

# LHC Phenomenology of SUSY multi-step GUTs

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W. Kilian, JR, PL **B642** (2006), 81; arXiv:0709.4202;  
work in progress (with F. Braam, K. Mallot, D. Wiesler)

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# Motivation for (SUSY) Unification

## Incompleteness/Theoretical Dissatisfaction

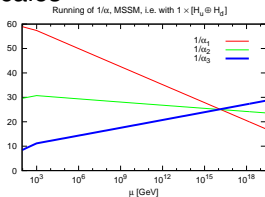
EWSB,  $H$ ,  $m_\nu$ , DM, hierarchy, . . . . ., reducible representation:

$$q(\mathbf{3}, \mathbf{2})_{\frac{1}{3}} \oplus \ell(\mathbf{1}, \mathbf{2})_{-1} \oplus u^c(\overline{\mathbf{3}}, \mathbf{1})_{-\frac{4}{3}} \oplus d^c(\overline{\mathbf{3}}, \mathbf{1})_{\frac{2}{3}} \oplus e^c(\mathbf{1}, \mathbf{1}) \oplus H(\mathbf{1}, \mathbf{2})_1$$

Supersymmetry: consistent extrapolation to high scales

- ⇒ unification quantitatively testable
- ⇒ two Higgs doublets  $H^u, H^d$
- ⇒ TeV-scale SM-superpartners

Bottom-Up Approach: *just MSSM*



Unification verification only with megatons? What about colliders?

- ▶ SPA: super precision accurately
- ▶ **Look for chiral exotics**
- ▶ Physics beyond MSSM provides handle to GUT scale

# The Doublet-Triplet Splitting Problem

MSSM Higgses included in  $\mathbf{5}_H \oplus \bar{\mathbf{5}}_H$

$$\mathbf{5}_H = (\mathbf{3}, \mathbf{1})_{-\frac{2}{3}} \oplus (\mathbf{1}, \mathbf{2})_1 : \begin{pmatrix} D \\ H_u \end{pmatrix} \quad \bar{\mathbf{5}}_H = (\bar{\mathbf{3}}, \mathbf{1})_{\frac{2}{3}} \oplus (\mathbf{1}, \mathbf{2})_{-1} : \begin{pmatrix} D^c \\ \epsilon H_d \end{pmatrix}$$

$D, D^c$  colored triplet Higgses with charges  $\pm \frac{1}{3}$  (EW singlet)  
colored Dirac fermion  $\tilde{D}$  with charge  $-1/3$  (EW singlet)

Unification requires omitting colored part of SU(5) Higgs  $\mathbf{5}_H, \bar{\mathbf{5}}_H$

- 1) **Doublet-triplet splitting problem** ( $m_H \sim 100 \text{ GeV}, m_D \sim 10^{16} \text{ GeV}$ )

Welcome, since  $SU(5)$ -symmetric Higgs interactions would read

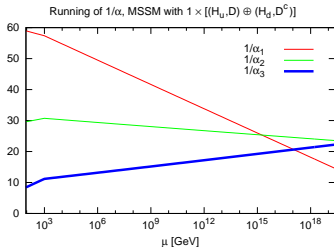
$$\begin{aligned} \bar{\mathbf{5}} \mathbf{10} \bar{\mathbf{5}}_H &= \ell H_d e^c + q \epsilon H_d d^c + q \epsilon \ell D^c + d^c u^c D^c \\ \mathbf{10} \bar{\mathbf{5}}_H \mathbf{10} &= q \epsilon H_u u^c + D u^c e^c + D q \epsilon q \end{aligned}$$

Generating SM masses  $\Rightarrow$  leptoquark *and* diquark coupl. for  $D, D^c$   
 $\Rightarrow$  triggers **rapid proton decay**

# Doublet-Triplet Splitting

Possible scenarios:

1. Colored singlets are heavy (GUT scale) = **doublet-triplet splitting**
  - ▶ enables exact unification near  $10^{16}$  GeV, excludes rapid proton decay
  - ▶ Proton decay may still be too fast (depends on superpotential)
  - ▶ Doublet-triplet splitting is not trivially available
2. Colored singlets are **light (TeV scale)**
  - ▶ Simple unification no longer happens near  $10^{16}$  GeV, nor elsewhere



- ▶ Proton-decay coupl. must be excluded: consistent with GUT symmetry?

