

Detecting Dark Matter in NUHM Models

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Universality

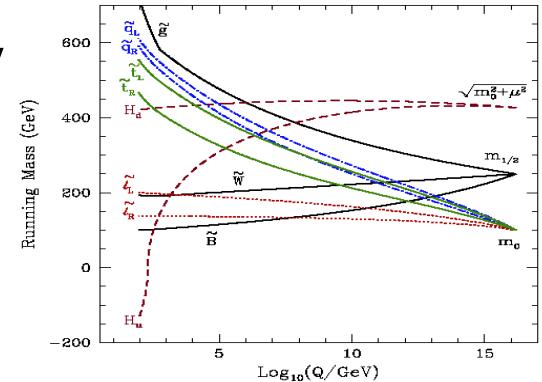
FCNC suppression: universality of matter fields that share quantum numbers (sfermion generations)

minimal Supergravity: Higgses+sfermions, gauginos

SUSY GUTs: varying degrees of universality

SO(10): Higgses, scalars, gauginos

SU(5): some masses equal



Mirage mediation: universality below GUT scale

Some string scenarios for SUSY breaking - maybe no universality at any scale!

Non-Universal Higgs Masses

- D. Matalliotakis and H. P. Nilles, Nucl. Phys. B 435 (1995) 115 [arXiv:hep-ph/9407251];
M. Olechowski and S. Pokorski, Phys. Lett. B 344, 201 (1995) [arXiv:hep-ph/9407404];
V. Berezinsky, A. Bottino, J. Ellis, N. Fornengo, G. Mignola and S. Scopel, *Astropart. Phys.* **5** (1996) 1, hep-ph/9508249; M. Drees, M. Nojiri, D. Roy and Y. Yamada, *Phys. Rev. D* **56** (1997) 276, [Erratum-ibid. D **64** (1997) 039901], hep-ph/9701219;
M. Drees, Y. Kim, M. Nojiri, D. Toya, K. Hasuko and T. Kobayashi, *Phys. Rev. D* **63** (2001) 035008, hep-ph/0007202; P. Nath and R. Arnowitt, *Phys. Rev. D* **56** (1997) 2820, hep-ph/9701301; J. R. Ellis, T. Falk, G. Ganis, K. A. Olive and M. Schmitt, *Phys. Rev. D* **58** (1998) 095002 [arXiv:hep-ph/9801445]; J. R. Ellis, T. Falk, G. Ganis and K. A. Olive, *Phys. Rev. D* **62** (2000) 075010 [arXiv:hep-ph/0004169];
A. Bottino, F. Donato, N. Fornengo and S. Scopel, *Phys. Rev. D* **63** (2001) 125003, hep-ph/0010203; S. Profumo, *Phys. Rev. D* **68** (2003) 015006, hep-ph/0304071; D. Cerdeño and C. Muñoz, *JHEP* **0410** (2004) 015, hep-ph/0405057;
- J. Ellis, K. Olive and Y. Santoso, *Phys. Lett. B* **539**, 107 (2002) [arXiv:hep-ph/0204192];
J. R. Ellis, T. Falk, K. A. Olive and Y. Santoso, *Nucl. Phys. B* **652**, 259 (2003) [arXiv:hep-ph/0210205].
- H. Baer, A. Mustafayev, S. Profumo, A. Belyaev and X. Tata, *Phys. Rev. D* **71**, 095008 (2005) [arXiv:hep-ph/0412059]. H. Baer, A. Mustafayev, S. Profumo, A. Belyaev and X. Tata, *JHEP* **0507** (2005) 065, hep-ph/0504001.

CMSSM → NUHM

CMSSM GUT-scale inputs:

m_0 , $m_{1/2}$ and A_0 and the sign of μ

Use electroweak vacuum conditions:

$$m_A^2(Q) = m_1^2(Q) + m_2^2(Q) + 2\mu^2(Q) + \Delta_A(Q)$$

$$\mu^2 = \frac{m_1^2 - m_2^2 \tan^2 \beta + \frac{1}{2} m_Z^2 (1 - \tan^2 \beta) + \Delta_\mu^{(1)}}{\tan^2 \beta - 1 + \Delta_\mu^{(2)}},$$

$$Q = (m_{\tilde{\tau}_R} m_{\tilde{\tau}_L})^{1/2} \quad m_A \equiv m_A(Q) \quad m_1^2(Q) = m_1^2 + c_1, \\ \mu \equiv \mu(m_Z) \quad m_2^2(Q) = m_2^2 + c_2, \\ \mu^2(Q) = \mu^2 + c_\mu,$$

CMSSM → NUHM

CMSSM – Higgs masses determined by $m_1^2(M_{GUT}) = m_2^2(M_{GUT}) = m_0^2$

NUHM1 – One extra free parameter: $m_1^2 = m_2^2$

$$\begin{aligned} m_1^2(\tan^2 \beta + 1 + \Delta_\mu^{(2)}) &= m_2^2(\tan^2 \beta + 1 - \Delta_\mu^{(2)}) + m_Z^2(\tan^2 \beta - 1) - 2\Delta_\mu^{(1)} \\ &\quad + (m_A^2 - (\Delta_A(Q) + c_1 + c_2 + 2c_\mu))(\tan^2 \beta - 1 + \Delta_\mu^{(2)}). \end{aligned}$$

$$m_1^2 = m_2^2 \tan^2 \beta + \mu^2(\tan^2 \beta - 1 + \Delta_\mu^{(2)}) + \frac{1}{2}m_Z^2(\tan^2 \beta - 1) - \Delta_\mu^{(1)}.$$

NUHM2 – No constraint at GUT scale:

$$\begin{aligned} m_1^2(1 + \tan^2 \beta) &= m_A^2(Q) \tan^2 \beta - \mu^2(\tan^2 \beta + 1 - \Delta_\mu^{(2)}) - (c_1 + c_2 + 2c_\mu) \tan^2 \beta \\ &\quad - \Delta_A(Q) \tan^2 \beta - \frac{1}{2}m_{BZ}^2(1 - \tan^2 \beta) - \Delta_\mu^{(1)} \end{aligned}$$

$$\begin{aligned} m_2^2(1 + \tan^2 \beta) &= m_A^2(Q) - \mu^2(\tan^2 \beta + 1 + \Delta_\mu^{(2)}) - (c_1 + c_2 + 2c_\mu) \\ &\quad - \Delta_A(Q) + \frac{1}{2}m_{BZ}^2(1 - \tan^2 \beta) + \Delta_\mu^{(1)}. \end{aligned}$$

DD-Relevant Parameters

Neutralino mass eigenstates:

$$\chi = Z_{\chi 1} \tilde{B} + Z_{\chi 2} \tilde{W} + Z_{\chi 3} \tilde{H}_1 + Z_{\chi 4} \tilde{H}_2.$$

$$(\tilde{W}^3, \tilde{B}, \tilde{H}_1^0, \tilde{H}_2^0) \begin{pmatrix} M_2 & 0 & \frac{-g_2 v_1}{\sqrt{2}} & \frac{g_2 v_2}{\sqrt{2}} \\ 0 & M_1 & \frac{g_1 v_1}{\sqrt{2}} & \frac{-g_1 v_2}{\sqrt{2}} \\ \frac{-g_2 v_1}{\sqrt{2}} & \frac{g_1 v_1}{\sqrt{2}} & 0 & -\mu \\ \frac{g_2 v_2}{\sqrt{2}} & \frac{-g_1 v_2}{\sqrt{2}} & -\mu & 0 \end{pmatrix} \begin{pmatrix} \tilde{W}^3 \\ \tilde{B} \\ \tilde{H}_1^0 \\ \tilde{H}_2^0 \end{pmatrix}$$

