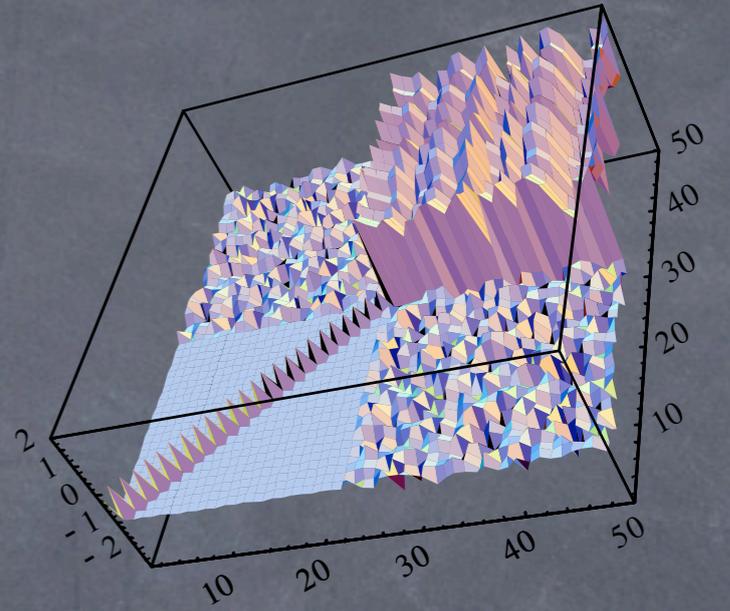
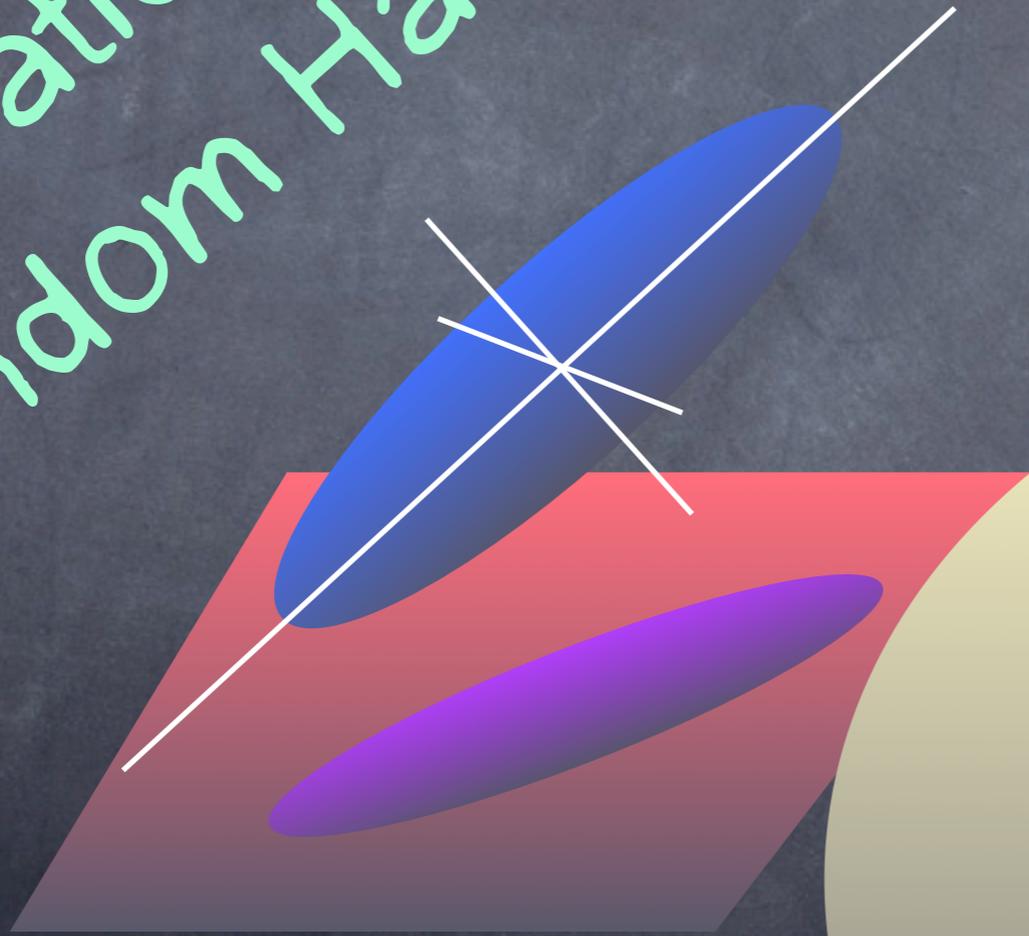
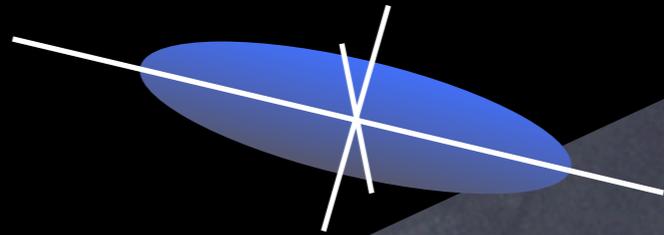


Grand Anarchy: unification patterns of random Hamiltonians



John Polston

The quantum mechanical coupling of a low energy system to a high energy one invariably leads to perturbative instability, naturally interpreted as dynamical instability.



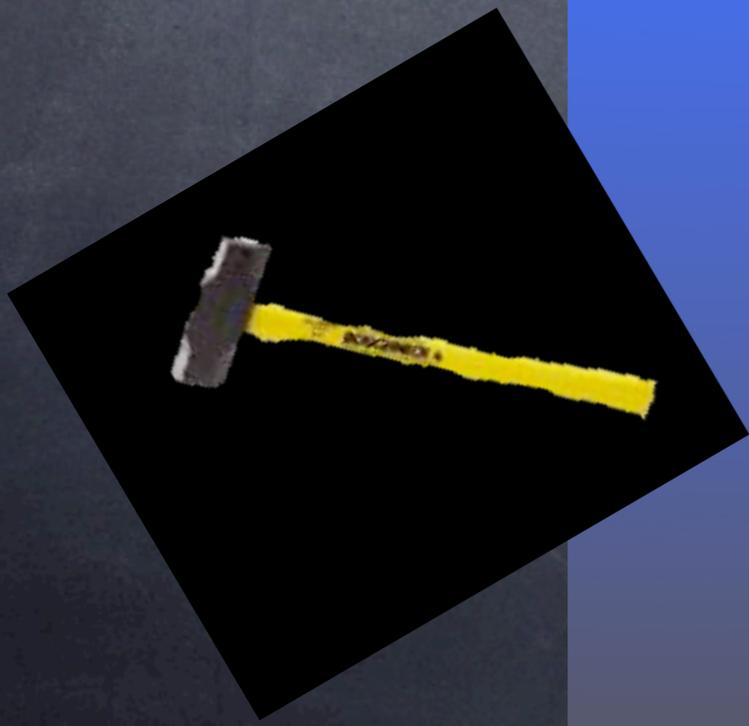
Let $\dim(H) = N \times N$

define random $H = U H_0 U^\dagger$,

with $U = \text{random} \in U(N)$, by Haar measure

case 1: H_0 has 1 "gutscale" eigenvalue,
 $N-1$ "small" eigenvalues $O(1)$

What are generic features of
random submatrix H_{PP} ?



