

A Pyramid Scheme for Particle Physics

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Outline

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Cosmological SUSY Breaking

- Cosmological constant an input parameter [Banks](#)
- LEFT $\Lambda = 0$ limit
 - Super-Poincaré symmetry
 - Discrete R -symmetry $\Rightarrow \Lambda = 0$
- LEFT $\Lambda \neq 0$ limit $\Rightarrow R$ -breaking operators
 - UV/IR mixing effect
 - Metastable SUSY violating state with $m_{3/2} = K\Lambda^{1/4}$,
 $K = \mathcal{O}(10)$
 - Constant term in superpotential such that c.c. = Λ
 - Tunneling probability of order $\mathcal{O}(e^{-\pi(RM_P)^2})$

\Rightarrow Constraints on LEFT

Metastable DSB

- Metastable dynamical SUSY breaking
 - DSB \Rightarrow Natural hierarchy of scales $M_{SUSY} \ll M_P$ [Witten](#)
 - More generic than DSB in stable states [Intriligator, Seiberg, Shih](#)
- Direct gauge mediation and MSSM $\Rightarrow G \times SU(1, 2, 3)$
 - No messenger sector
 - Solution of SUSY flavor problem
 - One loop gauge coupling unification
- SM gauge couplings in perturbative regime

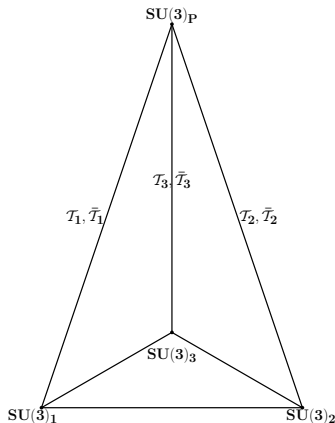
Trinification and Pyramid scheme

● Trinification Glashow

- $SU(3)^3 \times Z_3$
- Gauge bosons in $(8, 1, 1) \oplus \dots$
- 2 Higgs and 3 families in
- $(3, \bar{3}, 1) \oplus (\bar{3}, 1, 3) \oplus (1, 3, \bar{3})$

● Pyramid scheme

- Extra $SU(3)_P$
- Extra matter (trianons)
- $T_1 + \bar{T}_1 = (3, 1, 1, \bar{3}) \oplus (\bar{3}, 1, 1, 3)$
- $T_2 + \bar{T}_2 = (1, 3, 1, \bar{3}) \oplus (1, \bar{3}, 1, 3)$
- $T_3 + \bar{T}_3 = (1, 1, 3, \bar{3}) \oplus (1, 1, \bar{3}, 3)$
- One singlet $S \Rightarrow \mu$ -problem



$$\text{GUT } SU(3)_P \times SU(3)^3 \times Z_3 \rightarrow \text{LEFT } SU(3)_P \times SU(1, 2, 3)$$

Discrete R-symmetry

- Non-anomalous
- Yukawa terms, Higgs terms and trianon terms allowed
- Dimension 4 and 5 B and L violating terms forbidden (apart from neutrino seesaw operator)

$$W = \sum_{i=1}^3 (m_i + y_i S) T_i \bar{T}_i + \text{GUT terms} + g_\mu S H_u H_d \\ + \lambda_u H_u Q \bar{U} + \lambda_d H_d Q \bar{D} + \lambda_L H_d L \bar{E} + \frac{\lambda_\nu}{M} (L H_u)^2 + W_0$$

$SU(3)_P$ with $N_F = 9$ flavors \Rightarrow IR free

- Without ISS mass terms
 - Free theory with SUSY
 - Forbidden by CSB \Rightarrow Dynamical metastable SUSY violating state with $m_{3/2} = K\Lambda^{1/4}$
- With ISS mass terms
 - Two heavy masses \Rightarrow Assumed SUSY and R -symmetry breaking metastable state

Phenomenology suggests heavy $m_{1,3} > \Lambda_3$ with $m_2 \approx \Lambda_3$

Mass hierarchy $\Lambda_3 \approx m_2 \lesssim m_3 \ll m_1$

Spectrum

$$W = X(\det M/\Lambda_3 - P\bar{P} - \Lambda_3^2) + \text{trianon terms} + \dots$$