Squark Mixing: $diag(m_1^2, m_2^2) \equiv \eta M^2 \eta^{-1}$

$$\begin{pmatrix} \cos \theta_f & \sin \theta_f e^{i\gamma_f} \\ -\sin \theta_f e^{-i\gamma_f} & \cos \theta_f \end{pmatrix} \equiv \begin{pmatrix} \eta_{11} & \eta_{12} \\ \eta_{21} & \eta_{22} \end{pmatrix}$$

DD Processes

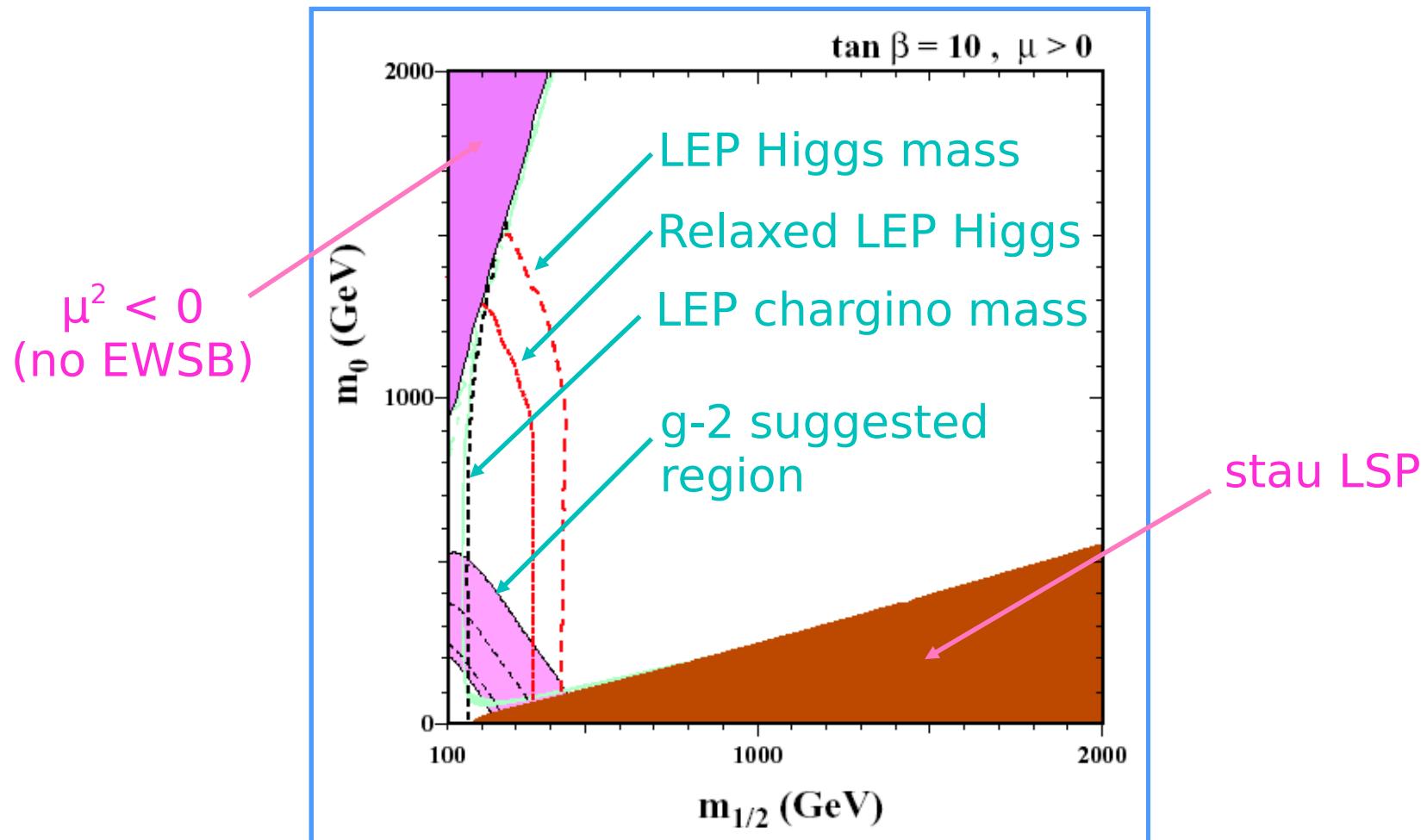
Neutralino-nucleon interaction Lagrangian:

$$\mathcal{L} = \alpha_{2i} \bar{\chi} \gamma^\mu \gamma^5 \chi \bar{q}_i \gamma_\mu \gamma^5 q_i + \alpha_{3i} \bar{\chi} \chi \bar{q}_i q_i$$

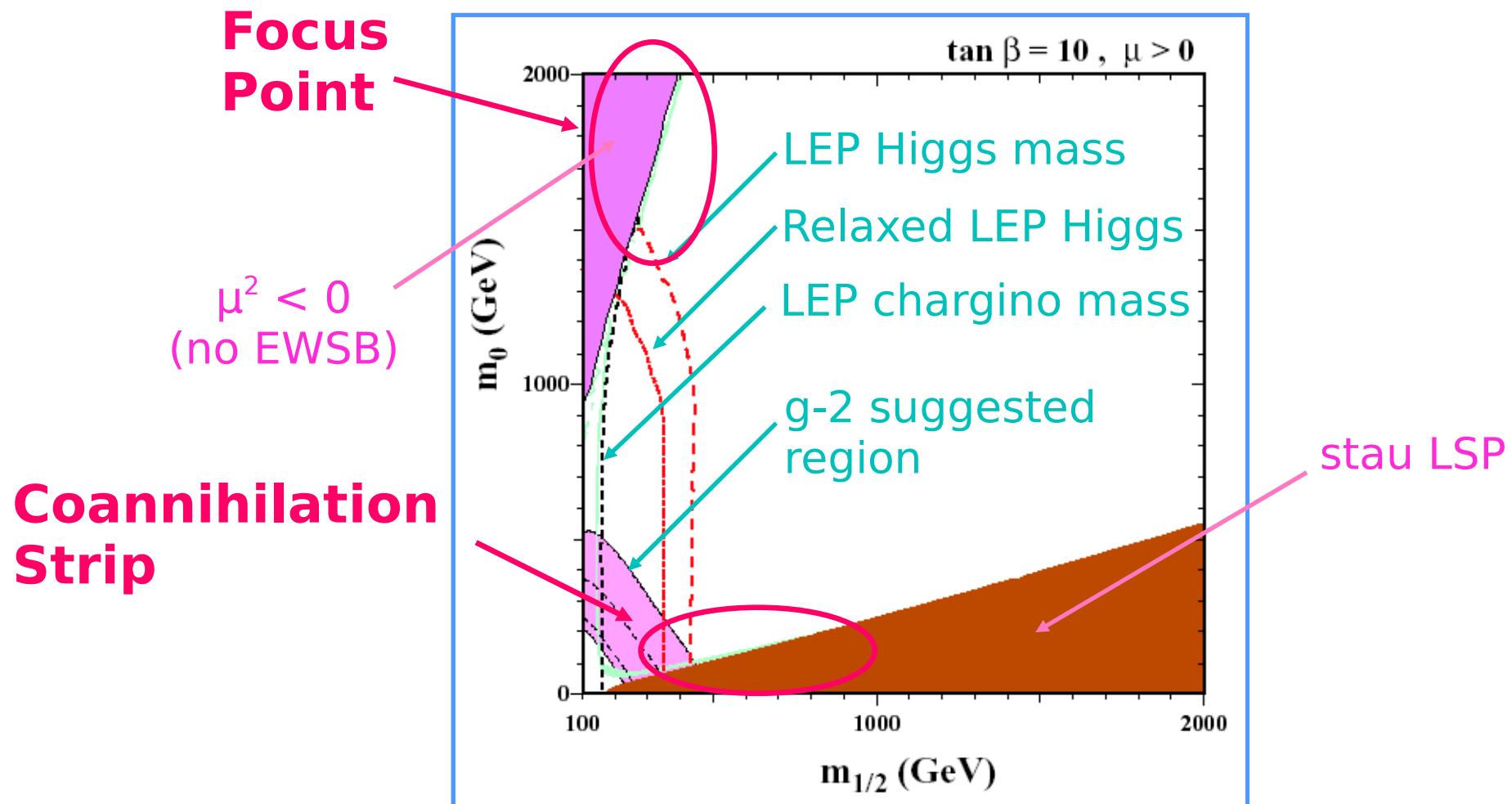
$$\begin{aligned} \alpha_{2i} &= \frac{1}{4(m_{1i}^2 - m_\chi^2)} [|X_i|^2 + |Y_i|^2] + \frac{1}{4(m_{2i}^2 - m_\chi^2)} [|W_i|^2 + |V_i|^2] \\ &\quad - \frac{g^2}{4m_Z^2 \cos^2 \theta_W} [|Z_{\chi 3}|^2 - |Z_{\chi 4}|^2] \frac{T_{3i}}{2} \end{aligned}$$

