

The WIMP Forest

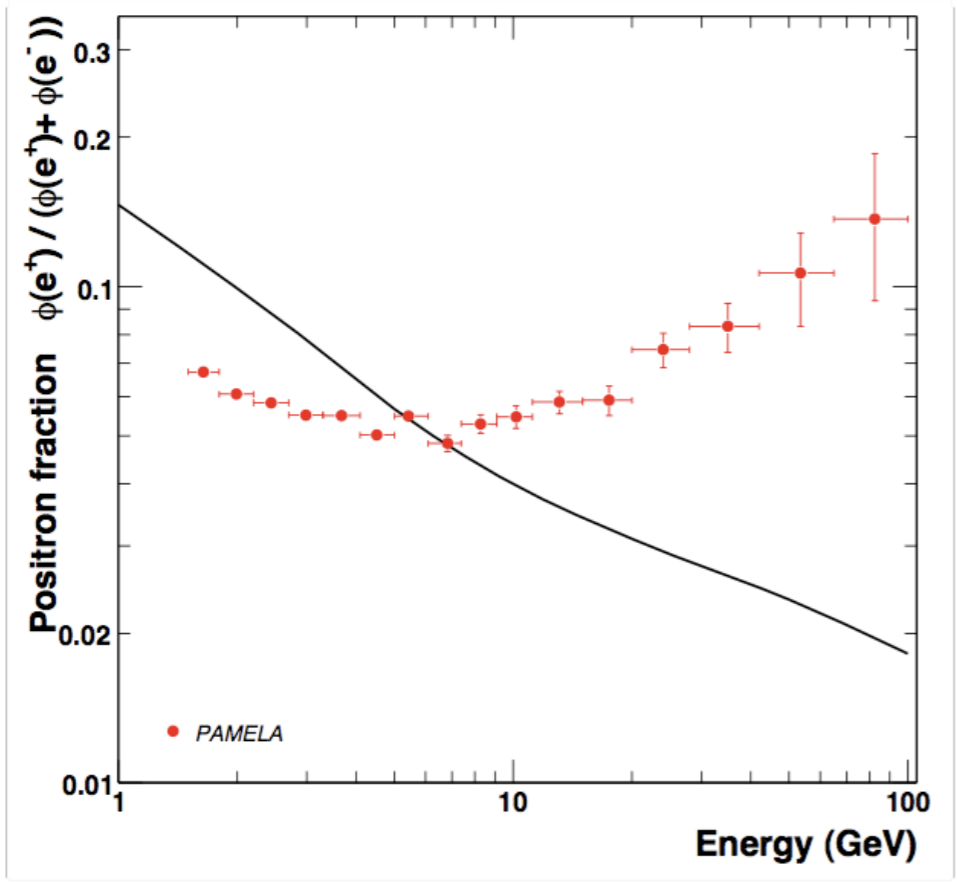
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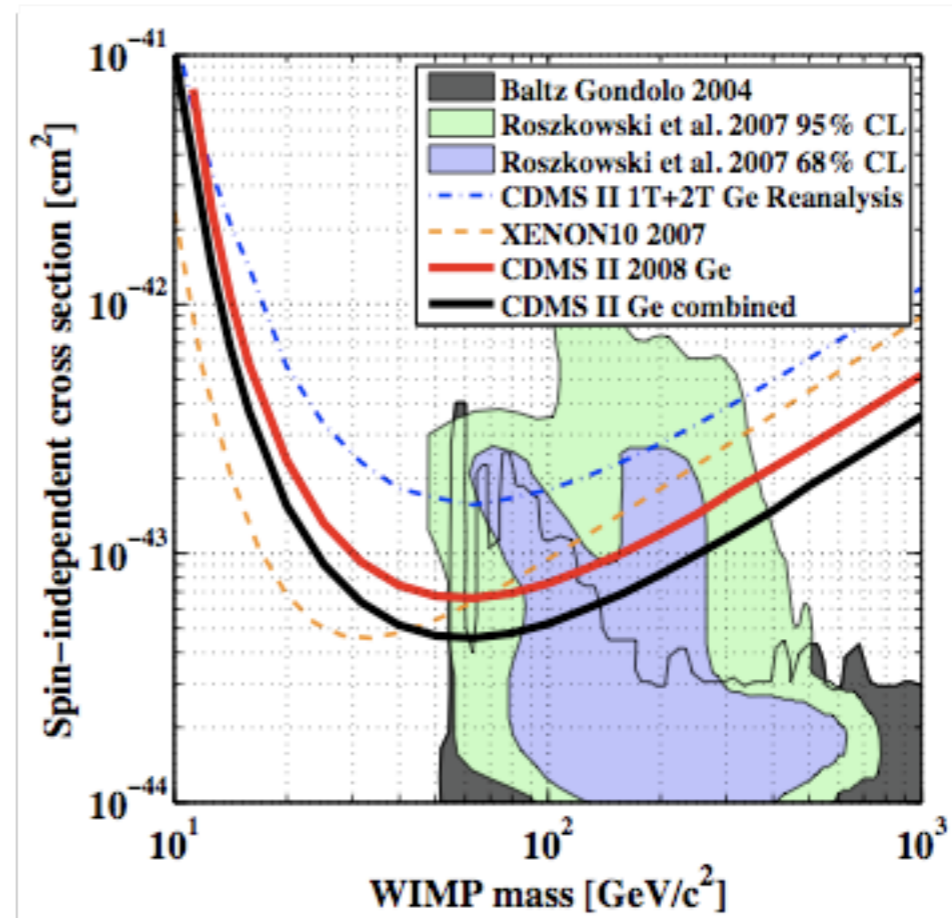
Based on arXiv:0904.1442



Cosmic Rays



Direct Detection



Colliders

LHC Alive! (?)