random projectors $P+Q=1$,

$P^2=P$, $Q^2=Q$, $PQ=QP=0$.

dimensions $\dim P + \dim Q = N$

Random subsystem "P" ($\dim P$)

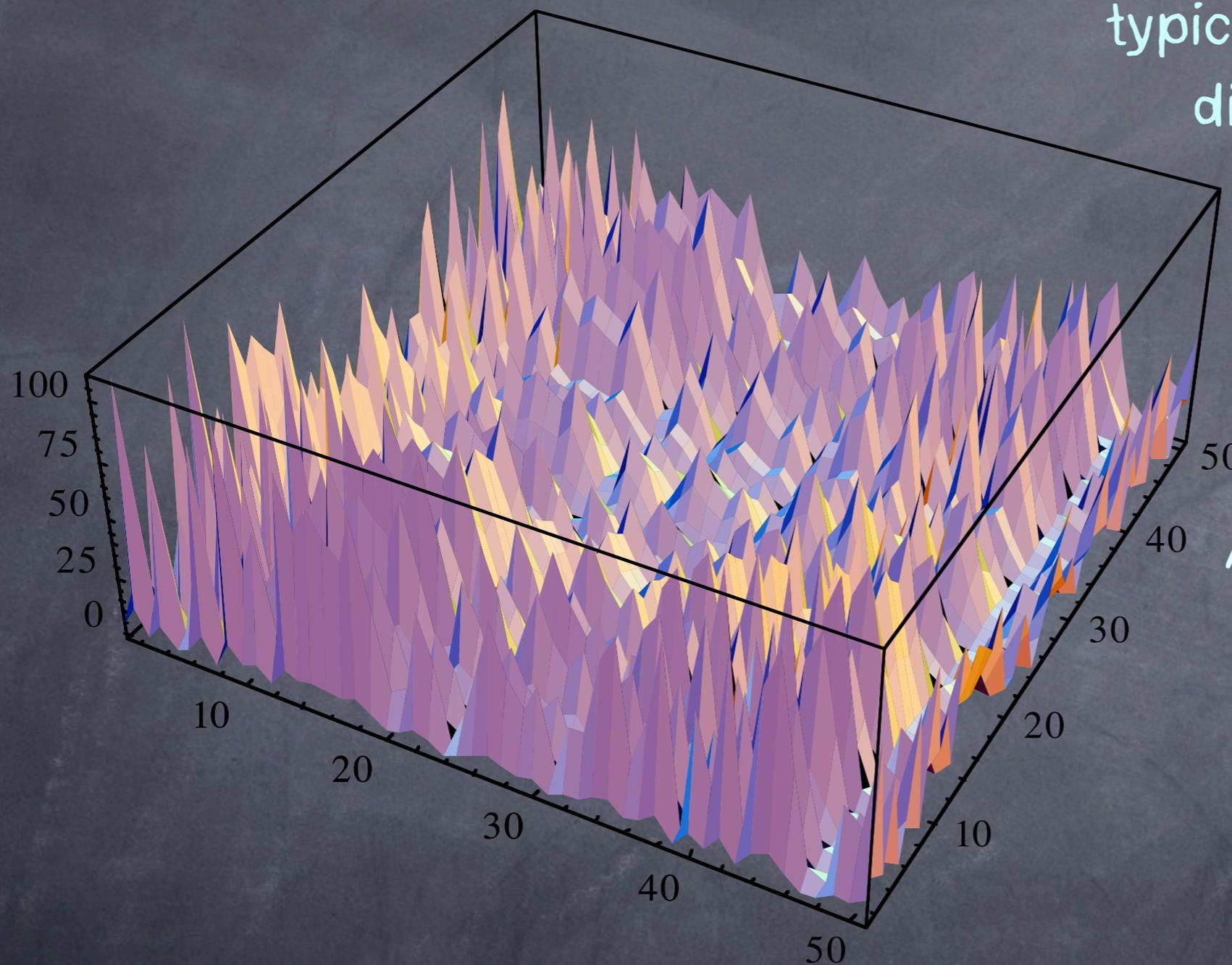
has Hamiltonian $H_{PP} = PHP$

In the diagonal frame of H_{PP} we have

$$H_{tot} \begin{pmatrix} E_1^{(0)} & 0 & 0 \dots & & \\ 0 & \dots & 0 \dots & & H_{PQ} \\ 0 & 0 & E_P^{(0)} \dots & & \\ & & & & \\ & H_{QP} & & & H_{QQ} \end{pmatrix} .$$

random gut-Hamiltonian in random frame

typical $H_{ij} \sim 0 \pm O(E_{\text{gut}}/N)$, $i \neq j$
diagonals $H_{ii} \sim O(E_{\text{gut}}/N)$



Abs[H_{ij}] shown;
 $E_{\text{gut}}=10^3$

Three dimensional plot of random matrix elements $|H_{ij}|$. One eigenvalue $E_{\text{gut}}=10^3$; remaining eigenvalues random $(0,1)$

Observation 1, case 1:

If H has 1 gutschale eigenvalue, ALMOST EVERY random submatrix H_{PP} has $P-1$ small eigenvalues, plus 1 large one

exact evals = 0.586, 0.231, 0.219, 0.161, 0.119, 0.109, 999

2 x 2 evals = 0.22073, 88.1463

80.7	-24.48	165.6	67.91	115.5	165.7	-21.58
-24.48	7.667	-50.22	-20.6	-35.01	-50.13	6.495
165.6	-50.22	340.8	139.8	237.6	341.	-44.55
67.91	-20.6	139.8	57.52	97.5	140.	-18.33
115.5	-35.01	237.6	97.5	165.9	237.8	-31.08
165.7	-50.13	341.	140.	237.8	341.7	-44.77
-21.58	6.495	-44.55	-18.33	-31.08	-44.77	6.148

4 x 4 evals = 0.300463, 0.183553, 0.118482, 486.088.

hierarchy !

Choose Any of an Infinity of Subspaces

If H has 1 gutscale eigenvalue, EVERY random submatrix H_{PP} has $P-1$ small eigenvalues, plus 1 large one

