

Electroweak and QCD corrections to W+jet production at hadron colliders

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1 Introduction

Importance of W+jet production: $pp \rightarrow l\nu_l + \text{jet} + X$

- large SM cross section ($\sim 1\text{nb}$ after basic cuts) → standard candle
- single-W production often shows additional jet activity
 ↪ relevance for W-mass determination at the LHC
- offers precision test for jet dynamics in QCD
- dominant channel for high- p_T leptons
 ↪ important background for various new-physics searches

Theoretical status:

- NLO QCD DYRAD [Giele et al. '93]; MCFM [Campbell/R.K.Ellis '02];
 Melnikov/Petriello '06; Catani et al. '09
- NLO EW for stable W bosons
 Kühn, Kulesza, Pozzorini, Schulze '07; Hollik, Kasprzik, Kniehl '07

Motivation for this work:

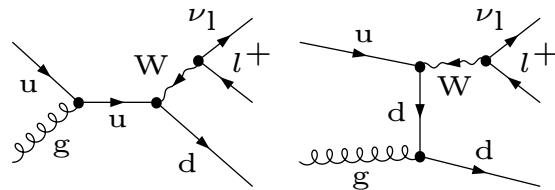
Inclusion of full off-shell effects in NLO QCD+EW predictions

↪ e.g. essential for W-mass determination

2 Calculation of NLO corrections

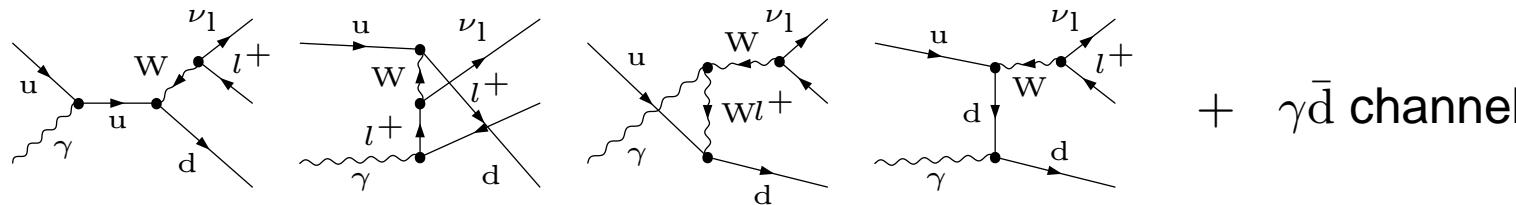
2.1 Lowest-order prediction

LO diagrams:



+ crossed graphs for $u\bar{d}$ and $g\bar{d}$ channels

Contributions from photon-induced processes



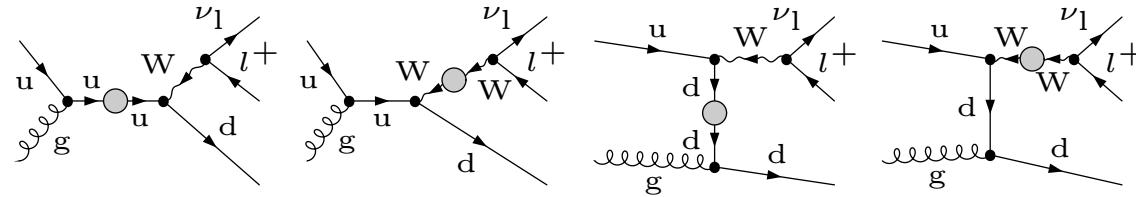
Features of the LO cross section:

- IR safety requires at least lower cut on $p_{T,\text{jet}}$
↪ apply jet algorithm for NLO cross section before cut on $p_{T,\text{jet}}$
- contributions from photon-induced processes generically small ($\sim 1\%$)
↪ inclusion of LO and NLO QCD as corrections

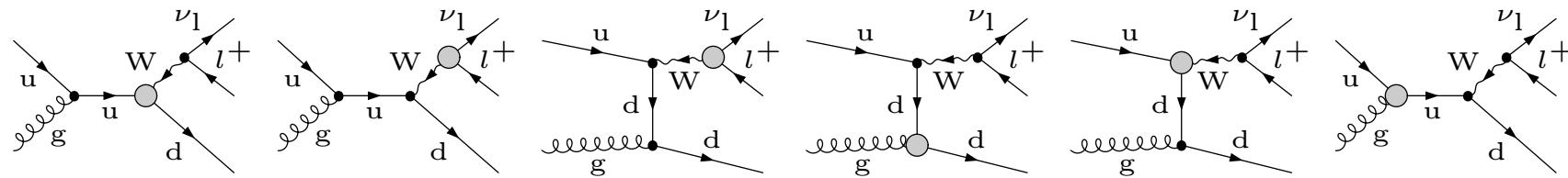
2.2 Survey of NLO EW+QCD corrections and technical details

1PI loop insertions in virtual corrections: $\mathcal{O}(100)$ EW loop diagrams per channel

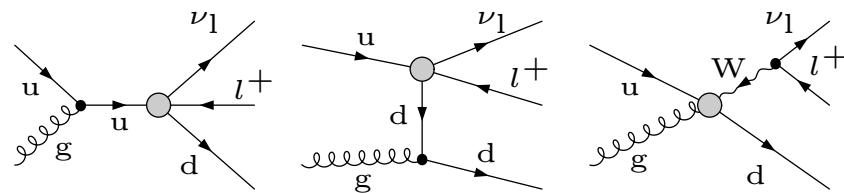
Self-energy insertions:



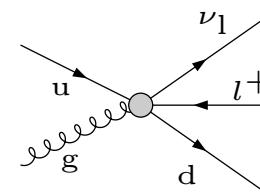
Vertex corrections:



Box corrections:

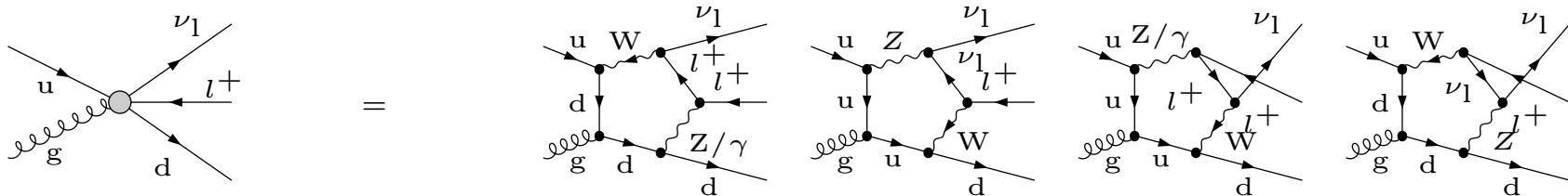


Pentagon graphs:



Most complicated parts of the loop calculation:

- pentagon graphs:



↪ stable reduction to box integrals without inverse Gram determinants

Denner, S.D. '02,'05 (similar to Binoth et al. '05)

- numerical instabilities in Passarino–Veltman reduction of tensor integrals
 - ↪ expansion about exceptional points

Denner, S.D. '05 (similar to Giele et al. '04; R.K.Ellis et al. '05)

- gauge-invariant treatment of **W resonances**

↪ “complex-mass scheme” Denner, S.D., Roth, Wieders '05

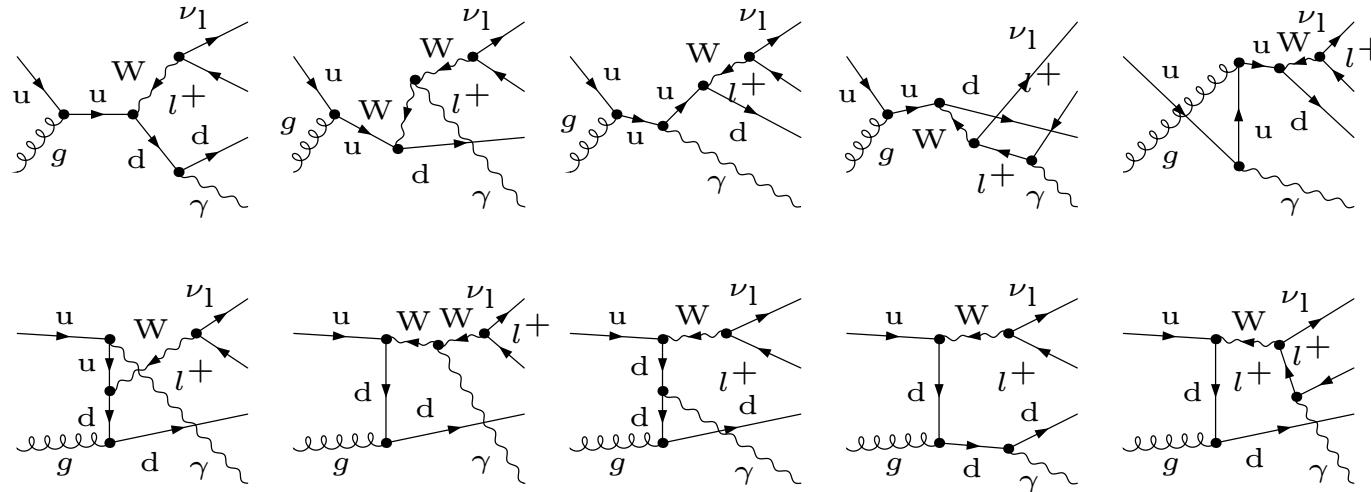
Real emission corrections

- **NLO QCD:**

- ◊ gluon bremsstrahlung: $u\bar{d} \rightarrow l^+ \nu_l gg$, $ug \rightarrow l^+ \nu_l dg$, $\bar{d}g \rightarrow l^+ \nu_l \bar{u}g$
- ◊ gg fusion: $gg \rightarrow l^+ \nu_l \bar{u}d$
- ◊ gluon splitting $g^* \rightarrow q\bar{q}$: $u\bar{d} \rightarrow l^+ \nu_l q\bar{q}$ + crossed variants

- **NLO EW:**

- ◊ photon bremsstrahlung: $u\bar{d} \rightarrow l^+ \nu_l g\gamma$, $ug \rightarrow l^+ \nu_l d\gamma$, $\bar{d}g \rightarrow l^+ \nu_l \bar{u}\gamma$



Note: $u\bar{d} \rightarrow l^+ \nu_l g\gamma$ contributes to both

- NLO **EW** corrections to $u\bar{d} \rightarrow l^+ \nu_l g$ and
- NLO **QCD** corrections to $u\bar{d} \rightarrow l^+ \nu_l \gamma$

Further technical details

- **Generic features**
 - ◊ two completely independent calculations of all ingredients
 - ↪ results in mutual agreement
 - ◊ dipole subtraction for QCD and photonic IR (soft and collinear) singularities
Catani, Seymour '96; S.D. '99; S.D., Kabelschat, Kasprzik '08
 - ↪ checked against phase-space slicing
 - **MPI/FR**
 - ◊ loop amplitude generation with **FEYNARTS 1** Böhm, Denner, Küblbeck '90
 - ↪ algebraic reduction with inhouse **MATHEMATICA** routines
 - ◊ tree amplitudes evaluated analytically via spinor formalism
 - ◊ phase-space integration via **VEGAS**
 - **PSI**
 - ◊ loop amplitude generation with **FEYNARTS 3** Hahn '00
 - ↪ algebraic reduction with **FORMCALC** Hahn '98-'09 and **POLE** Meier '05; Mück
 - ◊ real corrections generated with program **POLE**
 - ↪ automatic generation of tree amplitudes via spinor formalism, dipole subtraction terms, and multi-channel phase-space integration

