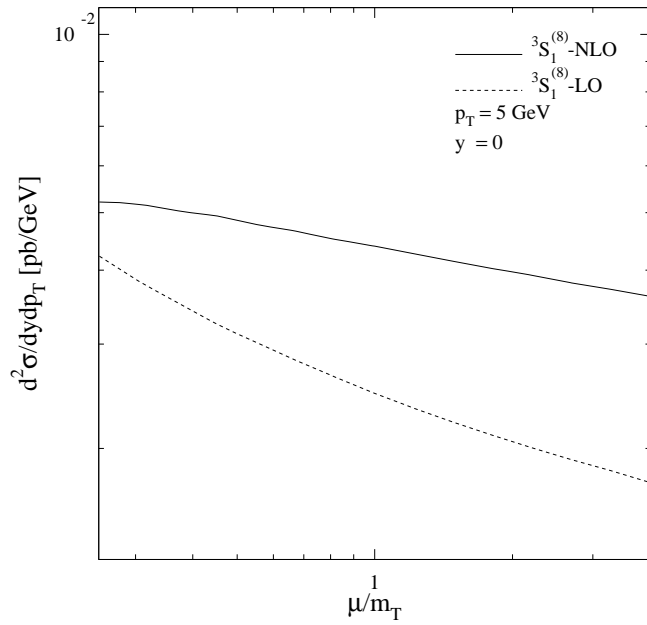
### Input:

- $e^+e^-$  linear collider with  $\sqrt{S} = 500$  GeV
- IS bremsstrahlung with  $\theta_{\max} = 25$  mrad
- Beamstrahlung with  $\Upsilon = 0.053$  (TESLA)
- $J/\psi$ ,  $\chi_{cJ}$ , and  $\psi'$  MEs from E. Braaten, B.A.K., J. Lee, Phys. Rev. **D62** (2000) 094005
- $\gamma$  PDFs from M. Glück, E. Reya, I. Schienbein, Phys. Rev. **D60** (1999) 054019
- FS prompt  $\gamma$ :  $p_T^\gamma > 3$  GeV,  $|y^\gamma| < 2.79$

