

# Properties of the Higgs boson

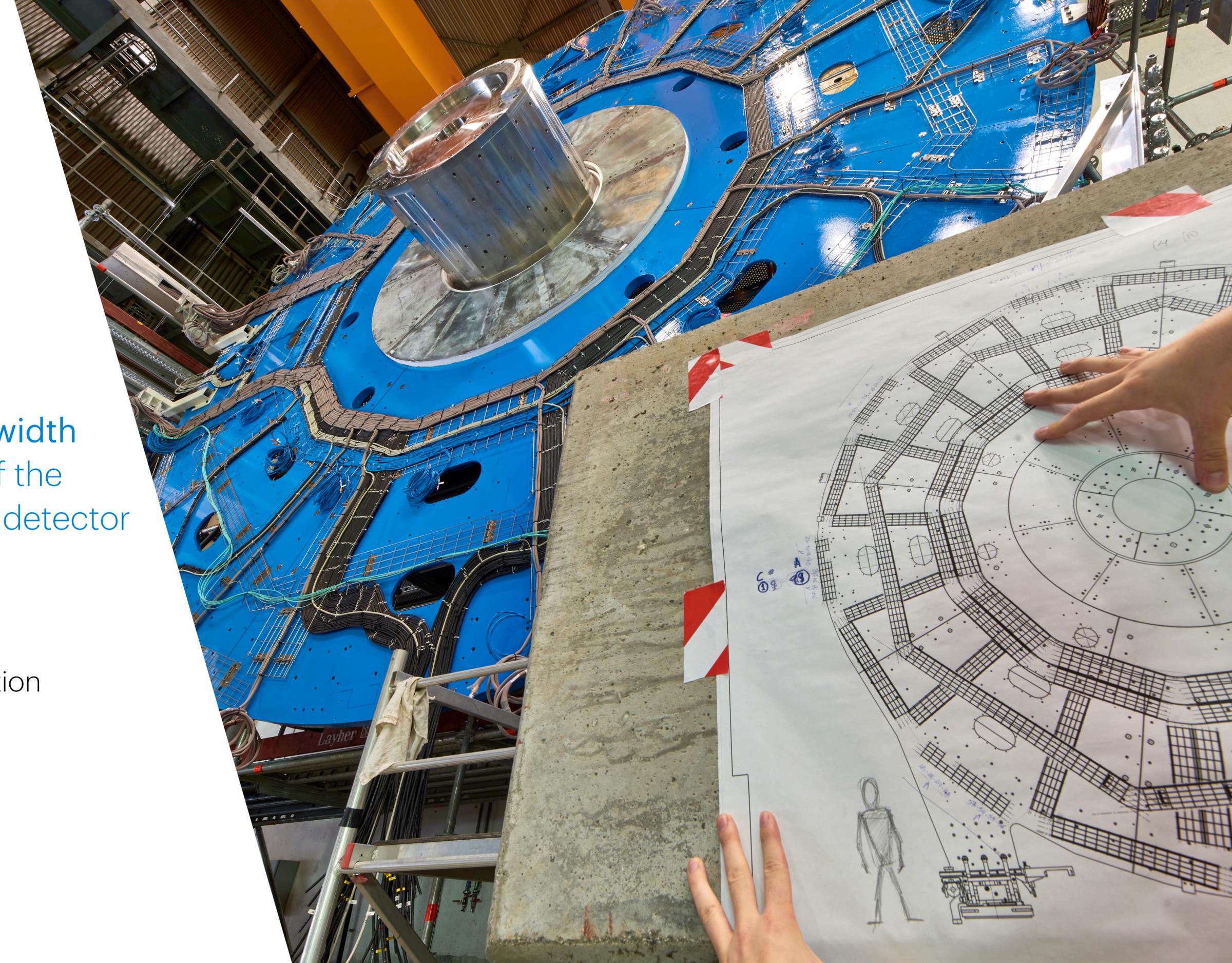
Measurements of the mass, width and coupling CP structure of the Higgs-boson with the ATLAS detector

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On behalf of the ATLAS collaboration

