High energy cosmic rays and the potential for new physics discovery

Haim Goldberg Northeastern University, Boston

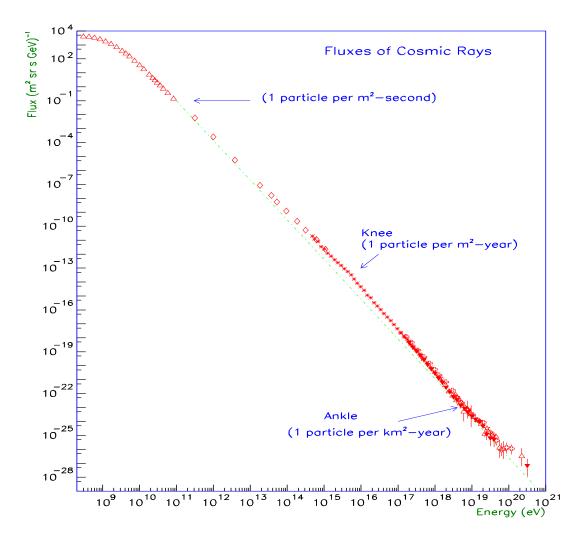
## May 16, 2006 Pheno2006

- Lightening review of cosmic rays: spectrum, sources, losses
- Examples of New Physics discovery:
  - Split supersymmetry
  - High energy neutralinos in cosmic rays (decay of relics)
  - Anomalous neutrino interactions at high energies (general treatment, TeV black holes, sphalerons)
  - Quantum decoherence

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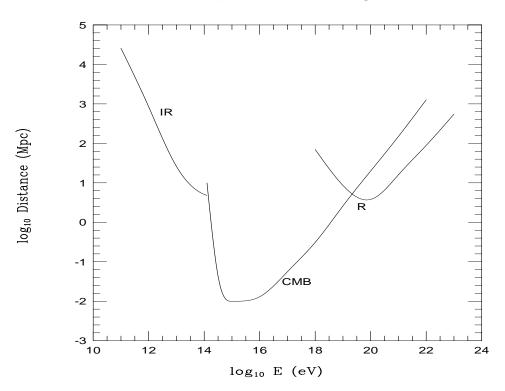
## The cosmic ray spectrum



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### Cosmic ray horizons

- protons pion photoproduction on CMB; GZK horizon at 50–100 Mpc
- nuclei photodisintegration on CMB;  $\sim$  same horizon
- photons pair production off CMB, IR, radio and B; window for distant UHE photons only if  $B < 10^{-12}$  G.



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# Sources

- AGN's (linear Fermi mechanism)
- Accretion shock waves
- Relativistic jets
- Magnetars
- Starburst galaxies (collective effects)
- GRB fireballs
- Top down mechanisms
- + many more

New physics: Split SUSY at Auger

Anchordoqui, Nunez, HG, PRD71:065014,2005

 Split Supersymmetry: fine-tune to small Higgs mass, scalar masses are large, gluino mass O(TeV), long-lived

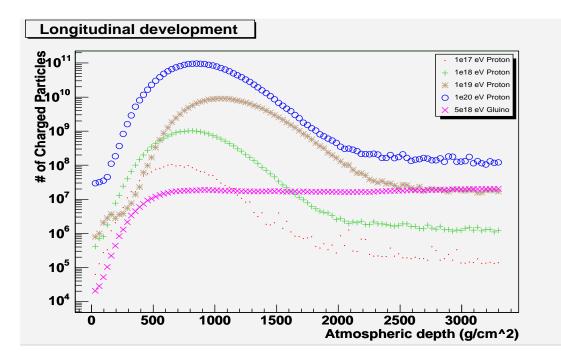
$$au_0 \simeq 25 \,\, {
m yr} \,\, \left(rac{{
m TeV}}{M_{ ilde g}}
ight)^5 \left(rac{M_{
m SUSY}}{10^{11}\,\,{
m GeV}}
ight)^4$$

Arkani-Hamed, Dimopoulos, JHEP 0506:073, 2005

- isotope relic abundance  $\rightarrow M_{
  m SUSY} \lesssim 10^{13} {
  m ~GeV}$
- bound from BBN: decaying  $\tilde{g}$ 's  $\rightarrow$  baryons which break apart <sup>4</sup>He  $\rightarrow$  excess D and Li Arvanitaki et al, hep-ph/0504210  $\rightarrow M_{\rm SUSY} < 10^{12} {
  m GeV}$

