

# Implications of SUSY Phenomenology for Baryogenesis

**Björn Garbrecht**

*Theoretical Nuclear, Particle, Astrophysics and Cosmology  
(NPAC) Group  
Department of Physics  
University of Wisconsin–Madison*

Work with D.J.H. Chung, M.J. Ramsey-Musolf and S. Tulin

**PHENO 09, Madison, May 11th**

# Baryogenesis and Phenomenology

Terrestrial  
Experiments:  
Colliders, EDM  
Searches,...

Theoretical  
Parameters  
(M)SSM

Cosmology:

$$\eta_B = \frac{n_B}{n_\gamma} \approx 6 \times 10^{-10},$$

dark matter, ...

This work: establish (new) connections  
& improve predictions for  $\eta_B$ .

## Leptogenesis:

Point out that  $\phi_\mu^{M_{1,2}}$ ,  $\phi_\mu^A$  lead to leptogenesis in  $\nu$ MSSM.

New connection of leptogenesis to  $CP$ -violating phenomenology.

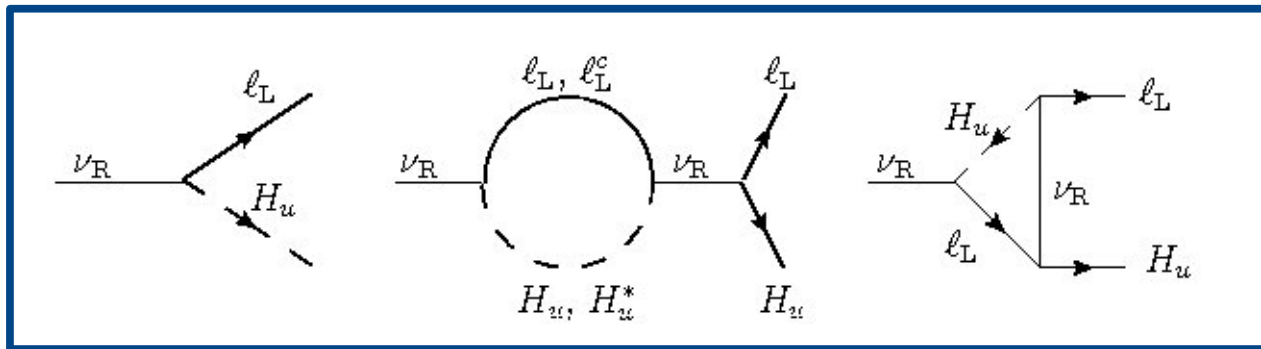
DJH Chung, BG, MJ Ramsey-Musolf,  
arxiv:0904:1591 [hep-ph]

## EWBG:

Besides  $CP$ -phases  $\phi_\mu^{M_{1,2}}$ ,  $\phi_\mu^A$ , also sparticle mass spectrum and  $\tan \beta$  are crucial.

DJH Chung, BG, MJ Ramsey-Musolf, S Tulin, PRL.(2009),  
[arXiv:0808.1144 [hep-ph]]

# Leptogenesis from Leptonic $CP$ -Phases



Yukawa interactions:

$$H_u \nu_R^c T \mathbf{Y}_\nu \ell_L + \text{h.c.}$$

Fukugita, Yanagida '86;  
Covi, Roulet, Vissani '96

→ each vertex suppressed by  $\mathbf{Y}_\nu$ ,  $CP$ -asymmetry by  $\mathbf{Y}_\nu^4$

• Type I see-saw:  $\tilde{\mathbf{m}} = \mathbf{Y}_\nu^T \mathbf{Y}_\nu \frac{v^2}{\mathbf{m}_{\nu_R}} \lesssim O(\text{eV}) \longrightarrow m_{\nu_R}^i \gtrsim 10^8 \text{ GeV}$   
Davidson, Ibarra '02

• **Loophole: Resonant Leptogenesis** Pilaftsis '97; Pilaftsis, Underwood '04

•  $CP$ -violating phases for leptogenesis in general unrelated to  $CP$ -violating observables at low energies. but: flavor effects: Pilaftsis, Underwood '05; Abada, Davidson, Josse-Michaux, Losada, Riotto '06; Nardi, Nir, Roulet, Racker '06; Pascoli, Petcov, Riotto '07

