

Z - H - η COUPLINGS IN LITTLE HIGGS MODELS

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- Little Higgs pseudo-axions
- Product Group v. Simple Group
- ILC pheno

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LITTLE HIGGS PSUEDO-AXIONS

Little Higgs models have pseudo-axions!

[Kilian, DR, Reuter, PRD(71) 015008]

→ corresponds to broken diagonal generator

In general, extra $U(1)$ s appear, but don't gauge:

A' tends to conflicts w/ EW data

Examples:

1. Minimal moose $[SU(3)]^4 \rightarrow SU(3)$: rank 8 → 2 ∴ 6 η;
gauged $SU(3) \times SU(2) \times U(1)$, rank 2 higher than EW,
∴ 2 η eaten by heavy VBs, 4 pseudo-axions remain
2. Simple group $[SU(4)]^3 \rightarrow [SU(3)]^2 \times SU(2)$: rank 9 → 5 ∴ 4 η;
gauged $SU(4) \times U(1)$, rank 2 higher than EW,
∴ 2 η eaten, 2 remain
3. Schmaltz μ -model: $SU(3) \rightarrow SU(2)$: rank 2 → 1 ∴ 1 η;
gauged $SU(2)$ is EW sector ∴ no η eaten, 1 remains

Littlest Higgs ($SU(5) \rightarrow SO(5)$) model pseudo-axion

NGBs parameterized by 5×5 matrix $\Sigma = (\exp \frac{2i}{F} \Pi) \Sigma_0$

$$\Sigma_0 = \begin{pmatrix} 0 & 0 & 1_{2 \times 2} \\ 0 & 1 & 0 \\ 1_{2 \times 2} & 0 & 0 \end{pmatrix}, \quad \Pi = \frac{1}{\sqrt{2}} \begin{pmatrix} \eta/\sqrt{10} & h & \phi \\ h^\dagger & -4\eta/\sqrt{10} & h^T \\ \phi^\dagger & h^* & \eta/\sqrt{10} \end{pmatrix}$$

Yukawa interactions also have η -dependent factors $\xi = \exp \frac{i}{\sqrt{5}F} \eta$

$$\chi_L = \begin{pmatrix} i\tau^2 \xi^{\beta_0} T_L & iq_L & 0 \\ -iq_L^T & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad \text{and} \quad iT_2^2 = \text{diag}(0, 0, -i\tau^{2*})/2$$

$$\begin{aligned} \mathcal{L}_Y^t &= \lambda_1 F \xi^{\beta_1} \bar{t}_R \text{Tr} [\Sigma^* (iT_2^2) \Sigma^* \chi_L] - \lambda_2 F \xi^{\beta_2} \bar{T}_R T_L + \text{h.c.} \\ \mathcal{L}_Y^b &= -\frac{\lambda_b}{\sqrt{2}} F \xi^{\beta_3} \bar{b}_R \text{Tr} [\Sigma Y_1 \chi_L] + \text{h.c.} \end{aligned}$$