**Properties of G air showers** 

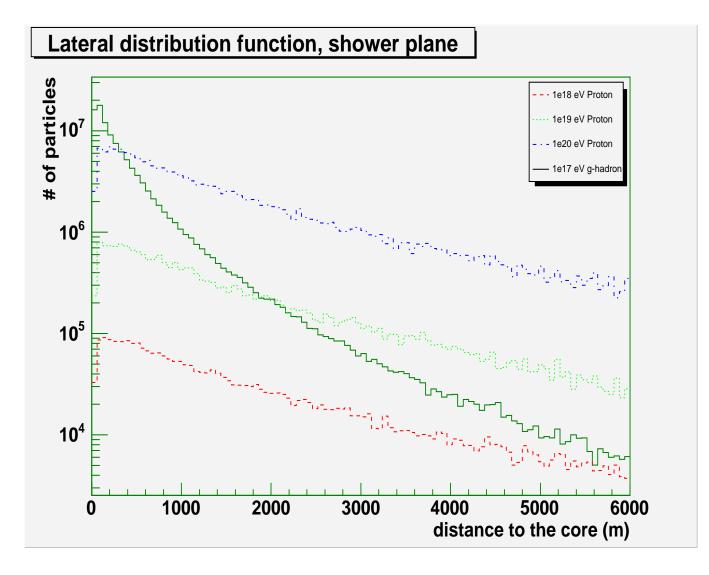
- *G*'s from astrophysical source, air shower in atmosphere. cp Hewett et al hep-ph/0408248: *G*'s created in atmosphere, flux limited for  $M_G > 200$  GeV
- inelasticity  $\sim (1~{
  m GeV}/M_G)$ , minishowers, no distinct  $X_{
  m max}$  J. Gonzalez et al, hep-ph/0504210, 0504260





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## Lateral Shower Profile





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• For one event/yr at PAO, require

$$\int_{E_{G,\,{
m min}}}^{E_{G,\,{
m max}}} J_G(E_G) \, dE_G \, pprox 1.4 imes 10^{-21} \, {
m cm}^{-2} \, {
m s}^{-1} \, {
m sr}^{-1}$$

• Production at source associated with neutrino production via pion decay

$$\int J_
u(E_
u) \ dE_
u = rac{2}{3} \, rac{\sigma_{
m inel}}{\sigma_{pp 
ightarrow G}(\hat{s}_{
m min})} \, rac{\langle N_
u 
angle}{N_G} \, \int J_G(E_G) \ dE_G$$

• Expected neutrino flux  $J_{\nu}(1.5 \times 10^{12} \text{ GeV}) \approx 4.4 \times 10^{-30} \text{ GeV}^{-1} \text{cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1}$ 

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Comparison with limits and a lower bound on  $M_{\rm SUSY}$ 

- Within a factor of 2 of limits set by RICE, and extrapolated cascade bounds (related to diffuse photon background set by EGRET
- If a few events observed in 10 yrs, and expect  $\sim 1~{
  m yr}^{-1}$  on basis of high energy neutrinos observed, can set a lower limit on the lifetime of  $M_G$

$$rac{E_G}{M_G} au_0 > H^{-1} ~
ightarrow ~ au_0 \gtrsim 100 {
m yr}$$

$$ightarrow M_{
m SUSY} > 6 imes 10^{10} {
m ~GeV}$$

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May 16, 2006 Pheno2006 **Detecting ultra-high energy neutralinos** 

Barbot, Drees, Halzen, Hooper, PLB563:132,2003; Anchordoqui, HG, Nath, PRD70:025014,2004

- Detect air showers produced by high energy neutralinos emitted in long-lived relic *X* particle decay
- Underlying process is  $ilde{\chi} + q(ar{q}) o ilde{q}(ar{ ilde{q}}) o ext{all}$
- Weak interaction strength → restrict events to large zenith angle to eliminate hadronic bkgd
- Elimination of neutrino bkgd to be described
- Require large acceptance  $\sim 2400~{\rm km^3}$  we akin to EUSO space-based fluorescence detector