3 x 3 evals = 0.155178, 0.286777, 428.725

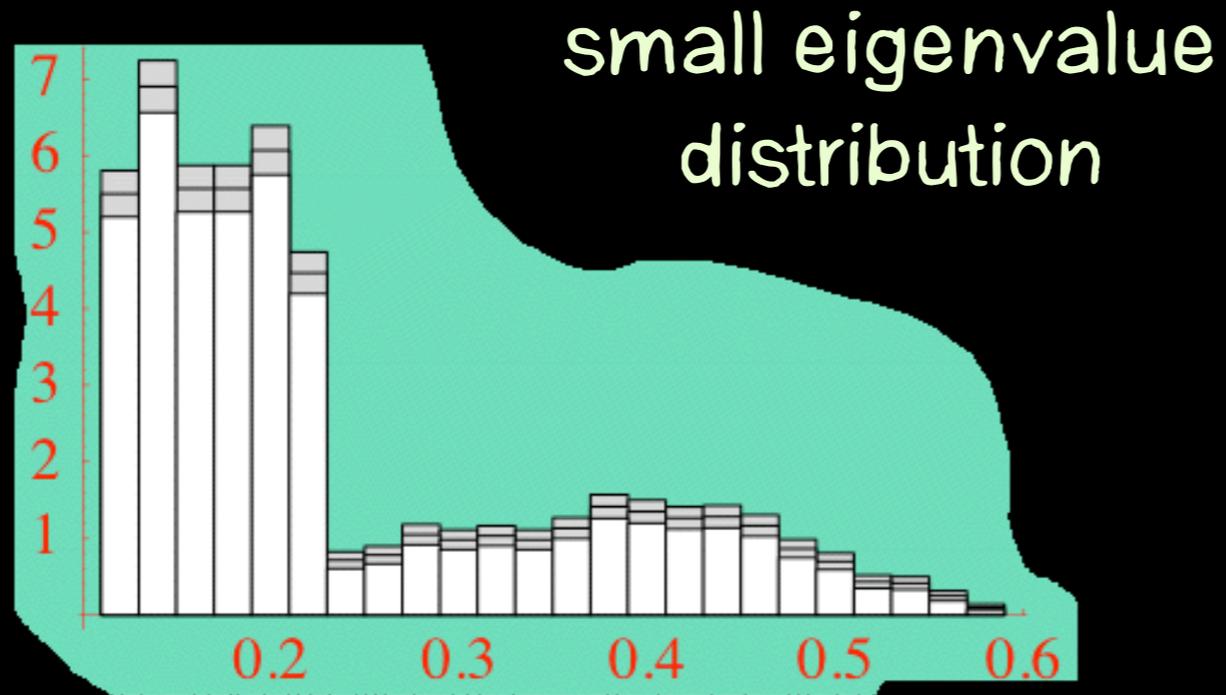
same hierarchy !

80.7	-24.48	165.6	67.91	115.5	165.7	-21.58
-24.48	7.667	-50.22	-20.6	-35.01	-50.13	6.495
165.6	-50.22	340.8	139.8	237.6	341.	-44.55
67.91	-20.6	139.8	57.52	97.5	140.	-18.33
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hierarchy !

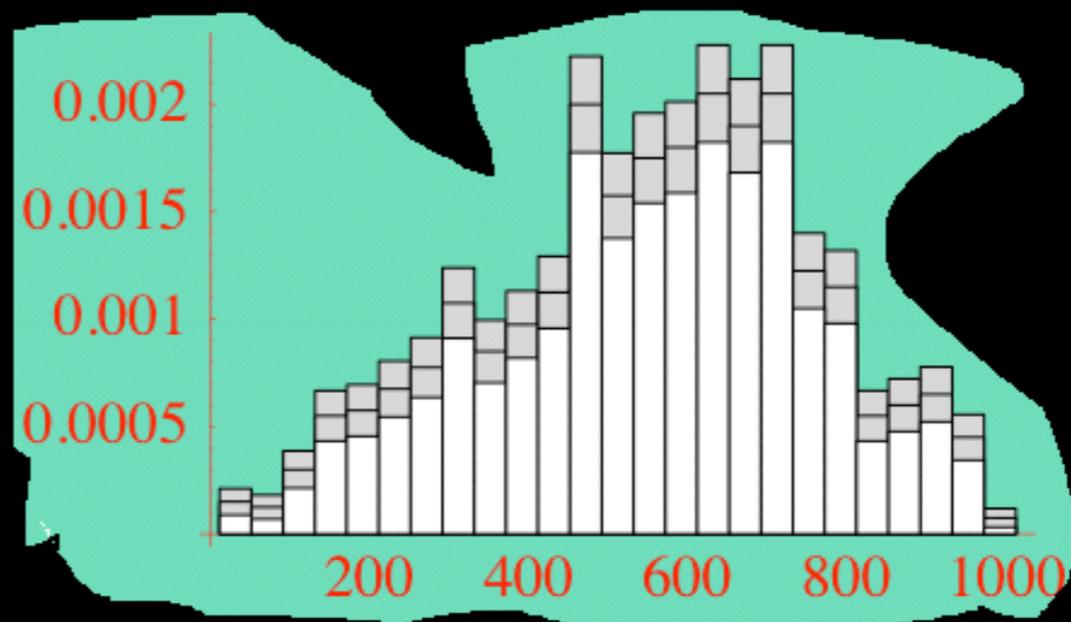
new 4 x 4 evals = 0.366832, 0.194968, 0.142922, 570.548

there's almost always a hierarchy:
highly bimodal distributions on any dimension



$$E_{\text{gut}} = 10^3$$

large eigenvalue
distribution



Why ?

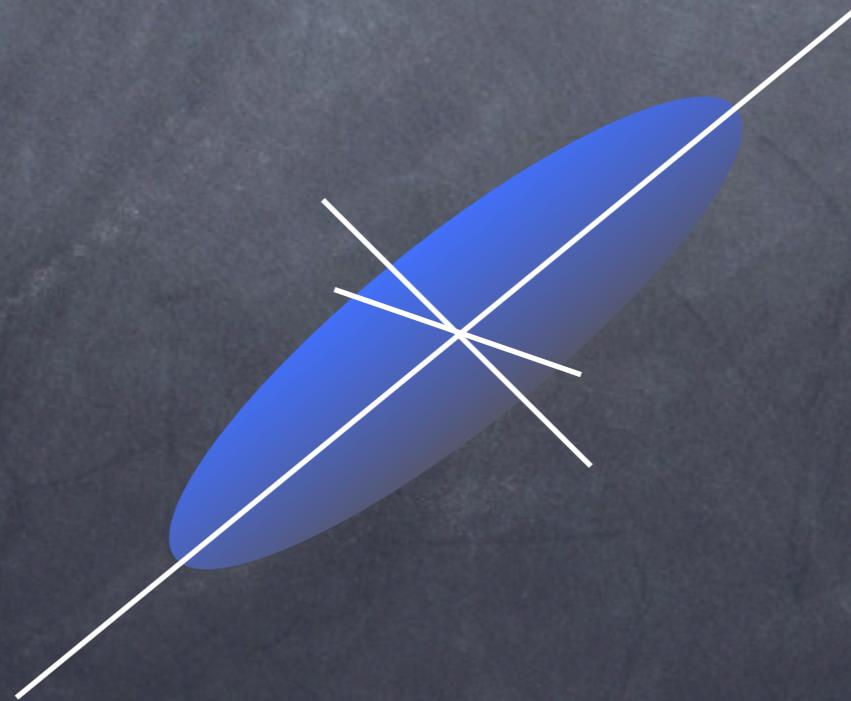
"entropy of the Hamiltonian"

$$H_r = H / \text{tr}(H)$$

$$S_H = -\text{tr}(H_r \log(H_r));$$

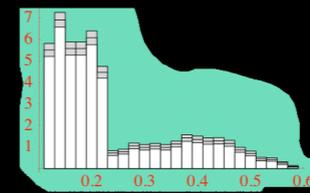
$$0 \leq S_H \leq \log(N)$$

$\exp(S_H) =$ effective dimension of the eigenspace

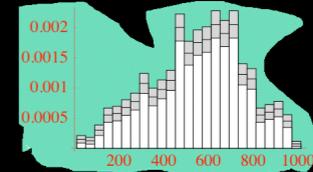


(an invariant
geometrical
measure)