$$\begin{aligned} \alpha_{3i} &= -\frac{1}{2(m_{1i}^2 - m_\chi^2)} \text{Re} [(X_i)(Y_i)^*] - \frac{1}{2(m_{2i}^2 - m_\chi^2)} \text{Re} [(W_i)(V_i)^*] \\ &\quad - \frac{gm_{q_i}}{4m_W B_i} \left\{ \left(\frac{D_i^2}{m_{H_2}^2} + \frac{C_i^2}{m_{H_1}^2} \right) \text{Re} [\delta_{2i} (gZ_{\chi 2} - g'Z_{\chi 1})] \right. \\ &\quad \left. + D_i C_i \left(\frac{1}{m_{H_2}^2} - \frac{1}{m_{H_1}^2} \right) \text{Re} [\delta_{1i} (gZ_{\chi 2} - g'Z_{\chi 1})] \right\}, \end{aligned}$$

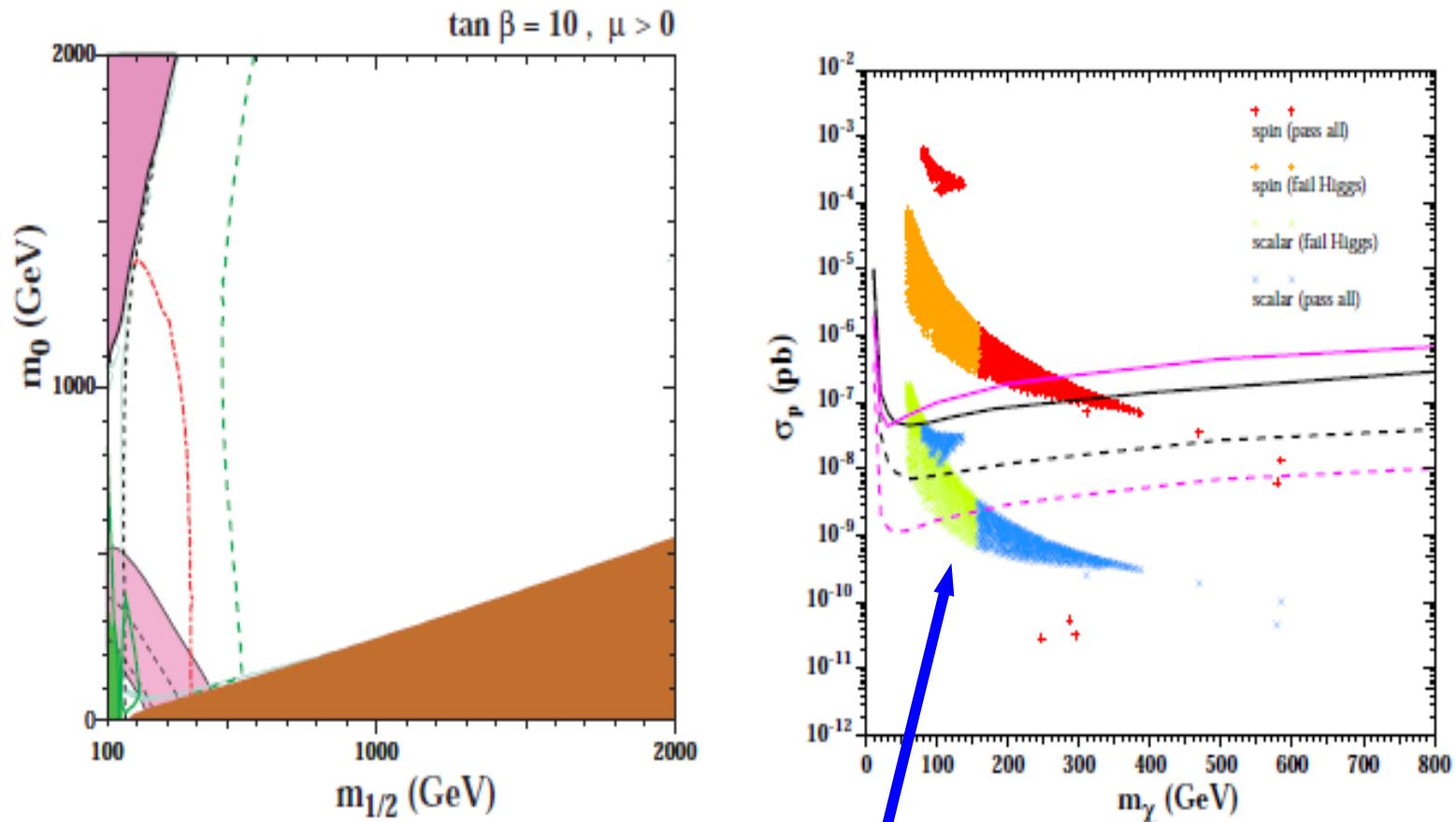
CMSSM



CMSSM

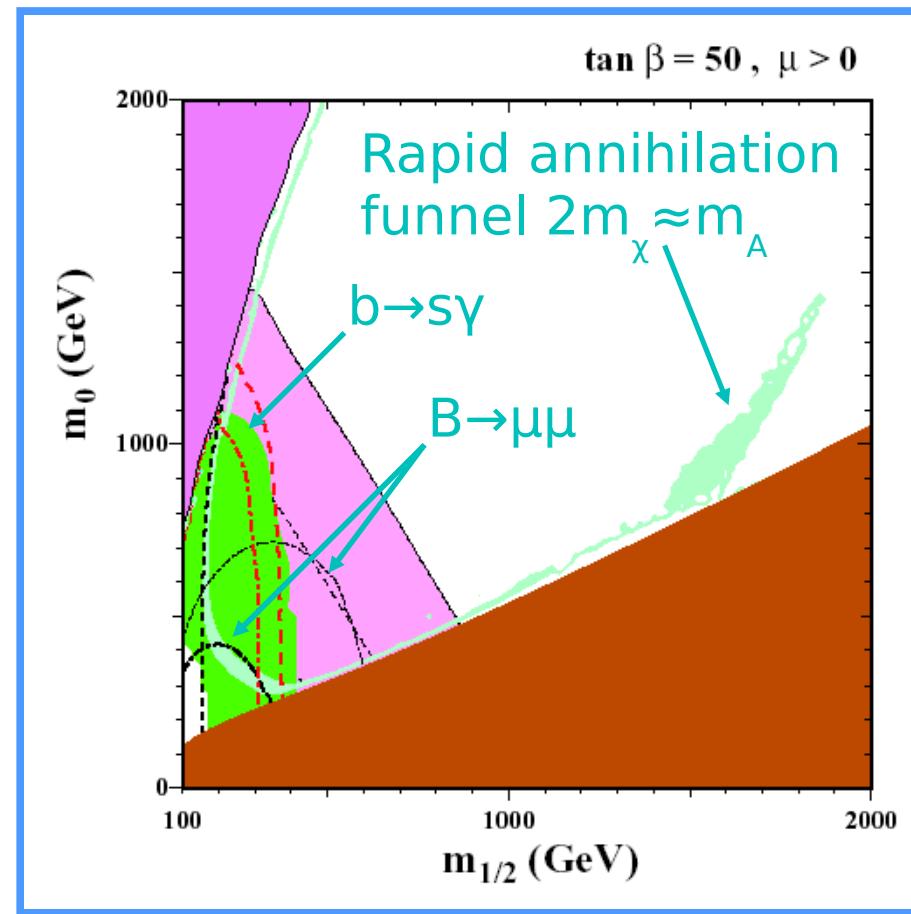


CMSSM

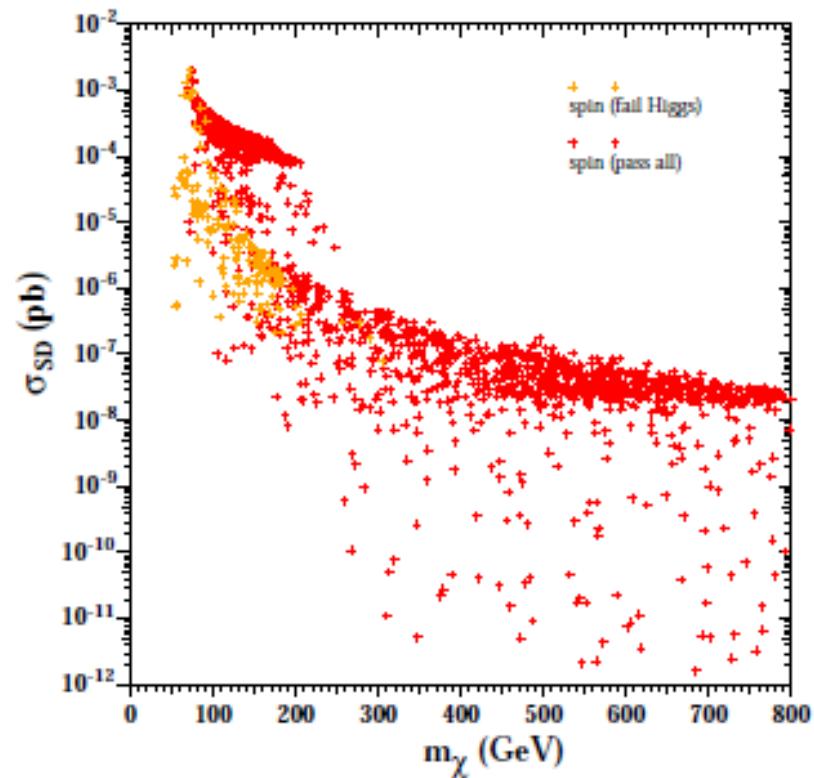
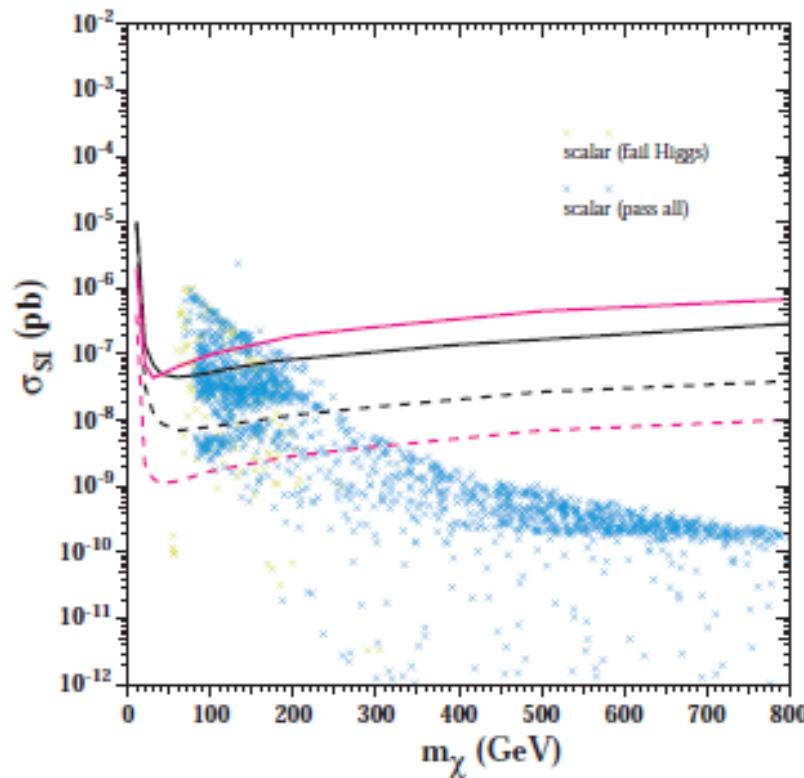


$$\left\{ \begin{array}{l} \text{If } \Omega_\chi < \Omega_{DM}^{WMAP} \text{ then} \\ \sigma = \sigma_0 \frac{\Omega_\chi}{\Omega_{DM}^{WMAP}} \end{array} \right\}$$

