

# Five-loop QED contributions for the anomalous magnetic moment of the muon

**Christian Sturm**

Brookhaven National Laboratory  
Physics Department  
High Energy Theory Group  
Upton, NY

- I. Introduction & Motivation
- II. Calculation
- III. Results & Discussion
- IV. Summary & Conclusion

In collaboration with: [P.A. Baikov](#) and [K.G. Chetyrkin](#)  
Nucl. Phys. **B183** (Proc.Suppl.) 2008

# Introduction & Motivation

Generalities; Anomalous magnetic moment of the muon:  $a_\mu$

- Magnetic moment of any system:
  - I.) Motion of el. charges
  - II.) Intrinsic mag. moments
- Dirac theory predicts for a lepton  $\ell = e, \mu, \tau$ :

$$\vec{\mu}_\ell = g_\ell \left( \frac{e}{2m_\ell} \right) \vec{S}, \quad g_\ell = 2$$

- **Quantum fluctuations:**  $\rightsquigarrow$  deviation from  $g_\ell = 2$ :

$$g_\ell = 2(1 + a_\ell) \quad \rightsquigarrow \text{precise test of QFT}$$

More formally:

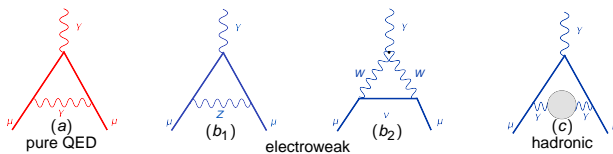
$$= \bar{u}(p') \left[ \gamma^\mu F_E(q^2) + i \frac{\sigma^{\mu\nu} q_\nu}{2m_\mu} F_M(q^2) \right] u(p)$$

- In the static limit ( $q^2 \rightarrow 0$ ):  $F_E(0) = 1$ ,  $F_M(0) = a_\mu$

# Introduction & Motivation

## Theory: Higher order corrections

- Higher order corrections are classified into 3 classes:

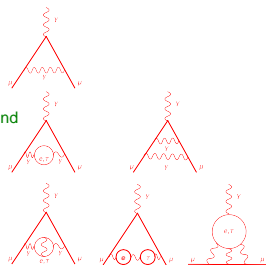


- The QED part is known to 4-loops (and leading terms in 5 loops!)  
(next slide)
- The EW part is known to 2-loops  
R. Jackiw, S. Weinberg; G. Altarelli et al.; I. Bars, M. Yoshimura; A. Czarnecki et al.
- The hadronic part is known but with limited accuracy  
Bouchiat, et al.; M. Gourdin, et al.; Brodsky, de Rafael; Hagiwara et al.; Alemany et al.; Davier et al.; Passera et al.  
Dominant theoretical uncertainties to the muon anomaly  
Effects at present not from first principles

# Introduction & Motivation

Theory: QED contributions to  $a_\mu$

$$\begin{aligned}
 a_\mu^{\text{QED}} &= \left(\frac{\alpha}{\pi}\right) 0.5 && \text{Schwinger} \\
 &+ \left(\frac{\alpha}{\pi}\right)^2 0.765857410(27) && \text{Sommerfield; Petermann; Suura \& Wichmann; Elend} \\
 &+ \left(\frac{\alpha}{\pi}\right)^3 24.05050964(87) && \text{Barbieri, Laporta, Remiddi,...; Kinoshita; Czarnecki, Skrzypek; Friot, Greynat, de Rafael} \\
 &+ \left(\frac{\alpha}{\pi}\right)^4 130.8055(80) && \text{Kinoshita, Lindquist; Kinoshita, Nio; Kinoshita, Nizic, Okamoto; Aoyama, Hayakawa, Kinoshita, Nio; Lautrup, de Rafael} \\
 &+ \left(\frac{\alpha}{\pi}\right)^5 663(20) \text{ In progress} && \text{Kinoshita et al.; Kataev; Laporta; Baikov et al.}
 \end{aligned}$$



growing coefficients

# Introduction & Motivation

## The value

The present experimental value is terrifically accurate!:

$$a_{\mu}^{\text{exp}} = 116592080(63) \cdot 10^{-11}$$

E821: Final Report: PRD73 (2006)