$0 < \text{eval} < 1$

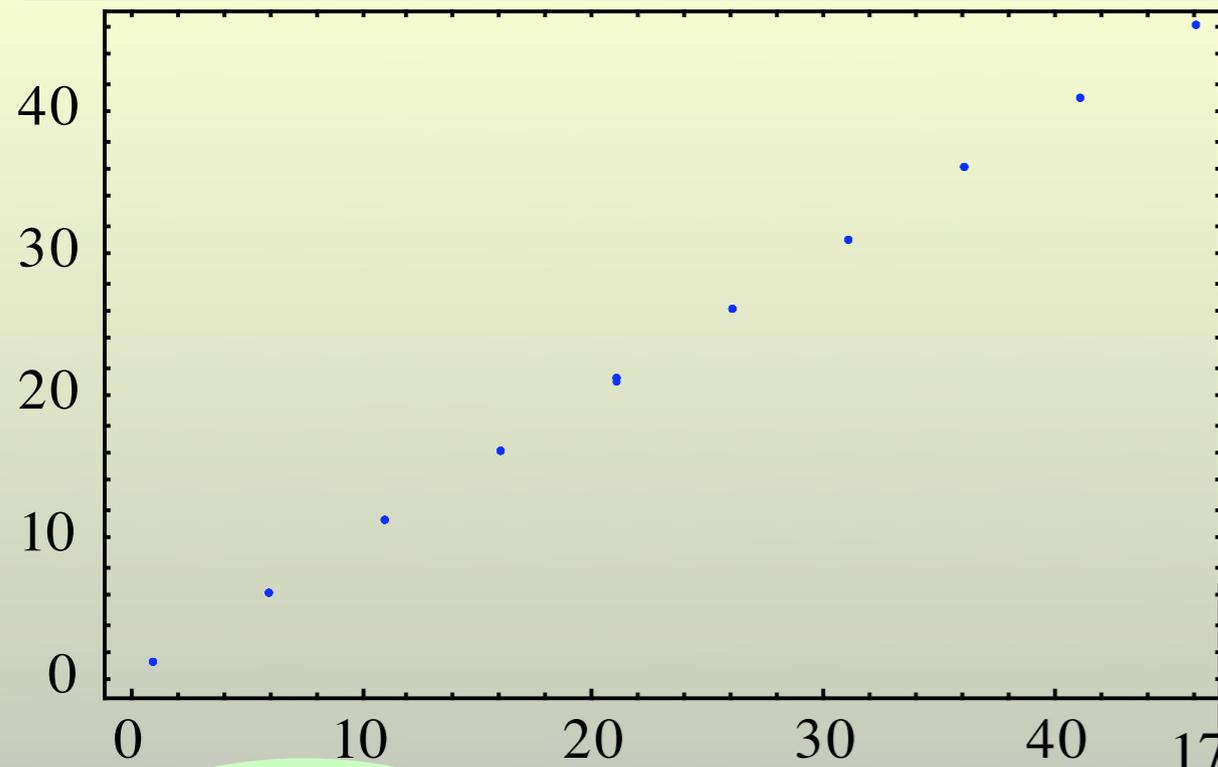


$100 > \text{eval} < 1000$



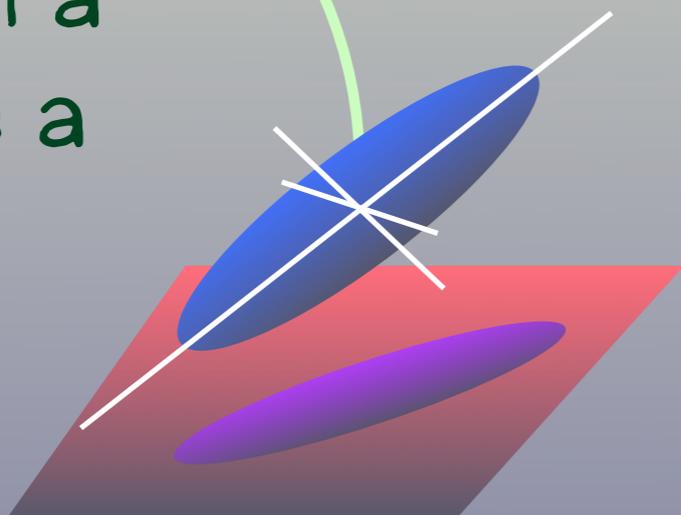
entropy S_{PP} scales like entropy S_H

$\exp(S_H) =$ effective dimension of the eigenspace



number of big evals

"almost all shadows of a needle are a needle"