CIPANP 2022

Hilton Orlando - Lake Buena Vista



04 July 2022: CERN celebrated its 10th anniversary of the Higgs boson discovery announcement.

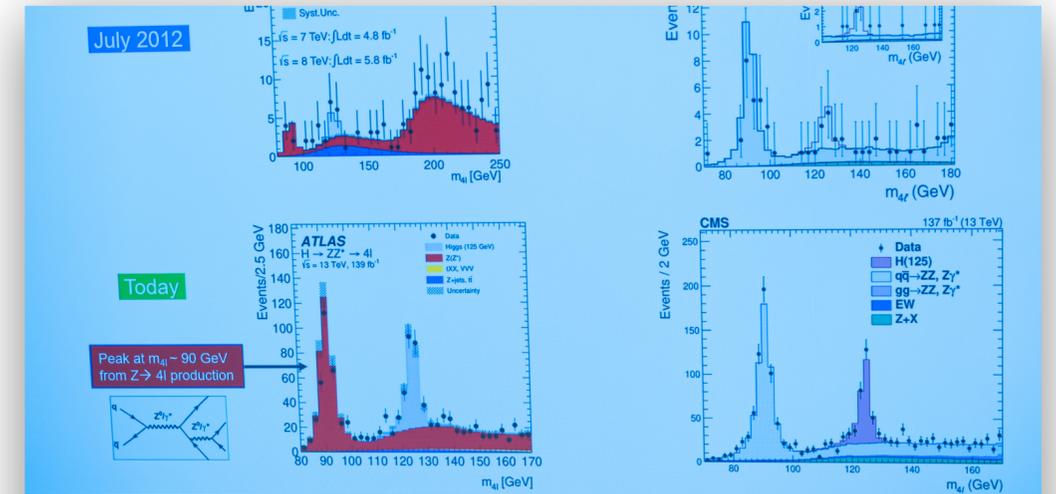


From J. Ellis et. al. 1976 [[paper](#)] on the phenomenological profile of the Higgs boson:

small. For these reasons we do not want to encourage big experimental searches for the Higgs boson, but we do feel that people performing experiments vulnerable to the Higgs boson should know how it may turn up.

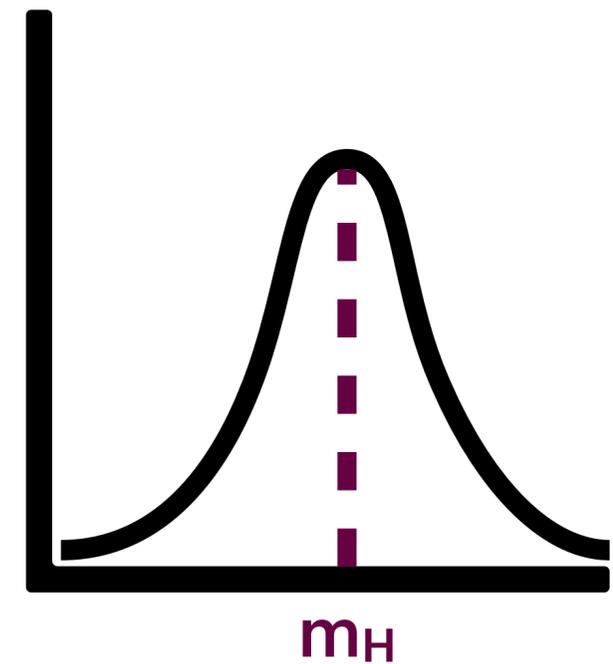
Today: Overview of current Higgs property measurements:

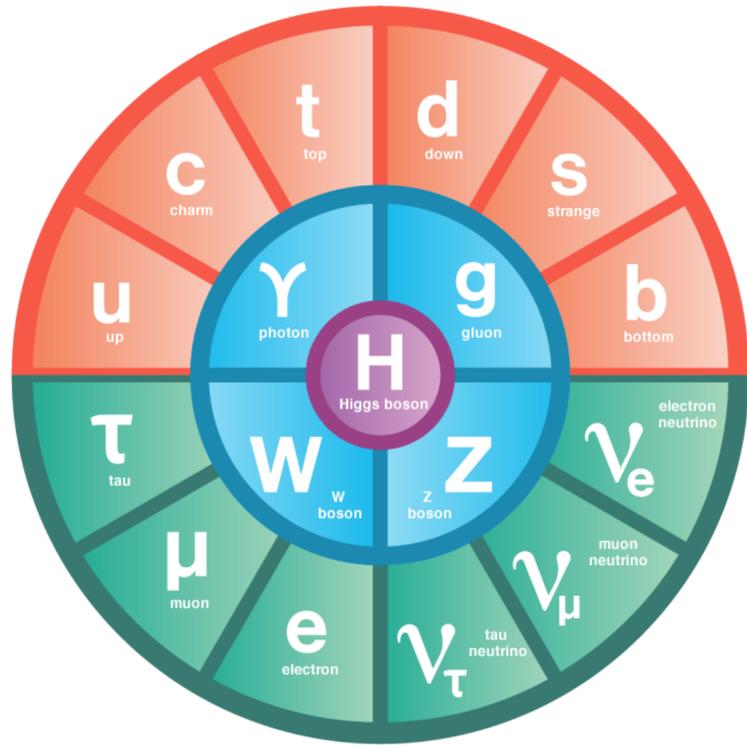
1. Mass of the Higgs boson
2. Width of the Higgs
3. Spin and Charge-Parity (CP) state of the Higgs



MASS

1.





Does the Higgs behave as expected by the SM?

## Why measure the Higgs mass?

Not predicted by the theory!

- Higgs couplings to SM particles are completely defined by the mass.
- **Shape of the Higgs potential fully known in the theory**, once the  $v_{ev}$  and  $m_H$  have been determined.

$$\mathcal{L} \subset V(\phi) = V_0 + \frac{1}{2} m_H^2 H^2 + \lambda v H^3 + \frac{1}{4} \lambda H^4$$

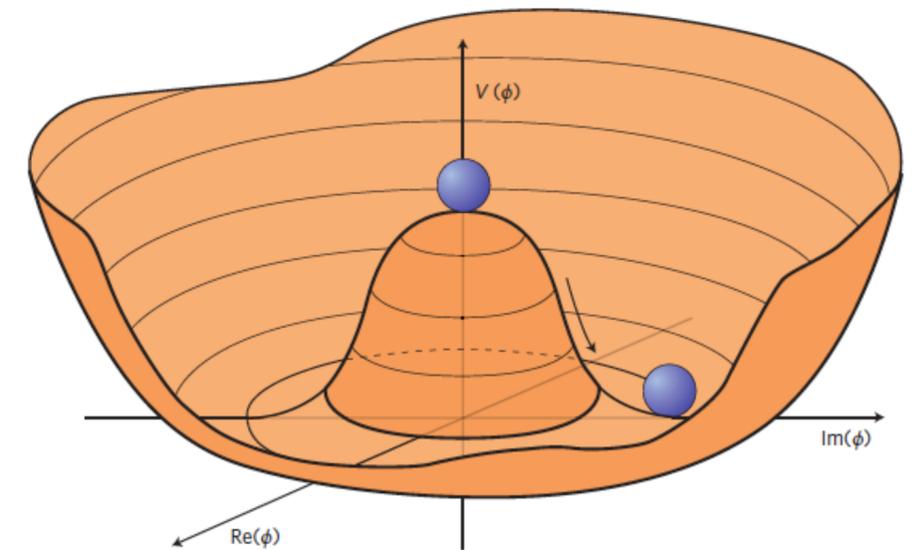
## How to measure the mass:

Precision:  $H \rightarrow ZZ^* \rightarrow 4\ell$  and  $H \rightarrow \gamma\gamma$ ,  $\ell = \text{electron } (e) \text{ or muon } (\mu)$

Despite low BR for each, profit from clean detector signature.

Precision greatly impacted by both statistics and experimental systematics.