- EUSO at present is planned to dock on JAXA platform on the ISS
- Riken laboratory has agree to pay for increasing eff of photomultipliers
- Most significantly, JAXA has proposed to use their HTV launcher in place of NASA shuttle
- Mature project with high promises

EUSO collaboration, Nucl. Phys. Proc. Suppl. 151, 401 (2006)

Constraints on neutralino flux

- Normalization: X particle density + homogenous population of astrophysical sources fits UHE proton flux
- EGRET bounds: photons resulting from *X* decay conform with EGRET limits on GeV gamma ray flux
- Dark matter: X density consistent with negligible contribution to WMAP dark matter fraction

Eliminating the neutrino background

- For a small range of squark masses, neutralino cross sections are
  - \* sufficiently smaller than neutrino cross sections so that restricting upcoming showers to pass through enough earth would screen out neutrinos
  - \* sufficiently large to give enough events in the atmosphere
- $\bullet$  For bino-type neutralinos,  $m_{\tilde{q}}\simeq 1~{\rm TeV}$  does the job

Event rate at EUSO

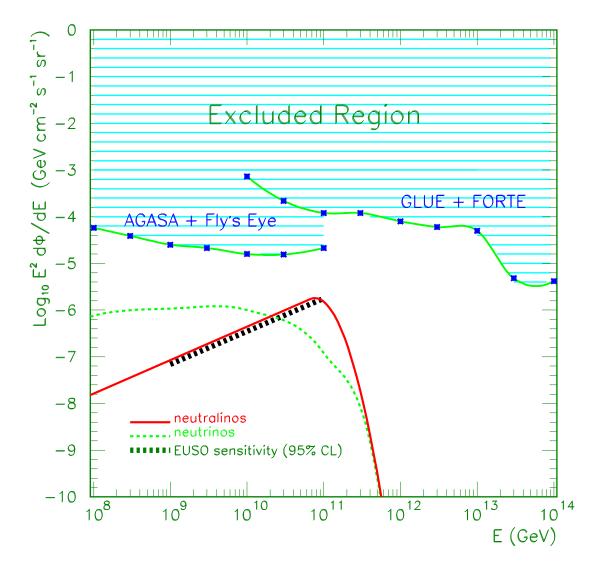
$$\mathcal{N} = \int_{E_{ ilde{\chi}}^{ ilde{\min}}}^{E_{ ilde{\chi}}^{ ext{max}}} dE_{ ilde{\chi}} \; N_A \; P \; rac{d\Phi}{dE_{ ilde{\chi}}} \, \sigma_{_{ ilde{\chi}N}} \, A \; \epsilon_{_{ ext{DC}}} \; t,$$

- Estimate neutrino bkgd is about 0.3 events in 3 years
- Require 3.09 neutralino events for significance at 95% CL
- For parameters given, find  $\sim$  4 or 5 events
- Results in figure to follow

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## EUSO sensitivity to neutralino flux



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IceCube and anomalous neutrino interactions

Anchordoqui, Feng, HG, PRL 96:021101,2006

UHE hadronic cosmic rays  $(E > 10^{10} \text{ GeV})$  expected to be accompanied by UHE neutrinos:

$$p+\gamma 
ightarrow n+\pi^+, \ \ \pi^+ 
ightarrow e^+ 
u_e \ 
u_\mu \ ar
u_\mu$$