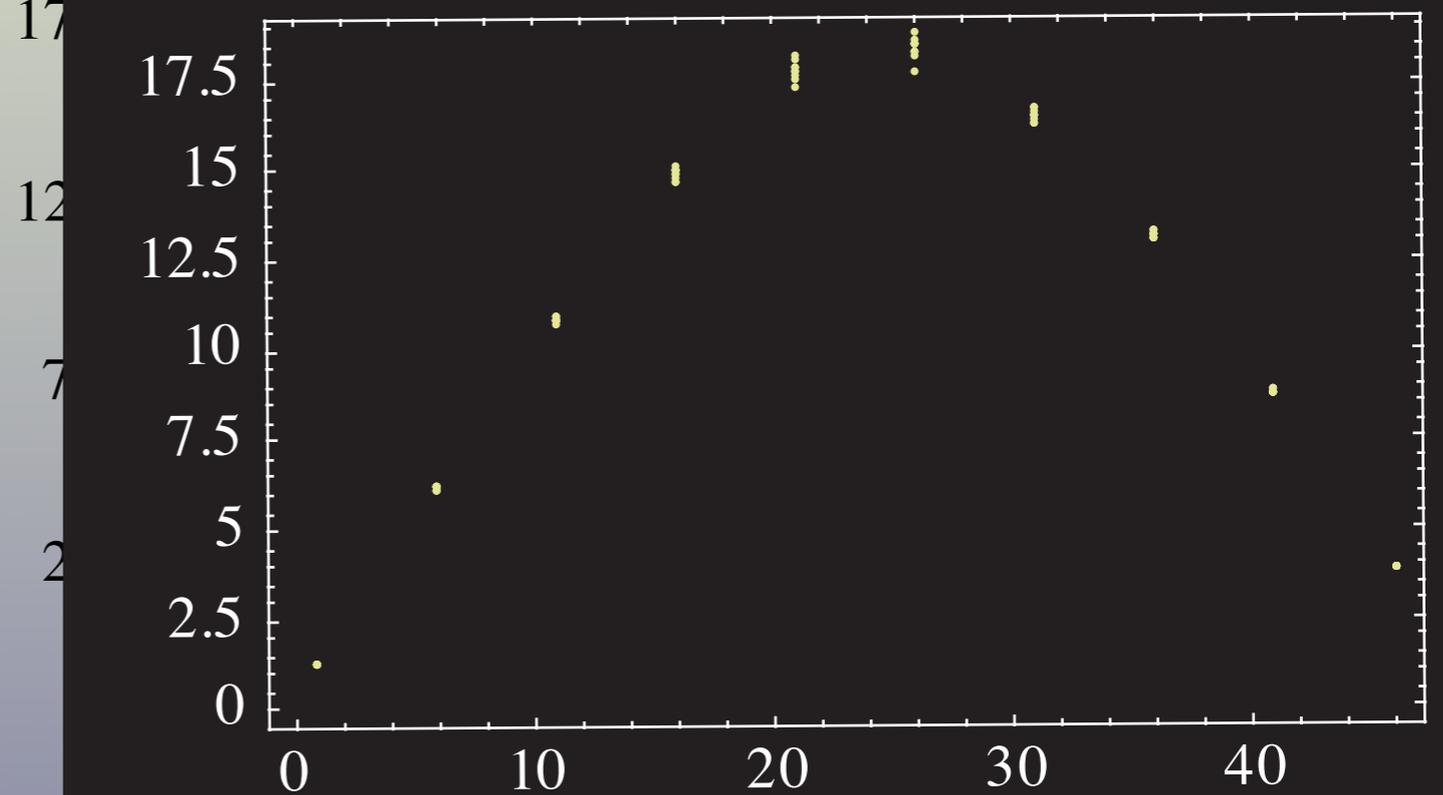
$$H_r = H / \text{tr}(H)$$

$$S_H = -\text{tr}(H_r \log(H_r))$$

$$S_{PP} = -\text{tr}(H_{rPP} \log(H_{rPP}))$$

$\exp(S_{pp}) =$ effective dimension of the eigenSUBspace of dim $N/2$

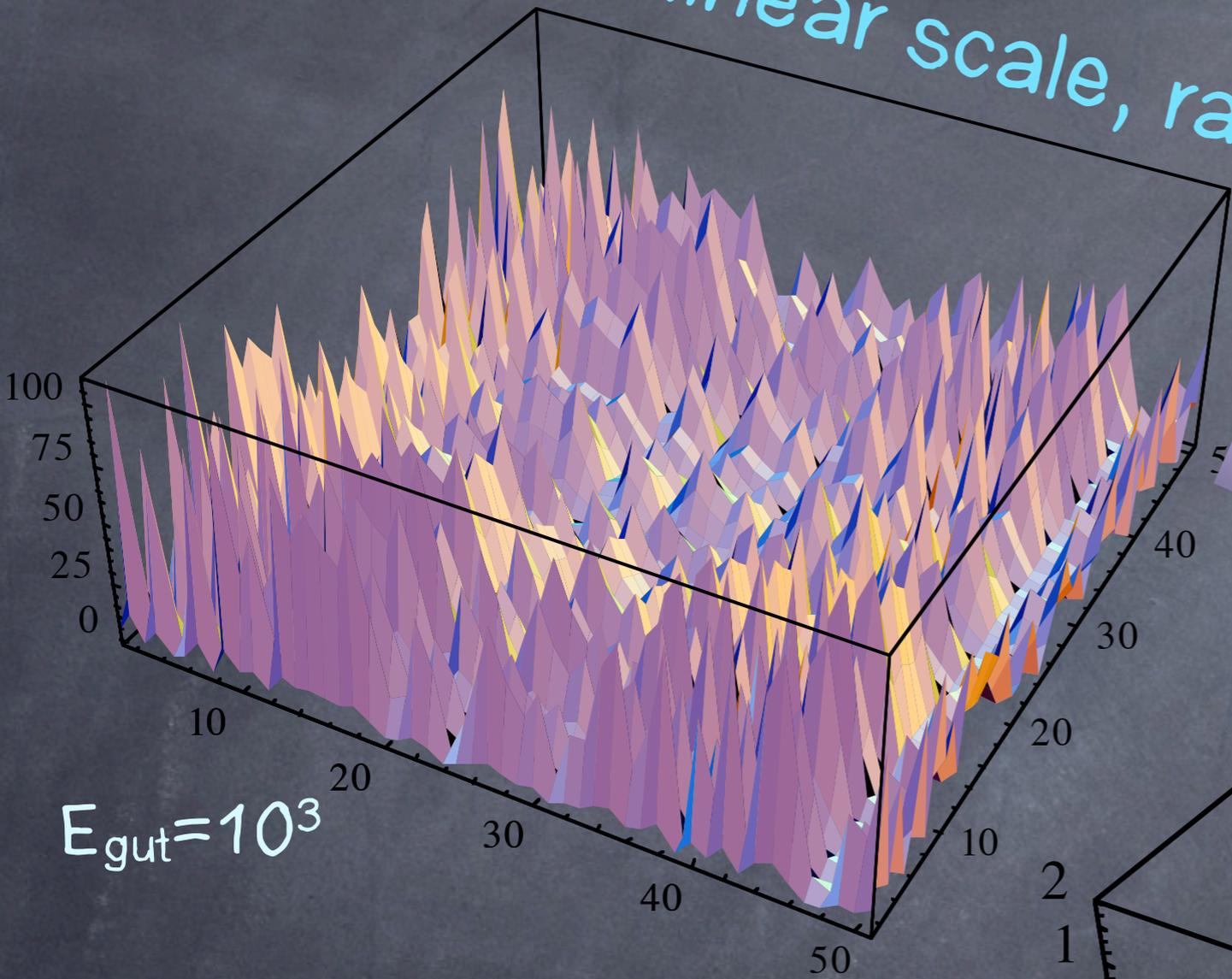
- hpp- sub entropy v big



number of big evals

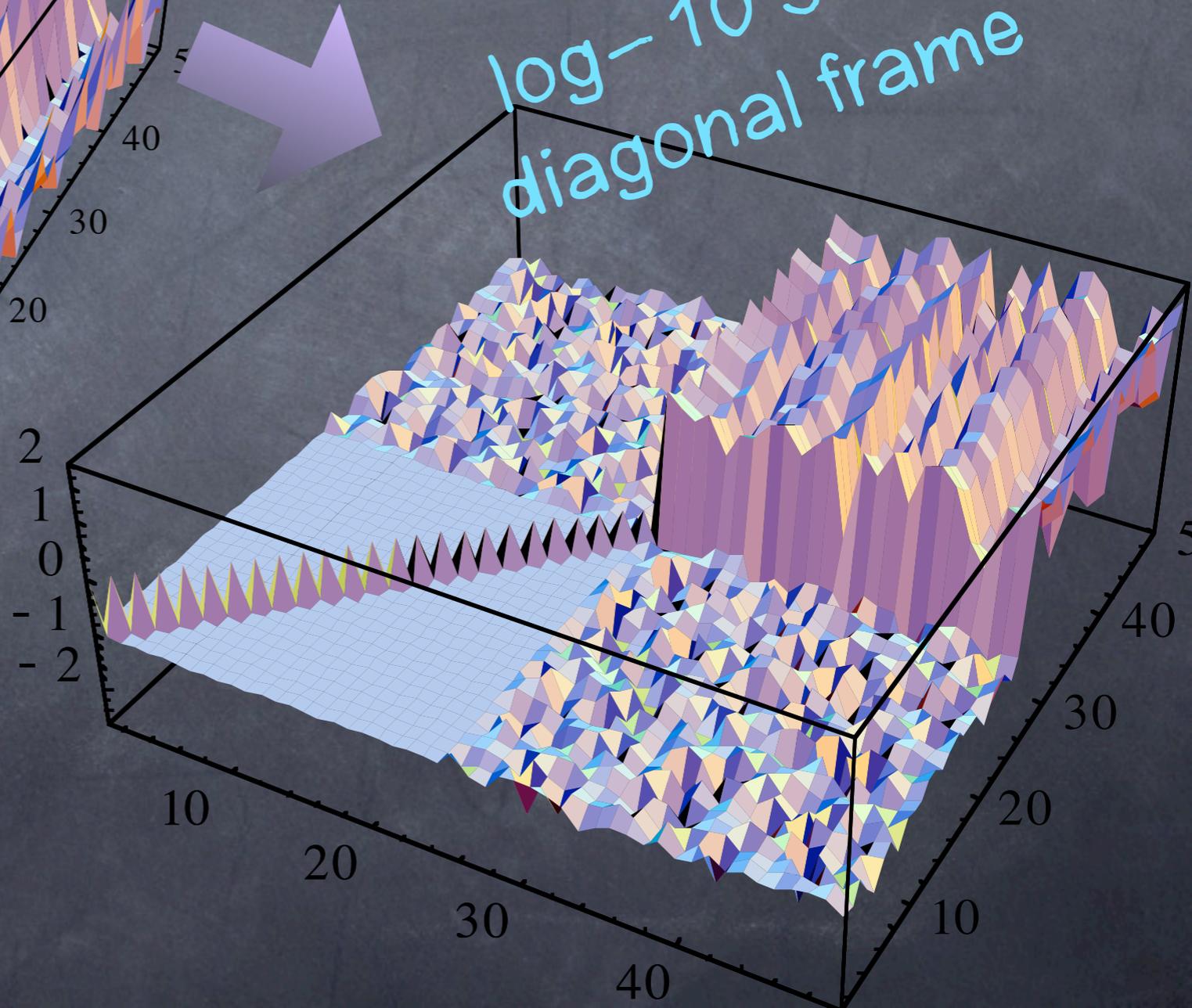