3 Numerical results for the LHC

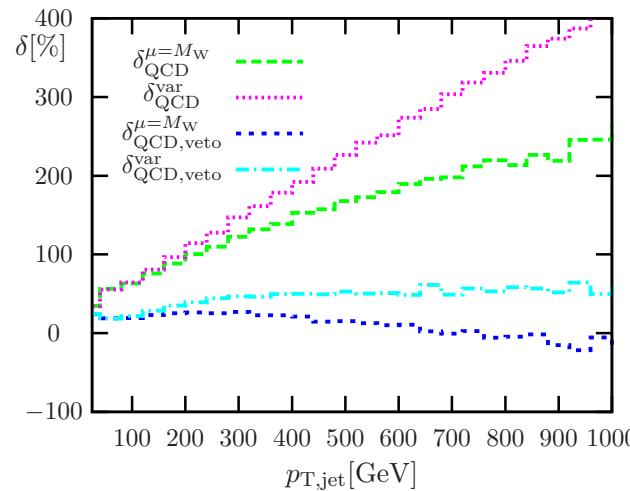
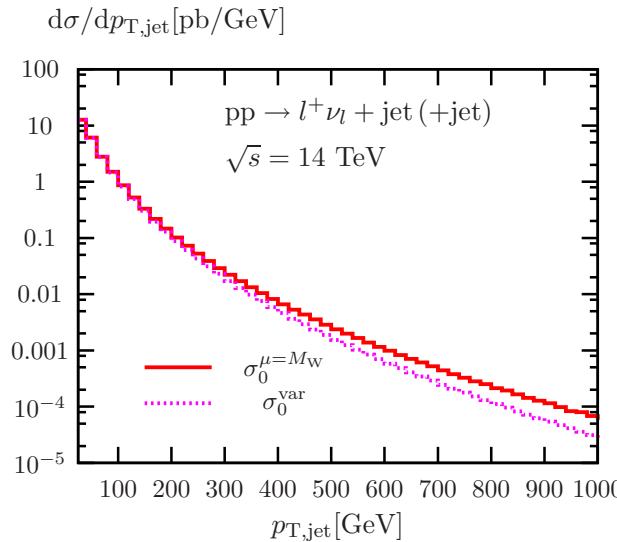
Setup and definition of observables

- PDF set MRST2004QED which includes NLO QCD+EW corrections
 - ↪ initial-state collinear QCD and photonic singularities removed via factorization
 - ↪ PDF set includes photon density
- two scale choices: $\mu = \mu_{\text{ren}} = \mu_{\text{fact}} = M_W$ (fixed) or $\mu = \sqrt{M_W^2 + (p_T^{\text{had}})^2}$ (variable)
- k_T -algorithm for jet definition
- basic cuts: $p_{T, \text{jet1}} / l / \text{miss} > 25 \text{ GeV}$, $|y_{\text{jet1}} / l| < 2.5$, $R_{l, \text{jet}} > 0.5$
- jet veto optionally applied to 2nd hard jet if $p_{T, \text{jet2}} > p_{T, \text{jet1}} / 2$
- two lepton identifications: “bare leptons” ($l = \mu^+$)
 - ↪ large corrections $\propto \alpha \ln(m_l^2 / Q^2)$or photon–lepton recombination if $R_{\gamma, l} < 0.1$ ($l = e^+$)
 - ↪ cancellation of all $\alpha \ln(m_l^2 / Q^2)$ terms a la KLN
- photon fragmentation function for photon–jet separation Glover, Morgan '94



Transverse-momentum distribution of the hardest jet

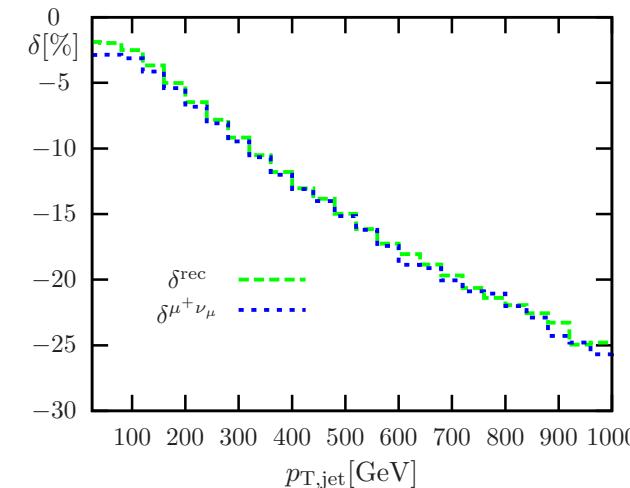
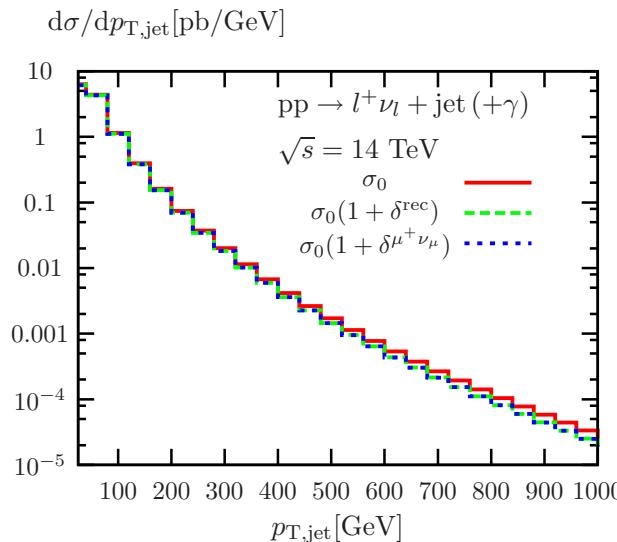
QCD corrections:



Large positive corrections
from W+2jets
(mainly back-to-back jets)

→ significant reduction of
corrections via jet veto

EW corrections:



Large neg. corrections
due to EW Sudakov logs

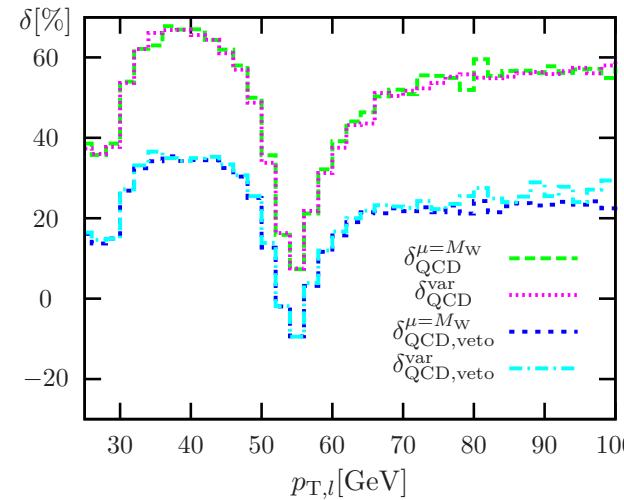
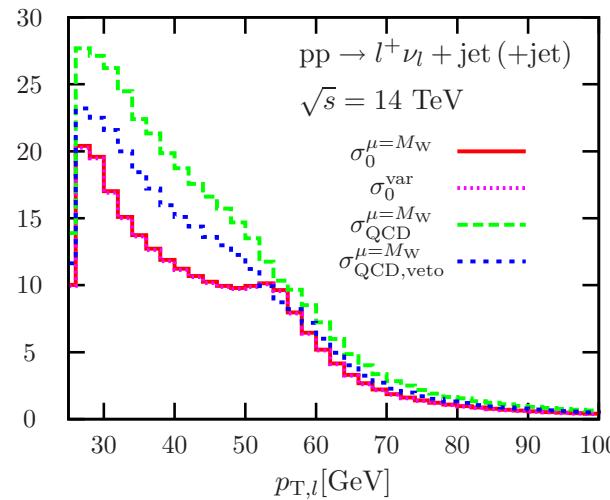
→ qualitative agreement
with previous results for
on-shell Ws [Kühn et al. '07](#)
[Hollik et al. '07](#)



Transverse-momentum distribution of the charged lepton

QCD corrections:

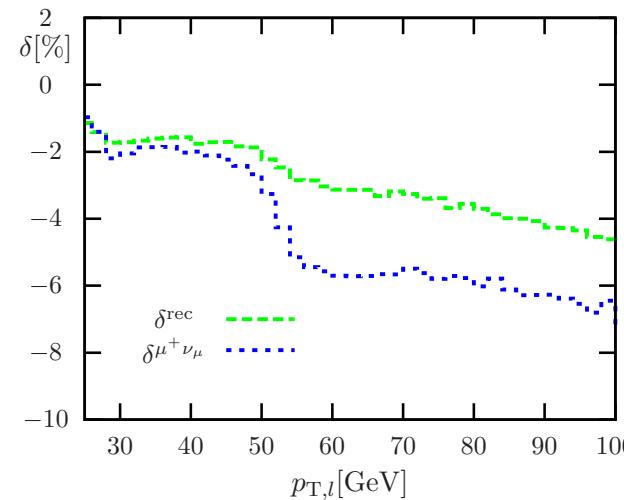
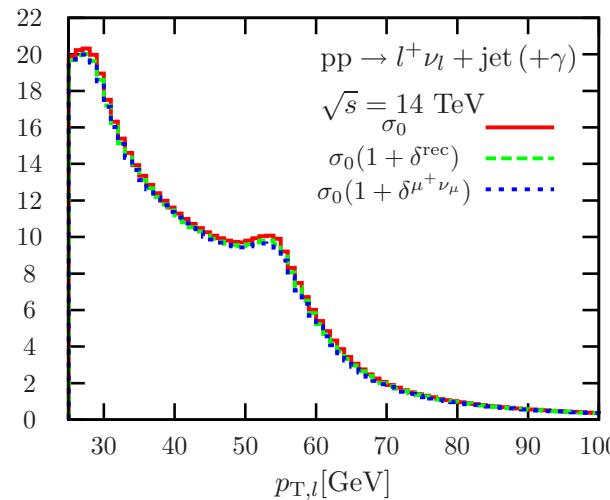
$d\sigma/dp_{T,l} [\text{pb}/\text{GeV}]$



- large corrections distorting the shape
- corrections reduced by jet veto

EW corrections:

$d\sigma/dp_{T,l} [\text{pb}/\text{GeV}]$

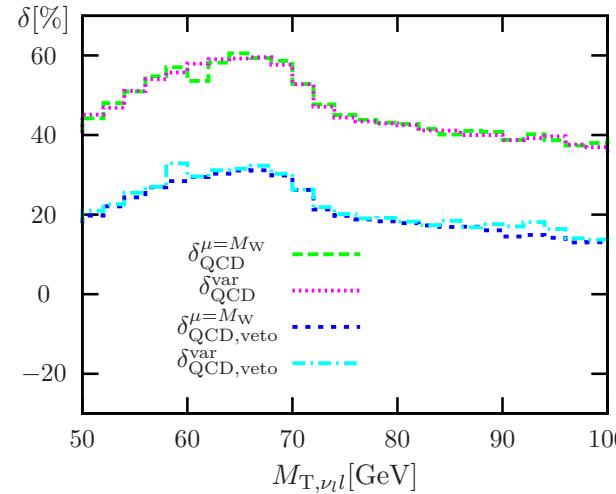
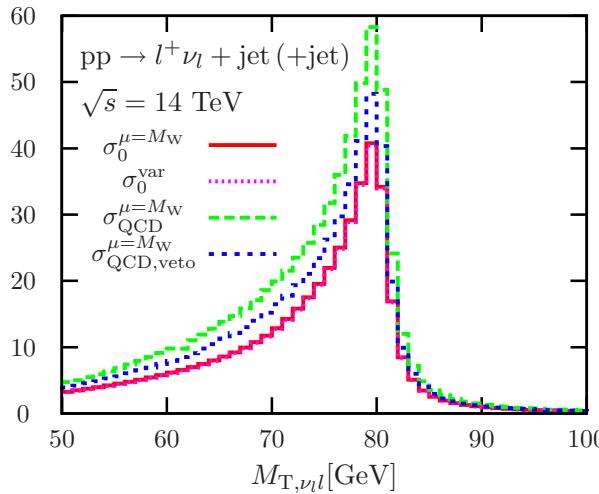


- moderate corrections of $\mathcal{O}(5\%)$
- Jacobian peaks slightly distorted

Transverse-mass distribution of the W boson

QCD corrections:

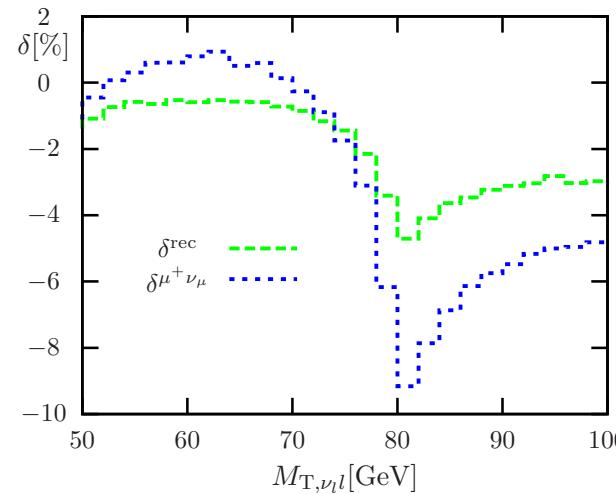
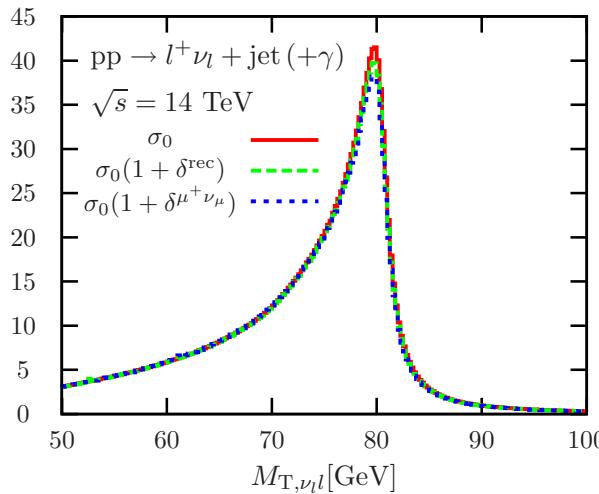
$d\sigma/dM_{T,\nu l l} [\text{pb}/\text{GeV}]$



- corrections smooth near Jacobian peak
- corrections reduced by jet veto

EW corrections:

$d\sigma/dM_{T,\nu l l} [\text{pb}/\text{GeV}]$



- corrections of $\mathcal{O}(5\text{--}10\%)$
- Jacobian peak significantly distorted



5 Conclusions

$W+jet$ production is a very important process at Tevatron and LHC.
(standard candle, W mass, background, etc.)

Our calculation provides

- recalculation of NLO QCD corrections
- first calculation of **NLO EW corrections including full off-shell effects**
(e.g. relevant for W -mass determination)
↪ building block for NNLO corrections of $\mathcal{O}(\alpha\alpha_s)$ for single- W production

Size of corrections:

- **EW corrections particularly large at high p_T** of leptons and jets
- photon-induced processes and EW–QCD interferences
phenomenologically unimportant

Outlook:

- low- p_T range should be further improved by **soft-gluon resummation**
or by merging with parton showers
- off-shell **Z+jet production** straightforward with the same approach



Backup slides



Photon–jet separation via photon fragmentation function $D_{q \rightarrow \gamma}$

Why?