or

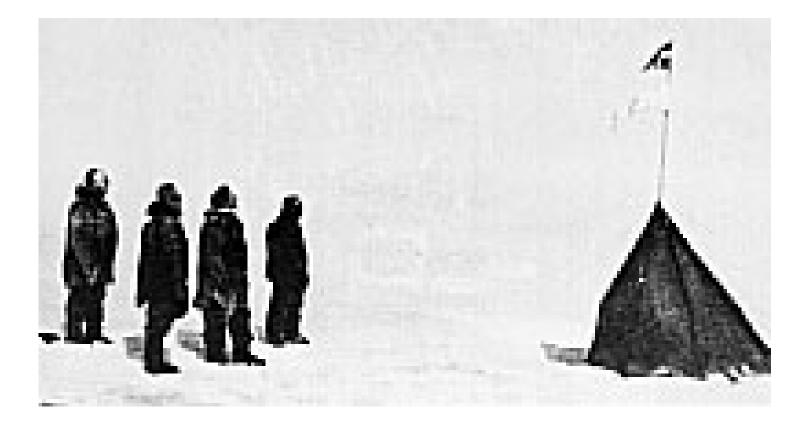
- $p+p 
  ightarrow {
  m nucleons} + \pi$ 's,  $\pi^{\pm} 
  ightarrow e^{\pm} + 
  u$ 's
- For optically thin source (neutrons escape), can relate  $\Phi^{\nu}$  to cosmic ray flux Waxman, Bahcall PRD 59 023002 (1999)
- Find (for *pp*)

$$E^2 \Phi^{\nu}_{\rm WB} \simeq 2 \times 10^{-8} \ {\rm GeV} \ {\rm cm}^{-2} \ {\rm s}^{-1} \ {\rm sr}^{-1}$$

### for each flavor.

• Specialize to  $10^7 \text{ GeV} < E_{\nu} < 10^{7.5} \text{ GeV}$  – minimize atmospheric background, have sufficient flux

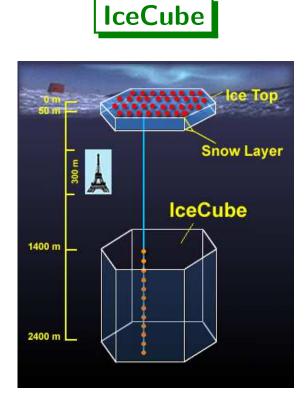
### Neutrino astronomy 1911: early assessment of South Pole site



#### **December 14th**

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#### Halzen, astro-ph/0311004

- Effective Area  $pprox 1 \ {
  m km}^2$
- $E_{\rm th} \approx 100 ~{\rm GeV}$
- 4800 PMT's on 80 strings
- $\mu$ -track angular resolution  $\Rightarrow 1^{\circ} \times 1^{\circ}$  bin
- Calibration ⇔ IceTop
  - $rac{1}{\sim}$  1 km<sup>2</sup> air-shower detector with 160 stations

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# **Tracks:** Cosmic muons and CC interactions of $\nu_{\mu}$ 's Angular resolution $1^{\circ} \times 1^{\circ}$

Showers: $\nu_e$  or  $\nu_{\tau}$  CC interactionsAll NC interactionsMuon bremsstrahlung near detectorAngular resolution  $10^\circ \times 10^\circ$ 

Salient observation: event rates for these have different dependence on cross section

A. Kusenko, T. Weiler, PRL 88,161101 (2002)

- Down-going shower event rate can be enhanced because of a large flux or anomalously large  $\sigma_{\nu N}$  or some combination of both
- Up-going rate also increases with increasing  $\phi^{\nu}$ ; however, because of absorption on passage through earth, event rate decreases with increasing  $\sigma_{\nu N}$
- Absorption effects will limit up-going events to those coming at just below the horizon Earth-skimming events J. L.Feng et al, PRL 88,161102(2002)

Working equations: downgoing

$$\mathcal{N}_{
m down} = C_{
m down} \; rac{\phi^
u}{\phi^
u_{
m WB}} \; rac{\sigma_{
uN}}{\sigma_{
m SM}}$$

- Constant  $C_{\text{down}}$  depends on exposure, acceptance, and varies according to neutrino flavor from experiment to experiment
- $\phi^{
  u}_{
  m WB}$  and  $\sigma_{
  m SM}$  serve only as normalization factors
- Specialize to electron showers, find

$$C_{
m down} = 4 \,\, {
m events}$$

in 15 yrs

Working equations: up-going

In this case

$$\mathcal{N}_{ ext{up}} = C_{ ext{up}} \; rac{\phi^
u}{\phi^
u_{ ext{WB}}} \; rac{\sigma^2_{ ext{SM}}}{\sigma^2_{
uN}} \quad \left(rac{\sigma_{
uN}}{\sigma_{ ext{SM}}} > 1
ight)$$

