

QCD CORRECTIONS TO HIGGS AND VECTOR-BOSON PRODUCTION

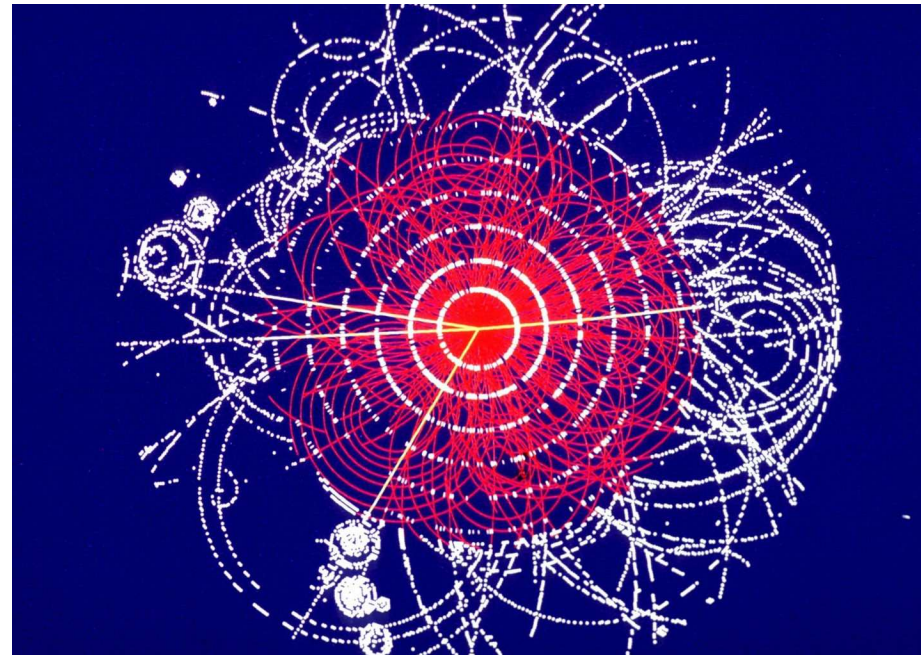
Carlo Oleari

Università di Milano-Bicocca, Milan

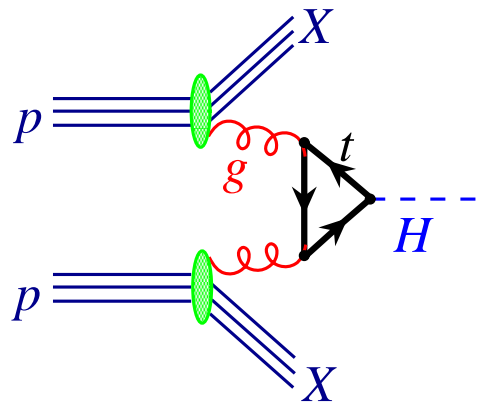
Pheno2006, Madison

16 May 2006

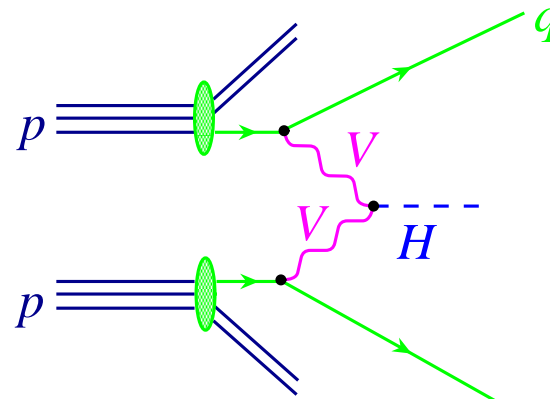
- Higgs production
- QCD corrections to vector-boson-fusion (VBF) production of
 - Higgs boson
 - W/Z
 - WW/ZZ
- Conclusions



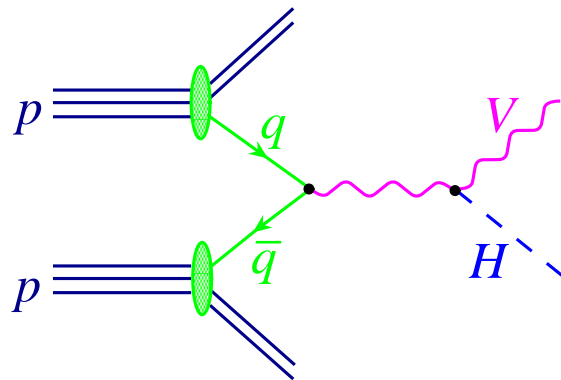
Production Modes



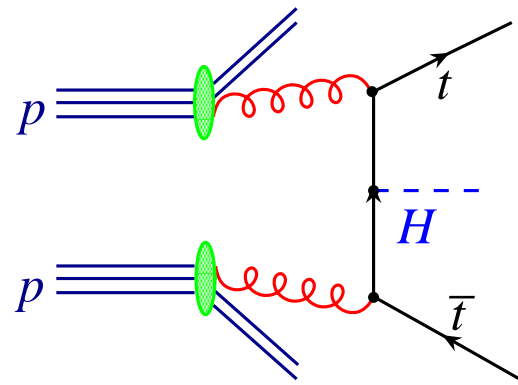
Gluon fusion



Weak-Boson Fusion



Higgs Strahlung



$t\bar{t}H$

Discovery is not the whole story!!

At least as important as the discovery, is the **detailed study** of the **properties** of the Higgs-like resonance: determination of all the quantum numbers and couplings of the state

- mass
- gauge couplings
- Yukawa couplings
- self couplings
- charge
- color
- spin
- CP quantum numbers
- ...

VBF: spontaneous symmetry breaking

$$\mathcal{L}_{\text{Higgs, kin}} = (D_\mu \Phi)^\dagger (D^\mu \Phi) \quad \Phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix}$$

$$D^\mu = \partial^\mu - igW_i^\mu \frac{\sigma^i}{2} - ig' \frac{Y(\Phi)}{2} B^\mu$$

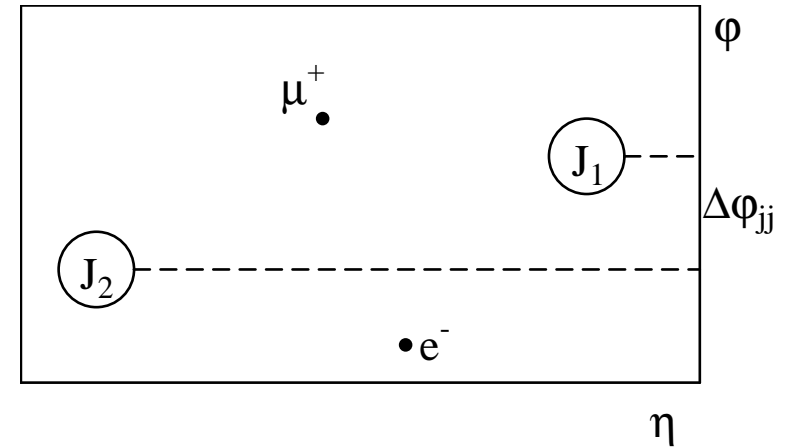
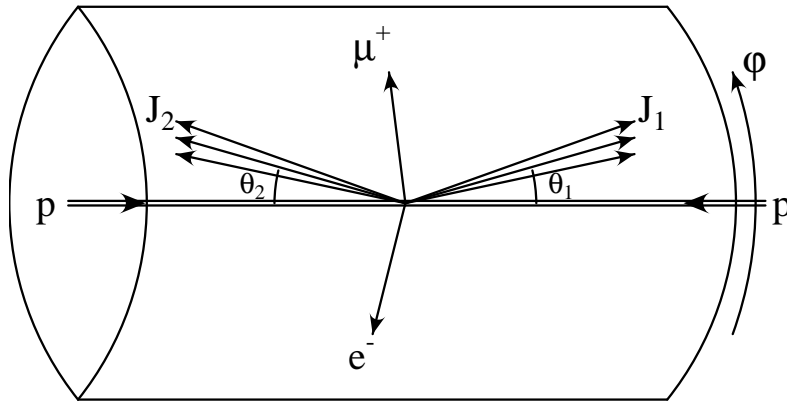
After spontaneous symmetry breaking $\Phi(x) \rightarrow \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v + H(x) \end{pmatrix}$

$$(D^\mu \Phi)^\dagger D_\mu \Phi = \frac{1}{2} \partial^\mu H \partial_\mu H + \left[\left(\frac{g}{2}\right)^2 W^{\mu+} W_\mu^- + \frac{1}{2} \frac{(g^2 + g'^2)}{4} Z^\mu Z_\mu \right] (v + H)^2$$