Seeing the Light... from Dark Matter

- One way to (indirectly) detect Dark Matter is via annihilations into photons.
 - Searches performed with Fermi/GLAST and/or ground-based ACT's.
 - Advantage of γ 's \rightarrow travel in straight lines w/o much energy loss... identify sources and/or trace dark matter distribution?
- Expected spectrum of γ 's from DM annihilation has two main components:
 - "Secondary" γ 's from annihilation into charged SM particles which then radiate γ 's OR hadronize and decay into γ 's ($\pi^0 \rightarrow \gamma\gamma$).

The result is a CONTINUUM of γ 's.

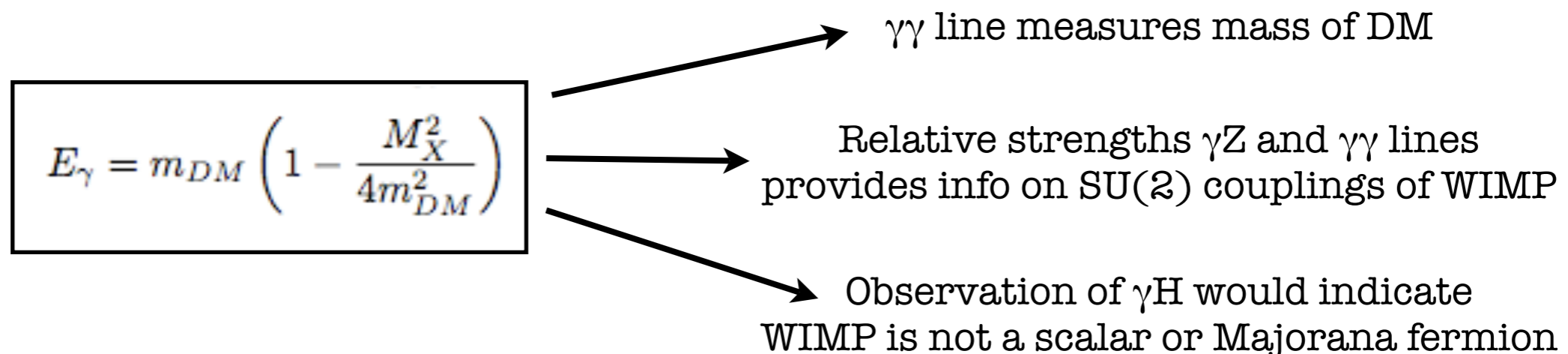
- "Primary" γ 's from LOOP-LEVEL annihilations into final states involving γ 's ($\text{DM} + \text{DM} \rightarrow \gamma + \text{X}$).

The result is (are) mono-energetic line(s) super-imposed onto continuum.

- Generally, lines are expected to be small (loop-suppressed) compared to the continuum.
- However, if observable, lines provide a nice discriminant against astrophysical backgrounds.

The WIMP Forest

- What if the nature of DM is such that continuum emission is suppressed while production of “direct” photons is via “large-ish” couplings?
 - The former can occur if DM annihilates mainly into “photon-unfriendly” states (such as Z’s or Higgs)
 - The latter can occur, for example, in models where the WIMP is related to the SM hypercharge gauge boson.
- The position/strength of lines can provide a wealth of information about DM.



- If there are other particles in the “dark sector” with masses appreciable to the WIMP mass (but $\leq 2m_{DM}$), you could possibly observe a SERIES of LINES... i.e., a WIMP FOREST!!!

Case Study: The Chiral Square

- Model of Universal Extra Dimensions (UEDs) with TWO extra dimensions (Dobrescu & Ponton, JHEP 0403 (2004) 071)

- Identify adjacent sides of square:

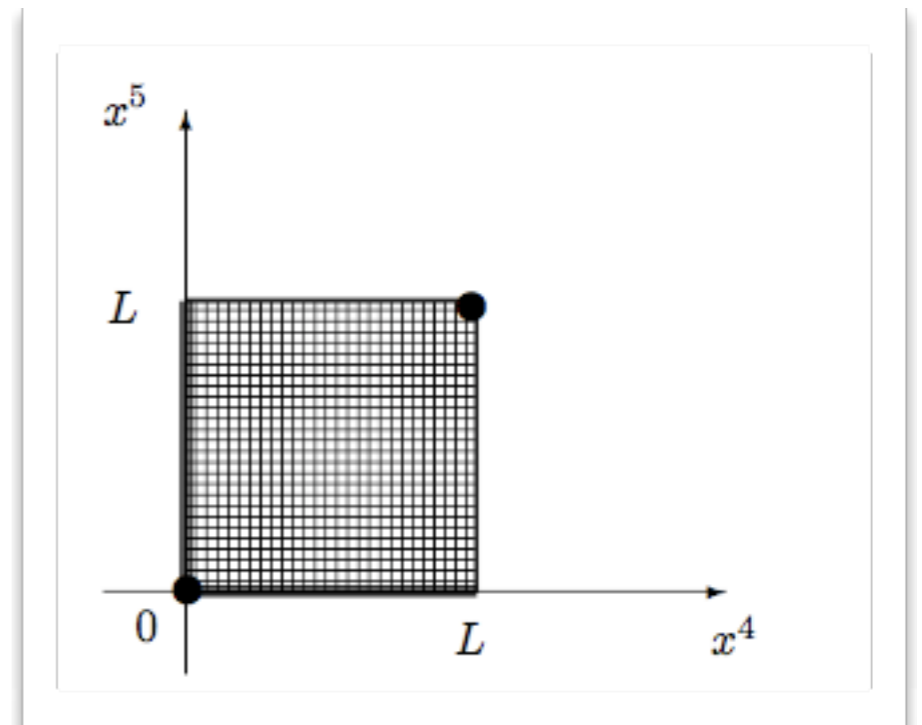
$$(y, 0) \equiv (0, y) \quad (y, L) \equiv (L, y)$$