• Small Yukawa couplings and/or large  $m_{\nu_R}$

→ leptogenesis effectively decoupled from (MS)SM.

for loopholes, see e.g.: Pilaftsis '04; Blanchet, Chacko, Granor, Mohapatra '09

# Leptogenesis from Non-Leptonic $CP$ -Phases

DJH Chung, BG, MJ Ramsey-Musolf, arxiv:0904:1591 [hep-ph]

- Avoid suppression by powers of neutrino Yukawa couplings.
- Take  $CP$ -violating phases from Higgs sector beyond the SM.
- Add  $\nu_R$  ( $m_{\nu_R}$  unrelated to  $CP$ -asymmetry), but no extra symmetries/particles/interactions/tuning in lepton sector.

## Worked examples:

- Non-SUSY two Higgs doublet model and  $CP$ -violation from 3<sup>rd</sup> generation Yukawa couplings.
- Leptogenesis in  $\nu$ MSSM
  - bino driven ← this talk
  - wino driven ← completely analogous to bino driven
  - $A_{t,b,\tau}$ -term driven ← very similar to bino driven scenario

# Bino-Driven Leptogenesis in the $\nu$ MSSM

The main ingredients:

$b$  -term provides mixing between  $H_u, H_d$

$$\begin{aligned} \mathcal{L} \supset & -(m_{H_u}^2 + \mu^2)H_u^\dagger H_u - (m_{H_d}^2 + \mu^2)H_d^\dagger H_d - b (H_u^T H_d + H_u^\dagger H_d^*) \\ & - \frac{1}{2}M_1 \bar{\Psi}_{\tilde{B}} \Psi_{\tilde{B}} - \mu \bar{\Psi}_{\tilde{H}^0} \Psi_{\tilde{H}^0} - \mu \bar{\Psi}_{\tilde{H}^+} \Psi_{\tilde{H}^+} \\ & - \frac{g_1}{\sqrt{2}} \left[ \bar{\Psi}_{\tilde{H}^+} \left( -H_d^{-*} P_L + e^{i\phi_\mu^{M_1}} H_u^+ P_R \right) \Psi_{\tilde{B}} \right. \\ & \quad \left. + \bar{\Psi}_{\tilde{H}^0} \left( -H_d^{0*} P_L - e^{i\phi_\mu^{M_1}} H_u^0 P_R \right) \Psi_{\tilde{B}} + \text{h.c.} \right] \end{aligned}$$

Have chosen  $\mu, M_1$  to be real.  
 $CP$ -violation resides in this interaction term.

$$\begin{aligned} \Psi_{\tilde{H}^+} &= \begin{pmatrix} \tilde{H}_u^+ \\ \tilde{H}_d^{-\dagger} \end{pmatrix}, & \Psi_{\tilde{H}^0} &= \begin{pmatrix} -\tilde{H}_u^0 \\ \tilde{H}_d^{0\dagger} \end{pmatrix}, \\ \Psi_{\tilde{B}} &= \begin{pmatrix} \tilde{B} \\ \tilde{B}^\dagger \end{pmatrix} \end{aligned}$$

# Higgs Mixing

$$\begin{pmatrix} H_u^+ \\ H_d^{-*} \end{pmatrix} = \begin{pmatrix} \cos \alpha & \sin \alpha \\ -\sin \alpha & \cos \alpha \end{pmatrix} \begin{pmatrix} H_1^+ \\ H_2^+ \end{pmatrix}, \quad \begin{pmatrix} H_u^0 \\ H_d^{0*} \end{pmatrix} = \begin{pmatrix} \cos \alpha & -\sin \alpha \\ \sin \alpha & \cos \alpha \end{pmatrix} \begin{pmatrix} H_1^0 \\ H_2^0 \end{pmatrix}$$