Tree-level HWW and HZZ couplings require vacuum-expectation value

Gauge interactions of non-VEV scalars are bilinear in Φ

VBF signature



$$\eta = \frac{1}{2} \log \frac{1 + \cos \theta}{1 - \cos \theta}$$

Characteristics and detector requirements:

- energetic jets in the **forward** and **backward** directions ($p_T > 20$ GeV)
- Higgs decay products **between** tagging **jets**
- Little gluon radiation in the central-rapidity region, due to **colorless** W/Z exchange (**central jet veto**: no extra jets with $p_T > 20$ GeV and $|\eta| < 2.5$)

NLO QCD corrections to Higgs production in VBF

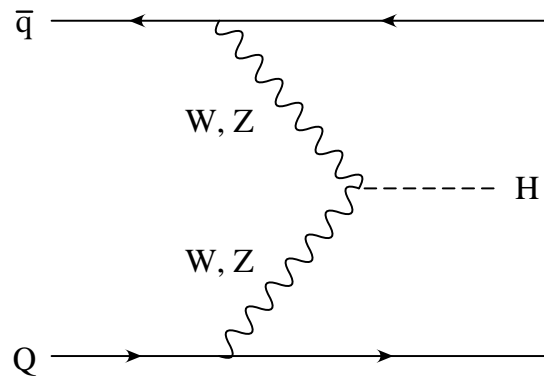
To extract **Higgs-boson coupling constants** with full experimental precision, a **theoretical prediction** of the SM production cross sections with error well **below 10% is required**, and this clearly entails knowledge of the next-to-leading order QCD corrections.

The question then arises whether the **K factors** (the ratio between the next-to-leading and the leading-order cross section) and the **scale dependence**, determined for the **inclusive production cross section**, are valid for less inclusive quantities.

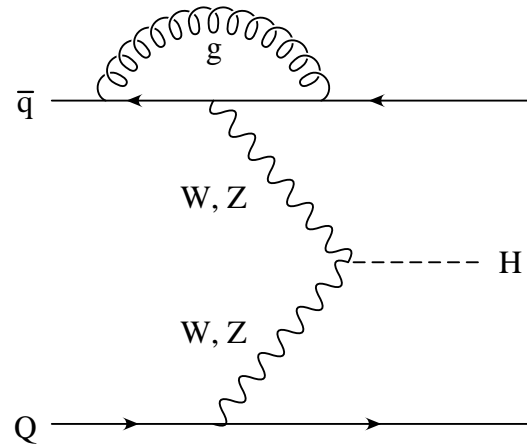
To address this question, we [Figy, C.O. and Zeppenfeld, hep-ph/0306109] have implemented the **QCD corrections** to VBF in a **fully flexible NLO parton-level Monte Carlo** program: **VBFNLO**.