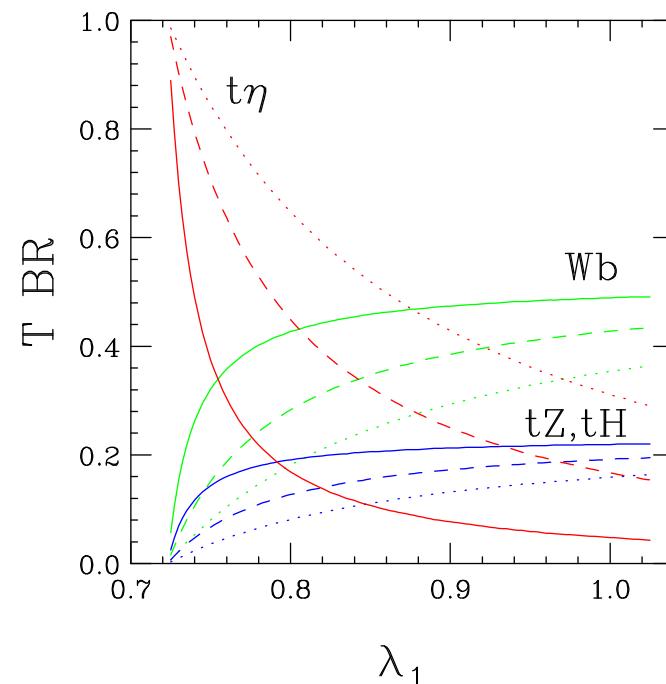
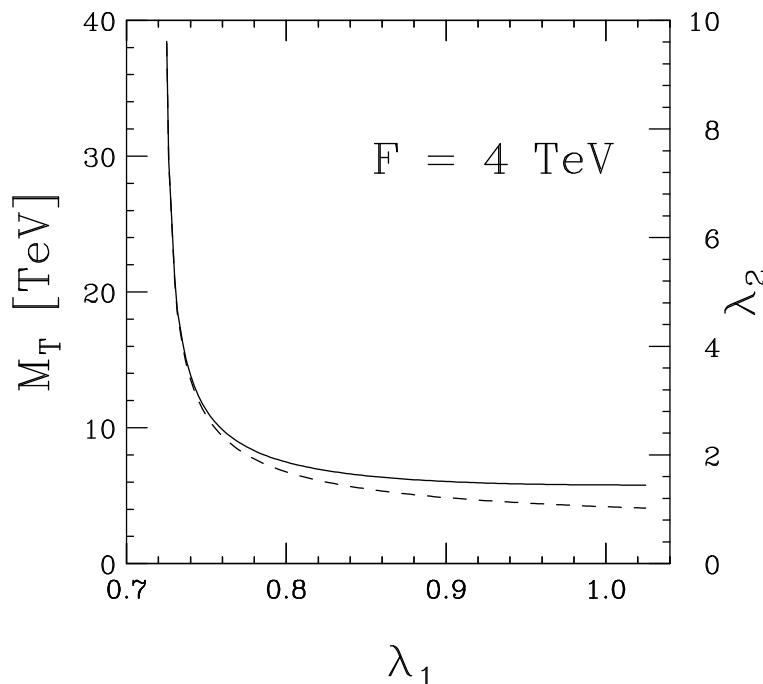
β_i are real #s, diff.'s of $U(1)_\eta$ fermion hypercharges

Littlest Higgs pseudo-axion – what do we know?

- mass is free parameter ($>$ few GeV to avoid P.-Q. axion limits)
- β_i are free parameters
- no tree-level gauge boson couplings; loop-induced $gg, \gamma\gamma$ decays
- t, T couplings to H, η complementary & alters T decays!

$\bar{T}TH$	$O(\frac{v}{F})$
$\bar{T}tH$	$O(1) \mathcal{P}_L + O(\frac{v}{F}) \mathcal{P}_R$
$\bar{t}tH$	$O(1)$

$\bar{T}T\eta$	$O(1) \gamma_5$
$\bar{T}t\eta$	$O(1) \mathcal{P}_R + O(\frac{v}{F}) \mathcal{P}_L$
$\bar{t}t\eta$	$O(\frac{v}{F}) \gamma_5$



The μ -model pseudo-axion

$SU(3) \times U(1)_Y \rightarrow SU(2)_L \times U(1)_Y$ via 2 non-linear sigma fields:

$$\Phi_1 = \exp\left[i\frac{F_2}{F_1}\Theta\right] \begin{pmatrix} 0 \\ 0 \\ F_1 \end{pmatrix}, \quad \Phi_2 = \exp\left[-i\frac{F_1}{F_2}\Theta\right] \begin{pmatrix} 0 \\ 0 \\ F_2 \end{pmatrix},$$

$$\Theta = \frac{1}{F} \left\{ \frac{\eta}{\sqrt{2}} + \begin{pmatrix} 0 & 0 & h \\ 0 & 0 & \\ h^\dagger & & 0 \end{pmatrix} \right\}, \quad F^2 = F_1^2 + F_2^2$$