CMSSM

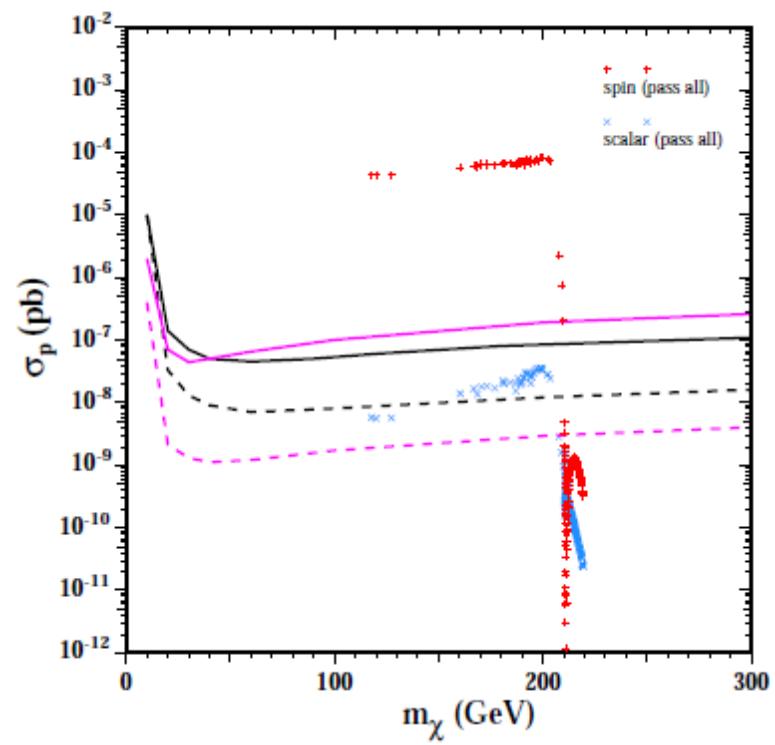
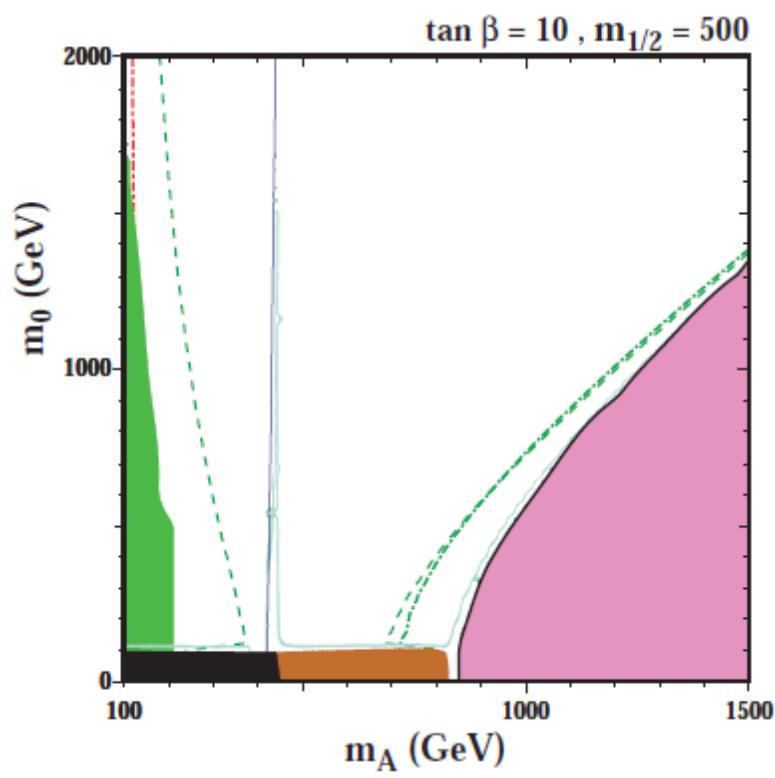


CMSSM

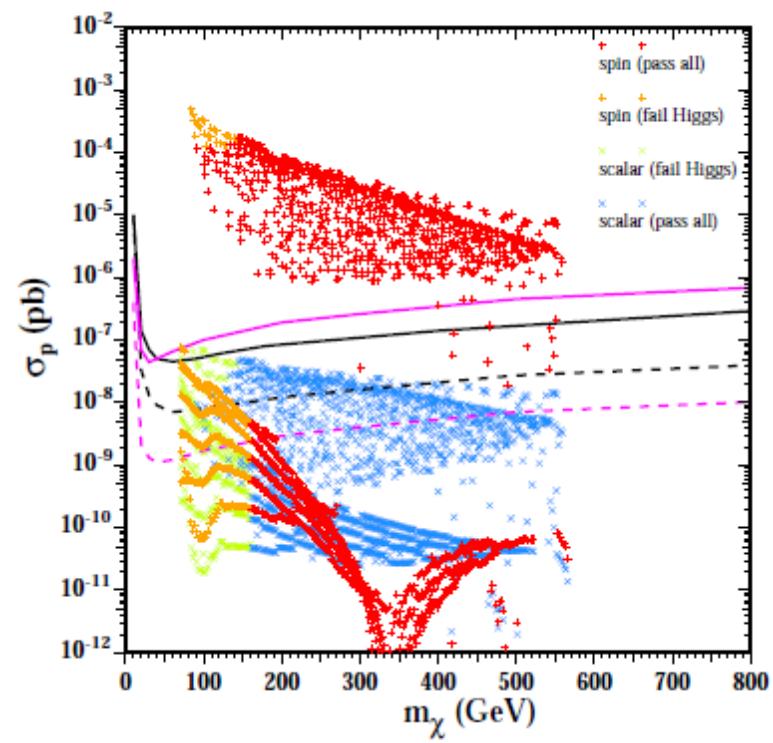
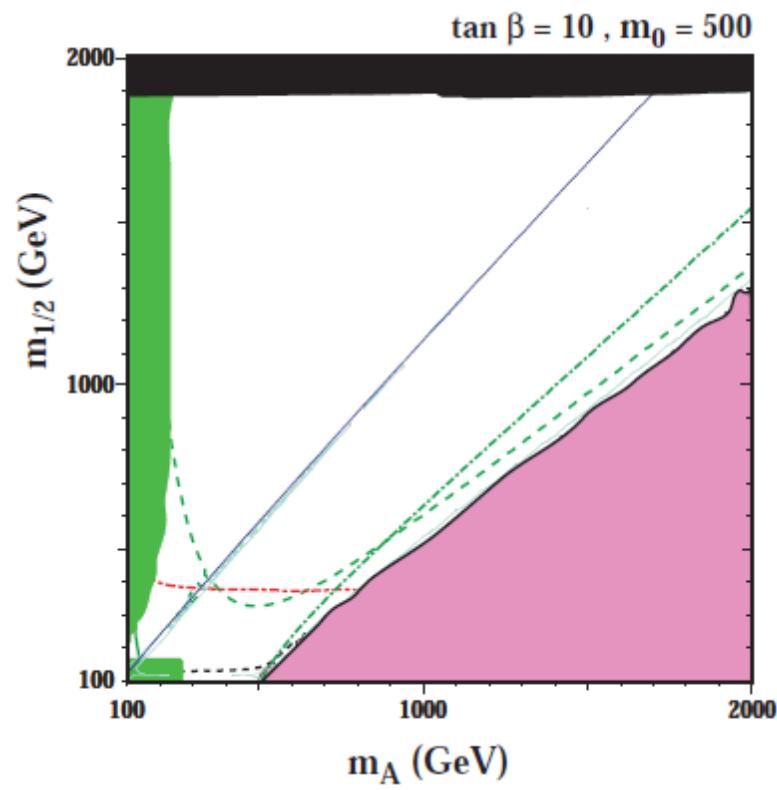


$5 \leq \tan \beta \leq 55, 0 \leq m_{1/2} \leq 2000 \text{ GeV}, 100 \text{ GeV} \leq m_0 \leq 2000 \text{ GeV},$
 $-3m_{1/2} \leq A_0 \leq 3m_{1/2}$ We consider $\mu < 0$ only for $\tan \beta < 30$.