NLO QCD correction diagrams

Leading order diagrams

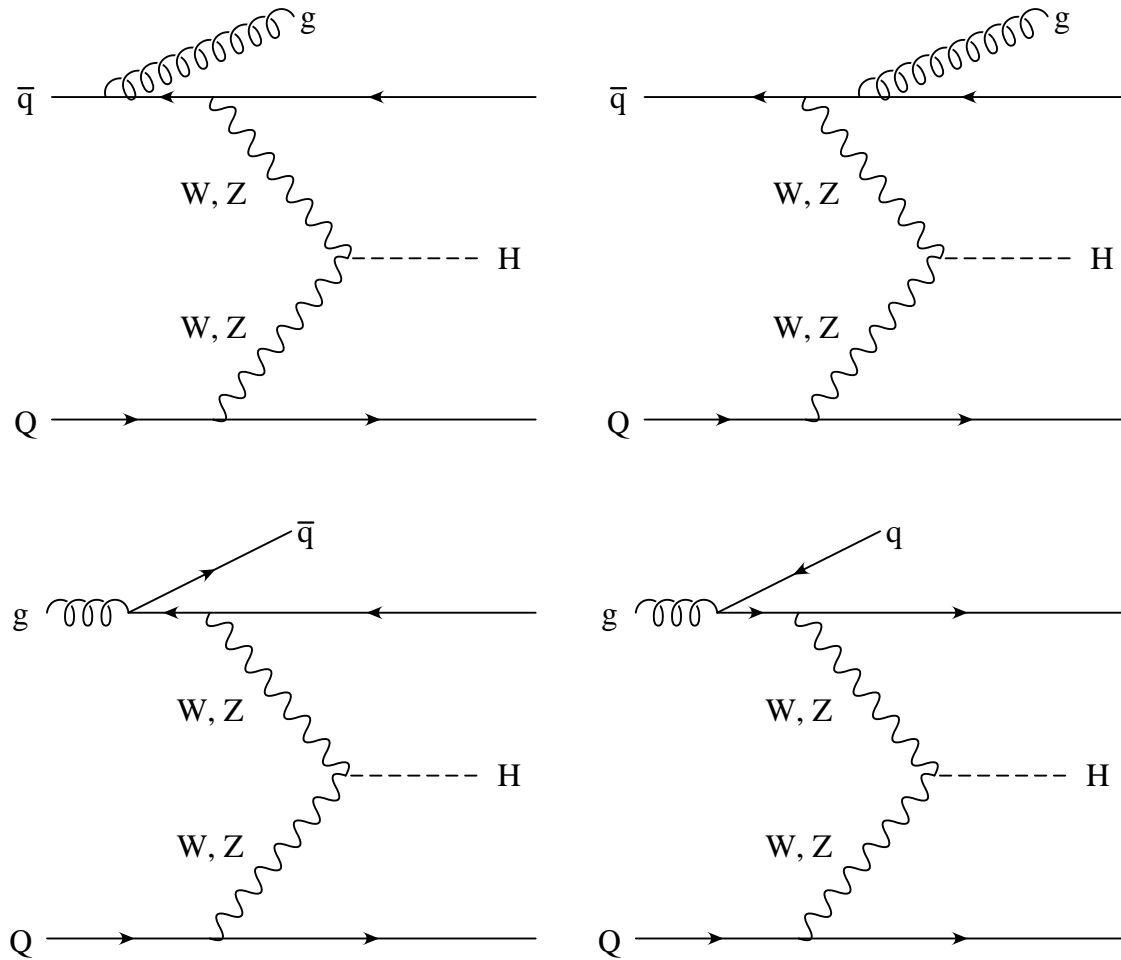


NLO: virtual diagrams



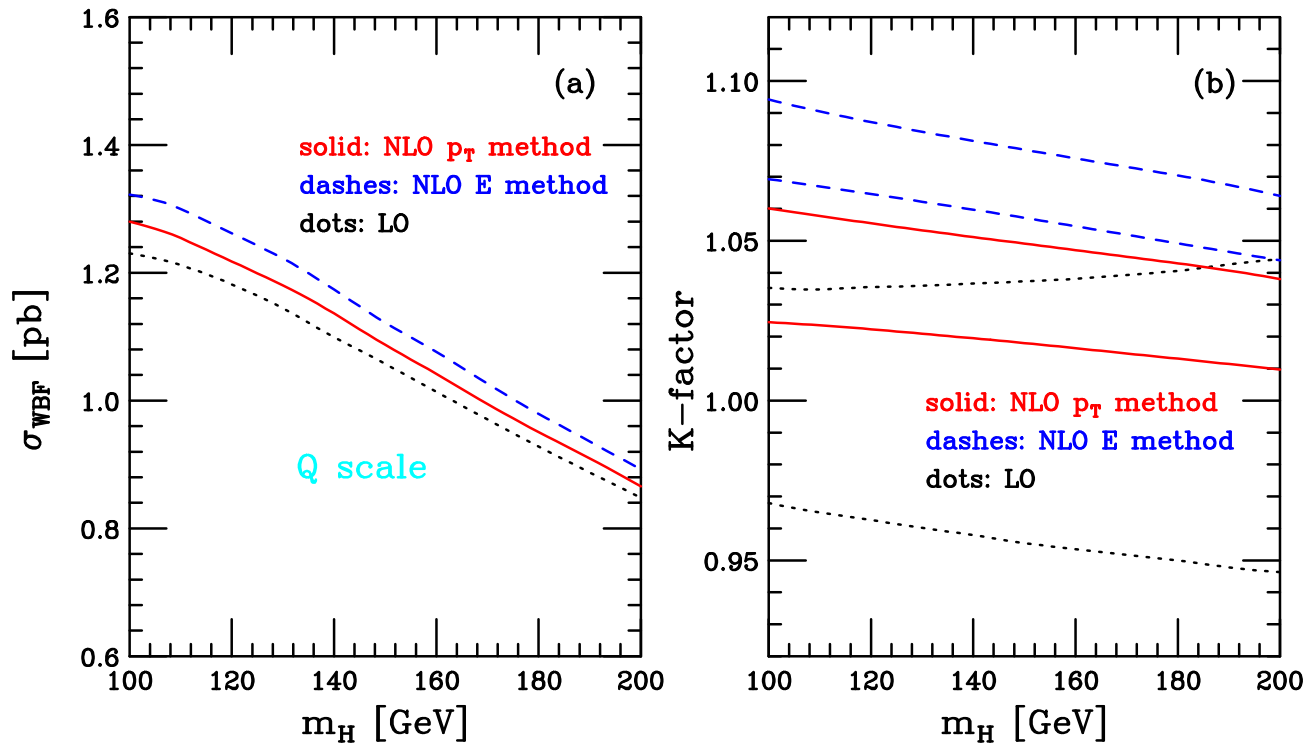
NLO QCD correction diagrams

NLO: real diagrams



plus **crossed processes**: $\bar{q} \rightarrow q$, and/or $Q \rightarrow \bar{Q}$

Total cross section (LHC)



[Figy, C.O. and Zeppenfeld, hep-ph/0306109]

$$K = \frac{\sigma(\mu_R, \mu_F)}{\sigma^{LO}(\mu_F = Q)}$$

To test **scale dependence** we use two different scales

$$\mu_F = \xi_F m_H \quad \mu_R = \xi_R m_H$$

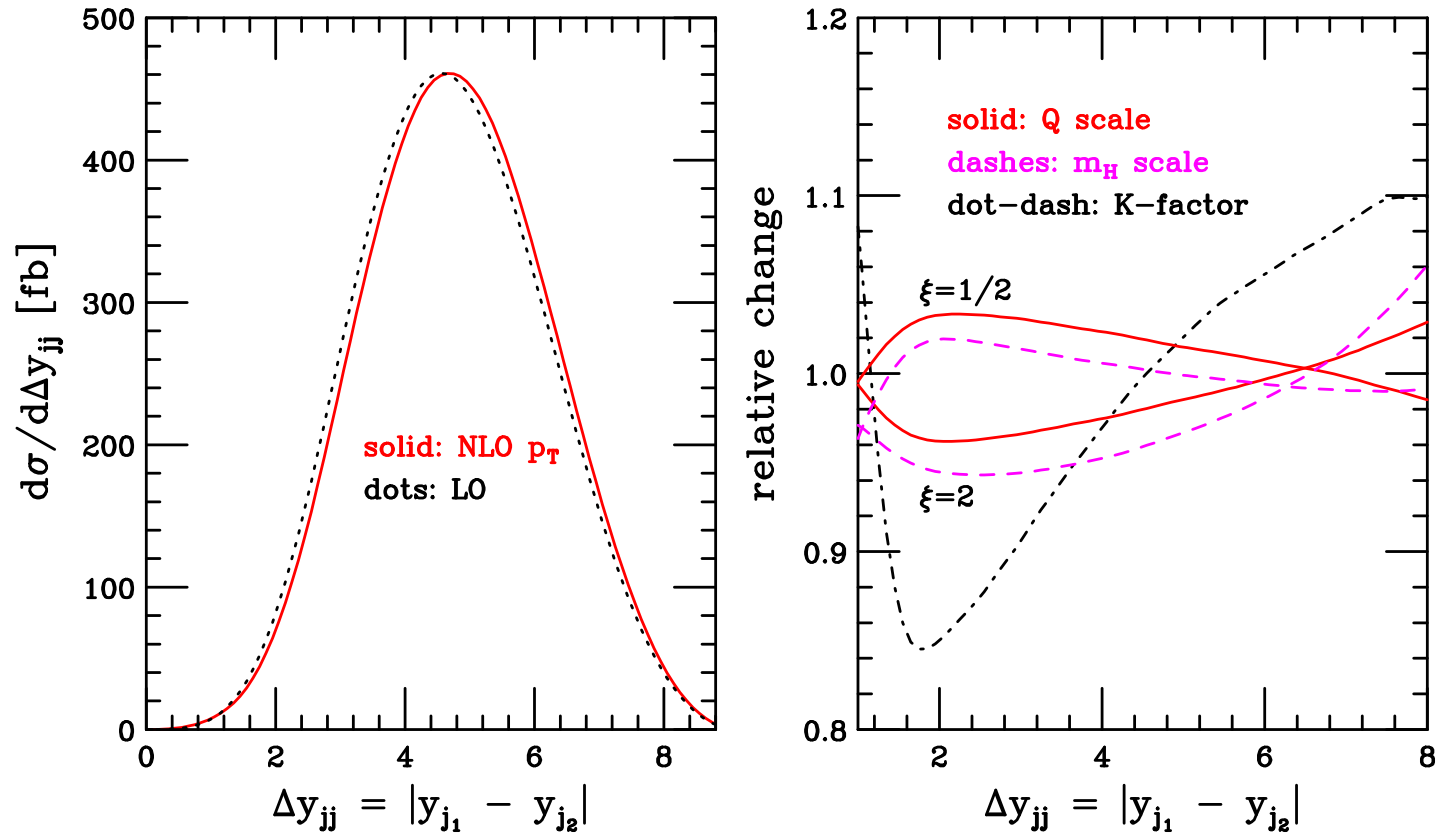
$$\mu_{Fi} = \xi_F Q_i \quad \mu_{Ri} = \xi_R Q_i$$

Q_i = virtuality of the exchanged weak boson (on upper or lower quark line)

The largest scale variations when we vary the **renormalization** and the **factorization scale** in the same direction: $\xi = \xi_R = \xi_F$ with $1/2 \leq \xi \leq 2 \implies \pm 2\%$ only

A uniform $\pm 3.5\%$ PDF uncertainty of the total cross section over the entire range of m_H .