## Further MSSM Issues

2)  **$\mu$  problem:** SUSY  $\mu$ -term  $\mu H_u H_d$ , not related to soft breaking

Why  $\mu \sim \mathcal{O}(100 \text{ GeV})$ , not  $\mathcal{O}(10^{16} \text{ GeV})$ ?

⇒ Possible extension as a solution: singlet Higgs  $S$  with superpotential

$$\lambda S H_u H_d \rightarrow \lambda \langle S \rangle H_u H_d = \mu H_u H_d$$

⇒ NMSSM:  $\langle S \rangle$  should be somehow related to soft-breaking

Large top Yukawa coupl. drives effective  $H_u$  mass squared negative:

This mechanism may also be responsible for  $\langle S \rangle$  in the NMSSM

- ▶ requires existence of vectorlike pair of chiral superfields
  - ▶ for instance,  $D$  and  $D^c$  (colored) with  $SDD^c$
  - ▶ ... as required by  $SU(5)$ , if  $S H_u H_d$  is present, gives Dirac mass to  $D$
- ▶ Without tree-level quartic coupl., the CW mechanism implies  $\langle S \rangle \sim 4\pi m_{\text{soft}}$ , so  $\langle S \rangle \gg \langle H \rangle$ .

3) **Higgs-matter unification:** Why only one family of Higgs matter?  $E_6$  unifies Higgs fields with SM matter... Possible scenarios:

1. Omit one bi-triplet  $D, D^c$  family ⇒ doublet-triplet splitting
2. Add one extra MSSM Higgs family ⇒ ESSM
3. Different unification pattern

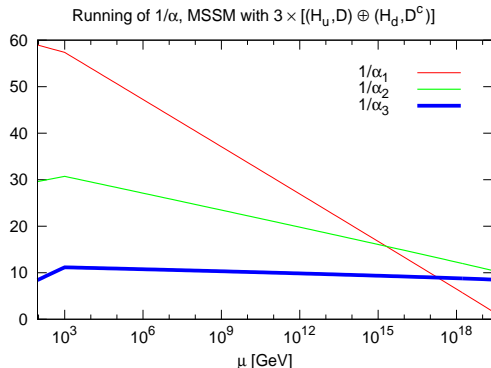
# Running Couplings: PSSSM

[Braam/JR/W.Kilian, 2006/2009]

- ▶ Additional particles spoil simple unification
- ▶ Gauge coupling unification below  $\Lambda_{Planck}$  due to intermediate Pati-Salam

$$SU(4) \times SU(2)_L \times SU(2)_R [\times U(1)_X]$$

symmetry at  $\sim 10^{16}$  GeV



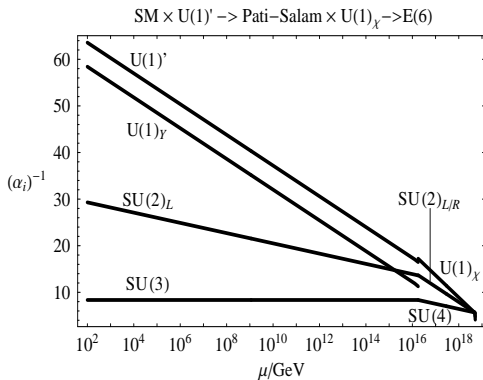
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- ▶  $SU(2)_R$  and  $SU(2)_L$ : identical particle content
- $\Rightarrow$  Identical running
- ▶ Crossing of  $SU(4)$  and  $SU(2)_{L/R}$  couplings determines  $E_6$  breaking scale