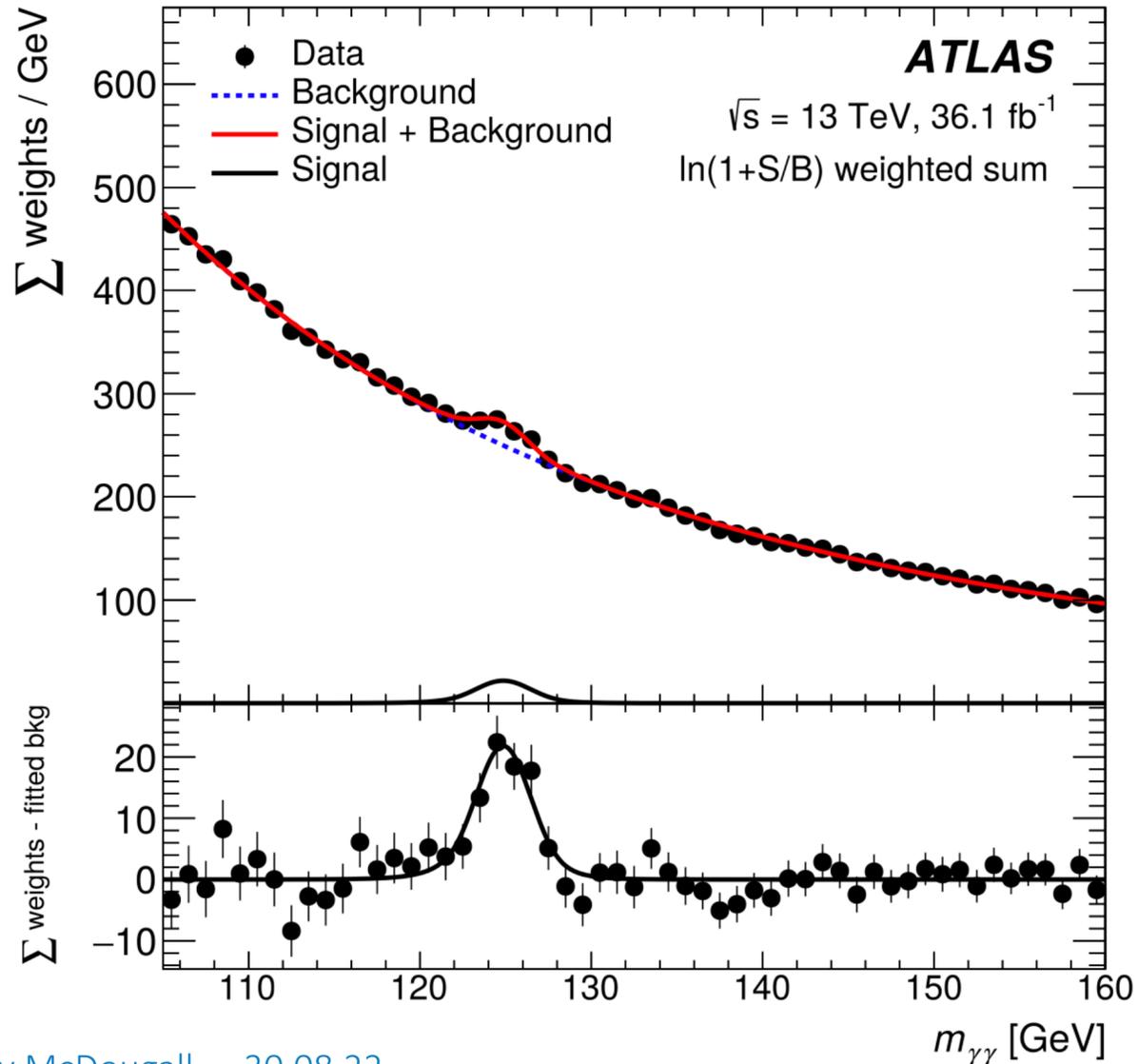
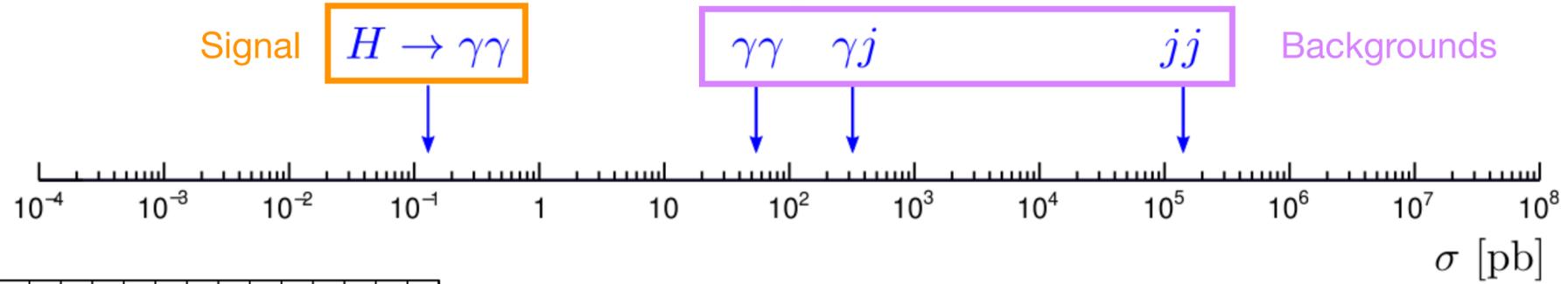
-> Require extremely good calibration of  $\gamma$  energy scale,  $\gamma$  energy resolution,  $e$  energy scale,  $\mu$  momentum scale, etc.