- collinear quarks and photons have to be recombined → quasiparticle otherwise corrections $\propto \ln(m_q^2/Q^2)$ → perturbative “IR instability”
 - quark and gluon jets cannot be distinguished event by event
↪ common recombination required for quarks/gluons with photons
- ⇒ $\underbrace{(\text{g}_{\text{hard}} + \gamma_{\text{soft}})}_{\text{EW corr. to } W+\text{jet}}$ and $\underbrace{(\text{g}_{\text{soft}} + \gamma_{\text{hard}})}_{\text{QCD corr. to } W+\gamma}$ both appear as 1 jet

Solution:

- exclude events with photon energy fraction $z_\gamma = \frac{E_\gamma}{E_{\text{jet}} + E_\gamma} > z_0$ for (jet + γ) quasiparticles (chosen value $z_0 = 0.7$)
 - subtract convolution of LO cross section with
- $$D_{q \rightarrow \gamma}^{\overline{\text{MS}}} (z_\gamma, \mu_{\text{fact}}) \Big|_{\text{mass. reg.}} = \frac{\alpha Q_q^2}{2\pi} P_{q \rightarrow \gamma}(z_\gamma) \left[\ln \frac{m_q^2}{\mu_{\text{fact}}^2} + 2 \ln z_\gamma + 1 \right] \quad \leftarrow \text{cancels coll. singularities}$$
- $$+ D_{q \rightarrow \gamma}^{\text{ALEPH}}(z_\gamma, \mu_{\text{fact}}) \quad \leftarrow \text{non-perturbative part fitted to ALEPH data}$$
- where $P_{q \rightarrow \gamma}(z_\gamma) = \frac{1+(1-z_\gamma)^2}{z_\gamma} = \text{quark-to-photon splitting function}$



Integrated cross section for various cuts on $p_{\text{T}, \text{jet}}$

$\text{pp} \rightarrow l^+ \nu_l \text{ jet} + X$ at $\sqrt{s} = 14 \text{ TeV}$

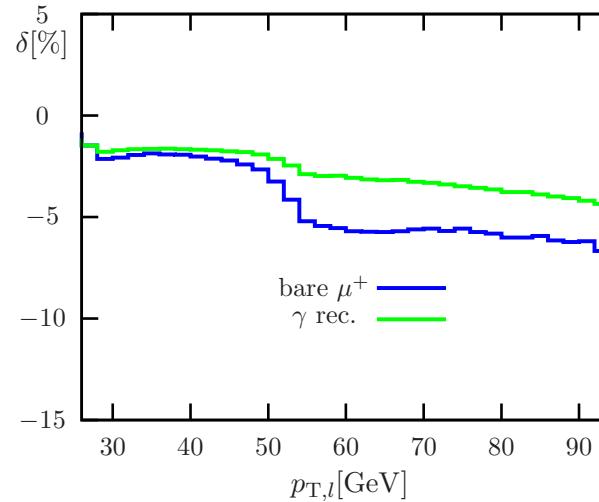
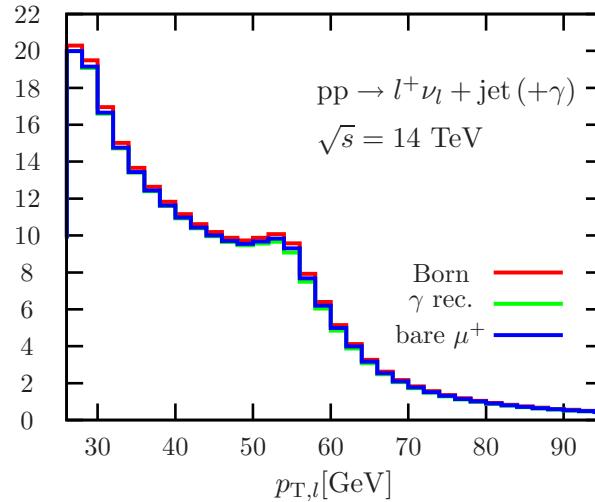
$p_{\text{T}, \text{jet}} / \text{GeV}$	$25 - \infty$	$50 - \infty$	$100 - \infty$	$200 - \infty$	$500 - \infty$	$1000 - \infty$	
$\sigma_{\text{Born}}^{\mu=M_W} / \text{pb}$	509.45(5)	182.74(2)	49.777(5)	8.1086(9)	0.3156(1)	0.0117(1)	
$\sigma_{\text{Born}}^{\text{var}} / \text{pb}$	502.66(5)	176.39(1)	45.382(4)	6.4990(6)	0.1850(1)	0.0048(1)	
$\delta_{\text{EW}}^{\mu^+ \nu_\mu, \text{var}} / \%$	-3.07(6)	-3.35(1)	-4.64(1)	-8.50(1)	-18.0(1)	-28.0(1)	EW Sudakov
$\delta_{\text{EW}}^{\text{rec, var}} / \%$	-2.07(1)	-2.55(2)	-4.18(1)	-8.37(3)	-17.8(1)	-27.9(1)	logs
$\delta_{\text{QCD}}^{\mu=M_W} / \%$	48.0(1)	64.8(1)	80.7(1)	115	188	270(1)	
$\delta_{\text{QCD}}^{\text{var}} / \%$	47.9(2)	65.4(1)	85.8(1)	135	270	494(1)	
$\delta_{\text{QCD, veto}}^{\mu=M_W} / \%$	21.5(1)	18.2(1)	22.5(2)	24.6(1)	5.3(1)	-26.4(2)	jet veto
$\delta_{\text{QCD, veto}}^{\text{var}} / \%$	22.3(1)	20.8(1)	29.8(2)	43.3(2)	52.4(1)	58.7(1)	
$\delta_{\gamma, \text{NLO}}^{\text{var}} / \%$	0.38	0.70	1.18	1.86	3.26	5.18(1)	γ -induced processes
$\delta_{\gamma, \text{NLO, veto}}^{\text{var}} / \%$	0.35	0.64	1.10	1.76	3.03	4.73(1)	
$\delta_{\text{IF}}^{\text{var}} / \%$	0.05	0.13	0.51	1.88	11.50	49.95	QCD-EW interferences
$\delta_{\text{IF, veto}}^{\text{var}} / \%$	0.01(1)	0.03	0.12	0.40	1.63	4.72	