# $E_6$ unification with intermediate Pati-Salam Symmetry

The quantum numbers of the **27** under the Pati-Salam gauge group are

$$\mathbf{Q}_R = ((u^c, d^c), (\nu^c, e^c)) = (\bar{\mathbf{4}}, \mathbf{1}, \mathbf{2})_{\frac{1}{2}}$$

$$\mathbf{Q}_L = ((Q, L)) = (\mathbf{4}, \mathbf{2}, \mathbf{1})_{\frac{1}{2}}$$

$$\mathbf{H} = (H^u, H^d) = (\mathbf{1}, \mathbf{2}, \mathbf{2})_{-1}$$

$$\mathbf{D} = (D, D^c) = (\mathbf{6}, \mathbf{1}, \mathbf{1})_{-1}$$

$$\mathbf{S} = (\mathbf{1}, \mathbf{1}, \mathbf{1})_2$$

▶ Lepton number becomes 4<sup>th</sup> color

$$\text{▶ } T_{SU(4)}^{15} \propto \frac{B-L}{2}$$

$$\text{▶ } Y = \frac{B-L}{2} + T_R^3$$

▶ Integrating out  $\nu^c$

⇒ appropriate breaking

⇒ see-saw mechanism  
(i.e. naturally small  
neutrino masses)



# Flavor Symmetry

## Proton decay?

- ▶ Triplets included: PS-symmetric superpotential contains leptoquark and diquark couplings  $DQ_RQ_R = \epsilon_{\alpha\beta\gamma}\epsilon_{jkl}D_\alpha(Q_R)_{\beta j}(Q_R)_{\gamma k}$

Possible solution: **extra flavor symmetry**  $SU(3)_F$  (or  $SO(3)_F$ )

$\Rightarrow$   $D$  diquark coupling with  $SU(2)_R, SU(3)_c, SU(3)_F$ :

$$DQ_RQ_R = \epsilon^{abc}\epsilon_{\alpha\beta\gamma}\epsilon_{jkl}D_\alpha^a(Q_R)_{\beta j}^b(Q_R)_{\gamma k}^c$$

Vanishes due to total antisymmetry  $\Rightarrow$  **no proton decay**

Analogous for  $\epsilon^{abc}\epsilon_{\alpha\beta\gamma}\epsilon_{jkl}(D^c)_\alpha(Q_L)_{\beta j}^b(Q_L)_{\gamma k}^c$

- ▶ **Leptoquark coupling of  $D$  not affected**

Eff. superpotential from (spontan.) breaking of LR and/or flavor symm.:

- ▶ Exclude spurions  $\propto \epsilon_{\alpha\beta\gamma}$  (color space)  $\Rightarrow$  diquark couplings absent
- ▶ **baryon number as low-energy symmetry, flavor symmetry not (necessarily)**

# The LHC phenomenology

[Braam/JR/Wiesler]

Next step: **Provide a viable low-energy spectrum**

## At LHC:

### 1) 1 – 3 pairs of scalar leptoquarks $D_L, D_R$ .

- ▶ probably heavy  $\lesssim 1$  TeV (but hierarchy is possible)
- ▶ pair-produced in  $gg$  fusion at LHC
- ▶ decay into  $\ell u$  and  $\nu d$ :
  - generation-diagonal, or just third-generation:  $\tau t$  and  $\nu b$  or
  - generation-crossed (flavor symmetry!):  $ec, eb, \mu d, te, t\mu \dots$
  - $gq \rightarrow D\ell$  production enhanced
  - or, if  $R$ -parity is violated, may mix with down-type squarks.

### 2) 1 – 3 fermionic leptoquarkinos $\tilde{D}$

- ▶ are probably heavy as well, but somewhat lighter than scalars (because  $m^2 = \lambda \langle S \rangle^2 + m_{\text{soft}}^2$ )
- ▶ are also pair produced (maybe singly if  $R$ -parity is violated)
- ▶ decay into  $\tilde{\ell} j$ , or  $\ell \tilde{q}$ , or  $\nu \tilde{q}$ 
  - rich signatures!
  - spin measurement distinguishes from ordinary squarks

# A little bit of Pheno

## 3) (non)"standard" MSSM Higgs

- ▶ Relaxed Higgs bounds (like in NMSSM)
- ▶ Possibly large invisible decay ratio ( $\tilde{\chi}^0, a$ )

## 4) 2 – 4 doublets of unhiggses

- ▶ probably only pair-produced: Drell-Yan, maybe Higgs decays (singlets involved)
- ▶ missing-energy signatures, unique identification could be difficult: ILC?