$\gamma\gamma$  decay channel: extremely low cross-section, but very clean signature.

Phys. Lett. B 784 (2018) 345

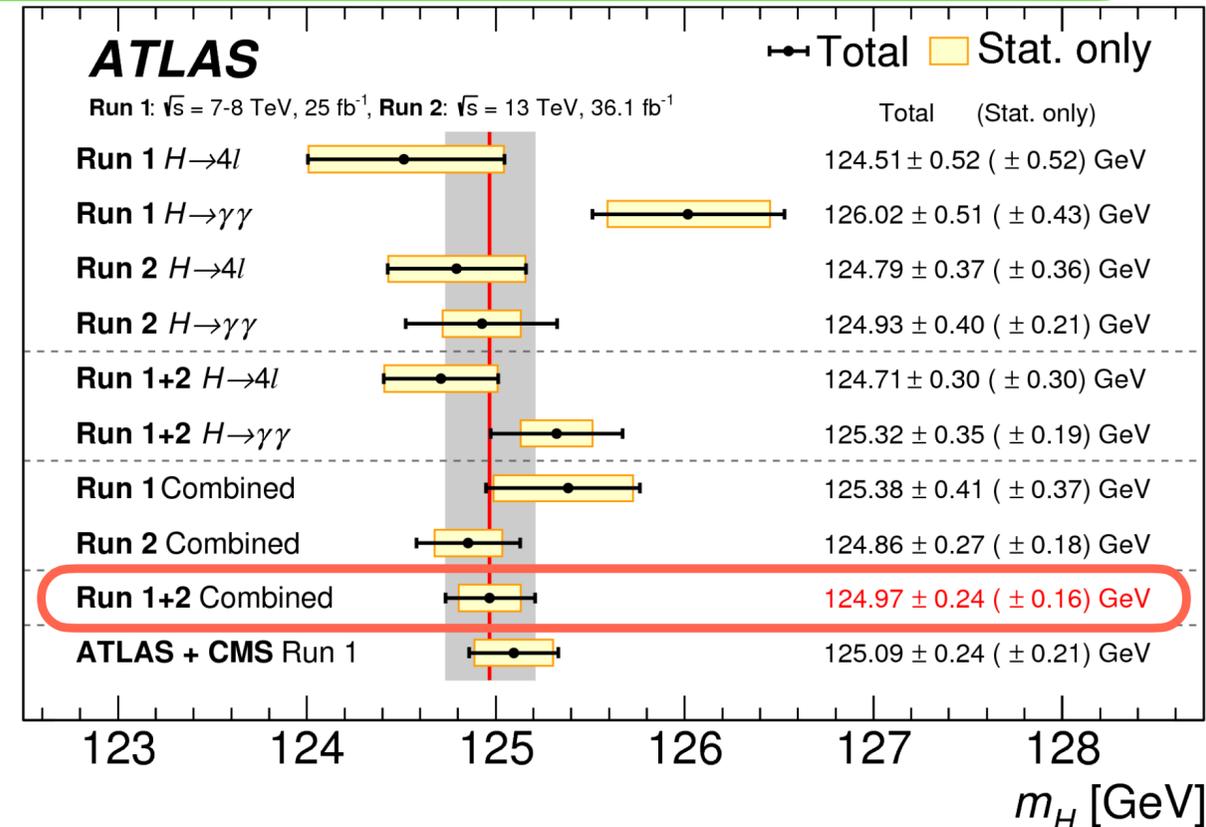
Dominant backgrounds: SM  $\gamma\gamma$  production,  $\gamma$ +jet, di-jet.