$$\tan 2\alpha = \frac{2b}{m_{H_u}^2 - m_{H_d}^2}$$

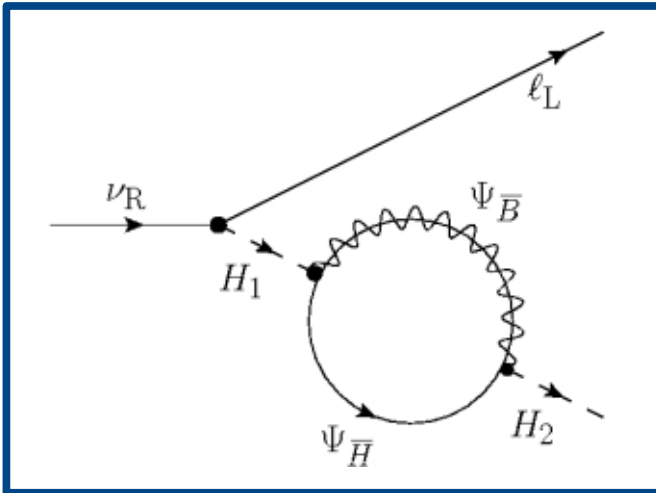
Recall:  $b$ -term is related to  $\tan \beta$  in the MSSM:  $b = m_A^2 \sin \beta \cos \beta$

$$\text{For } m_A \gg m_Z: \alpha \approx \frac{1}{\tan \beta}$$

# CP-Asymmetry

Assume  $m_{H_1} < m_{\tilde{B}} + m_{\tilde{H}} < m_{H_2}$

$$\varepsilon = \frac{\Gamma_{\nu_R \rightarrow \ell_L + H_2} - \Gamma_{\nu_R \rightarrow \ell_L^c + H_2^*}}{\Gamma_{\nu_R \rightarrow \ell_L + H_2} + \Gamma_{\nu_R \rightarrow \ell_L^c + H_2^*} + \Gamma_{\nu_R \rightarrow \ell_L + H_1} + \Gamma_{\nu_R \rightarrow \ell_L^c + H_1^*}}$$

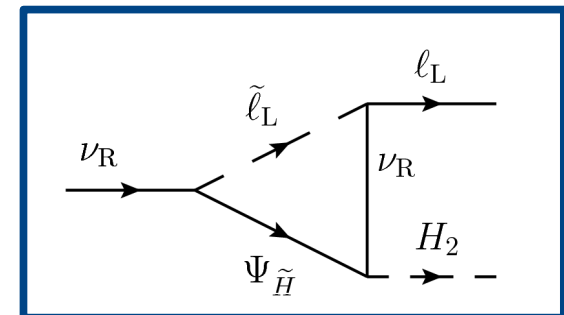


$$\varepsilon \approx \frac{\cos \alpha \sin \alpha \sin \phi_\mu^{M_1}}{8\pi} \frac{g_1^2 \mu M_1}{m_{H_2}^2 - m_{H_1}^2}$$

This vanishes if either  $\mu \rightarrow 0$ ,  $M_1 \rightarrow 0$ , as it should.

Thermal corrections will be sizable.

Also vertex correction can contribute.



# Strong Washout

- Take  $m_{\nu_R} = O(\text{TeV})$ , assume that the  $\nu_R$  were in thermal equilibrium before their decay.
- Suppression of lepton asymmetry due to inverse decays.