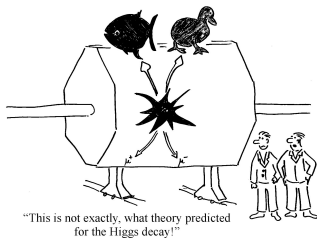
## 5) 1 – 3 singlet scalars + pseudoscalars

- ▶ masses, properties?

## 6) and all associated neutralinos ( $\leq 11$ ) and charginos ( $\leq 4$ )

- ▶ large and complicated chargino/neutralino mixing matrices. Decay chains at LHC become difficult to understand.

## 7) Either heavy $Z'$ (gauged NMSSM) or light pseudo-axion(s) $\eta$ corresponding to extra $U(1)$



**Conclusion:** LHC phenomenology rich

...

and confusing

# Investigation Of The Parameter Space

- ▶ # free parameters  $\sim \mathcal{O}(100)$ , additional assumptions:
  - Unified SSB-parameters
  - Flavor structure
 ⇒ Limitation on 14 free parameters
- ▶ Constraints:
  - (1) Experimental lower bounds on masses of new particles
  - (2) Running parameters perturbative up to  $\Lambda_{E_6}$
  - (3) Scalar (non-Higgs) mass terms have to remain positive  
(⇔ No unwanted symmetry breaking)

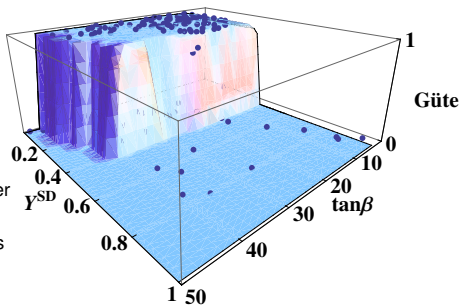
- ▶ 14-dim parameter space

⇒ grid scanning  $\rightarrow 10^{28}$  points

- ▶ Investigation per point (RGE, Higgs potential minimization, calculation of masses)  $\sim 5$  s

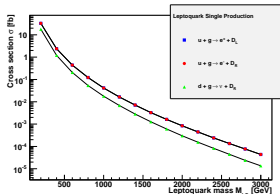
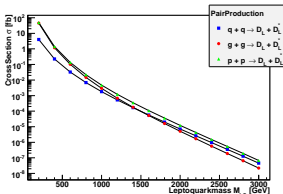
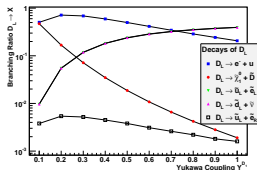
Sol: Monte-Carlo Markov-Chain through parameter space

⇒ Effective search for sensible parameter tuples



# Predictions for Collider Experiments

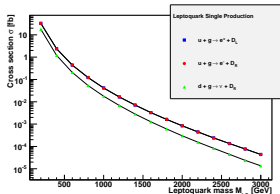
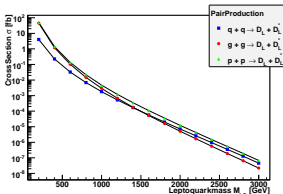
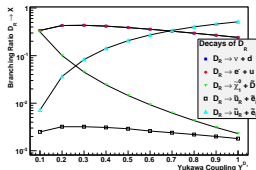
- ▶ Collider phenomenology with event generator WHIZARD  
[Kilian/Ohl/JR]
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- ▶ **First analyses:** BRs, cross sections for scalar leptoquarks, S/B
- ▶ In progress: leptoquarkino pheno