things are different in the diagonal frame of H_{pp}

linear scale, random frame



$E_{\text{gut}} = 10^3$

log-10 scale
diagonal frame

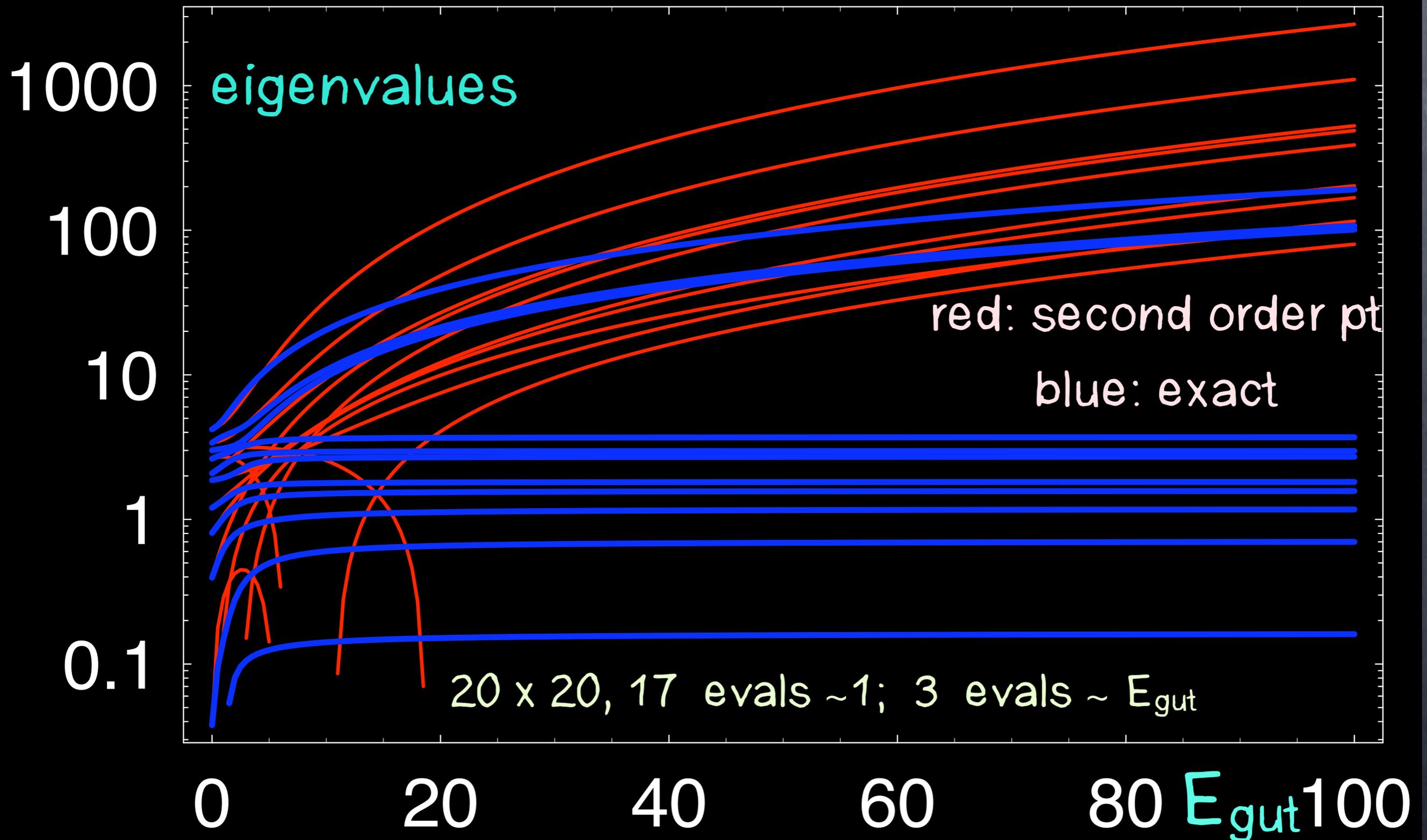


“decoupling
anarchy”

it's non-perturbative:

not available
from ordinary
perturbation theory

Rayleigh–Schroedinger perturbation theory fails with a hierachy



Brillouin–Wigner Perturbation Theory...