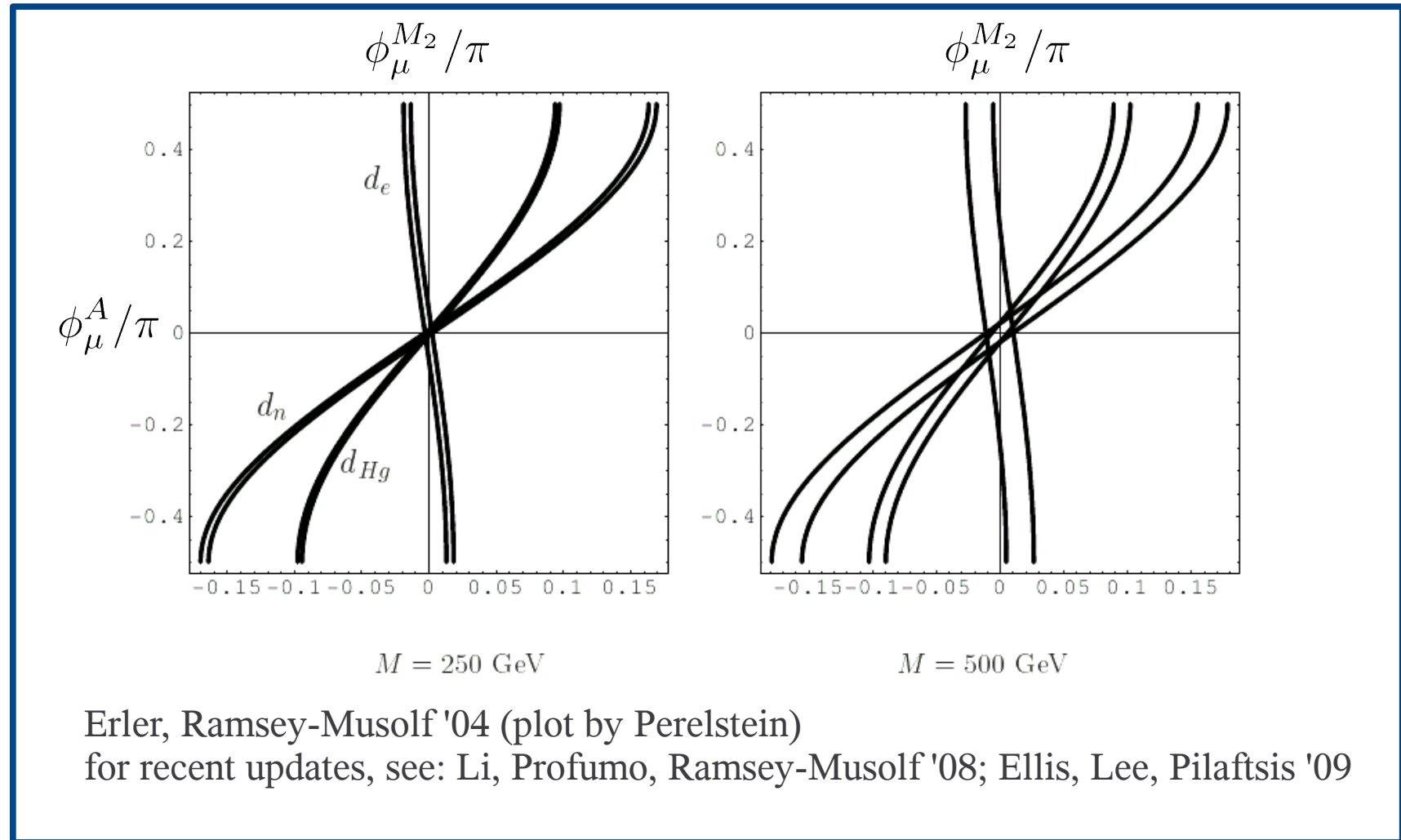
For analytical calculation, see Buchümller, Di Bari, Plümacher '04

Necessary  $CP$ -phase to explain the observed baryon asymmetry:

$$\sin \phi_{\mu}^{M_1} \approx -2 \times 10^{-6} \frac{\tilde{m}}{10^{-3} \text{eV}} \frac{m_{H_2}^4}{bg_1^2 \mu M_1}.$$



# Electric Dipole Moments



Higgs Sector/Collider Pheno: Pilaftsis, Wagner '99; Pilaftsis, Ellis, Wagner '00; Kittel '09;...

**A very testable scenario for baryogenesis.**

# Sterile Neutrinos as Drain for Lepton Number

- An alternative to conventional washout scenarios.
- If  $m_{\nu_R} < m_{H_2}$ ,  $CP$ -violating decays  $H_2^* \rightarrow \ell_L + \nu_R$ ,  $H_2 \rightarrow \ell_L^c + \nu_R$
- Yuakwa interaction  $y_\nu$  is out of equilibrium above  $T_{EW} = 100\text{GeV}$ , if

$$y_\nu \ll 10^{-7} \Leftrightarrow \frac{m_{\nu_R}}{10^3 \text{eV}} \frac{\tilde{m}}{10^{-3} \text{eV}} \ll 10^9$$

- Washout from inverse decays is completely negligible
- Cosmic exclusion window for light  $\nu_R$

Dodelson, Widrow '94; Akhmedov, Rubakov, Smirnov '98; Shaposhnikov et.al.; Kusenko et. al.