↪ Improvement of a factor of 14 compared to the classic CERN experiment

- QED four-loop contributes as much as  $380.8 \cdot 10^{-11}$   
(compared to the exp. uncertainty of  $\sim 60 - 70 \cdot 10^{-11}$ )
- <sup>1</sup>  $\Delta a_{\mu} = a_{\mu}^{\text{exp}} - a_{\mu}^{\text{theo.}} = 292(86) \cdot 10^{-11}$

The current theory prediction shows an

"interesting but not yet conclusive discrepancy" of  $\sim 3\sigma$

(or  $\sim 1\sigma$  if one uses  $\tau$ -data to describe the hadronic effects)

---

<sup>1</sup> A. Höcker, W. Marciano, PDG

# Calculation

## logarithmically enhanced contributions

There are 2 **sources** of numerically leading enhanced logs of large ratio  $\frac{M_\mu}{m_e} = 206.7682838$  pure QED contributions:

$$\begin{array}{cc}
 \begin{array}{c} \text{LBL diagram: a circle with a photon loop and external photon lines} \\ \sim \left(\frac{\alpha}{\pi}\right)^3 \left(\frac{2\pi^2}{3} \ln\left(\frac{M_\mu}{m_e}\right) + \dots\right) \\ \text{LBL: light by light scattering} \end{array} &
 \begin{array}{c} \text{VP diagram: a triangle with a photon loop and external photon lines} \\ \sim \left(\frac{\alpha}{\pi}\right)^2 \left(\frac{1}{3} \ln\left(\frac{M_\mu}{m_e}\right) + \dots\right) \\ \text{VP: vacuum polarization} \end{array}
 \end{array}$$

→ We will consider mixed VP contributions: photon propagator composed from electron loops and photon exchanges only!

$$a_\mu = \frac{\alpha}{\pi} \int_0^1 dx (1-x) \left[ d_R \left( \frac{-x^2 M_\mu}{1-x m_e}, \alpha \right) - 1 \right] \quad \text{B. Lautrup, de Rafael '74}$$

with  $d_R(q^2/m_e^2, \alpha) = 1/(1 + \alpha \Pi^{\text{OS}}(q^2/m_e^2, \alpha^{\text{OS}}))$   
 $\Pi^{\text{OS}}$  proper photon VP in OS-scheme

# Calculation

## Vacuum polarization function

### Required ingredients for the calculation:

- Vacuum polarization function

- High-energy limit  $\rightarrow$  massless propagators

The value of  $\overline{\Pi}(Q^2, m = 0, \overline{\alpha})$  is known at 4-loops with **Baicer** (P. Baikov, 2000-2008) in  $\overline{\text{MS}}$ -scheme

P.A. Baikov, K.G. Chetyrkin, J.H. Kühn

Traditionally in calculations of  $a_\ell$  everybody uses the classical OS-scheme:  $\alpha$  and all lepton masses are on-shell and:

$$\Pi^{\text{OS}}(Q = 0, m, \alpha^{\text{OS}}) = 0$$

$\rightsquigarrow$  Transform massless propagator from  $\overline{\text{MS}}$   $\rightarrow$  on-shell scheme

- Need:  $\overline{\text{MS}} \leftrightarrow$  On-shell relation at 4-loop

for fine structure constant conversion:  $\overline{\alpha} \rightarrow \alpha^{\text{OS}}$