...the alternative at strong coupling

$$E \sim E_1 + \sum_{mn} V_{pn} \frac{1}{E - E_n^{(0)}} V_{np} + O(V^3)$$

...implicit, transcendental equations for E...but not exact

Our procedure is related...but exists EXACT

“gap equation identity”

(new: see preprint):

$$\det(E - H) = (-1)^P E^P \det(E - K_P(E)) \det(E - H_{QQ})$$

$$K_p(E) = H_{PP} + H_{PQ} \frac{1}{E - H_{QQ}} H_{QP}$$

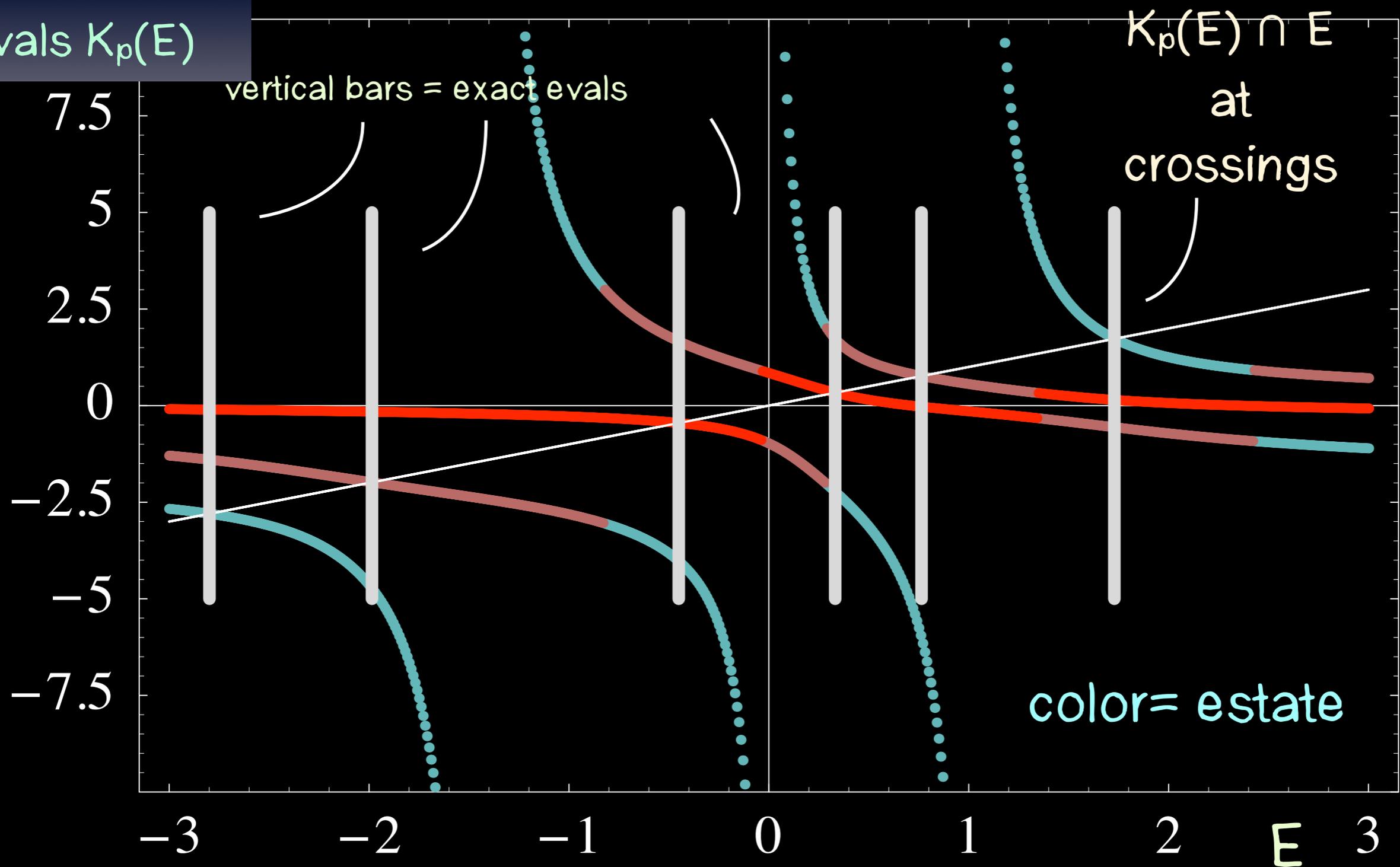
on subspace P, $\{K_p(E) \cap E\}$ has the same spectrum as

H on space $P \oplus Q$

fano; feshbach

$P+Q = 6$ exact eigenvalues on a $P = 3$ dimensional subsystem

evals $K_p(E)$



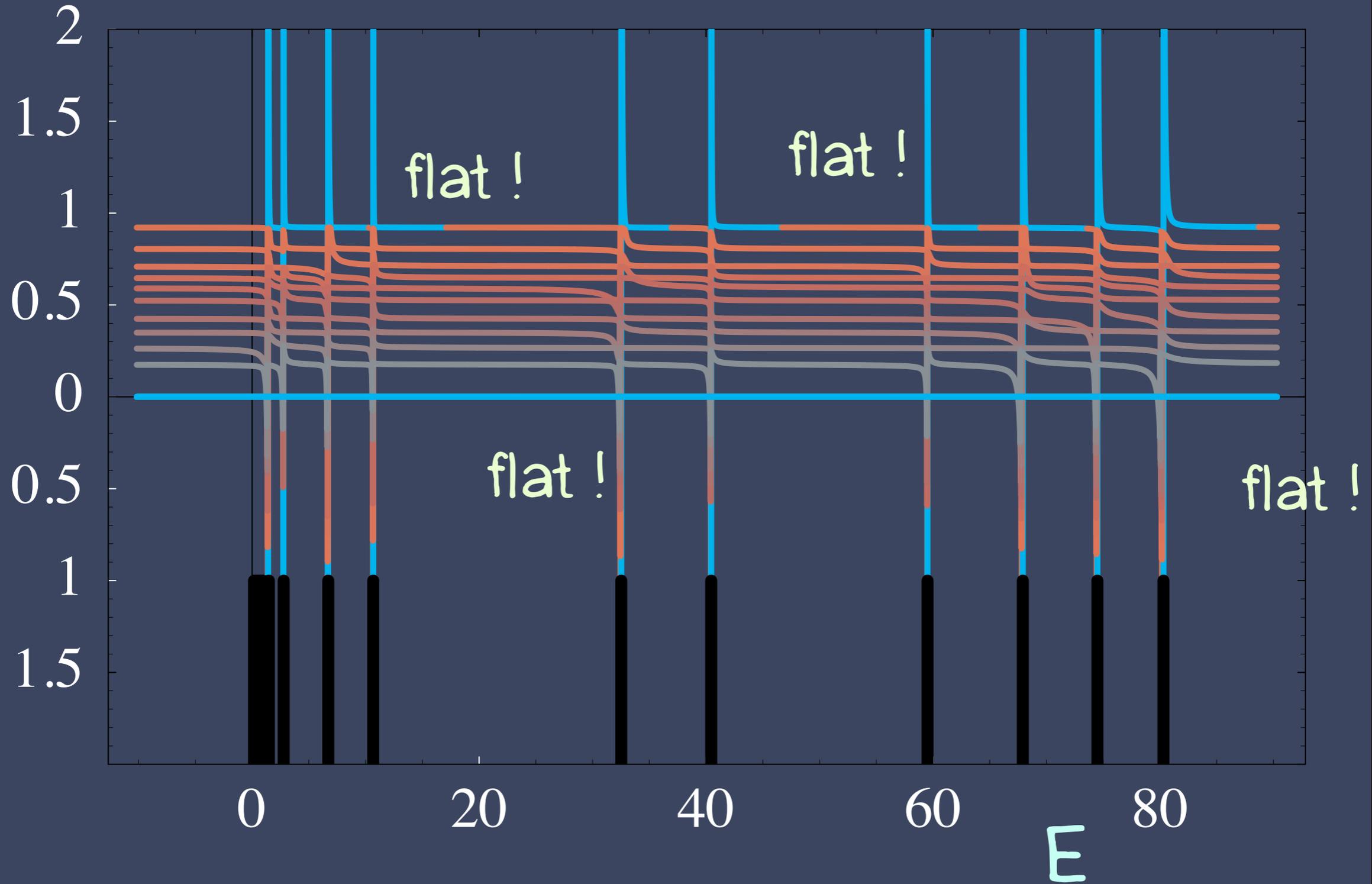
Solving the nonlinear eigenvalue equation $K_p(E) = E$. An $N=6$ dimensional system has been partitioned into direct sums of orthogonal $P=3$ and $Q=3$ dimensional subspaces