$$20 \text{ keV} \lesssim m_{\nu_R} \lesssim 1 \text{ GeV}$$

OK, keV neutrinos may serve as warm dark matter

Overclosure/BBN conflict

OK, decay before BBN

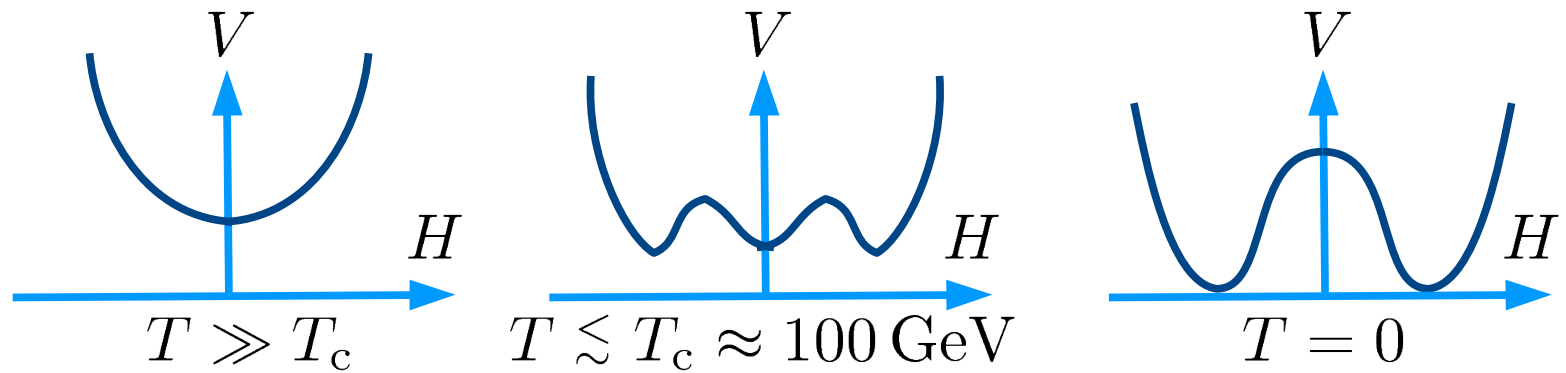
- Predicted  $CP$ -phase:

$$\sin \phi_\mu^{M_1} = -2 \times 10^3 \frac{m_{H_2}^4}{bg_1^2 \mu M_1} \frac{m_{H_2}}{100 \text{ GeV}} \frac{10^3 \text{ eV}}{m_{\nu_R}} \frac{10^{-3} \text{ eV}}{\tilde{m}}$$