# Calculation

On-shell  $\leftrightarrow$   $\overline{\text{MS}}$ -scheme for  $\alpha$

In QED MS in OS-schemes are related through:

$$\frac{\alpha^{\text{OS}}}{1 + \Pi^{\text{OS}}(\mathbf{Q}, m, \alpha^{\text{OS}})} = \frac{\bar{\alpha}}{1 + \overline{\Pi}(\mathbf{Q}, \bar{m}, \bar{\alpha})}$$

scheme invariant  
concept of the  
invariant charge

Required:  $\overline{\Pi}(\mathbf{Q} = 0, \bar{m}, \bar{\alpha})$

$$\rightsquigarrow \bar{\alpha} = \alpha^{\text{OS}} \left( 1 + \sum_{i \geq 1} \mathbf{C}_{\bar{\alpha}\alpha}^{(i)} \left( \frac{\alpha^{\text{OS}}}{\pi} \right)^i \right)$$

- IBP, 11 master integrals
- $\overline{\text{MS}} \leftrightarrow$  On-shell relation at 3-loop for mass conversion  
Steinhauser, Chetyrkin; Melnikov, van Ritbergen



# Results & Discussion

$\overline{\alpha}$ -OS-relation for conversion of fine structure constant

$$\overline{\alpha} = \alpha^{\text{OS}} \left( 1 + \sum_{i \geq 1} c_{\overline{\alpha}\alpha}^{(i)} \left( \frac{\alpha^{\text{OS}}}{\pi} \right)^i \right)$$

$$c_{\overline{\alpha}\alpha}^{(4)} = \frac{14327767}{9331200} + \frac{8791}{3240} \pi^2 + \frac{204631}{259200} \pi^4 - \frac{175949}{4800} \zeta_3 + \frac{1}{24} \pi^2 \zeta_3 + \frac{9887}{480} \zeta_5 - \frac{595}{108} \pi^2 \ln 2$$

$$- \frac{106}{675} \pi^4 \ln 2 + \frac{6121}{2160} \pi^2 \ln^2 2 - \frac{32}{135} \pi^2 \ln^3 2 - \frac{6121}{2160} \ln^4 2 + \frac{32}{225} \ln^5 2 - \frac{6121}{90} a_4 - \frac{256}{15} a_5$$

$$+ \ell_{\mu m} \left[ -\frac{383}{31104} + \frac{23}{108} \pi^2 - \frac{41}{144} \zeta_3 - \frac{2}{9} \pi^2 \ln 2 \right] + \frac{43}{144} \ell_{\mu m}^2 + \frac{13}{108} \ell_{\mu m}^3 + \frac{1}{81} \ell_{\mu m}^4,$$

$$\ell_{\mu m} = \ln \frac{\mu}{m}, \quad a_n = \text{Li}_n \left( \frac{1}{2} \right)$$

**Important:**  $\Pi(q^2/m^2, \alpha) = \Pi^\infty(q^2/m^2, \alpha) + \mathcal{O}(m^2/q^2)$

then the resulting error in  $a_\mu$  will be:

$$a_\mu = \frac{\alpha}{\pi} \int_0^1 dx (1-x) \left[ d_R^\infty \left( \frac{-x^2 M_\mu}{1-x m_e}, \alpha \right) - 1 \right] + \mathcal{O} \left( \frac{m_e}{M_\mu} \right)$$

of order  $m_e/M_\mu$  with  $d_R^\infty = 1/(1 + \alpha \Pi^\infty)$

# Results & Discussion

## Analytical 5-loop contributions to $a_\mu$

The resulting contributions to  $a_\mu$  coming from 4-loop terms in the photon propagator read :