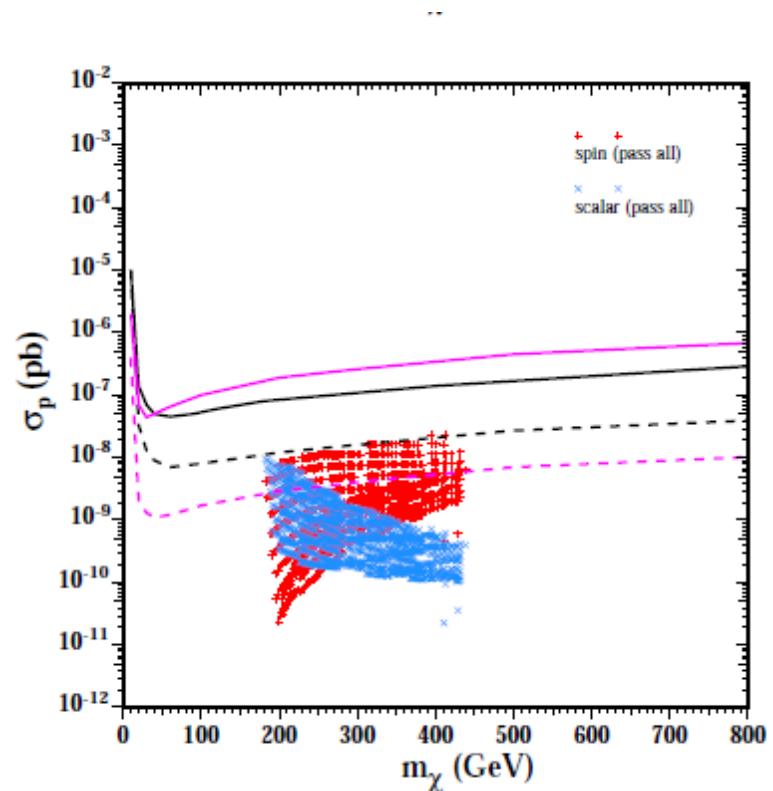
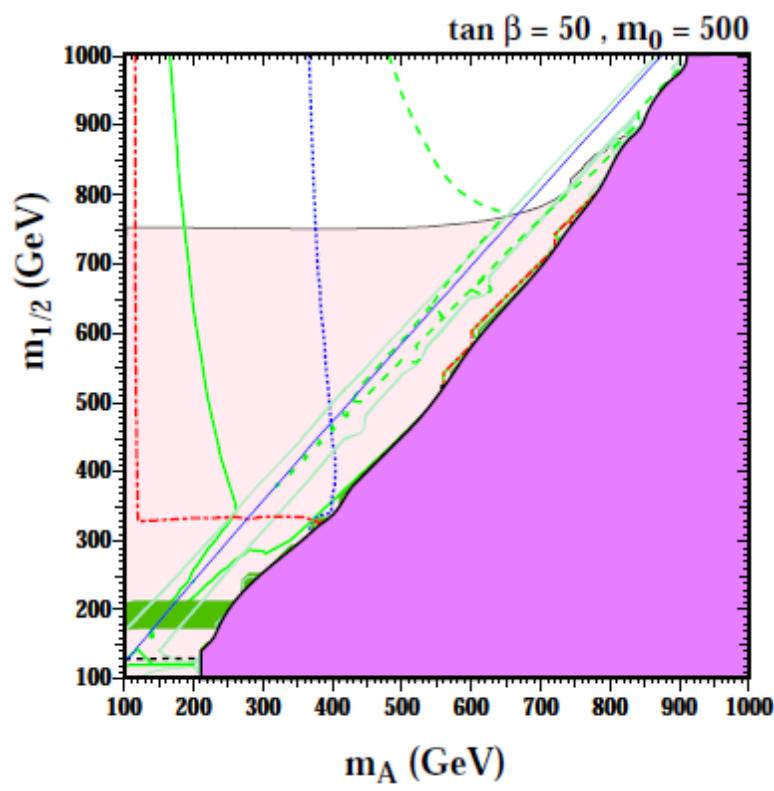
NUHM1



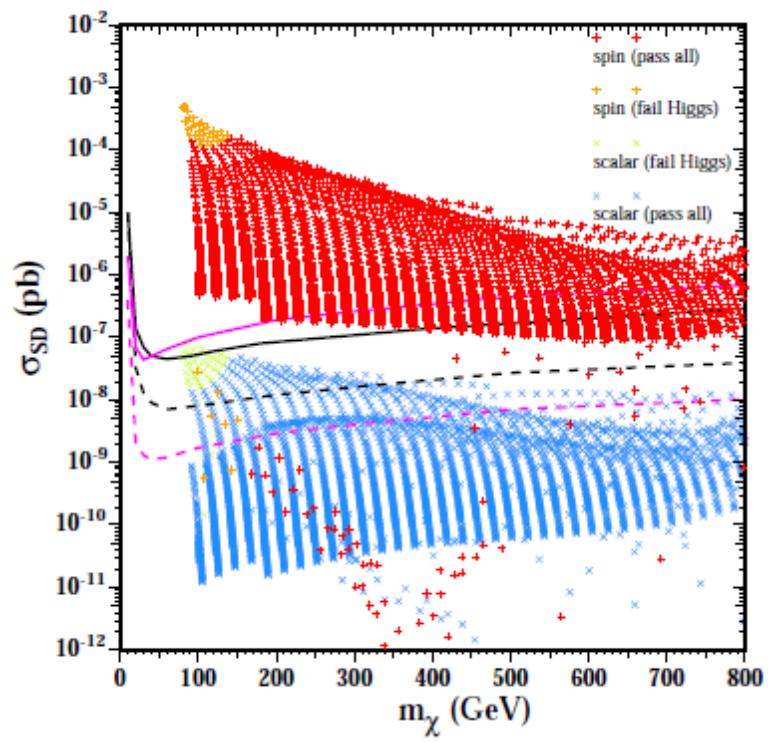
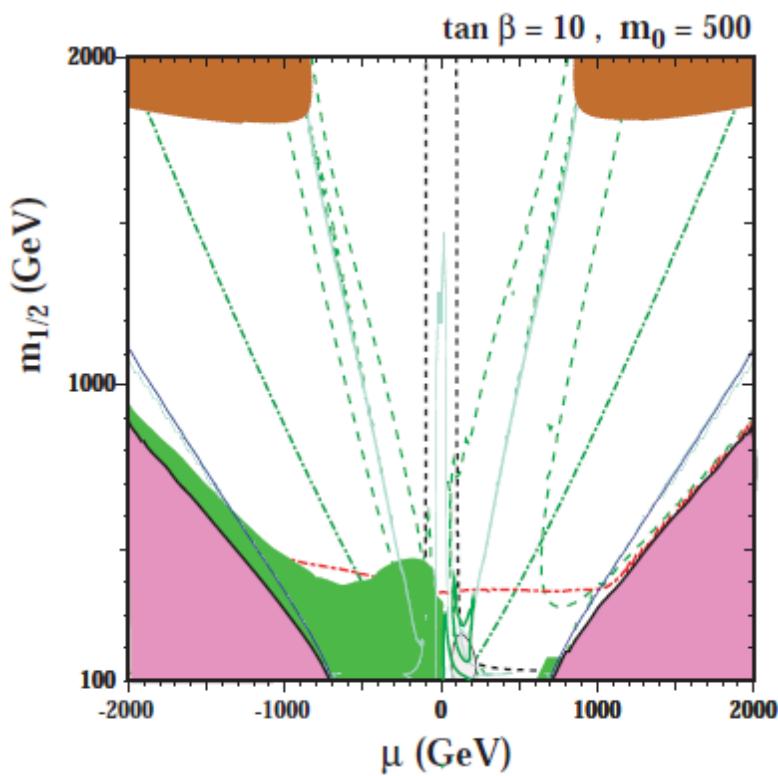
NUHM1