Does the rapidity gap survive?



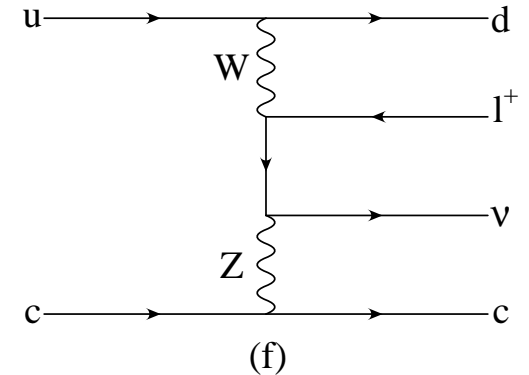
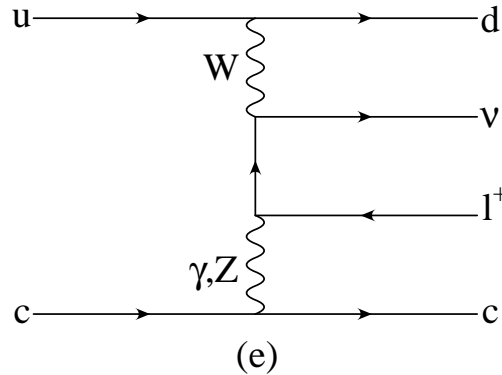
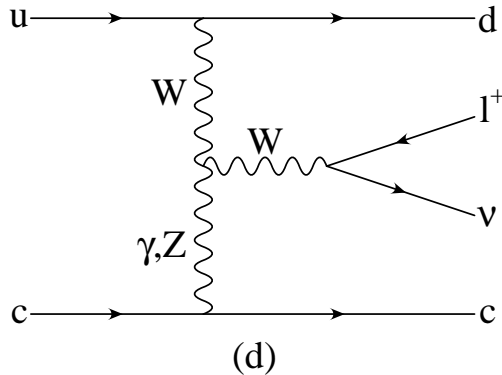
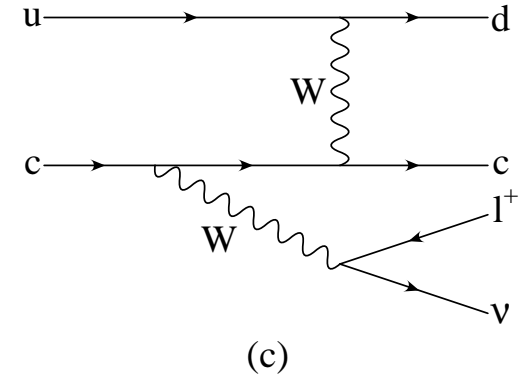
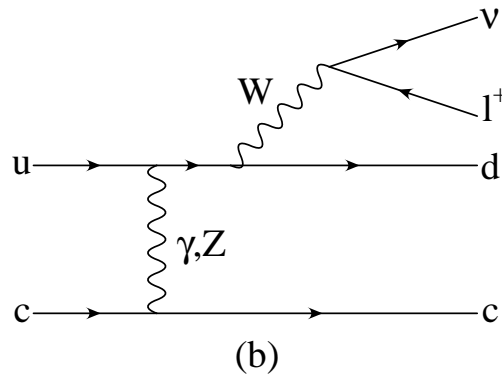
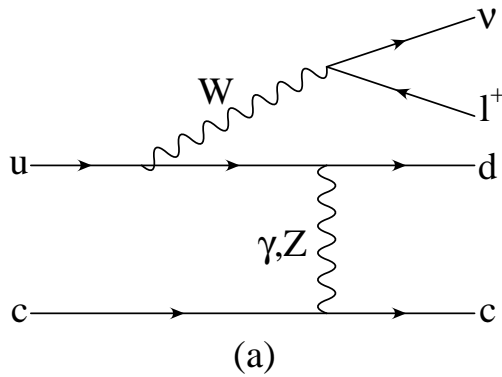
Tagging jets are slightly more forward at NLO than at LO

NLO QCD corrections to W and Z production via VBF

The production of $\ell\nu_\ell jj$ or $\ell^+\ell^- jj$ is another **important background** to the Higgs boson search in vector-boson fusion (VBF) at the LHC.

- $\tau^+\tau^- jj$ is a background to $H \rightarrow \tau^+\tau^-$ and $H \rightarrow W^+W^-$, when W 's and τ 's decay leptonically.
- $\ell\nu_\ell jj$ final state with an unidentified charged lepton, or $\nu_\ell\bar{\nu}_\ell jj$ events from $Z \rightarrow \nu_\ell\bar{\nu}_\ell$ decay, form a background to **invisible Higgs boson decay**
 - Higgs to the lightest **neutralinos** or **gravitinos** (in some region of the parameter space of SUSY models, these branching rates are large)
 - in **large extra dimensions**, where the Higgs boson mixes with scalar fields arising from gravity

Leading order diagrams

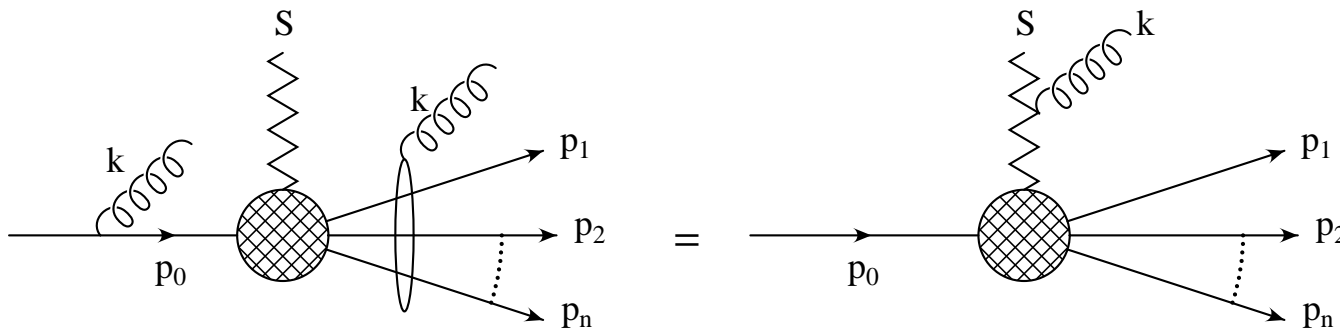


Neglect **annihilation** and **conversion diagrams** (where both the two bosons are time-like): **very suppressed** by VBF cuts.

Why compute NLO QCD corrections?