Stability of Procedure Under Hierarchy

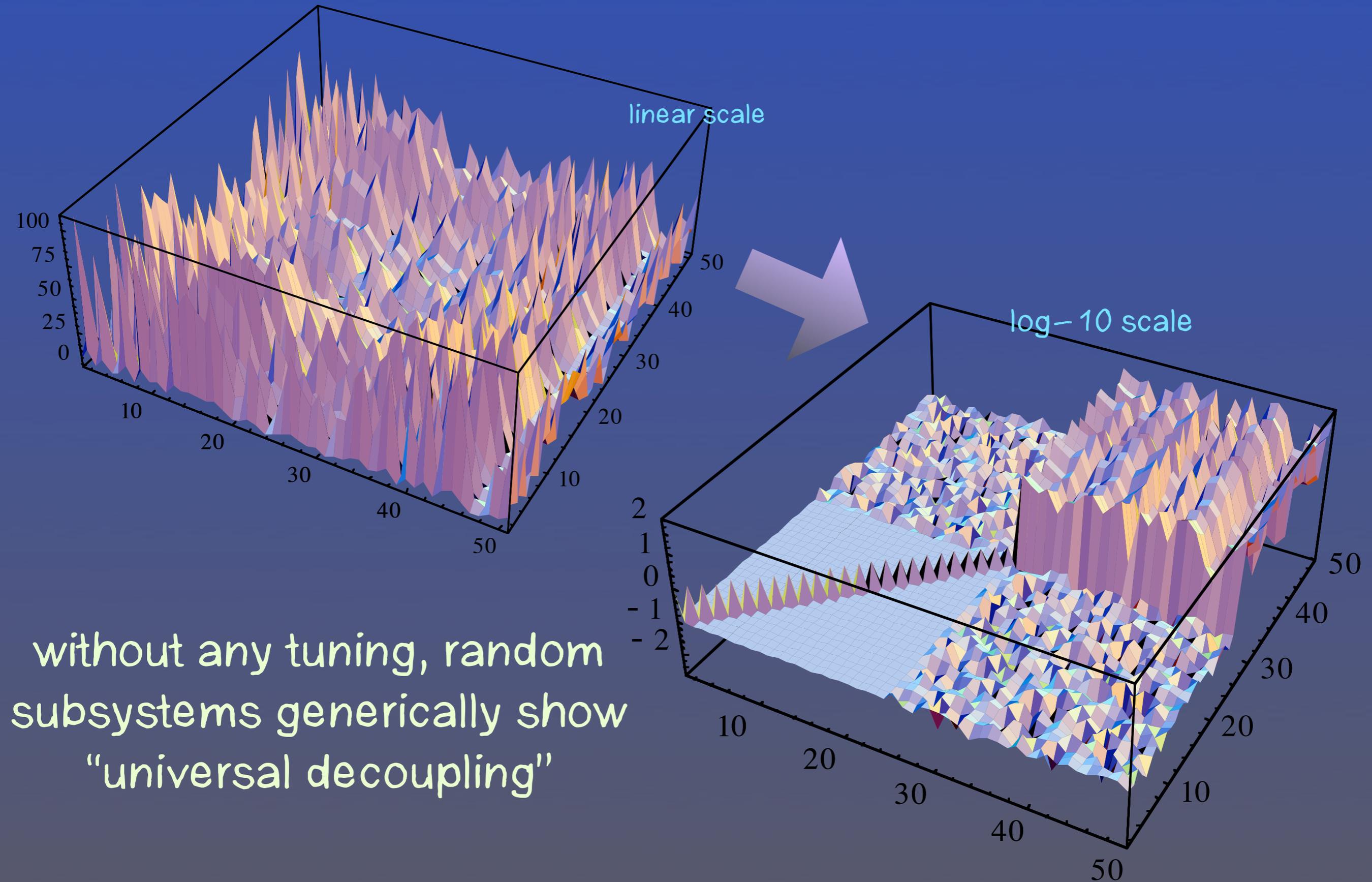
black bars are exact eigenvalues;

singularities are poles of H_{QQ}

evals $K_p(E)$



the perception of dynamical instability is deceptive.



without any tuning, random
subsystems generically show
"universal decoupling"

Grand Anarchy

Given a big hierarchy,
there's almost always a hierarchy
on a random subspace.

Diagonalizing on most
random subspaces,
Rayleigh Schroedinger
perturbation theory signals a
deceptive instability

...yet the Exact Gap Equation is Stable:

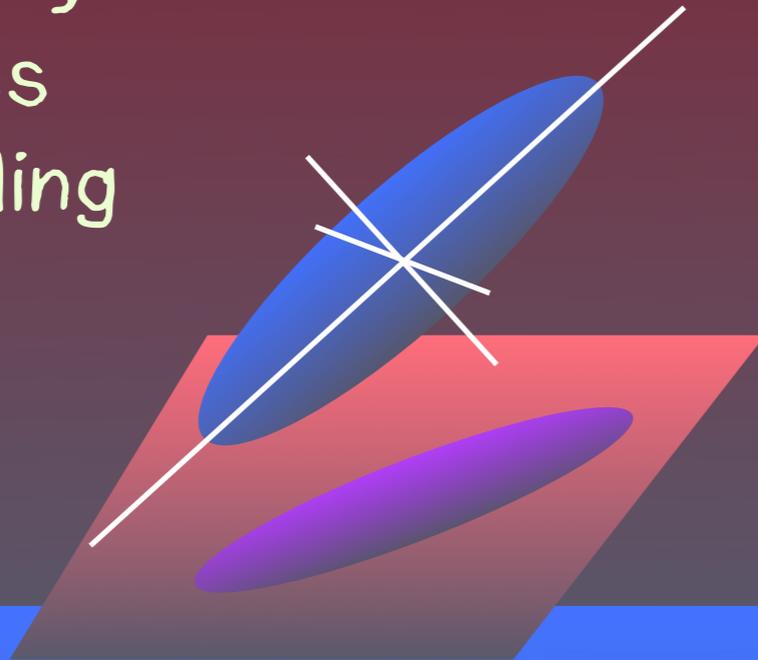
$$K_p(E) = H_{pp} + H_{pq} \frac{1}{E - H_{qq}} H_{qp}$$

$$K_p(E) |\psi_p\rangle = \lambda(E) |\psi_p\rangle$$

with $\lambda(E) \rightarrow E$

Grand Anarchy

geometry
implies
decoupling



equivalent to cancellations
by symmetries?
perhaps!

..yet the Exact Gap Equation is Stable:

$$K_p(E) = H_{pp} + H_{pq} \frac{1}{E - H_{qq}} H_{qp}$$

$$K_p(E) |\psi_p\rangle = \lambda(E) |\psi_p\rangle$$

with $\lambda(E) \rightarrow E$