- This holds for  $L^\ell \ll L^
  u < R_\oplus$ , corresponding to  $E > 10^7~{
  m GeV}$
- $\sigma_{\nu N}$  is defined with cuts so that shower energy fraction of same order as that of the CC SM process.
- Specialize to tau showers (distinctive topologies), find for 15 yr

$$C_{
m up}=20\,\,{
m events}$$

J.J Tseng et al; J.Jones et al; S.I.Dutta et al

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May 16, 2006 Pheno2006 With estimates of  $C_{\rm down}$  and  $C_{\rm up}$ , can determine sensitivities of IceCube to  $\phi^{\nu}$  and  $\sigma_{\nu N}$ 

- For a given set of observed rates  $\mathcal{N}_{up}^{obs}$  and  $\mathcal{N}_{down}^{obs}$ , two curves obtained in 2-D parameter space by setting  $\mathcal{N}_{up}^{obs} = \mathcal{N}_{up}$  and  $\mathcal{N}_{down}^{obs} = \mathcal{N}_{down}$
- Curves intersect at a point, yielding the most probable values of  $\phi^{\nu}$  and  $\sigma_{\nu N}$  for the given observations

Approximate likelihood analysis

• Fluctuations about this point define contours of constant  $\chi^2$  in an approximation to a multi-Poisson likelihood analysis

$$\chi^2 = \sum_i^{
m down,\,up} 2 \left[ \mathcal{N}^i - \mathcal{N}^i_{
m obs} + \mathcal{N}^i_{
m obs} \, \ln\left(rac{\mathcal{N}^i_{
m obs}}{\mathcal{N}^i}
ight) 
ight]$$

S. Baker and R. D. Cousins, Nucl. Instrum. Meth. A221,437(1984)

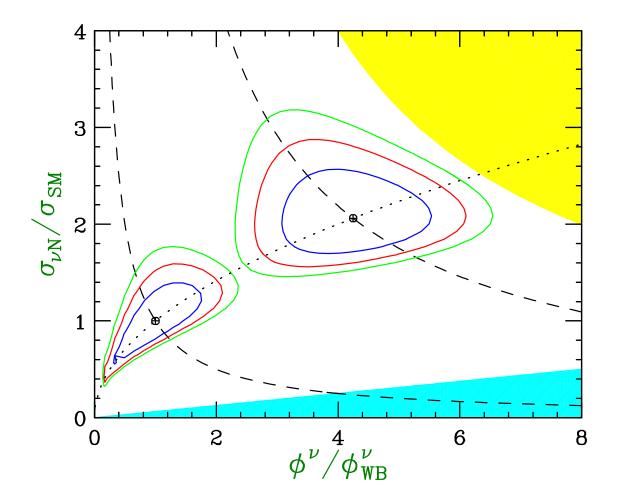
• Two illustrative cases:

 $(\mathcal{N}_{
m down}^{
m obs}, \ \mathcal{N}_{
m up}^{
m obs}) = (4, 20),$ conforming to SM/WB expectations  $(\mathcal{N}_{
m down}^{
m obs}, \ \mathcal{N}_{
m up}^{
m obs}) = (35, 20),$ indicating deviation from  $\sigma_{
m SM}$  and/or  $\phi_{
m WB}^{\nu}$ 

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90%, 99% and 99.9% CL contours for two cases in previous slide. Lower are deviation bounds, upper signal  $5\sigma$  new physics discovery (for any flux).

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- multi-KK graviton exchange in non-warped TeV scale gravity models Emparan, Masip, Rattazzi
- BH production in these models Feng and Shapere;

Anchordoqui et al; Ahn, Cavaglia, Olinto

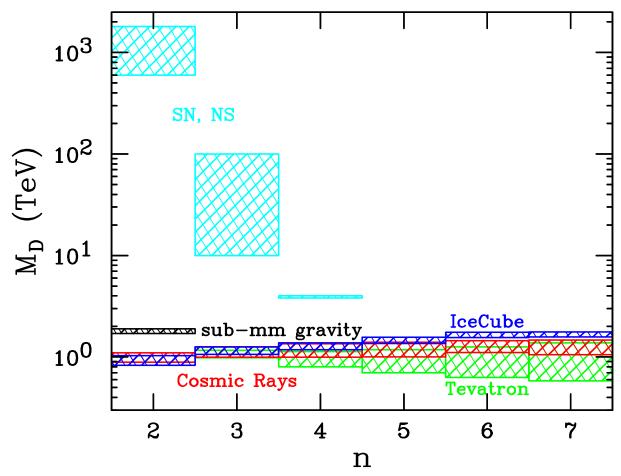
• cross section enhancements due to electroweak instantons Ahlers, Ringwald, Tu