- Little Higgs sym. broken radiatively (Yukawas)
→ global $U(1)_\eta$ symmetry left intact
- $U(1)_\eta$ breaking via μ -term: $V = -\mu^2 \Phi_1^\dagger \Phi_2 + \text{h.c.}$
- m_η predicted! $\Rightarrow m_\eta = \sqrt{\kappa} \mu$, $\kappa = \frac{F_1}{F_2} + \frac{F_2}{F_1} \geq 2$
- Higgs mass related: $m_H^2 = -2(\delta m^2 + m_\eta^2)$
→ δm^2 from 1-loop Coleman-Weinberg potential

μ -model has additional coupling: $ZH\eta$

Arises from kinetic terms for fields $\Phi_{1,2}$:

$$\mathcal{L} = (D_\mu \Phi_1)^\dagger (D^\mu \Phi_1) + (D_\mu \Phi_2)^\dagger (D^\mu \Phi_2)$$

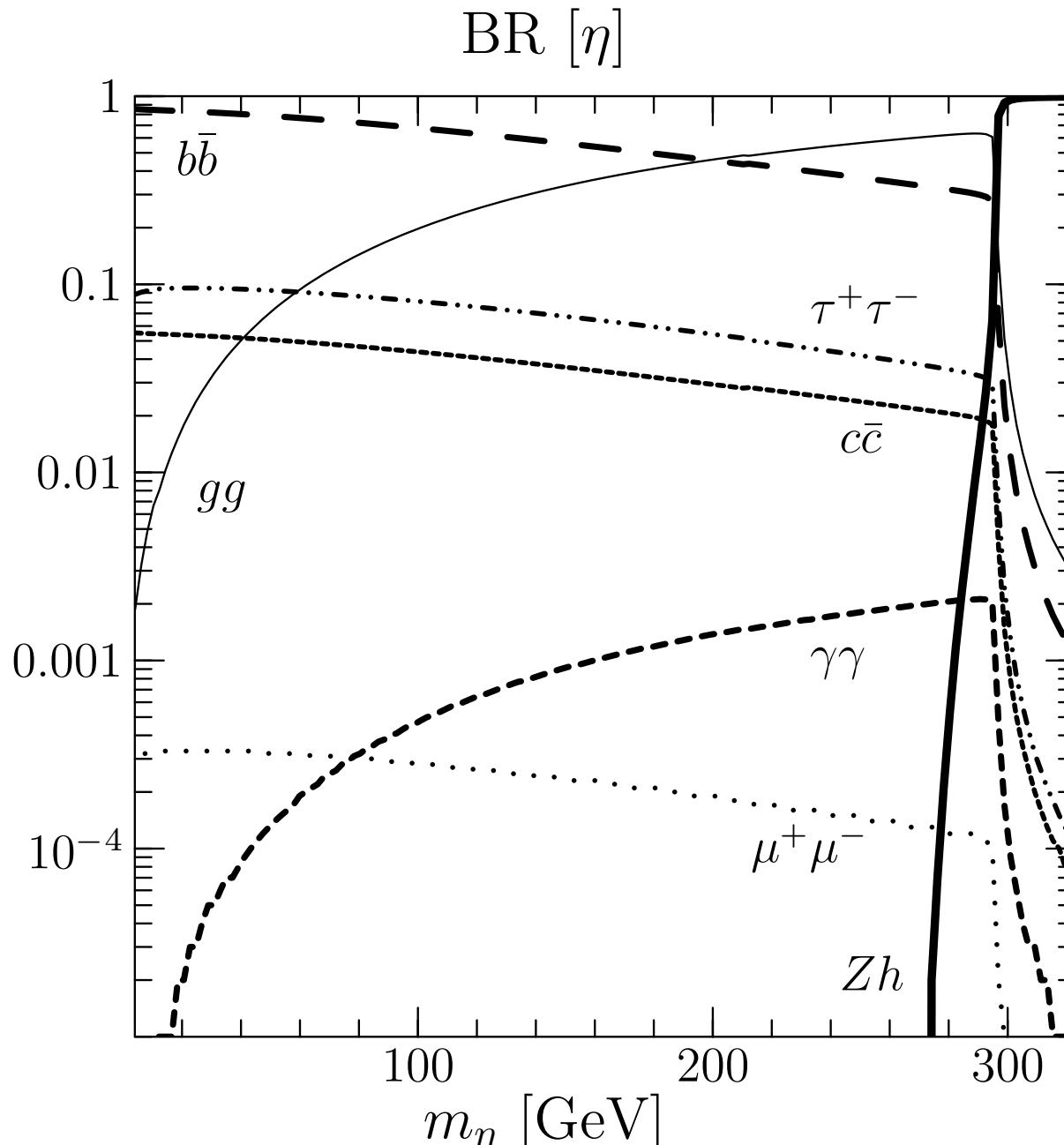
Leads to $ZH\eta$ derivative coupling:

$$\mathcal{L}_{ZH\eta} = \frac{m_Z}{\sqrt{2}F} N_2 Z_\mu (\eta \partial^\mu H - H \partial^\mu \eta), \quad N_2 = \frac{F_2^2 - F_1^2}{F_1 F_2}$$

- ▶ coupling forbidden for unbroken EW symmetry
(here, $SU(3)$), so must be proportional to v/F ,
but N_2 can be large!
- ▶ drastically changes pheno if kinem. accessible

μ -model η decay patterns

“Golden Point”, varying μ ($10 \lesssim \mu \lesssim 150$ GeV)



PRODUCT GROUP MODELS

V.

SIMPLE GROUP MODELS

Are these two models general observations? If so,

- existence of Z - H - η coupling distinguishes models!

Obvious collider tests:

- Drell-Yan $H\eta$ production
- inclusive $H \rightarrow Z\eta$ or $\eta \rightarrow ZH$ (depending on masses)
- $T \rightarrow t\eta \rightarrow tZH$ or $T \rightarrow tH \rightarrow tZ\eta$ (“)

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Theorem 1:

Product Group models cannot have a Z - H - η coupling

Theorem 2:

Simple Group models must have a Z - H - η coupling if $F_1 \neq F_2$

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$$\mathcal{L}_{\text{kin.}} = \sim f^2 (\partial_\mu \xi) \xi^\dagger \text{Tr} \left[(D^\mu \Sigma)^\dagger \Sigma \right] + \text{h.c.} \sim (\partial_\mu \eta) \text{Im} \text{Tr} \left[(D^\mu \Sigma)^\dagger \Sigma \right]$$

$$D_\mu \Sigma = \partial_\mu \Sigma + W_{1,\mu}^a (T_1^a \Sigma + \Sigma (T_1^a)^T) + W_{2,\mu}^a (T_2^a \Sigma + \Sigma (T_2^a)^T)$$

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5. B fields: $U(1)$ generators in global group also be traceless

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4. Two Φ fields: coupling $\propto \frac{F_2}{F_1} - \frac{F_1}{F_2}$ ($\rightarrow 0$ if $F_1 = F_2$, e.g. T-parity)