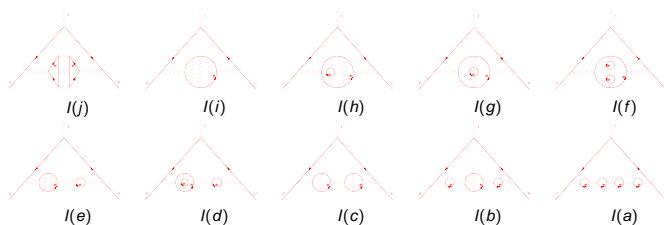
$$a_\mu^{\text{asyp.}} = \sum_{i \geq 2} a_\mu^{\text{asyp.},(i)} \left(\frac{\alpha}{\pi}\right)^i$$

$$\begin{aligned} a_\mu^{\text{asyp.},(5)} = & -\frac{296496193}{41990400} + \frac{45709}{58320} \pi^2 + \frac{212701}{518400} \pi^4 - \frac{4488523}{259200} \zeta_3 + \frac{35}{144} \pi^2 \zeta_3 + \frac{4}{3} \zeta_3^2 + \frac{10909}{720} \zeta_5 \\ & + \frac{35}{8} \zeta_7 - \frac{55}{24} \pi^2 \ln 2 - \frac{53}{675} \pi^4 \ln 2 + \frac{6121}{4320} \pi^2 \ln^2 2 - \frac{16}{135} \pi^2 \ln^3 2 - \frac{6121}{4320} \ln^4 2 \\ & + \frac{16}{225} \ln^5 2 - \frac{6121}{180} a_4 - \frac{128}{15} a_5 + \ell_{\mu e} \left[ \frac{1416095}{279936} + \frac{41}{972} \pi^2 - \frac{1855}{432} \zeta_3 - \frac{10}{3} \zeta_5 - \frac{2}{9} \pi^2 \ln 2 \right] \\ & + \ell_{\mu e}^2 \left[ -\frac{1507}{1944} + \frac{8}{81} \pi^2 + \frac{4}{3} \zeta_3 \right] - \frac{83}{243} \ell_{\mu e}^3 + \frac{8}{81} \ell_{\mu e}^4 + \mathcal{O}\left(\frac{m_e}{M_\mu}\right), \quad \ell_{\mu e} = \ln \frac{M_\mu}{m_e}, \quad a_n = \text{Li}_n\left(\frac{1}{2}\right) \end{aligned}$$

Numerically:  $\left(\frac{\alpha}{\pi}\right)^5 a_\mu^{\text{asyp.},(5)} = \left(\frac{\alpha}{\pi}\right)^5 62.2667 = 0.42105 \cdot 10^{-11}$   
 (compared to  $\left(\frac{\alpha}{\pi}\right)^5 663(20)$ )

# Comparison

Analytical  $\longleftrightarrow$  Numerical results



Subset	analytical	numerical	
$I(j)$	-1.21429	-1.24726(12)	✓
$I(i)$	+0.25237	—	
$I(g) + I(h)$	+1.50112	+1.56070(64)	✓
$I(f)$	+2.89019	+2.88598(9)	✓
$I(c)$	+4.81759	+4.74212(14)	✓
$I(d)$	+7.44918	+7.45270(88)	✓
$I(e)$	-1.33141	-1.20841(70)	✓
$I(b)$	+27.7188	+27.69038(30)	✓
$I(a)$	+20.1832	+20.14293(23)	✓

-Numerics from: T. Aoyama, M. Hayakawa, T. Kinoshita, M. Nio, N. Watanabe (2008) M. Nio, T. Aoyama, M. Hayakawa, T. Kinoshita (2007,2008) and T. Kinoshita, M. Nio (2006)

-Remaining differences should come from power suppressed corrections to the asymptotic result of  $\mathcal{O}(m_e/M_\mu)$

-Agreement with Kataev, where available

# Summary & Conclusion

- Analytical and numerical methods help in computing the QED contribution to  $(g - 2)_\mu$
- The conversions formula for  $\alpha^{OS}/\alpha^{\overline{MS}}$  is evaluated to four-loops  $\Rightarrow$  one could reexpress any (QED!)  $\mathcal{O}(\alpha^5)$  result in terms of the running  $\alpha^{\overline{MS}}$  or vice versa
- Asymptotic contributions to the VP part of  $(g - 2)_\mu$  in order  $\alpha^5$  are computed and support the numerical result of the group of Kinoshita