## All of these are shower producing

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Example of sensitivity; Large extra dimensional BH

If  $\mathcal{N}_{up}^{obs} = C_{up}$  and  $\mathcal{N}_{down}^{obs} = C_{down}$ , can bound BH cross section and thus scale of TeV gravity  $M_D$ 

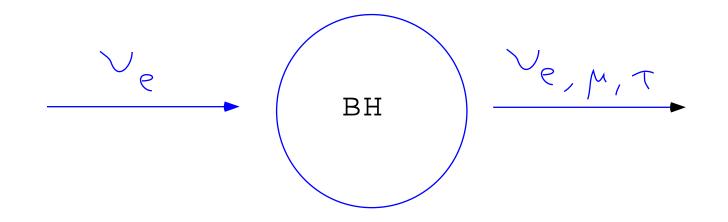




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Quantum Decoherence via neutrinos



**Density matrix description** 

$$egin{array}{rcl} 
ho^lpha(t) &=& |
u_lpha(t)
angle \langle
u_lpha(t)| \ P_{
u_lpha
ightarrow
u_eta} &=& {
m Tr}\left[
ho_lpha(t)
ho_eta
ight] \end{array}$$

**Evolve**  $\rho^{\alpha}(t)$  using modified Liouville equation

$$rac{\partial 
ho}{\partial t} = -i[H, \, 
ho] + \mathcal{D}[
ho]$$

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Quantum Decoherence (cont'd)

$$\mathcal{D}[
ho] = -rac{1}{2}\sum\limits_{j} \left( [b_j, \ 
ho \ b_j^\dagger] + [b_j \ 
ho, \ b_j^\dagger] 
ight)$$

- This conserves total probability  $Tr(\rho)$
- Can be obtained as interaction with heat bath via  $a^{\dagger}b + b^{\dagger}a$
- Monotonic increase of von Neumann entropy  $S = -\mathrm{Tr}(\rho \ln \rho) \to \mathbf{b} \text{ hermitian}$
- Expand in Gell-Mann matrices,  $\gamma_i = ext{eigenvalues} ext{ of } b^\dagger b$
- Dealing with distant sources  $\rightarrow$  only  $\gamma_3, \gamma_8$ survive oscillation, take  $\gamma_3 = \gamma_8$



$$egin{array}{rll} P_{\overline{
u}_lpha
ightarrow \overline{
u}_eta} &=& rac{1}{3} + \mathrm{e}^{-\overline{\gamma}\,d} \; \left[ rac{1}{2} \; (U_{lpha1}^2 - U_{lpha2}^2) (U_{eta1}^2 - U_{eta2}^2) \ &+& rac{1}{6} \; (U_{lpha1}^2 + U_{lpha2}^2 - 2 U_{lpha3}^2) (U_{eta1} + U_{eta2}^2 - 2 U_{eta3}^2) 
ight] \end{array}$$

Expect  $\gamma$  to be energy dependent

$$egin{array}{rcl} \overline{\gamma} &=& \kappa_n & (E_
u/{
m GeV})^n, & n=-1,0,1,2\dots \ &=& (E_
u/M_{
m QG})^{n-1} & E_
u \end{array}$$

- n = -1 is analogue of oscillation, n = 0 like decay
- $n \geq 2$  involves large scale  $M_{
  m QG}$