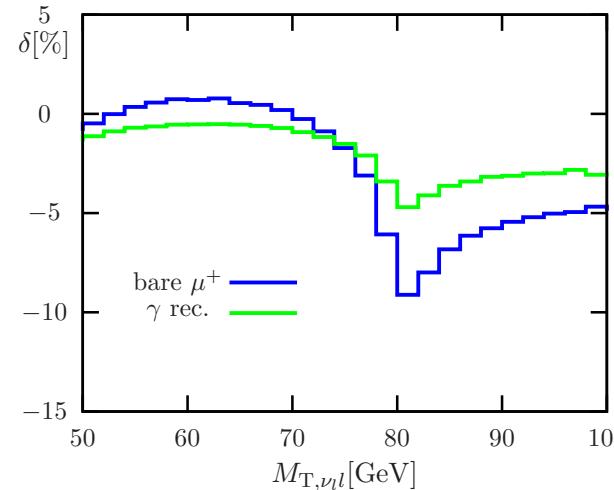
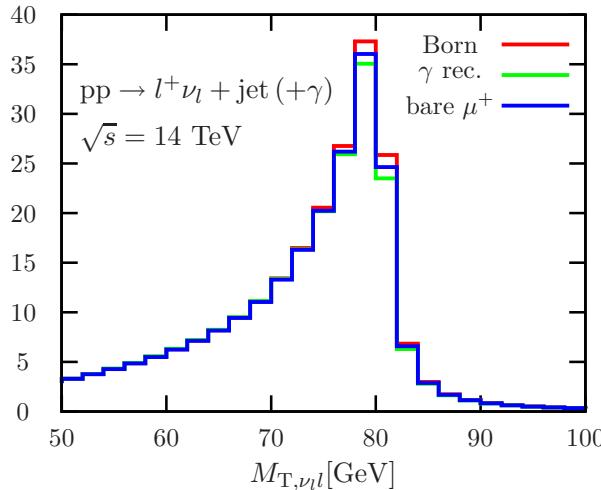
Comparison of EW corrections to W+jet and single (jet-inclusive) W production

→ interesting for W-mass determination via single-W production

$d\sigma/dp_{T,l} [\text{pb}/\text{GeV}]$



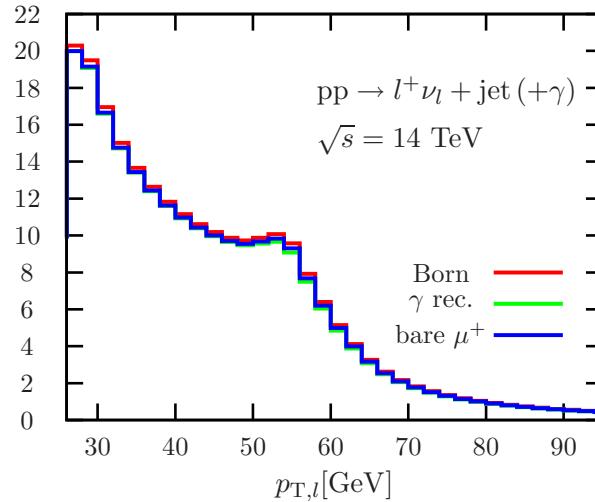
$d\sigma/dM_{T,\nu_l l} [\text{pb}/\text{GeV}]$



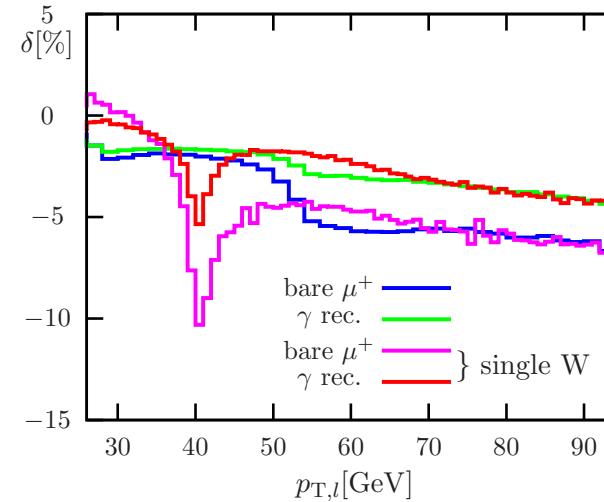
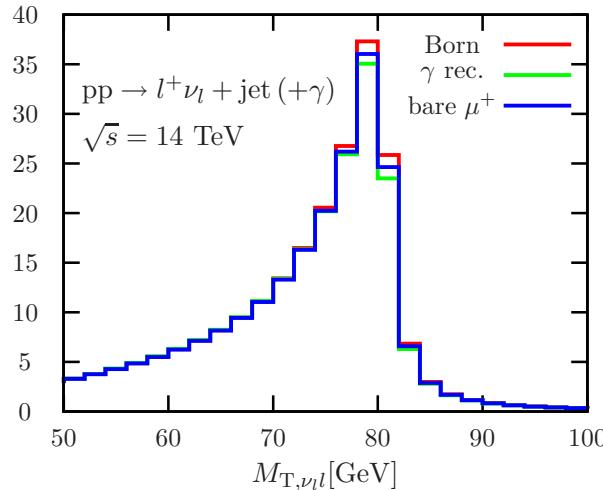
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$d\sigma/dp_{T,l} [\text{pb}/\text{GeV}]$