ILC PHENOMENOLOGY

(VERY PRELIMINARY)

$$\underline{e^+ e^- \rightarrow Z^* \rightarrow \eta H}$$

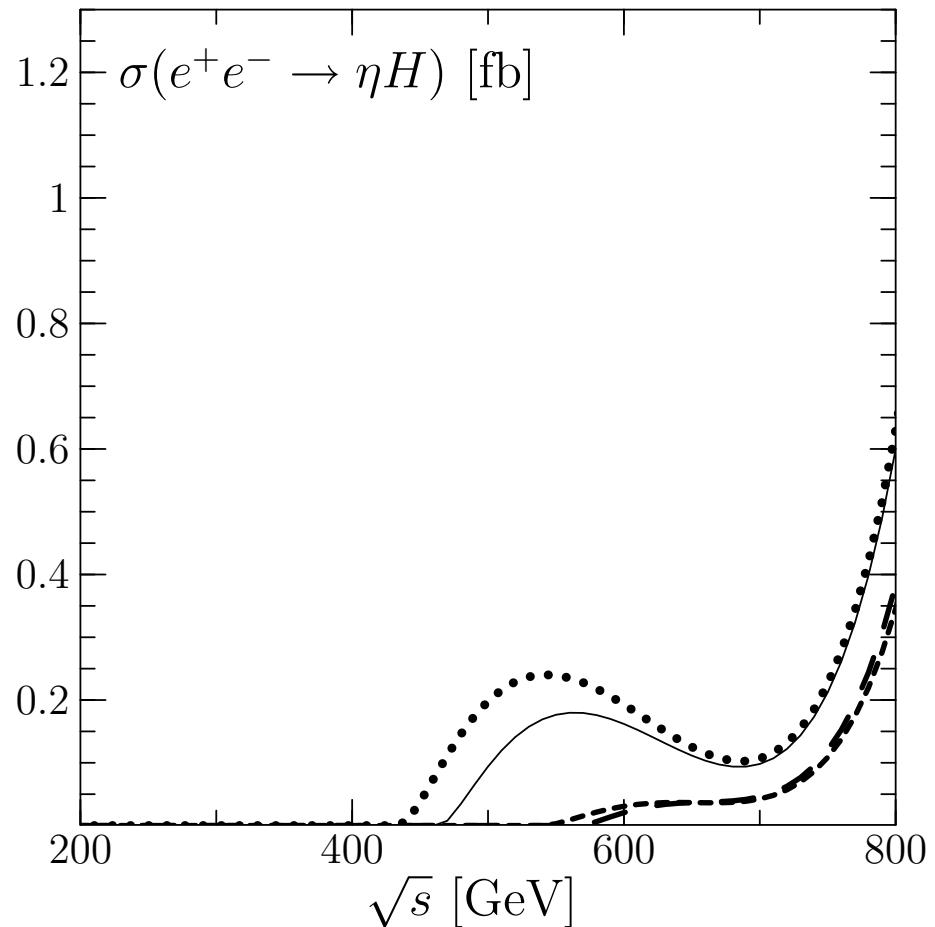
Foreword: Z' - H - η coupling also exists,
but assume the resonance is well above ILC energy

Lots of parameter space, so choose some representative points:

μ [GeV]	m_H [GeV]	m_η [GeV]
154	119	318
150	156	310
97	368	200
49	442	100

→ very small rates!

► strong destructive
interference with Z'



$$\overline{e^+ e^- \rightarrow Z^* \rightarrow \eta H \rightarrow ZHH} \quad (\eta \text{ decays})$$

- “Golden Point” of Schmaltz model: $F_{1,2} = 0.5, 2 \text{ TeV}$:
 $m_T = 1 \text{ TeV}, m_{Z'} = 1.11 \text{ TeV}$