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## Limits on $\gamma$

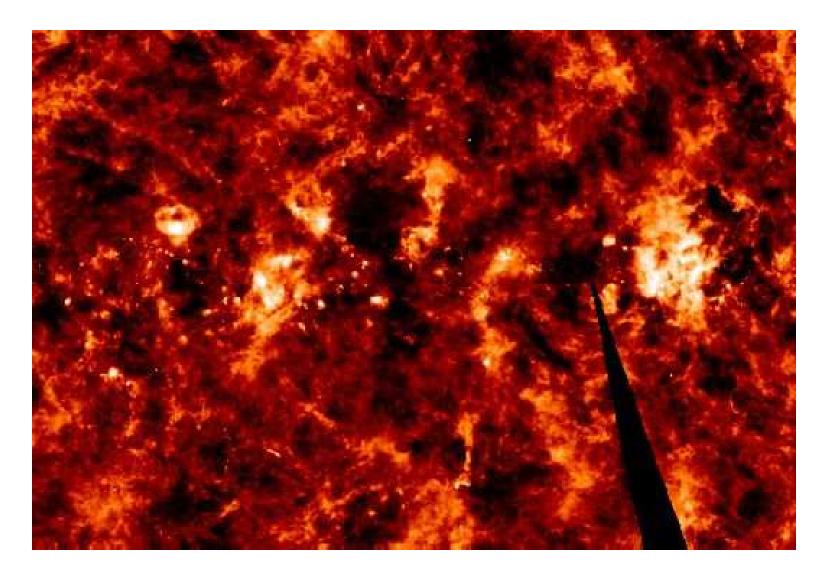
E. Lisi, A. Marrone, D. Montanino, hep-ph/0002053

- use atmospheric neutrinos  $\nu_{\mu}$  at SK and K2K
- can probe decoherence length  $\gamma^{-1}$  of order the oscillation length  $\sim 10^3$  km to obtain limits see later
- can obtain much stronger limits if there is source of neutrinos at large distance ( $\gtrsim kpc$ ) with observed flavor mix  $\neq 1:1:1$
- In what follows, IceCube is used as neutrino detection facility

## **Example: Neutrinos from Neutrons from Cygnus OB2**

- In PRD72:065019,2005 Anchordoqui, Gonzalez-Garcia, HG, Halzen, Hooper, Sarkar, Weiler use reported (at the time) anisotropy from direction of Cygnus OB2 region  $(d \approx 1.7 \text{ kpc})$
- Anisotropy identified as neutrons from photodisintegration of Fe
- $\overline{\nu}_e$ 's from neutron decay evolve to  $\overline{\nu}_e: \overline{\nu}_\mu: \overline{\nu}_\tau = 2.5: 1: 1$  after large distance – definitely not 1:1:1
- $\overline{\nu}$  fluxes fixed by n flux fixed by anisotropy

# Cygnus OB2



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# Limits on $\gamma$

- Obtain number of tracks and showers expected from atmospheric background
- Introduce an additional background of tracks and showers due to possible 1:1:1 source of neutrinos from nearby HEGRA  $\gamma$  ray source
- Calculate theoretically expected number of tracks and showers as a function of the decoherence parameter  $\overline{\kappa}_n$  and the HEGRA background
- Marginalize on HEGRA background to obtain bounds on  $\overline{\kappa}_n$  for a given observed number of tracks and showers

## Results for $\mathcal{N}_{\mathrm{obs}} = \mathcal{N}_{\mathrm{ATM}}$

Table 1: 90% CL limits on decoherence

	Cygnus OB2		SK & K2K	
n	$\kappa_{ m max}$ (GeV)	$M_{ m QG}^{ m min}$ (GeV)	$\kappa_{ m max}$ (GeV)	$M_{ m QG}^{ m min}$ (GeV)
-1	$1.0 imes10^{-34}$	_	$2.0 imes10^{-21}$	_
0	$3.2 imes10^{-36}$	—	$3.5 imes10^{-23}$	—
2	$2.0 imes10^{-44}$	$5.0 imes10^{43}$	$9.0 imes10^{-28}$	$1.1 imes 10^{27}$
3	$3.0 imes10^{-47}$	$1.8 imes10^{23}$	—	$3.3  imes 10^{13^{ *}}$

\* Obtained from limit on  $\gamma$ , with  $\langle E_{\nu} \rangle \approx 100 \text{ GeV}$ 

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For a wide variety of fundamental questions in particle physics

- Split Supersymmetry
- long-lived superheavy relics
- anomalous neutrino interactions at very high energies
- quantum decoherence

the new generation of detectors (Auger, IceCube, EUSO) can provide important constraints on the new physics, if not discovery.