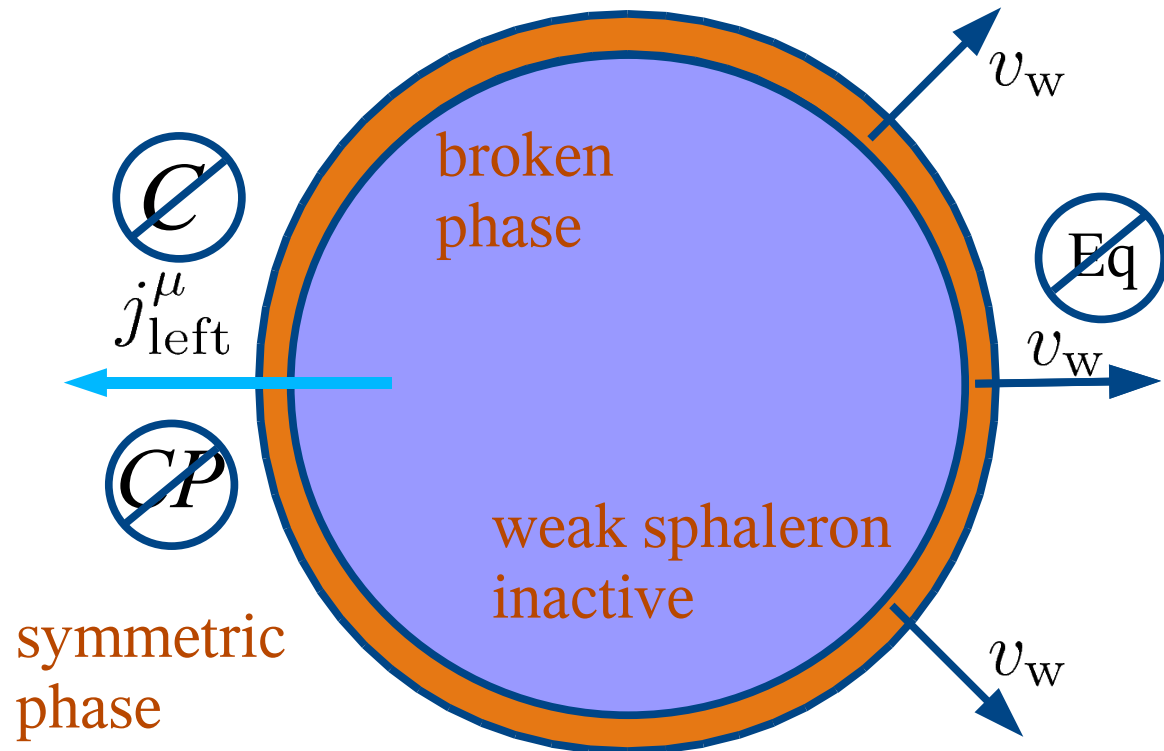
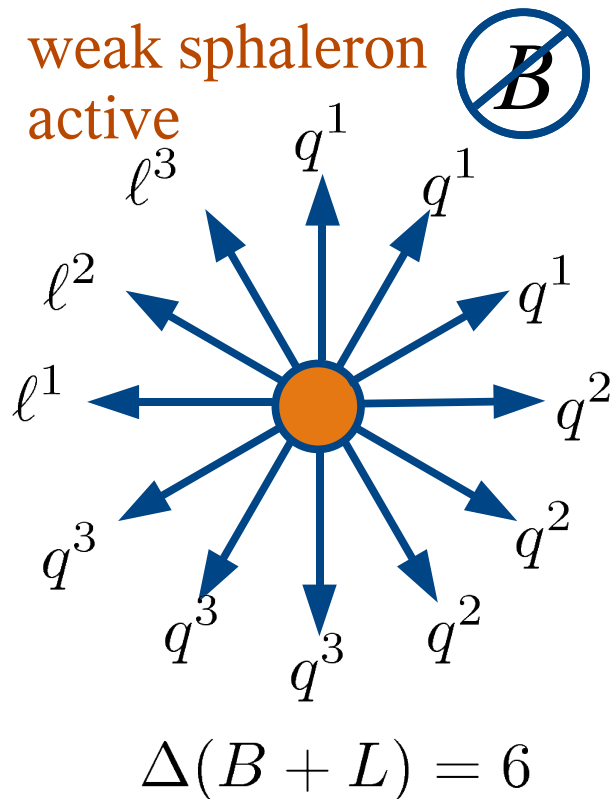
# Electroweak Baryogenesis

[Kuzmin, Rubakov, Shaposhnikov (1985); McLerran, Shaposhnikov, Turok, Voloshin (1990), Cohen, Kaplan, Nelson (1990)]