- Two kinds of messengers

- $N_F = N_C = 3$ theory moduli

- $M = Z_a \lambda^a$, $P = i\Lambda_3 e^{(q+p)/\Lambda_3}$ and $\bar{P} = i\Lambda_3 e^{(q-p)/\Lambda_3}$

- Heavy trianons

- SUSY breaking vacuum $\Rightarrow \langle P \rangle = \langle \bar{P} \rangle \neq 0$

- R -symmetry breaking vacuum $\Rightarrow \langle M \rangle \propto \text{Id}$

$$\Rightarrow V_{\text{non-SUSY}} \approx K_{M^\dagger M}^{-1} |m_2 \Lambda_3|^2$$

M messengers

$$m_{3/2} = X_g \frac{m_2 \Lambda_3}{m_P} = K \Lambda^{1/4}$$

$$m_{1/2}^i = 3X_i \frac{\alpha_i}{4\pi} m_2$$

- M field quartic scalar couplings of order $\mathcal{O}(m_2^2/\Lambda_3^2) \Rightarrow m_2 < \sqrt{4\pi} \Lambda_3$
- Chargino mass bound of 160 GeV $\Rightarrow X_2 \frac{m_2}{\text{TeV}} > 19.7$

\Rightarrow Example:

- Assumption: $m_2 = 1.7\Lambda_3$ such that $\frac{m_2^2}{4\pi\Lambda_3^2} \approx \frac{1}{4}$
- $X_2 > 4.2\sqrt{X_g/K} \Rightarrow \Lambda_3 = 5.1 \text{ TeV}$ and $m_2 = 8.6 \text{ TeV}$

$\mathcal{T}_{1,3}$ messengers

- $m_3 \gg \Lambda_3$

- Gluino effective couplings $\int d^2\theta (W_\alpha^3)^2 f(M/m_3, P/m_3, \bar{P}/m_3)$

- $m_3 \gtrsim \Lambda_3$ due to gluino constraint

- CW approximation breaks down
- Chiral perturbation theory not suitable

⇒ Gluino/chargino and squark/slepton mass ratios suppressed compared to usual gauge mediation

- No large contributions to Higgs potential ⇒ Good for little hierarchy problem

Cosmology

- Gauge mediation \Rightarrow LSP gravitino thus no WIMP dark matter candidate
- Hidden baryon-like states as dark matter [Banks, Mason, O'Neil](#)
- Pyramid scheme \Rightarrow 3 accidental baryon number-like symmetries
 - 2 unbroken $U_{B_{1,3}} \Rightarrow$ Observed DM density without asymmetry through non-thermal production with $T_{\text{reh}} < \Lambda_3$
 - 1 spontaneously broken U_{B_2}

Assumption: Negligible primordial asymmetries with low T_{reh}

- U_{B_3} -charged particles as DM
 - QCD-like interactions with dynamical scale Λ_3
 - Energy-independent annihilation cross section of order $\mathcal{O}(\Lambda_3^{-2})$
 - Annihilation to PNGB of spontaneously broken U_{B_2} with high multiplicity
 - PNGB of spontaneously broken U_{B_2} (pymion)
 - Leading U_{B_2} breaking operator $\int d^2\theta S(\det T_2)/M_{\text{GUT}} \Rightarrow$ Light (MeV range)
 - Stellar cooling rates bound satisfied
 - Colorless constituents \Rightarrow Decay to e^+e^- , photons and neutrinos from operators like $\alpha_2^2 \partial_\mu p J^\mu / \Lambda_3 \sim \alpha_2^2 m_e p e^+ e^- / \Lambda_3$
- \Rightarrow Positron excess with DM annihilation cross section $\sigma_0 = A/\Lambda_3^2$

Features and problems

Features

- Based on trinification \Rightarrow No Landau poles
- Heavy trianons needed for metastable SUSY violating state \Rightarrow colored sparticle mass suppression
- 2 unbroken baryon number-like symmetries \Rightarrow Non-thermal DM candidate compatible with experiments
- Pymion mass \Rightarrow Not produced in ordinary stars and positron excess compatible with experiments

Problems

- Existence of metastable SUSY breaking state \Rightarrow Non-zero meson VEV