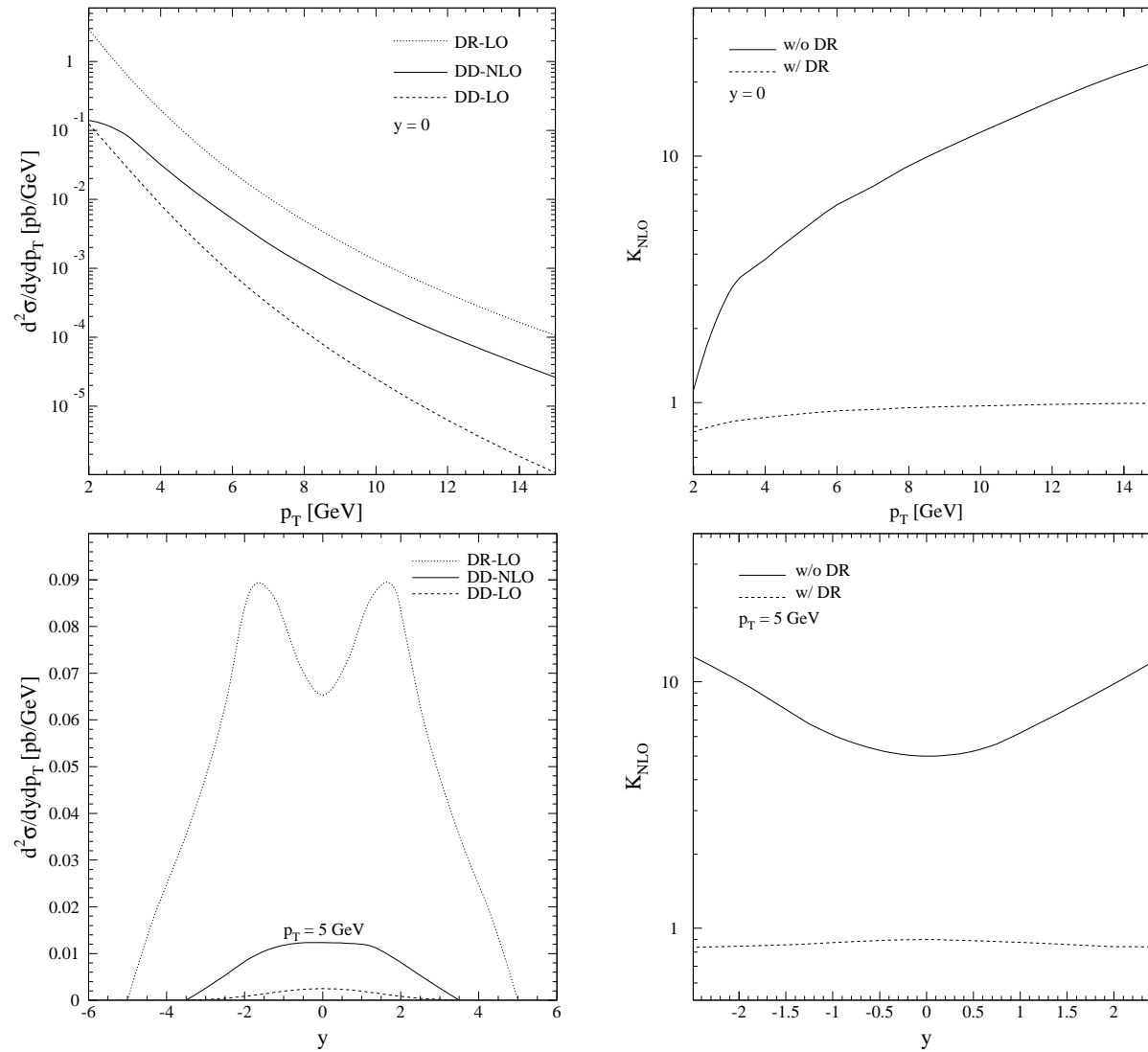
## (a) $\gamma\gamma \rightarrow J/\psi + j + X$