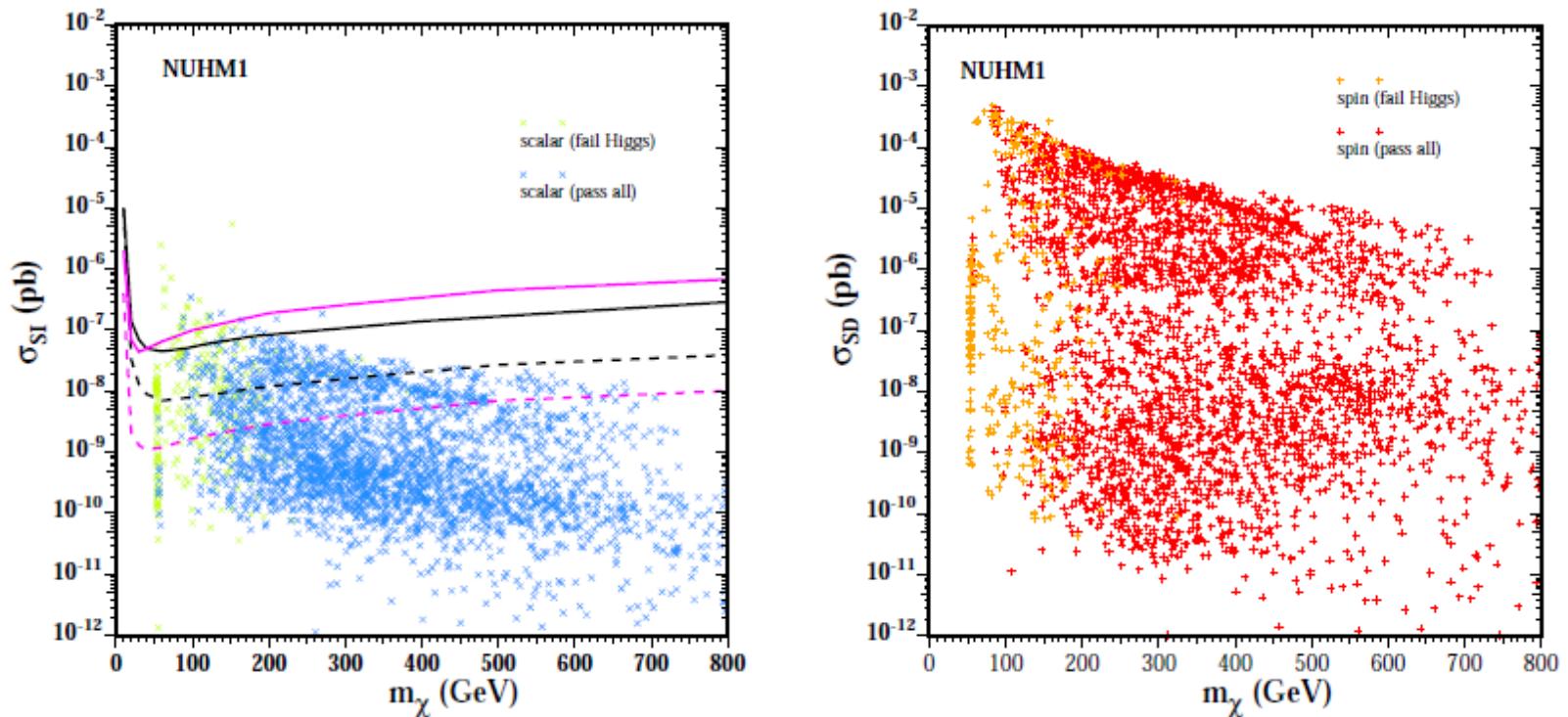
NUHM1



NUHM1

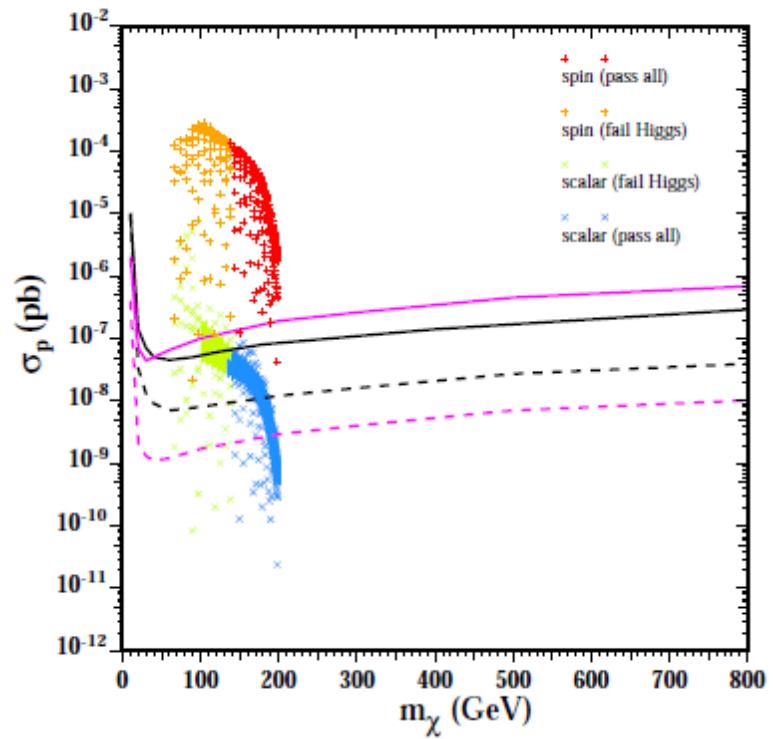
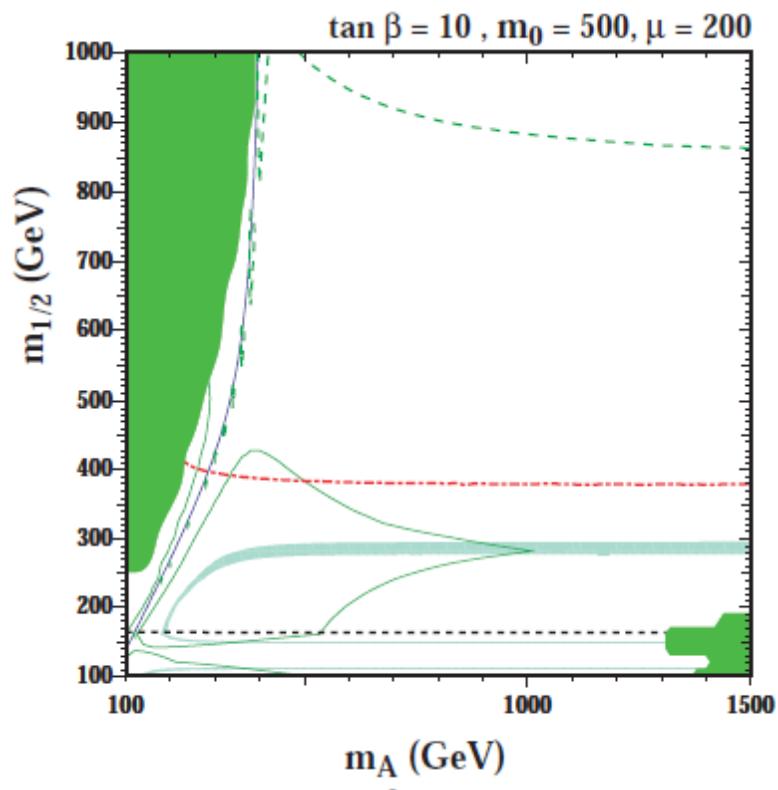