$d\sigma/dM_{T,\nu_l l} [\text{pb}/\text{GeV}]$



relative EW corrections
completely different

For single-W in NLO EW, see

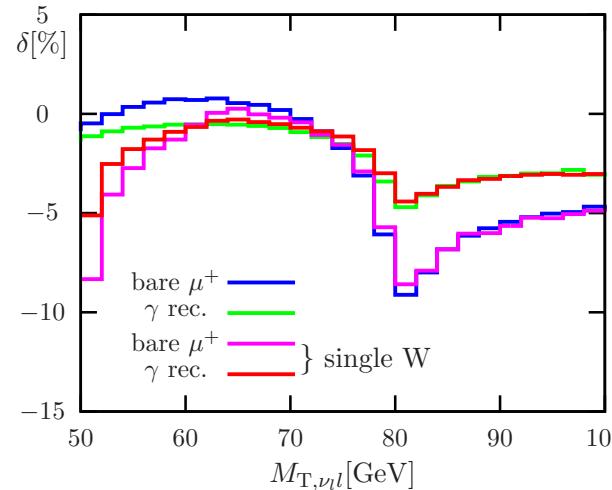
Baur et al. '98/'04

S.D./Krämer '01

Arbuzov et al. '06

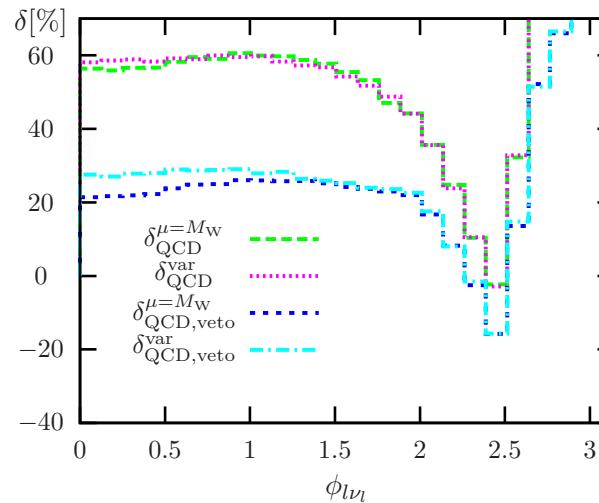
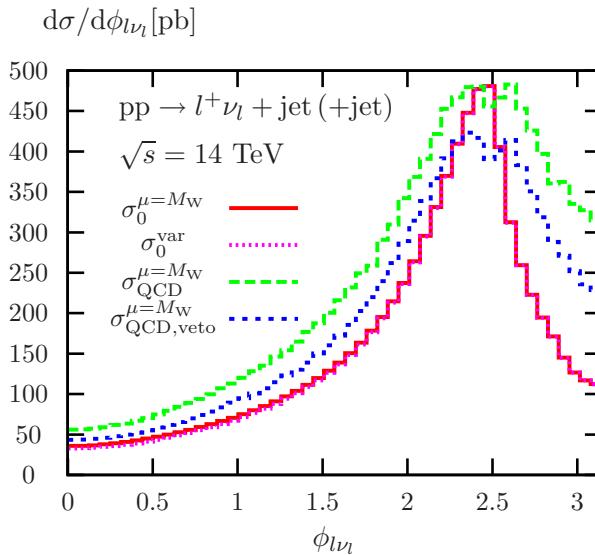
Carloni Calame et al. '06

relative EW corrections
practically identical
near Jacobian peak



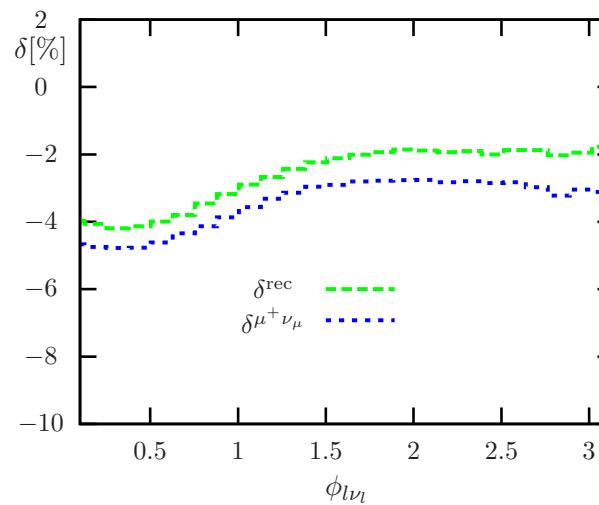
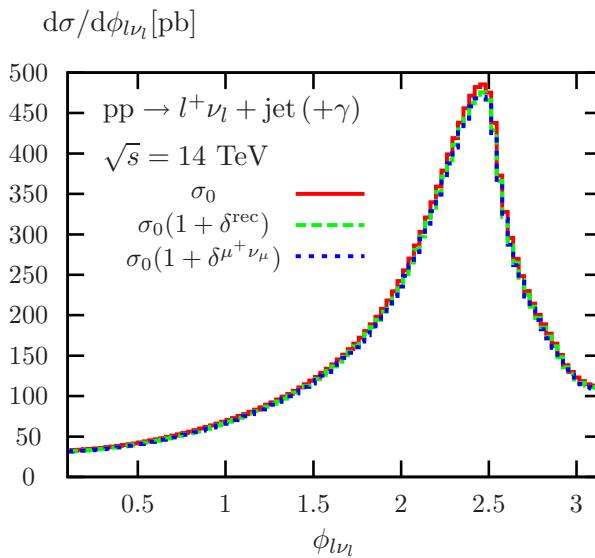
Distribution in the lepton–neutrino (p_T) azimuthal angle in the transverse plane

QCD corrections:



- significant distortion of shape
- corrections reduced by jet veto

EW corrections:



corrections small and relatively smooth