- Kaluza-Klein (KK) modes identified by TWO indices (j,k)

$$k \geq 0, \quad j \geq 1 - \delta_{k,0}$$

- Mass eigenvalues:

$$M_{(j,k)}^2 = M_0^2 + \pi^2 \frac{j^2 + k^2}{L^2},$$

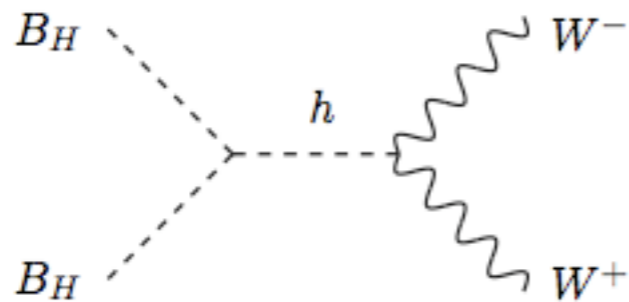


- Important for us: the lightest (non-SM) modes are the (1,0) modes, while the next-to-lightest modes are the (1,1) with masses $\approx \sqrt{2}$ times heavier!
(Compare to 5-d model where level-2 modes are TWICE as heavy)

The WIMP Candidate

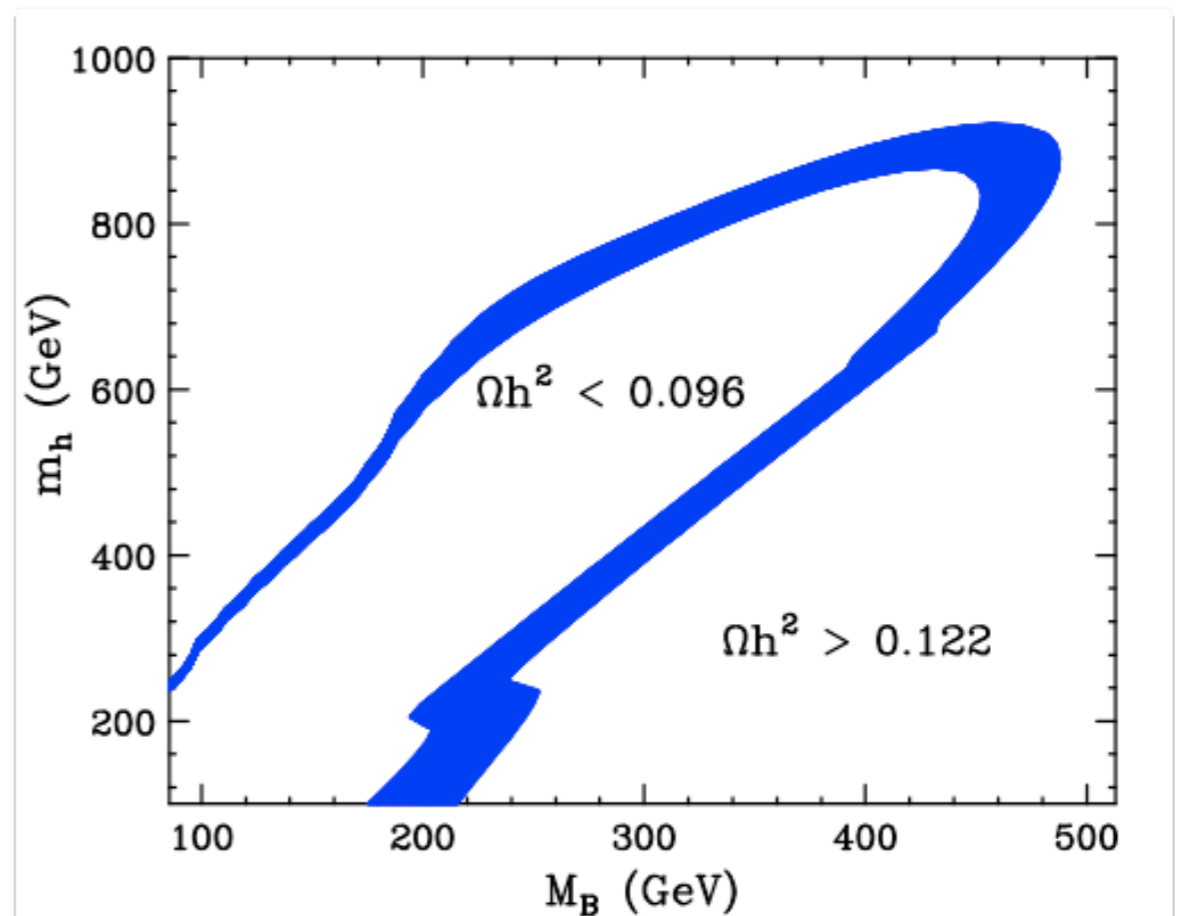
(Dobrescu et al., JCAP 0710:012,2007)

- Residual spacetime symmetry \rightarrow lightest KK particle (LKP) is stable
($j + k = \text{odd (even)} \rightarrow$ state is odd (even) under symmetry)
- LKP expected to be SCALAR partner of the hypercharge gauge boson
(a.k.a. “spinless photon”) $B^{(1,0)} \equiv B_H$
- Scalar LKP = suppressed annihilation to fermions...
annihilates mainly into W’s, Z’s and Higgs.
- Thermal relic abundance very sensitive to B_H mass... as well as Higgs mass...