NUHM1

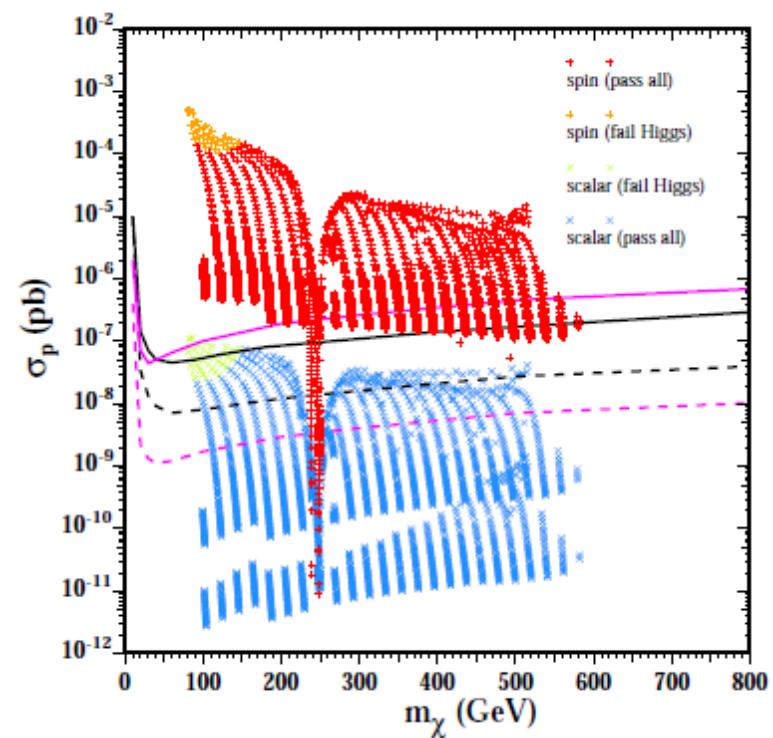
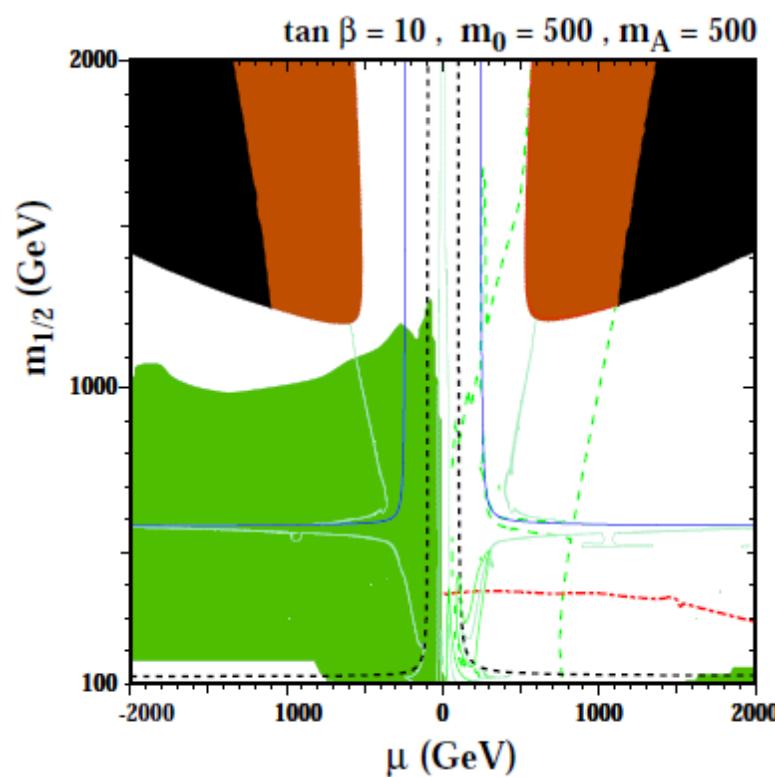


In both plots, we scan $5 \leq \tan \beta \leq 55$, $0 \leq m_{1/2} \leq 2000$ GeV, 100 GeV $\leq m_0 \leq 2000$ GeV, and $-3m_{1/2} \leq A_0 \leq 3m_{1/2}$. The common GUT-scale value of $m_1 = m_2$ is in the range $(-2000, 2000)$ GeV.

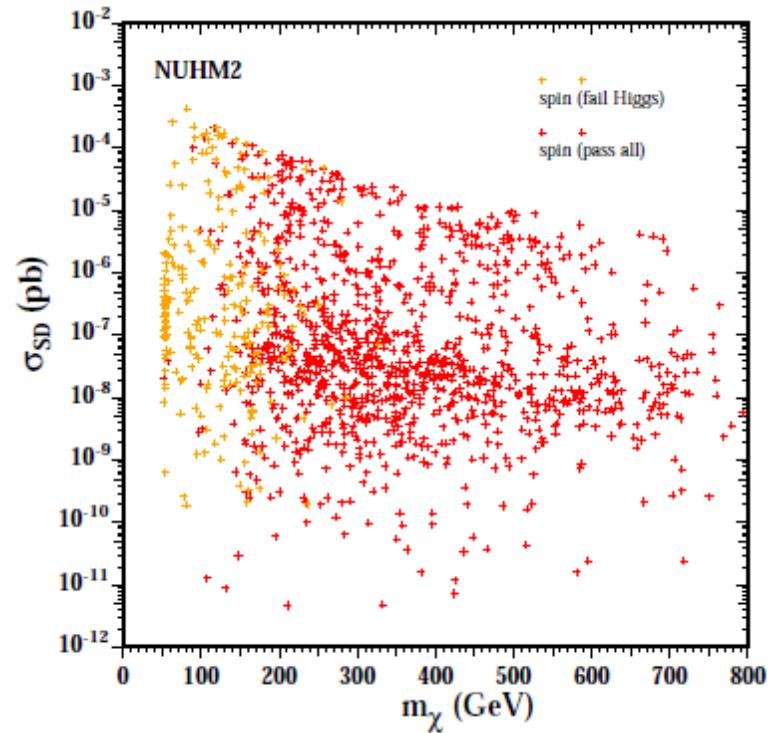
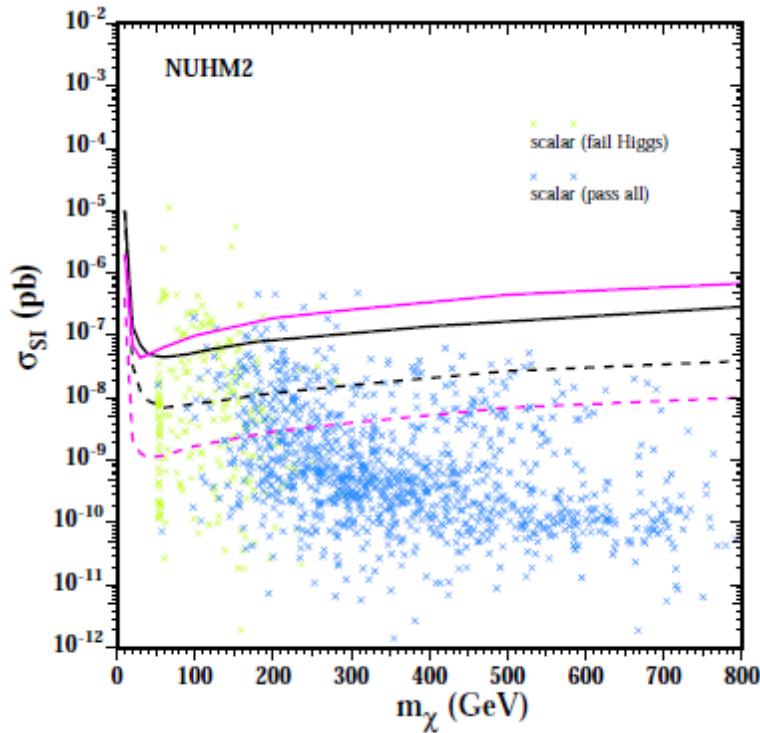
NUHM2



NUHM2



NUHM2



In both scans, we scan $5 \leq \tan \beta \leq 55$, $0 \leq m_{1/2} \leq 2000$ GeV, 100 GeV $\leq m_0 \leq 2000$ GeV, $-3m_{1/2} \leq A_0 \leq 3m_{1/2}$, and the GUT-scale values of m_1 and m_2 are each in the range $(-2000, 2000)$ GeV.

Summary

- Additional cosmologically-viable regions of parameter space lead to broader ranges of DM scattering rates than are possible in the CMSSM
 - Larger cross sections at larger mass (crossover regions with substantial Higgsino component)
 - Smaller cross sections at smaller mass (relic density suppressed by some mechanism?)
- If we observe SI scattering by heavy WIMP (not CMSSM), could be in crossover region...
- If LHC establishes light LSP but we don't see SI scattering, could mean Higgsino-like LSP