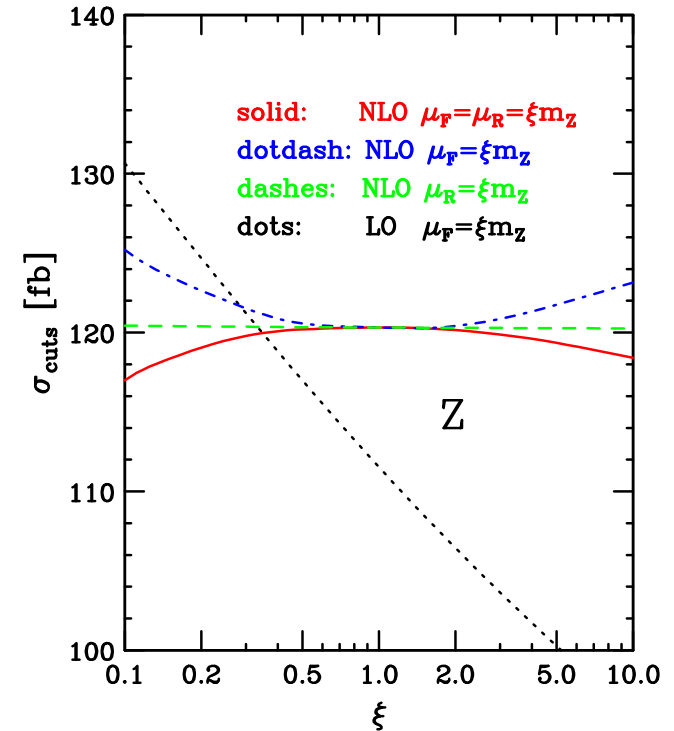
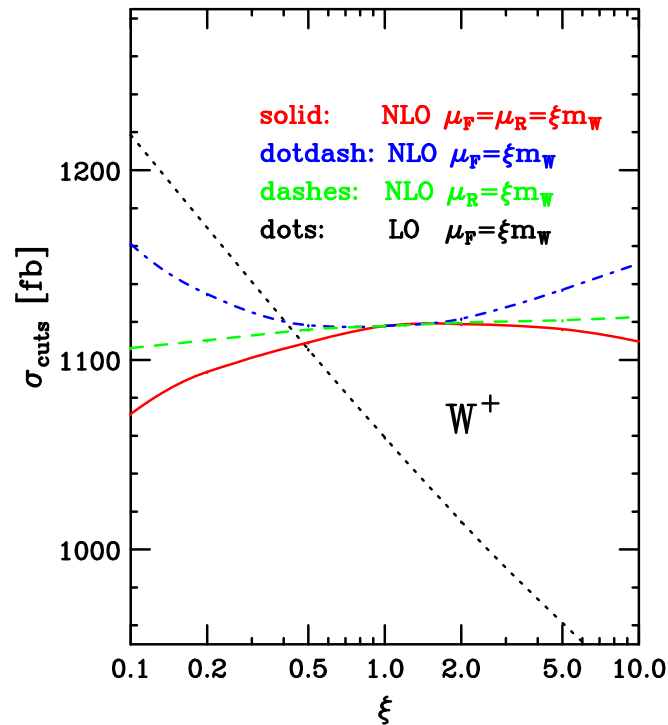
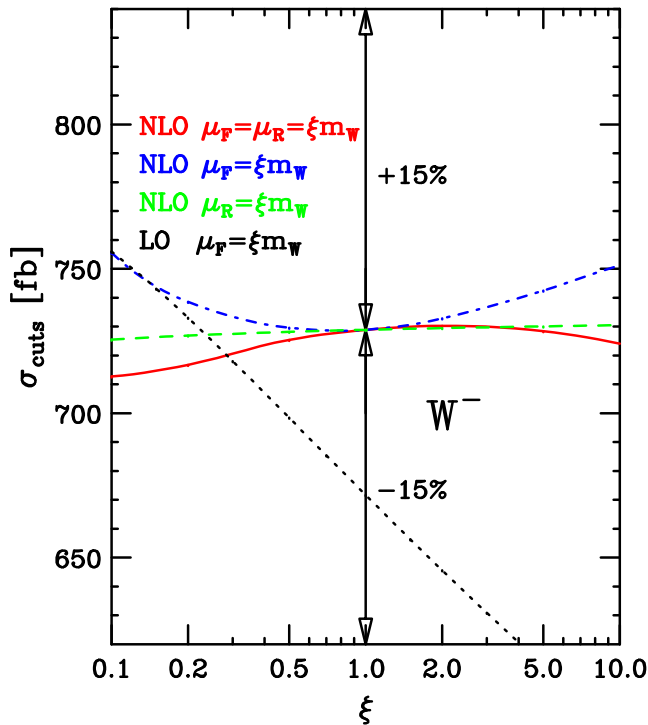
To exploit W and Z production via VBF as **calibration processes** for Higgs boson production

- as a tool to **understand the tagging of forward jets**
- to investigate **veto** of additional **central jets** in VBF: no extra jets with $p_T > 20$ GeV and $|\eta| < 2.5$



Color coherence and **angular ordering**

Total cross sections (LHC)



$$\mu_F = \xi_F m_V$$

$$\mu_R = \xi_R m_V$$

$$\mu_{Fi} = \xi_F Q_i$$

$$\mu_{Ri} = \xi_R Q_i$$

Considering the range $0.5 < \xi < 2$, the NLO cross sections change by **less than 1%** in all cases.

[C.O. and Zeppenfeld, hep-ph/0310156]

Results

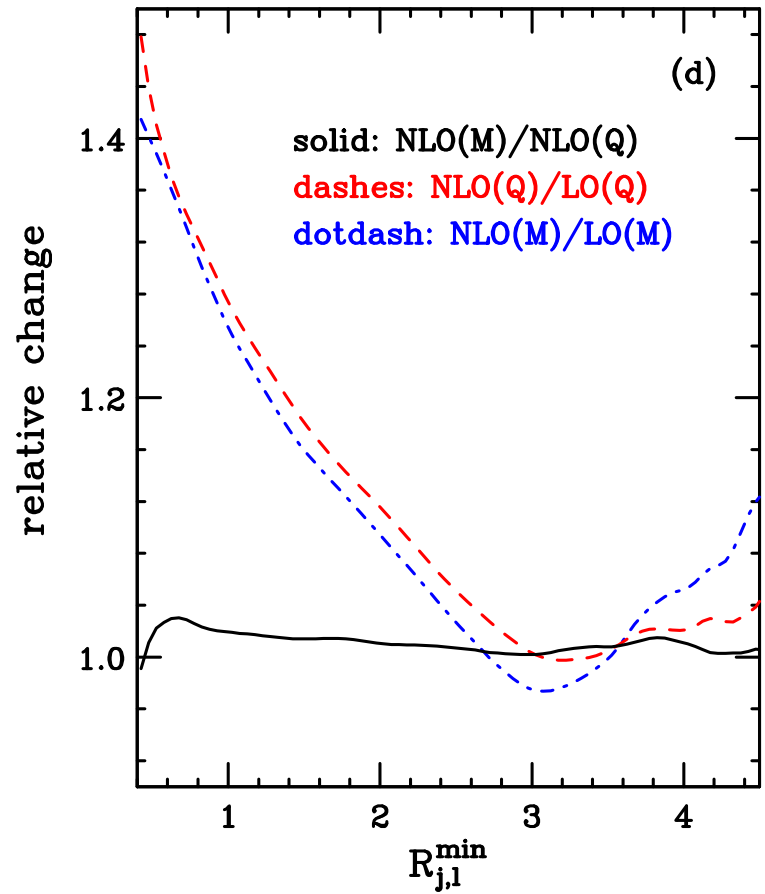
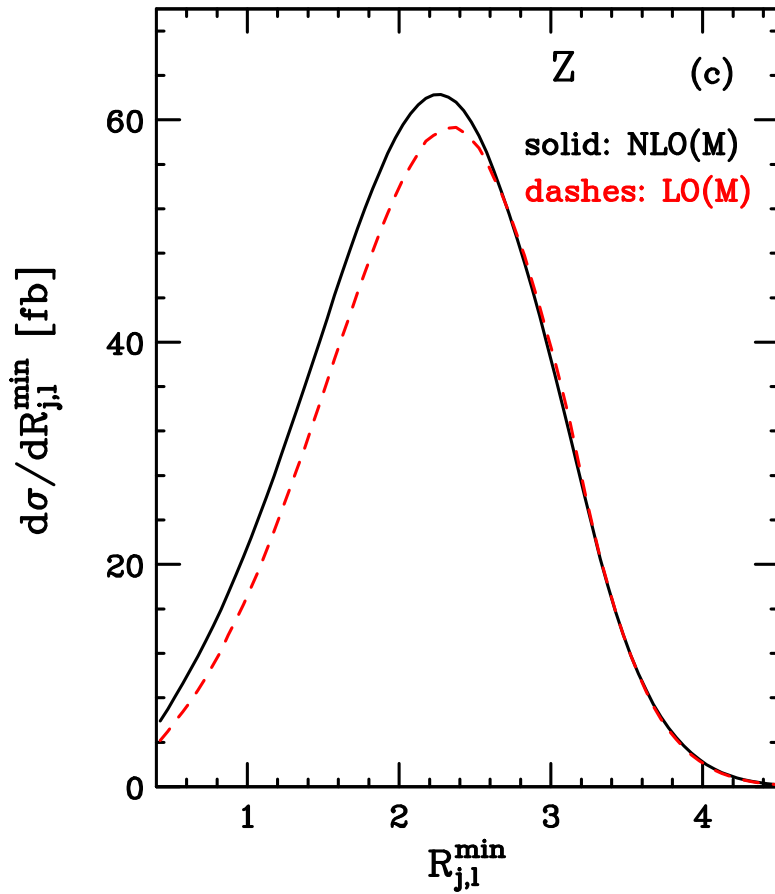
- PDF scale dependence of the total cross section within VBF cuts less than $\pm 4\%$