- WMAP constraints:

$$200 \text{ GeV} \lesssim M_B \lesssim 500 \text{ GeV}$$

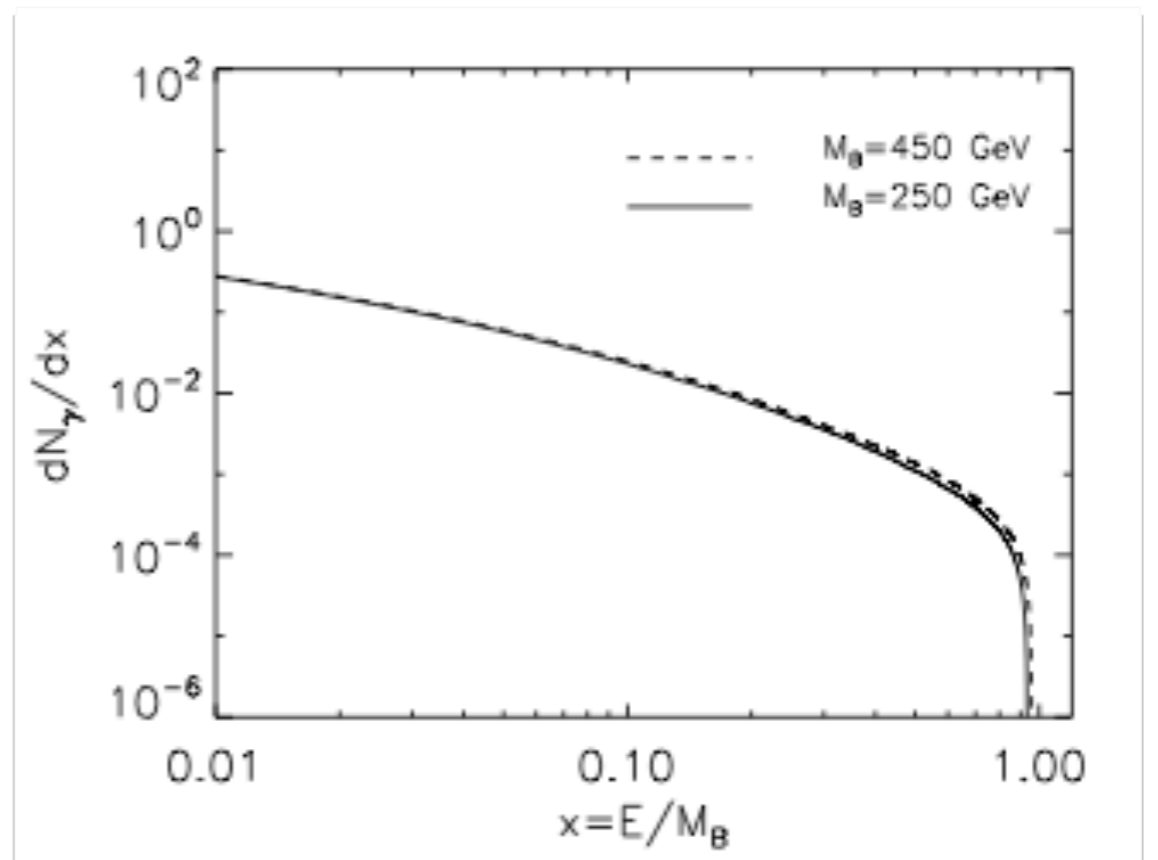


Gamma Ray Continuum

- In general, continuum consists of two components:
 - quark fragmentation/decay ($\pi^0 \rightarrow \gamma\gamma$): featureless, SOFT spectrum
 - final-state radiation: harder component of the form...

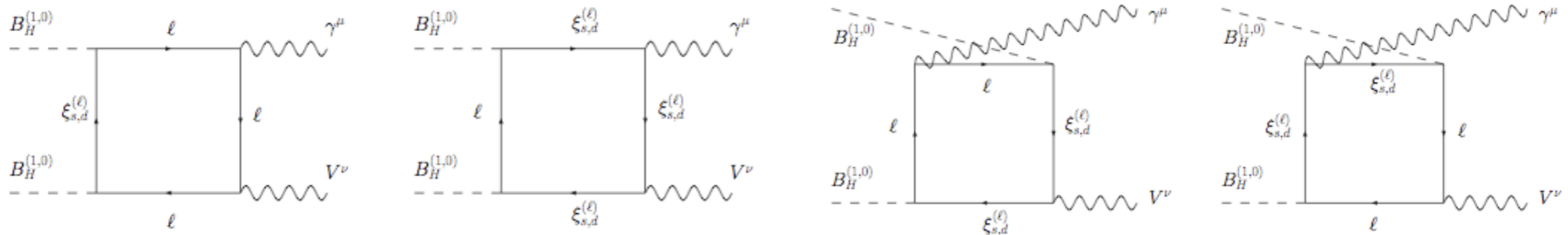
$$\frac{dN_{XX}}{dx} \approx \frac{\alpha Q_X^2}{\pi} \mathcal{F}_X(x) \log\left(\frac{s(1-x)}{m_X^2}\right) \longleftrightarrow \begin{aligned} \mathcal{F}_{\text{fermion}}(x) &= \frac{1 + (1-x)^2}{x} \\ \mathcal{F}_{\text{boson}}(x) &= \frac{1-x}{x} \end{aligned}$$

- Scalar nature of LKP suppresses continuum...
- Annihilation into “photon-unfriendly” states (ZZ + HH) 50% of the time.
- Continuum spectrum reminiscent of neutralino



Calculation of the “Lines”

- Annihilation to $\gamma + V$ final states proceed via box diagrams:



where $V = \gamma, Z$ or $B^{(1,1)}$ mode.