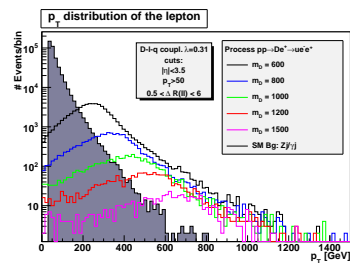
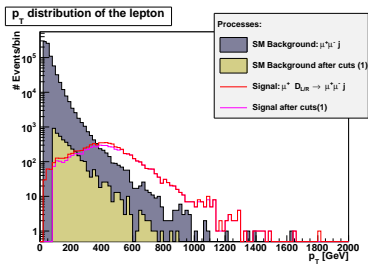
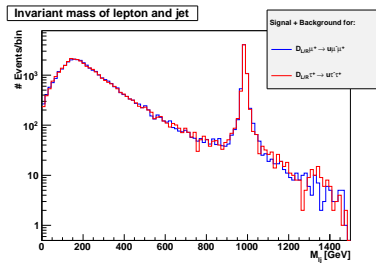
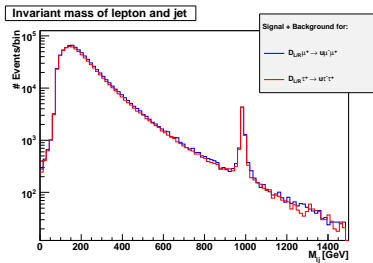
Cuts		Background	$m_D = 0.6 \text{ TeV}$		$m_D = 0.8 \text{ TeV}$		$m_D = 1.0 \text{ TeV}$	
$p_T$	$M_{\ell\ell}$	$N_{BG}$	$N_1$	$S_1/\sqrt{B}$	$N_2$	$S_2/\sqrt{B}$	$N_3$	$S_3/\sqrt{B}$
50	10	413274	64553	<b>93</b>	14823	<b>23</b>	4819	<b>7</b>
100	150	3272	40749	<b>194</b>	10891	<b>92</b>	3767	<b>45</b>
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# Dark Matter

MSSM Higgses:  $H_u^f, H_d^f$  with  $f = 1, 2, 3$

- \* VEV selects single direction (taken as  $f = 3$ ) in family space
- ⇒ 1 gen. MSSM Higgses, 2 gen. “unhiggses”

Ellis et al., 1985; Campbell et al., 1986

(2 bi-doublets = 8 charged and 8 neutral scalars + fermion superpartners)

In gauge interactions, unhiggses are pair-produced, thus suppressed in precision data, ... .. but also Yukawa interactions

- 1) FCNC
- 2) resonant single production in  $q\bar{q}$  or  $e^+e^-$  annihilation

Unhiggses very heavy *or* artificially aligned *or* suppressed

⇒ (approximate?)  $H$  parity: odd for unhiggses, even otherwise

*And why not? Flavor symmetry removes the need for  $R$  parity anyway.*

If  $H$  parity is exact:

- ▶ lightest unhiggs:  $H$  parity protected dark matter
- ▶ Pair production of unhiggses/unhiggsinos, cascade decays

... *and  $R$  parity is exact:*

- ▶ dark matter mix: interesting relic abundance  
(relaxes all neutralino bounds!)



# Summary/Outlook

## 3 independent building blocks for exotic SUSY phenomenology

### Color-triplet leptoquark/inos are present in the low-energy spectrum

- ▶ multi-step unification (PSSSM:  $E_6$  with PS symmetry)
- ▶ favoring PS symmetry above the R-neutrino mass scale

### Flavor symmetry prohibits proton decay

- ▶ instead of (or in addition to)  $R$  parity
- ▶ Superpotential terms are due to GUT- and flavor-breaking
- ▶ ... therefore do not exhibit GUT relations

### Higgs sector is flavored

- ▶ Unhiggses (1st and 2nd generation) carry conserved quantum number
- ▶ Unhiggses dark matter candidates
- ▶ Ordinary MSSM stuff might decay via  $R$ -parity violation

Investigation of the parameter space  $\Rightarrow$  Input for WHIZARD

Leptoquark phenomenology could hint toward GUT

### Further Studies:

- ▶ PS-symmetry breaking,  $U(1)$ -mixing, threshold corrections
- ▶ More LHC pheno (leptoquarkinos, Higgs, weak inos,  $Z'$  ...)