- *K* factors

$$K(x) = \frac{d\sigma_{NLO}/dx}{d\sigma_{LO}/dx}$$

flat for all the distributions we have checked, and QCD corrections affect distributions for less than a few percent.

Angular correlations of leptons and jets



$$R_{jl} = \sqrt{\Delta\eta_{jl}^2 + \phi_{jl}^2}$$

Additional **parton emission** at **NLO** reduces lepton **isolation**.

NLO QCD corrections to W^+W^- and ZZ production via VBF

An important **irreducible background** to Higgs searches at the LHC, in particular to the search for $H \rightarrow VV$ ($V = W, Z$) decays in VBF, is caused by continuum VV production, or more precisely, by

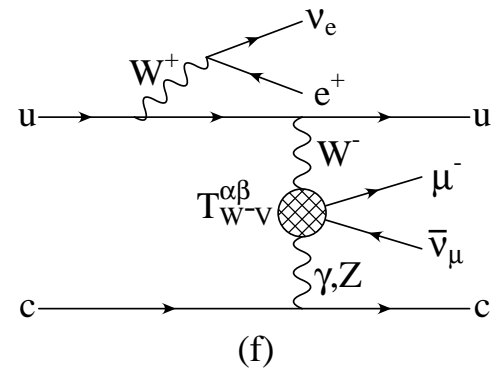
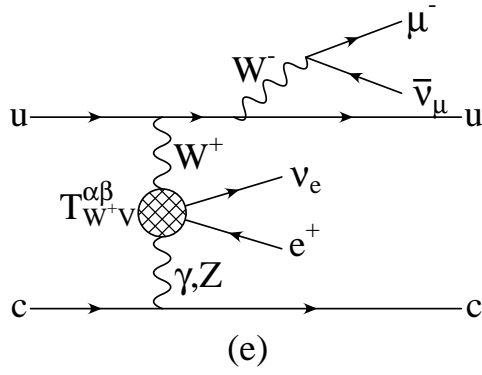
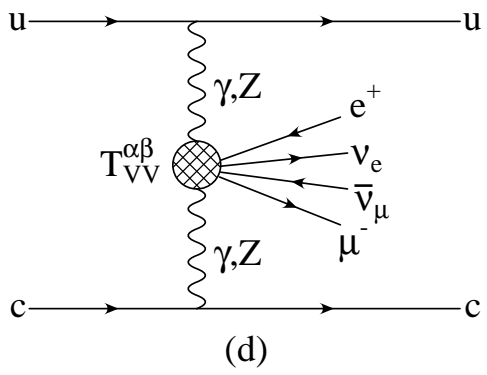
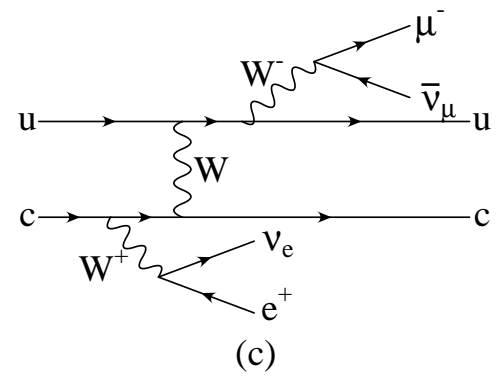
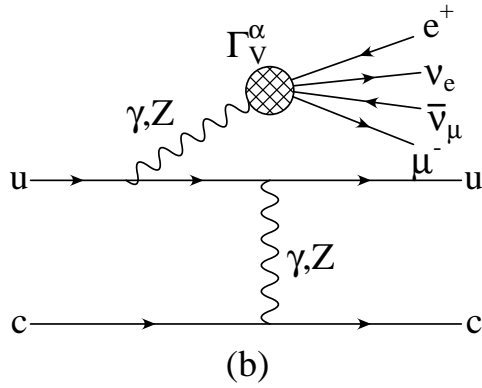
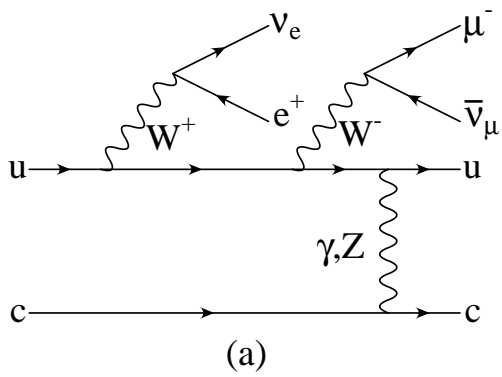
$$pp \rightarrow e^+ \nu_e \mu^- \bar{\nu}_\mu + 2 \text{ jets} \quad \text{and} \quad pp \rightarrow e^+ e^- \mu^+ \mu^- + 2 \text{ jets} \quad \text{or} \quad pp \rightarrow e^+ e^- \nu_\mu \bar{\nu}_\mu + 2 \text{ jets}$$

- ✗ similar **tagging-jet** and **leptonic distributions**
- ✗ suppression of **gluon radiation** in the central region (due to the t -channel color-singlet exchange nature of the VBF process)

Compute **QCD corrections** [Jäger, C.O. and Zeppenfeld, hep-ph/0603177 and hep-ph/0604200] since:

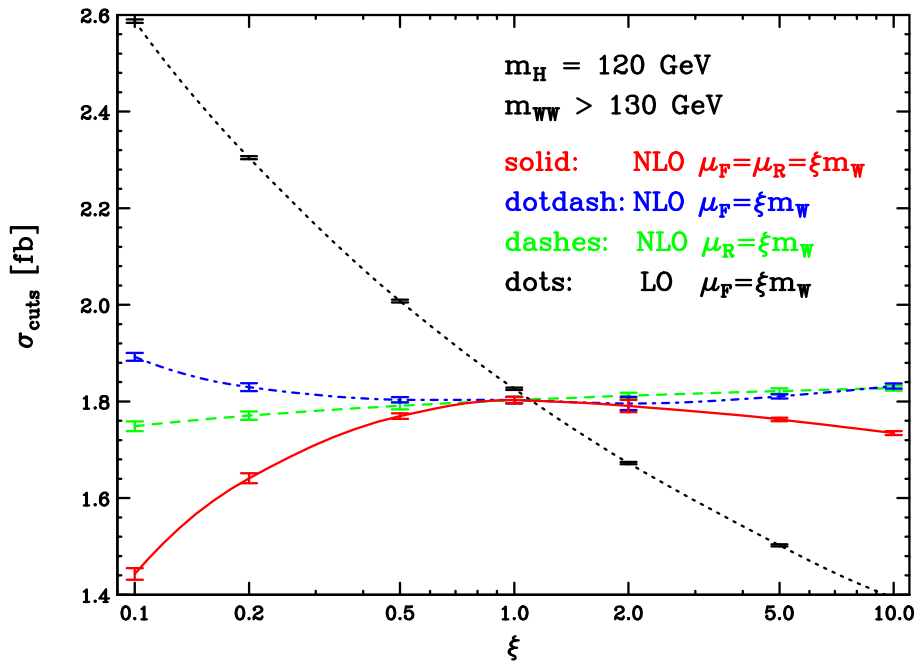
- ✓ W^+W^- contribution ranges between 15% and 3.5% of the Higgs signal, for Higgs boson masses between 115 and 160 GeV \implies important in the extraction of **Higgs boson couplings**
- ✓ one can use weak-boson scattering processes $VV \rightarrow V'V'$, and more precisely the **absence** of **strong enhancements** in these cross sections, as a **probe** for the **existence** of a **light Higgs boson**.

Leading order diagrams



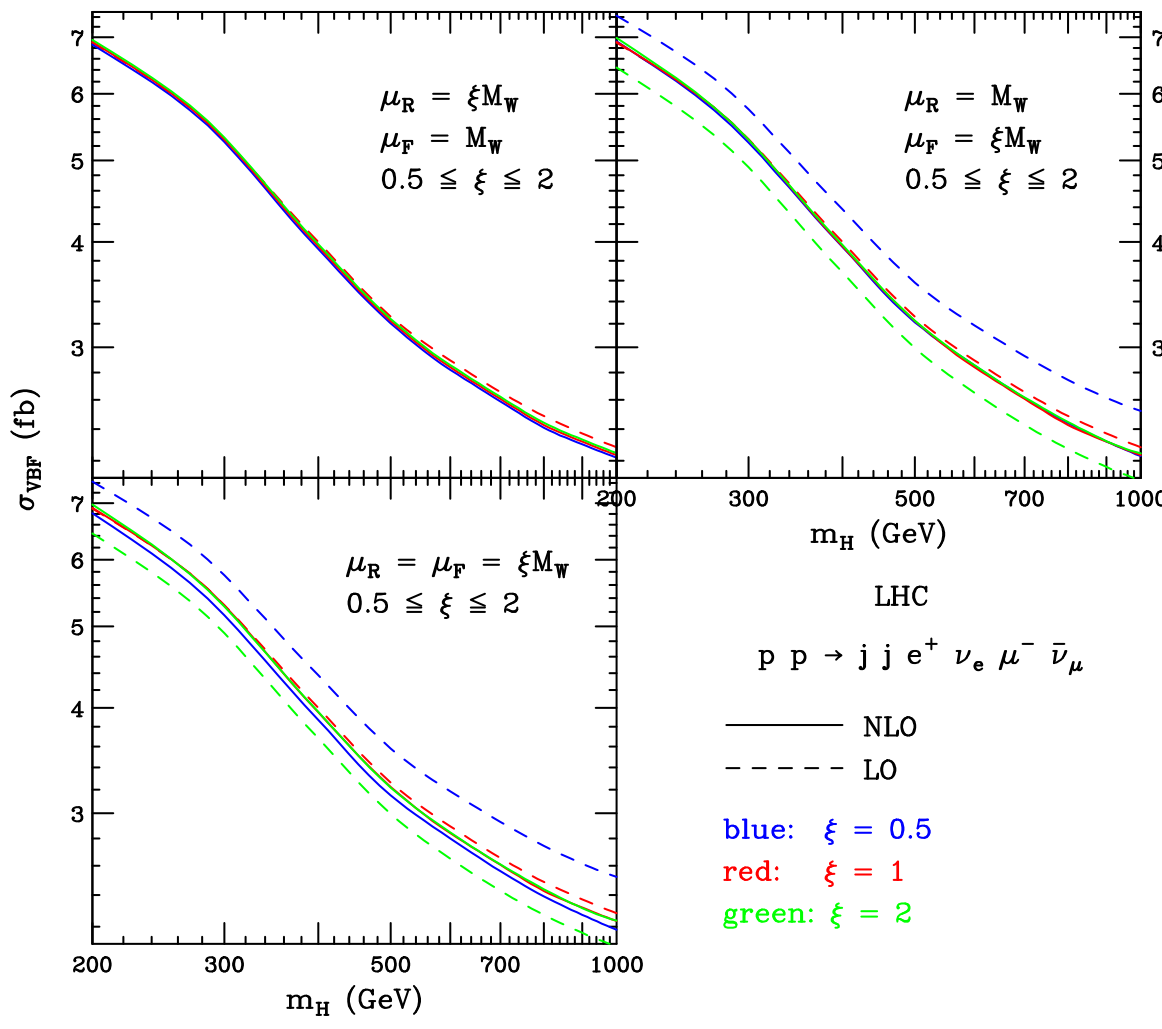
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Total cross section (LHC)