- In general, the amplitude takes the form:

$$\mathcal{M} = \epsilon_A^{\mu*}(p_A) \epsilon_B^{\nu*}(p_B) \mathcal{M}^{\mu\nu}(p_1, p_2, p_A, p_B) \longrightarrow \mathcal{M}^{\mu\nu} = A_1 g^{\mu\nu} + B_1 p_1^\mu p_1^\nu + B_2 p_2^\mu p_2^\nu + B_3 p_1^\mu p_2^\nu + B_4 p_1^\nu p_2^\mu + B_5 p_A^\nu p_B^\mu + B_6 p_1^\mu p_A^\nu + B_7 p_1^\nu p_B^\mu + B_8 p_2^\mu p_A^\nu + B_9 p_2^\nu p_B^\mu.$$

- Use the fact that WIMPs are non-relativistic (NR): $p_1 \simeq p_2 \simeq p \equiv (M_B, \mathbf{0})\dots$ and other tricks... A_1 is the dominant contribution.
- “NR-ness” of WIMPs causes problems when computing the loops... specifically the Passarino-Veltman (PV) TENSOR coefficients.
- Reason: PV coefficients depend INVERSELY on Gram Determinant (GD):

$$\text{GD} = \det(p_i \cdot p_j)$$

Line Cross Sections

- Used technique developed by R. Stuart (Comput. Phys. Commun. 48, 367 (1988))

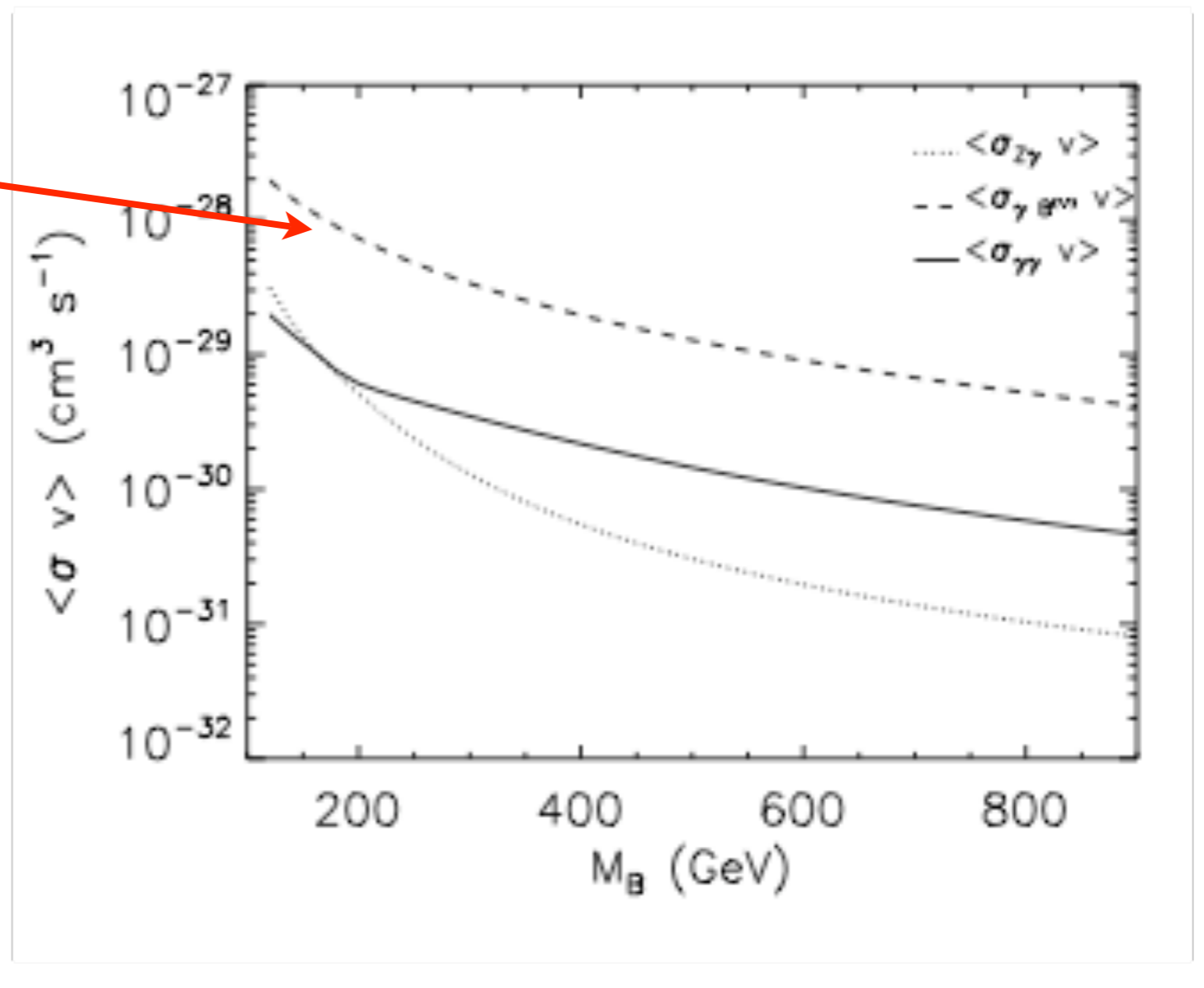
$$D_{27} = \alpha_{123}C_{24}(123) + \alpha_{124}C_{24}(124) \\ + \alpha_{134}C_{24}(134) + \alpha_{234}C_{24}(234),$$

and α 's are solved for assuming GD exactly vanishes.

Note relative size
of $B^{(1,1)}$ cross section!