$\delta_i, \delta_f$  dependences



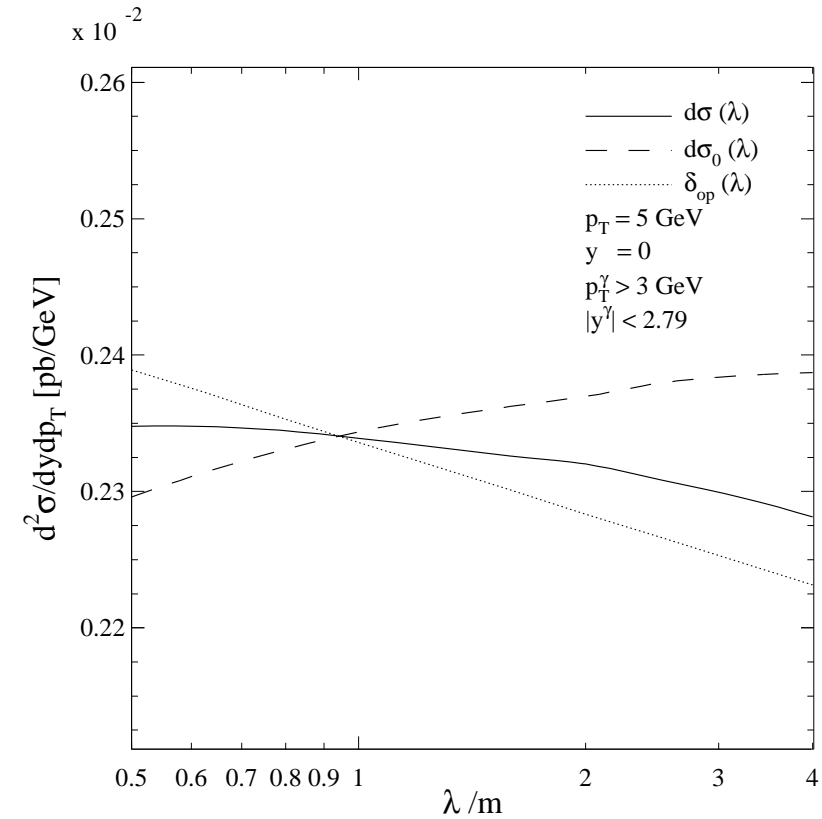
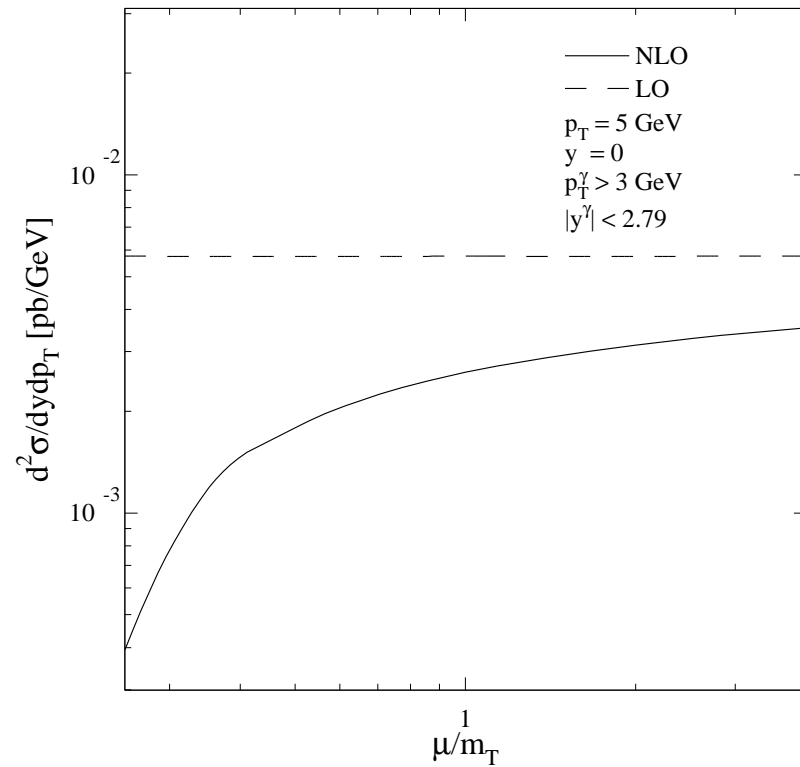
$\mu, \lambda, M$  dependences