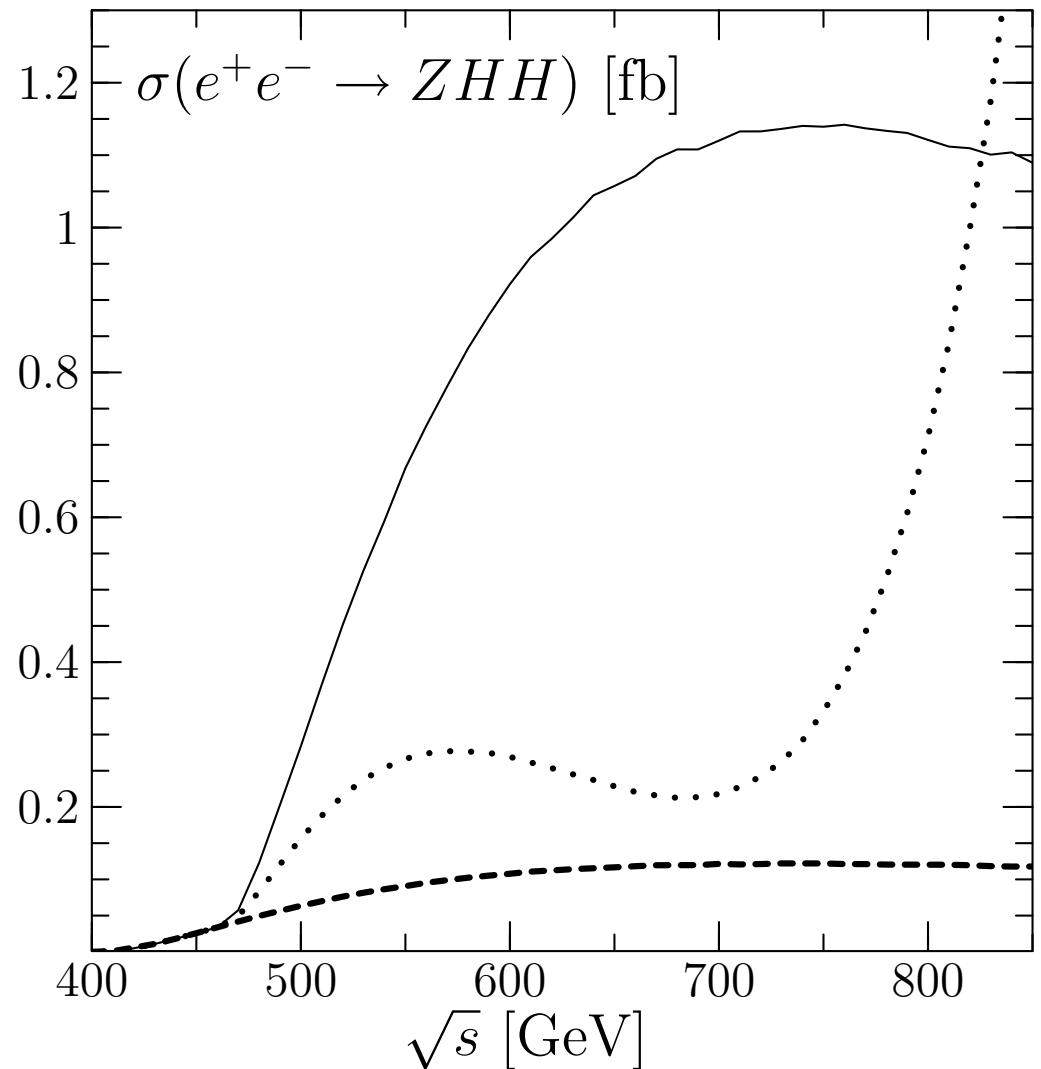
- for $m_\eta > 290 \text{ GeV}$,
 $\eta \rightarrow ZH$ dominates

dashed = SM

dotted = LH

solid = no Z' interf.

- SM ZHH marginal,
but LH has η resonance,
so would be observable



$$\underline{e^+ e^- \rightarrow Z^* \rightarrow \eta H \rightarrow H jj \quad (\eta \text{ decays to } gg)}$$