# Backup: $E_6$ Particle Content

- ▶ Unifies Higgs and matter fields
- ▶ Contains  $SU(3)_c \times SU(2)_L \times U(1)_Y [\times U(1)'] \Rightarrow$  additional  $Z'$
- ▶ **Ansatz:** all new particles in the spectrum at TeV scale
- ▶ SM  $\times U(1)'$  quantum numbers of the fundamental **27**:

**Notation:**  $(SU(3), SU(2))_{Y, Q'}$

$$Q_L = (\mathbf{3}, \mathbf{2})_{\frac{1}{6}, Q'_Q}$$

$$u^c = (\bar{\mathbf{3}}, \mathbf{1})_{-\frac{2}{3}, Q'_u}$$

$$d^c = (\bar{\mathbf{3}}, \mathbf{1})_{\frac{1}{3}, Q'_d}$$

$$H^u = (\mathbf{1}, \mathbf{2})_{\frac{1}{2}, Q'_{H^u}}$$

$$H^d = (\mathbf{1}, \mathbf{2})_{-\frac{1}{2}, Q'_{H^d}}$$

$$S = (\mathbf{1}, \mathbf{1})_{0, Q'_S} \neq 0$$

$$L_L = (\mathbf{1}, \mathbf{2})_{-\frac{1}{2}, Q'_L}$$

$$\nu^c = (\mathbf{1}, \mathbf{1})_{0, Q'_\nu=0}$$

$$e^c = (\mathbf{1}, \mathbf{1})_{1, Q'_e}$$

$$D = (\mathbf{3}, \mathbf{1})_{-\frac{1}{3}, Q'_D}$$

$$D^c = (\bar{\mathbf{3}}, \mathbf{1})_{\frac{1}{3}, -Q'_D}$$

# Backup: Sample Implementation

## Toy Model (no dynamics!)

Extend  $E_6 \times SU(3)_F$  to  $E_8$

... by implementing  $N = 2$  supersymmetry:

- ▶ We have: matter  $27_3$  and gauge  $78_1 + 1_8$ .
- ▶ Add: mirror matter  $\overline{27}_3$
- ▶ supersymmetrize by adding *matter*  $78_1 + 1_8$  and *gauge*  $27_3 + \overline{27}_3$ .

Decomposition of reps. in  $E_8 \rightarrow E_6 \times SU(3)_F$ :

$$248 = 27_3 \oplus \overline{27}_3 \oplus 78_1 \oplus 1_8$$

Result: matter 248 and gauge 248 (fundamental = adjoint)

# Backup: Sample Implementation — Top Down

## 1. Somewhat below $M_{\text{Planck}}$

- ▶  $N = 2 \rightarrow 1$  breaking removes mirror matter:  $\langle (\mathbf{27}_3)_i^a (\overline{\mathbf{27}}_3)_j^b \rangle = \delta^{ab} \delta_{j,i+1}$
- ▶  $E_6$  zero mode of chiral matter  $\mathbf{27}_3$ , maybe adjoint matter  $\mathbf{78}_1$  and  $\mathbf{1}_8$
- ▶ Flavor  $SU(3)$  on the zero modes (would be anomalous) is broken by colorless spurions, e.g., condensate  $\langle \mathbf{1}_8 \rangle$ .
- ▶  $E_6$  is broken to  $G_{\text{PS}}$  by colorless spurions, e.g., bilinear = Higgs “ $\mu$  term”  $\langle \overline{\mathbf{1}}_{2,2} \overline{\mathbf{1}}_{2,2} \rangle$  in the  $\overline{\mathbf{27}}_3 \overline{\mathbf{27}}_3$  mirror representation
- ▶ Additional allowed spurion = Singlet  $\langle \overline{\mathbf{1}}_{1,1} \rangle = \langle \overline{S} \rangle$  (3. gen.)