Due to less cancellation
at the amplitude level

Z and $B^{(1,1)}$ particles
decay (broadening of line)



Putting it All Together

- The differential flux (@ angle ψ w.r.t. GC) :

$$\frac{d\Phi_\gamma}{d\Omega dE}(\psi, E) = \frac{r_\odot \rho_\odot^2}{4\pi M_{BH}^2} \frac{dN_\gamma}{dE} \int_{\text{l.o.s.}} \frac{ds}{r_\odot} \left[\frac{\rho[r(s, \psi)]}{\rho_\odot} \right]^2 \longrightarrow \frac{dN_\gamma}{dE} = \sum_f \langle \sigma v \rangle_f \frac{dN_\gamma^f}{dE},$$

- Separate particle physics from astrophysics:

$$J \equiv \int_{\text{l.o.s.}} \frac{ds}{r_\odot} \left[\frac{\rho[r(s, \psi)]}{\rho_\odot} \right]^2 \longrightarrow \bar{J}(\Delta\Omega) = \frac{1}{\Delta\Omega} \int_{\Delta\Omega} J(\psi) d\Omega.$$

- Consider two profiles: Navarro-Frenk-White (NFW) and “Adiabatic”

Model	$\bar{J} (10^{-5})$
NFW	1.5×10^4
Adiabatic	4.7×10^7

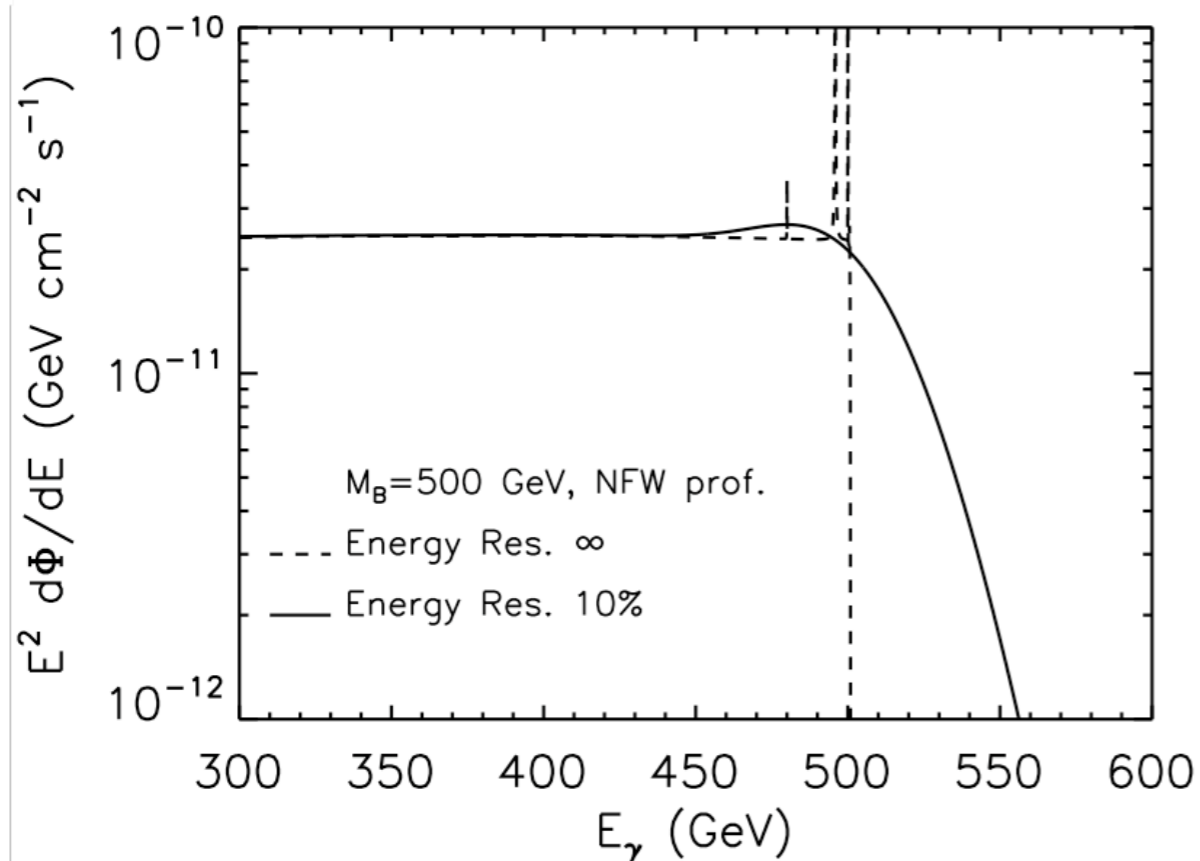
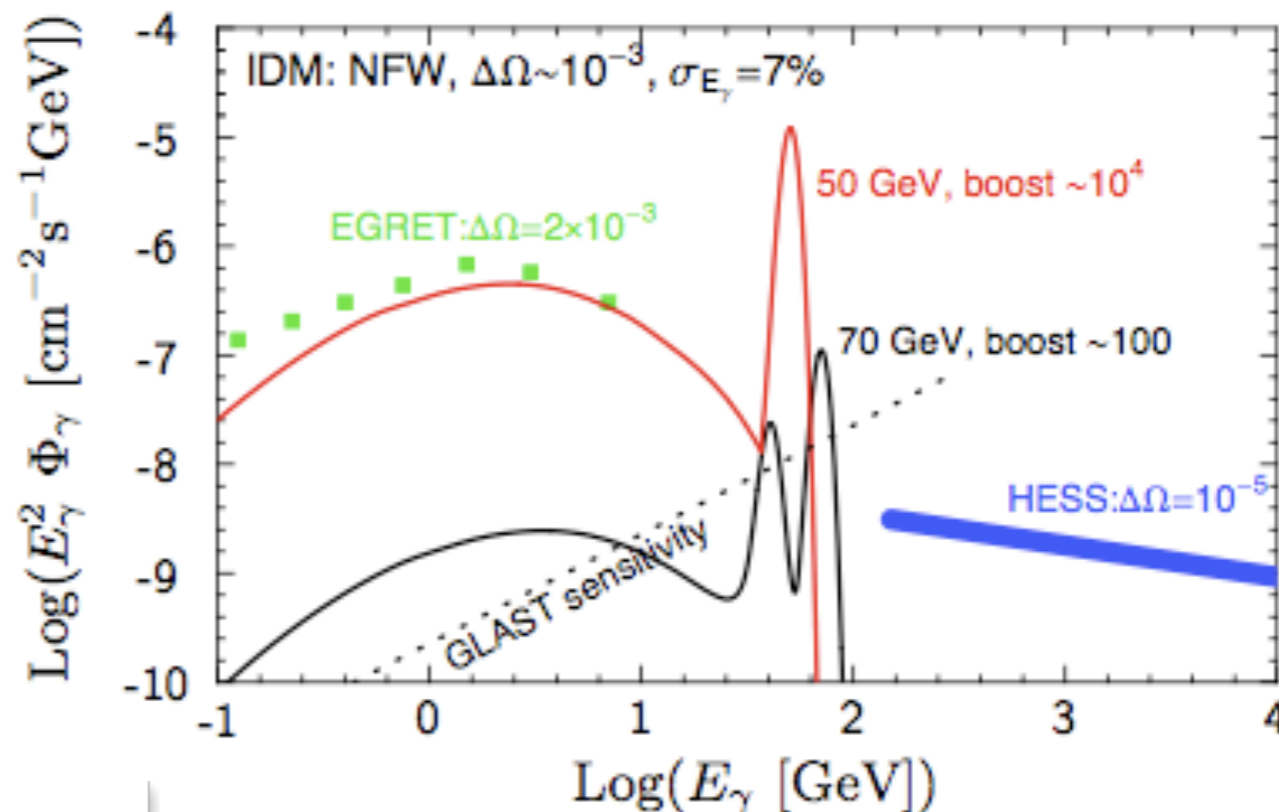
- Good news: signals like the ones considered here will “trace” DM distribution!
- To account for detector resolutions, convolve “signal” with Gaussian:

$$S_M(E_0) = \int dE G(E, E_0) S(E). \longrightarrow G(E, E_0) = \frac{1}{\sqrt{2\pi} E_0 \sigma} \exp \left[-\frac{(E - E_0)^2}{2\sigma^2 E_0^2} \right]$$

Setting the Stage

“The Inert Doublet Model”
(Gustafsson et al., PRL99: 041301, 2007)

WIMP = scalar
Lines = $\gamma\gamma$ and γZ



“5-d UED Model”
(Work in progress)

WIMP = vector gauge boson
Lines = $\gamma\gamma$, γZ and γH

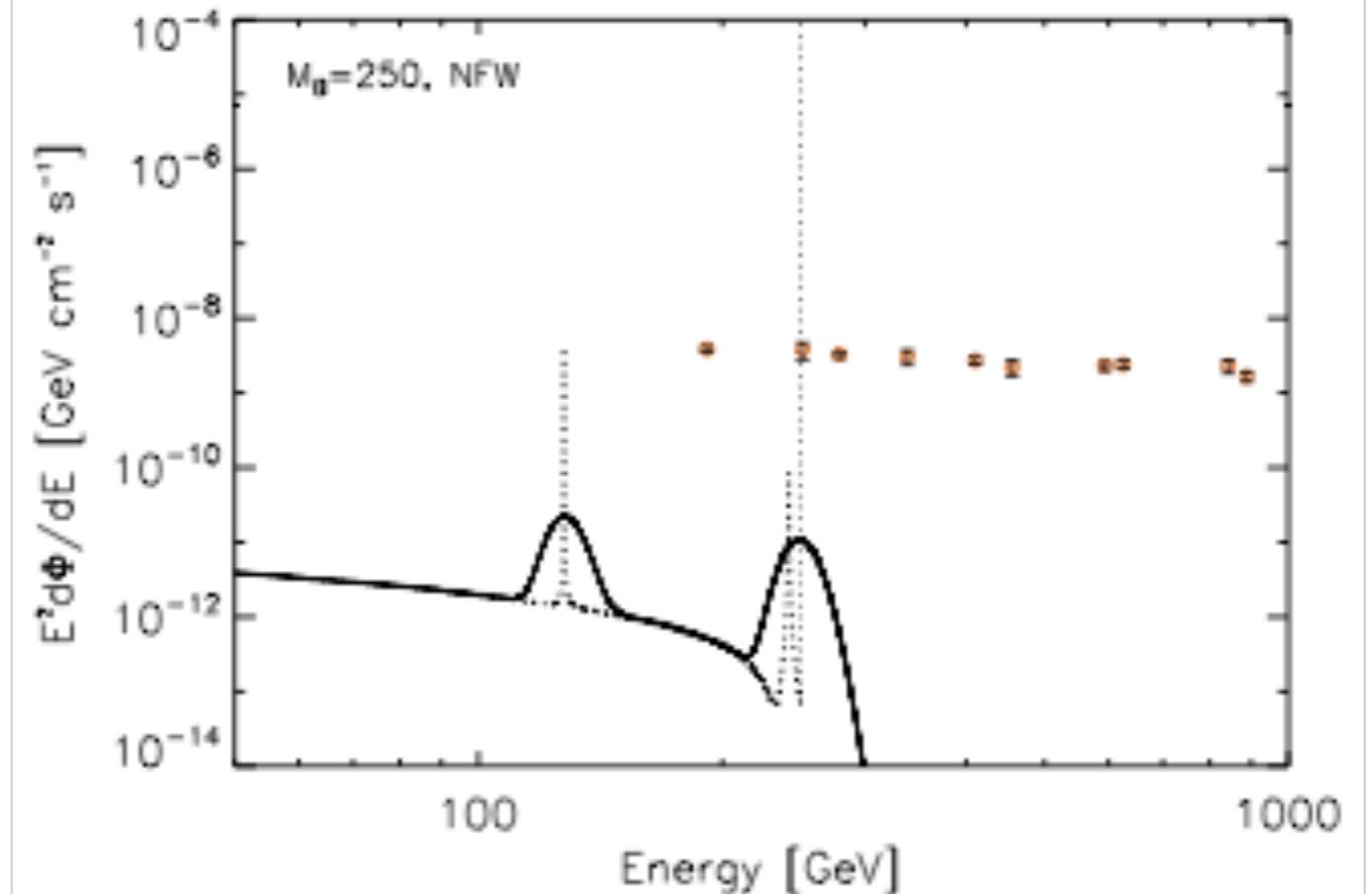
Results for the Chiral Square

- “WIMP Forest” characterized by 3 lines (and suppression of continuum)
- $\gamma\gamma$ line (@ $E_\gamma = M_B$), $Z\gamma$ (very close to M_B) and $B^{(1,1)}\gamma$ line
- Detector resolution ($\Delta E/E = 10\%$) effects smear $\gamma\gamma$ and $Z\gamma$ lines into one “bump”
- Distinctive feature of the Chiral Square model is the $B^{(1,1)}\gamma$ line!
- Since $B^{(1,1)}$ mass is comparable to WIMP mass...