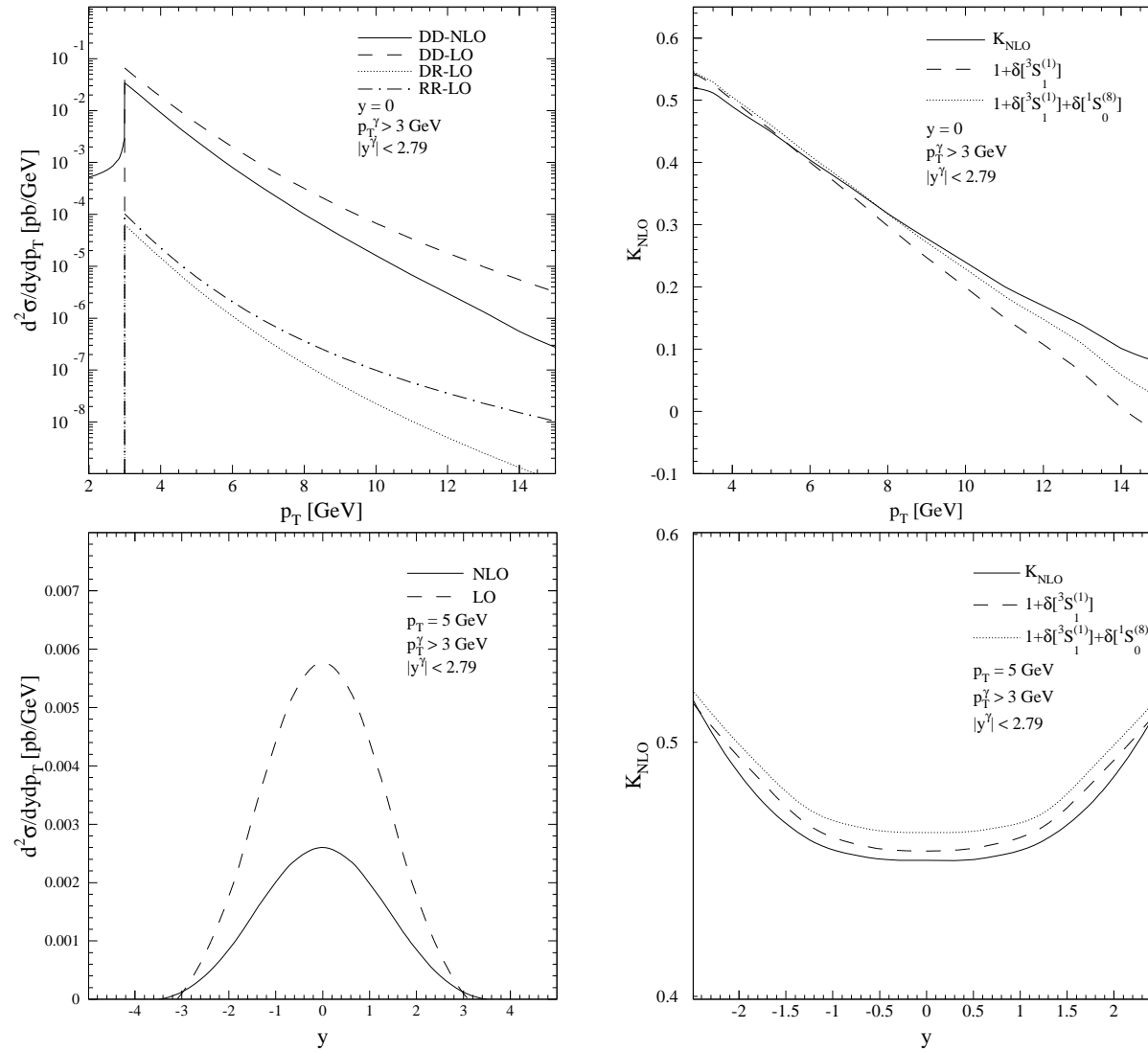




## (b) $\gamma\gamma \rightarrow J/\psi + \gamma + X$

$\mu, \lambda$  dependences





## 5. Summary and outlook

- $\gamma\gamma \rightarrow J/\psi + j + X$  w/ direct  $\gamma$ 's and prompt  $J/\psi$  for  $p_T > 0$  at NLO in NRQCD
  - Sizeable  $K$  factor due to opening of  $g \rightarrow c\bar{c} \left[ {}^3S_1^{(8)} \right]$  fragmentation-prone subprocesses  $\gamma\gamma \rightarrow c\bar{c} \left[ {}^3S_1^{(8)} \right] g$  and  $\gamma\gamma \rightarrow c\bar{c} \left[ {}^3S_1^{(8)} \right] q\bar{q}$  (proportional to  $e_q^2$ )
  - $M$  dependence of LO cross section (dominantly single-resolved) considerably reduced
- $\gamma\gamma \rightarrow J/\psi + \gamma + X$  at NLO in NRQCD
  - Direct photoproduction dominant
  - Unscreened  $\mu$  dependence
  - No  $M$  dependence  $\rightsquigarrow$  formally independent of single-resolved photoproduction
  - Small  $K$  factor from virtual corrections
- Extend analysis to  $ep$  photoproduction and hadroproduction
- Test NRQCD factorization by global NLO fit to high-statistics data from HERA II, Tevatron Run II, LHC, ILC, . . .