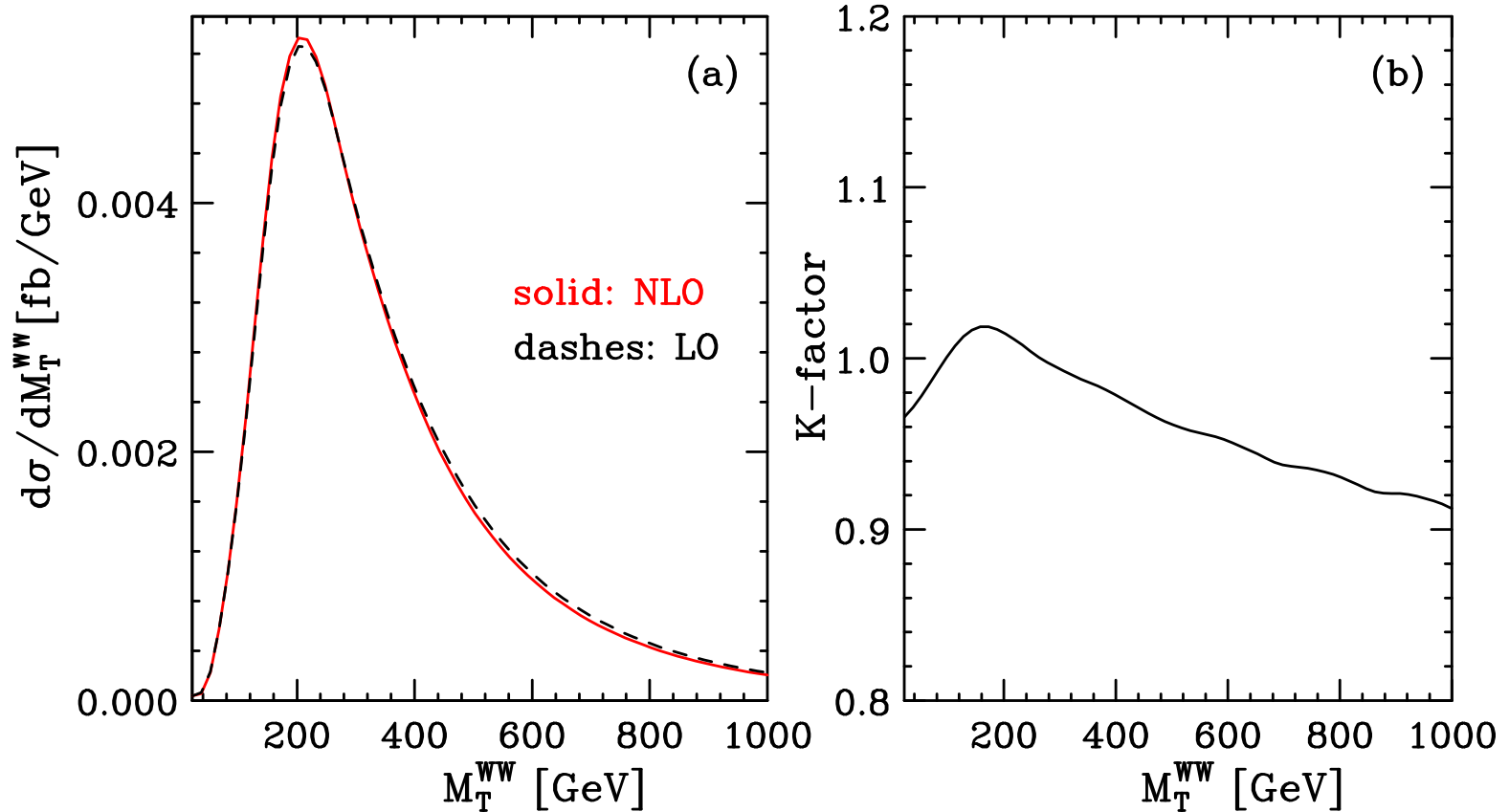


Scale variation less than 2%

$$m_{WW} = \sqrt{(p_e + p_\mu + p_{\nu_e} + p_{\nu_\mu})^2}$$

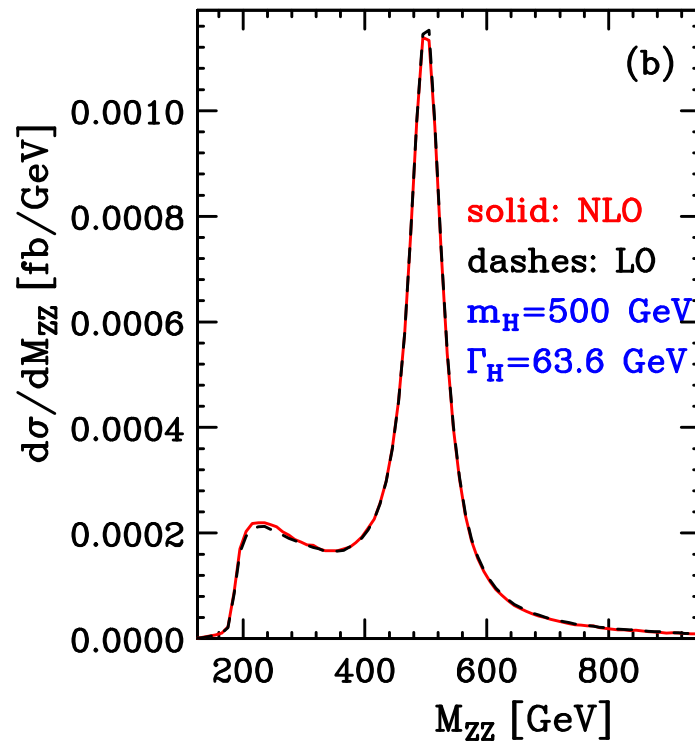
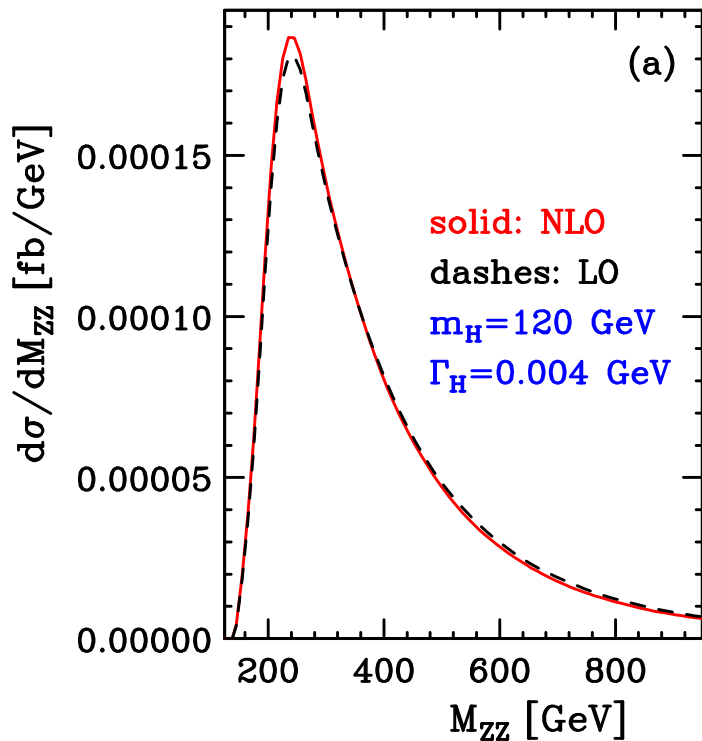


Transverse mass distribution for $e^+ \nu_e \mu^- \bar{\nu}_\mu$



$$M_T^{WW} = \sqrt{(\cancel{E}_T + E_{T,\ell})^2 - (\mathbf{p}_{T,\ell} + \mathbf{p}_T)^2} \quad E_{T,\ell} = \sqrt{\mathbf{p}_{T,\ell}^2 + m_{\ell}^2} \quad \cancel{E}_T = \sqrt{\mathbf{p}_T^2 + m_{\nu\nu}^2} \approx \sqrt{\mathbf{p}_T^2 + m_{\ell}^2}$$

Lepton-mass distribution for $e^+e^- \mu^+ \mu^-$



- $\mu_F = \mu_r = m_Z$
 - $\mu_F = \mu_r = m_{ZZ}$
 - $\sigma^{\text{LO}} / \sigma^{\text{LO}} = 1.8$
 - $\sigma^{\text{NLO}} / \sigma^{\text{NLO}} = 1.13$
- at $M_{ZZ} = 1.5$ TeV

$$M_{ZZ} = \sqrt{(p_{e^+} + p_{e^-} + p_{\mu^+} + p_{\mu^-})^2}$$

- ✓ **very sensitive** to a light Higgs boson: **pronounced resonance** behavior for $m_H \lesssim 800$ GeV
- ✓ for $m_H \sim 1$ TeV, the **peak is diluted** ($\Gamma_H \approx 500$ GeV) and the signal is distributed over a wide range in M_{ZZ}

Conclusions

- Once the Higgs boson has been found and its mass determined, the measurement of its **couplings to gauge bosons and fermions** will be of main interest. Here **vector-boson fusion** will be of **central importance** since it allows for independent observation in the $H \rightarrow \tau\tau$, $H \rightarrow WW$ and $H \rightarrow \gamma\gamma$.
- These measurements can be performed at the **LHC** with **statistical accuracies** on the measured cross sections times decay branching ratios, $\sigma \times B$, of **order 10% or even better**.
- This clearly requires knowledge of the **next-to-leading order QCD** corrections for **signal** and **backgrounds**. These **corrections**, in the case of H , W/Z and WW/ZZ production in VBF processes, are in general **small**, but distributions at LO and NLO have **different shapes!**
- **Jet veto** and **forward-jet tagging** still need **more investigation**.