$$E_\gamma = m_{DM} \left(1 - \frac{M_X^2}{4m_{DM}^2} \right)$$

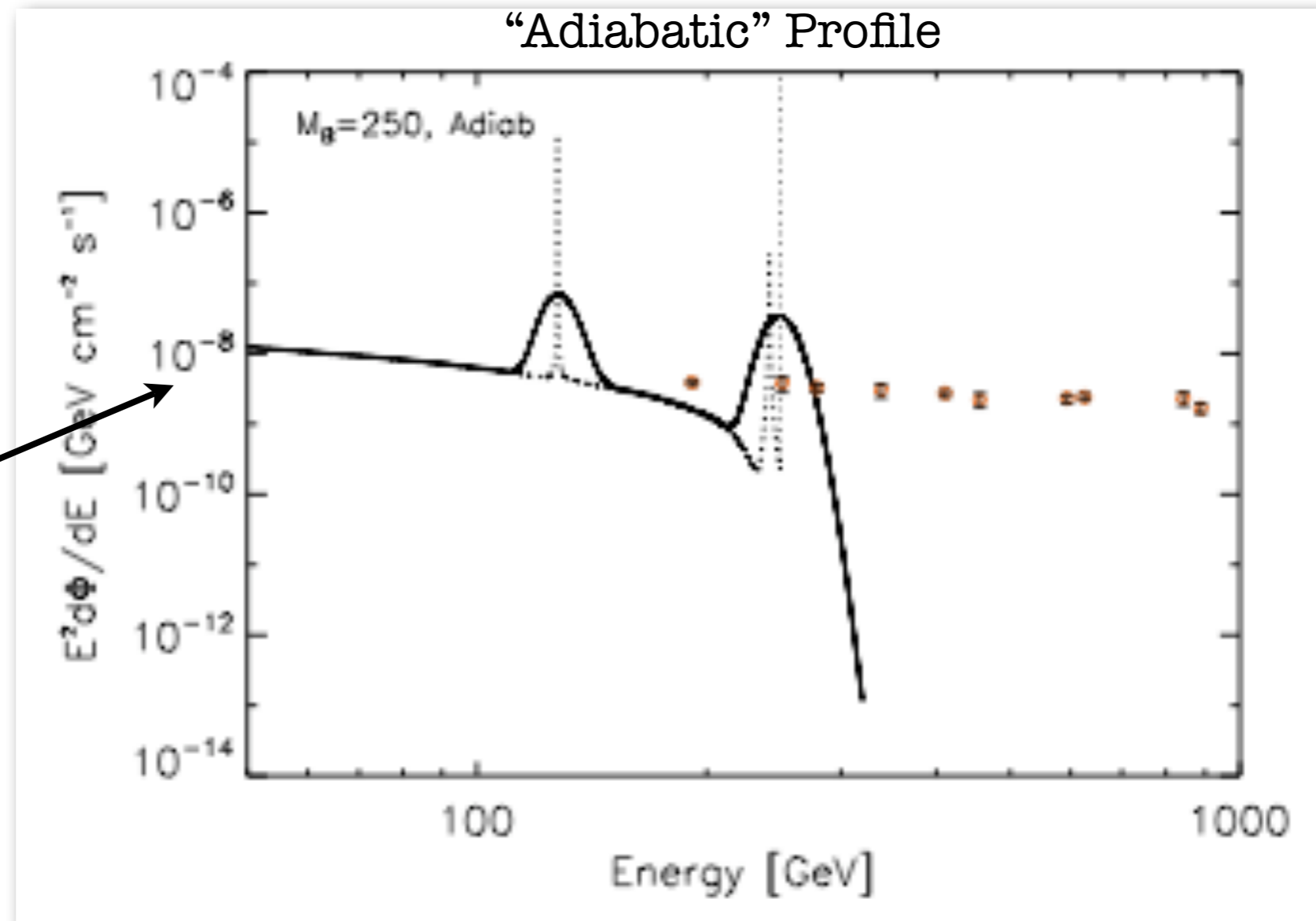
is well-separated from $\gamma\gamma + Z\gamma$ bump and distinguishable... even after detector resolution effects are accounted for!

NFW Profile



Data = HESS point source J1745-290
(foreground for our signal)

Large Boost Factors?



Note change
in scale

Conclusions

- I can't believe I made it this far in 13 minutes!
- Search for DM via annihilation into photons:
 - Particle physics: information on DM properties(?)
 - Astrophysics: trace out DM distribution!
- “WIMP Forest”: if there are other particles in the “dark sector” with masses appreciable to the WIMP mass (but $\leq 2m_{\text{DM}}$), you could possibly observe a SERIES of lines!
- Case study: 6-d UED “Chiral Square”
 - WIMP = scalar (“spinless”) photon
 - Relic Density: $200 \text{ GeV} \approx M_{\text{B}} \approx 500 \text{ GeV}$
- Gamma-ray Spectrum:
 - Suppressed continuum spectrum
 - VERY DISTINCTIVE TWO-BUMP FEATURE!