strong first order phase transition  
not with Standard Model



weak sphaleron active



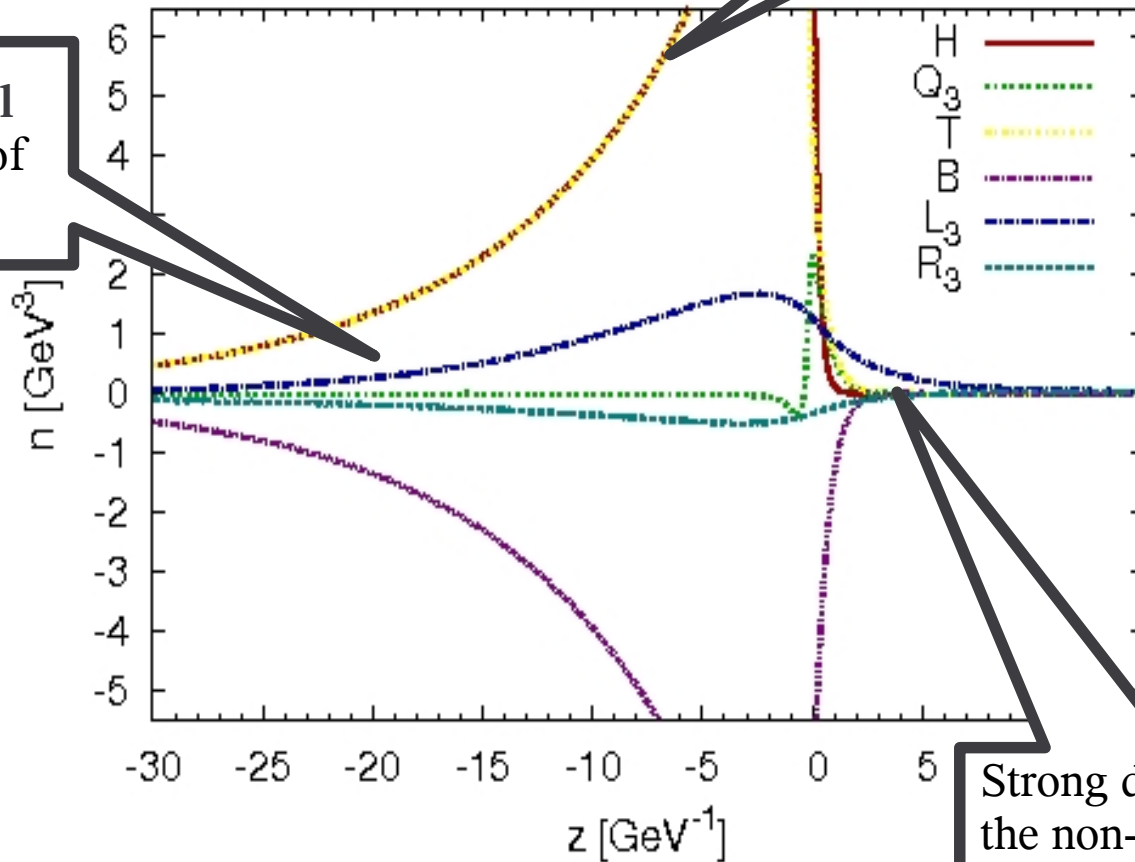
# Diffusion at the Bubble Wall

DJH Chung, BG, MJ Ramsey-Musolf, S Tulin,  
PRL.(2009),[arXiv:0808.1144 [hep-ph]]

Add particle and sparticle densities:

$$X = x + \tilde{x}$$

Diffusion & chemical  
equilibration ahead of  
the wall



Generation of Higgsino  
density from  $CP$ -  
violating local (?)  
source

Strong damping due to  
the non-zero Higgs  
VEV in the broken  
phase

$v_w$

# Sparticle Masses and Yukawa Couplings

Particle-sparticle equilibrium:

Introduce common charge densities and mass-dependent  $k$ -factors, e.g.

$$Q = q + \tilde{q} = t_L + b_L + \tilde{t}_L + \tilde{b}_L \quad k_Q = k_q + k_{\tilde{q}} \quad Q = \frac{T^2}{6} k_Q \mu_Q$$

$$T = t_R + \tilde{t}_R \quad k_T = k_{t_R} + k_{\tilde{t}_R} \quad T = \frac{T^2}{6} k_T \mu_T$$

...

...

...

Yukawa & Triscalar  
Equilibrium:

$$\mu_Q + \mu_H = \mu_T$$

$$\mu_Q - \mu_H = \mu_B$$

$$\frac{T}{k_T} + \frac{B}{k_B} - 2\frac{Q}{k_Q} = 0$$

Local baryon number  
conservation (WS slow):

$$T + B + Q = 0$$

$$\text{if } m_{\tilde{t}} = m_{\tilde{b}} \Rightarrow k_T = k_B \\ \Rightarrow Q = 0$$

Conversely, the sign of the baryon asymmetry depends on whether  
 $m_{\tilde{b}} > m_{\tilde{t}}$  or  $m_{\tilde{t}} > m_{\tilde{b}}$