Note: all spurions so far break flavor as well

## Result:

- ▶ PS symmetry
- ▶ all MSSM superpotential terms allowed, subject to PS symmetry/flavor constraints (no quark mixing):

$$\mathbf{27}_3 \mathbf{27}_3 \mathbf{27}_3 = 0 \quad (\mathbf{27}_3 \mathbf{78}_1 \overline{\mathbf{27}}_3), (\mathbf{78}_1 \mathbf{78}_1 \mathbf{78}_1), (\mathbf{27}_3 \mathbf{1}_8 \overline{\mathbf{27}}_3), \mathbf{1}_8 \mathbf{1}_8 \mathbf{1}_8$$

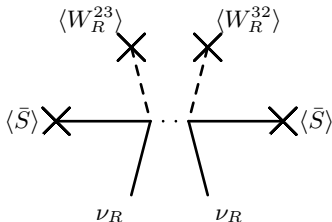
- ▶ Flavor dynamics in higher-dim. superpotential due to  $\mathbf{1}_8$  matter exchange

## Backup: Sample Implementation

only potentially dangerous term for proton decay:  $\mathbf{78}_1 \mathbf{78}_1 \mathbf{78}_1$ ,  
 because inserting (colorless) condensates into  $\mathbf{27}_3 \mathbf{78}_1 \overline{\mathbf{27}}_3$ ,  
 integrating out  $\mathbf{78}_1$   
 color-triplet leptoquark self-coupling  $XXX = 0$  (antisymmetry)

### 2. At $10^{15}$ GeV

Condensate in adjoint matter representation:  $\langle \mathbf{78}_1 \rangle = \langle W_R^{23} \rangle$   
 + higher-dimensional terms  $(\mathbf{27} \mathbf{78} \overline{\mathbf{27}})^2$



$\Rightarrow \nu_R$  Majorana mass

$\Rightarrow$  PS symmetry broken to SM

$\Rightarrow$  Leptoquark couplings possible for  $D, D^c$ , but no diquark couplings

# Backup: Sample Implementation

## 3. At $10^3$ GeV

Soft-breaking terms (hidden sector) induce radiative symmetry breaking  $\langle S \rangle$  via  $D/D^c$  loops

⇒  $\mu_D$ -term  $D^c \langle S \rangle D$  (Dirac masses)

⇒  $\mu_H$ -term  $H_u \langle S \rangle H_d$

⇒  $Z'$  mass if the extra  $U(1)$  broken by  $\langle S \rangle$  was gauged

... with flavor mixing

## 4. At $10^2$ GeV

Soft-breaking + effective  $\mu$ -term induce radiative symmetry breaking  $\langle H_u \rangle$  via  $t/t^c$  loops

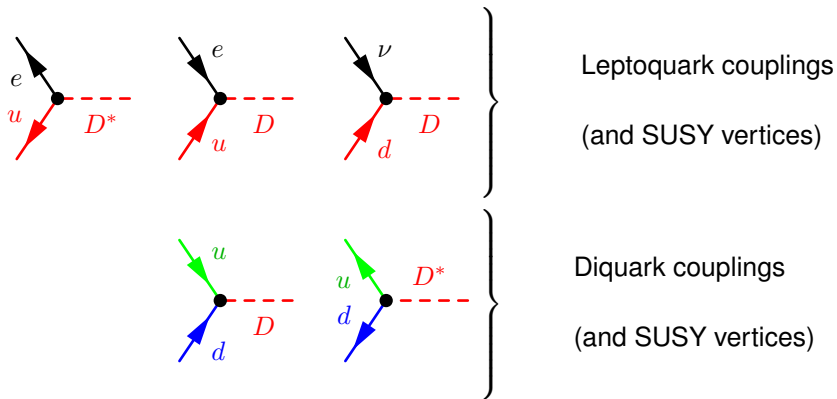
⇒  $\langle H_d \rangle$  due to Higgs superpotential + soft-breaking terms

⇒ Dirac masses for all charged MSSM matter

⇒ Majorana masses (see-saw) for  $\nu_L$

... again, with flavor mixing

# Backup: Interactions



Vector bosons induce e.g.  
 decay  $p \rightarrow e^+ \pi^0$   
 experimentally:  $\tau(p) > 10^{33}$  yrs