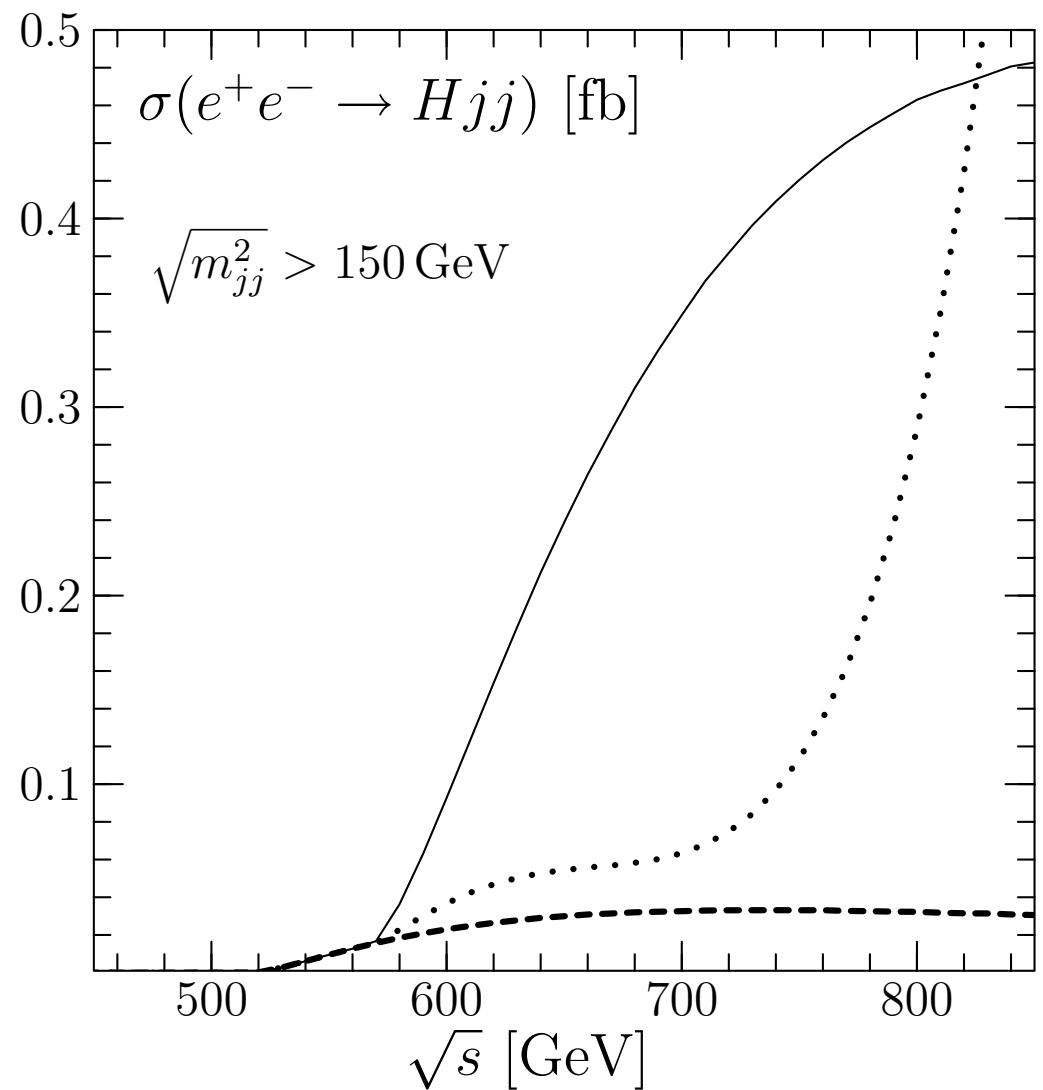
- “Golden Point” again
- for $m_\eta < 290$ GeV,
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- *MUCH* tougher case,
but perhaps possible



SUMMARY

- Little Higgs models generally have extra pseudoscalars
- these pseudo-axions *are not* ruled out by astro
- Product Group and Simple Group models differ:
the latter have Z - H - η , Z' - H - η couplings
(unless T-parity or other reason for $F_1 = F_2$)
- low rates at ILC, but clean environment promising
- LHC phenomenology on the way! (many more options)
 - $gg \rightarrow H \rightarrow Z\eta$
 - $gg \rightarrow \eta \rightarrow ZH$
 - $Z' \rightarrow Z\eta$
 - $T \rightarrow t\eta$