# Sparticle Mass and $\tan \beta$ Dependences

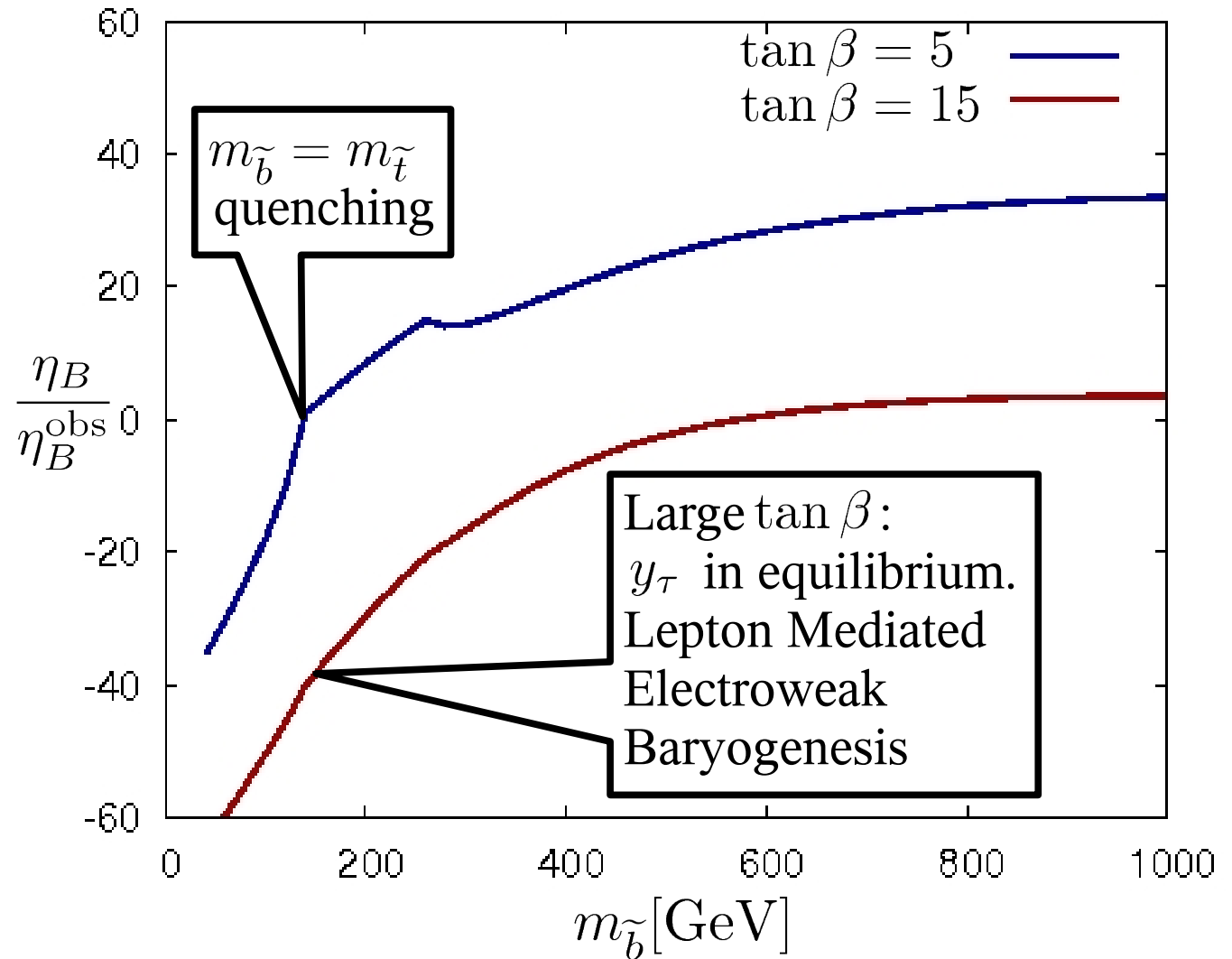
$$m_{\tilde{t}} = 102 \text{ GeV}$$

$$m_{\tilde{\tau}} = 104 \text{ GeV}$$

$$\mu = M_2 = 200 \text{ GeV}$$

All other sparticles heavy.

Assume fixed CP-violating source.



# Summary

## Leptogenesis:

- New scenario for leptogenesis with right handed neutrinos of TeV scale or below.
- Works within a conservative model:  $\nu$ MSSM.
- New connections of leptogenesis to phenomenology: EDMs, probes of Higgs sector.

## Electroweak Baryogenesis:

- Improved treatment of diffusion for electroweak baryogenesis.
- Magnitude and sign of baryon asymmetry are completely dependent on sparticle mass spectrum.
- Strong sensitivity also on size of Yukawa couplings and hence  $\tan \beta$ .