$H \rightarrow \gamma\gamma$  ( $36.1 \text{ fb}^{-1}$ ):

$$m_H^{\gamma\gamma} = 124.93 \pm 0.21 \text{ (stat)} \pm 0.40 \text{ (sys)} \text{ GeV}$$

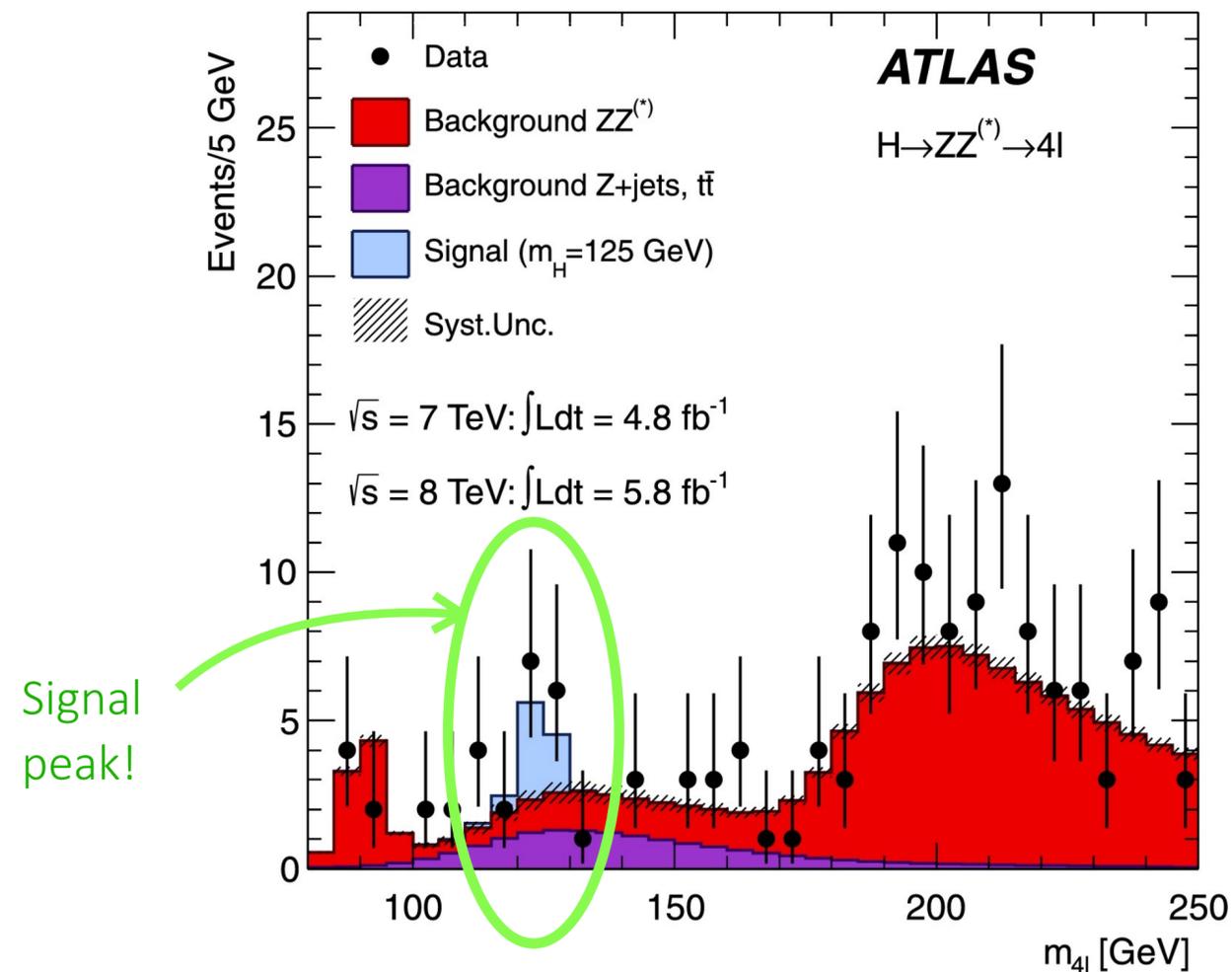
- Uncertainty dominated by experimental systematics.
- Two leading sources of experimental systematics are related to calibrating the EM calorimeter:  $\mathcal{O}(250) \text{ MeV}$



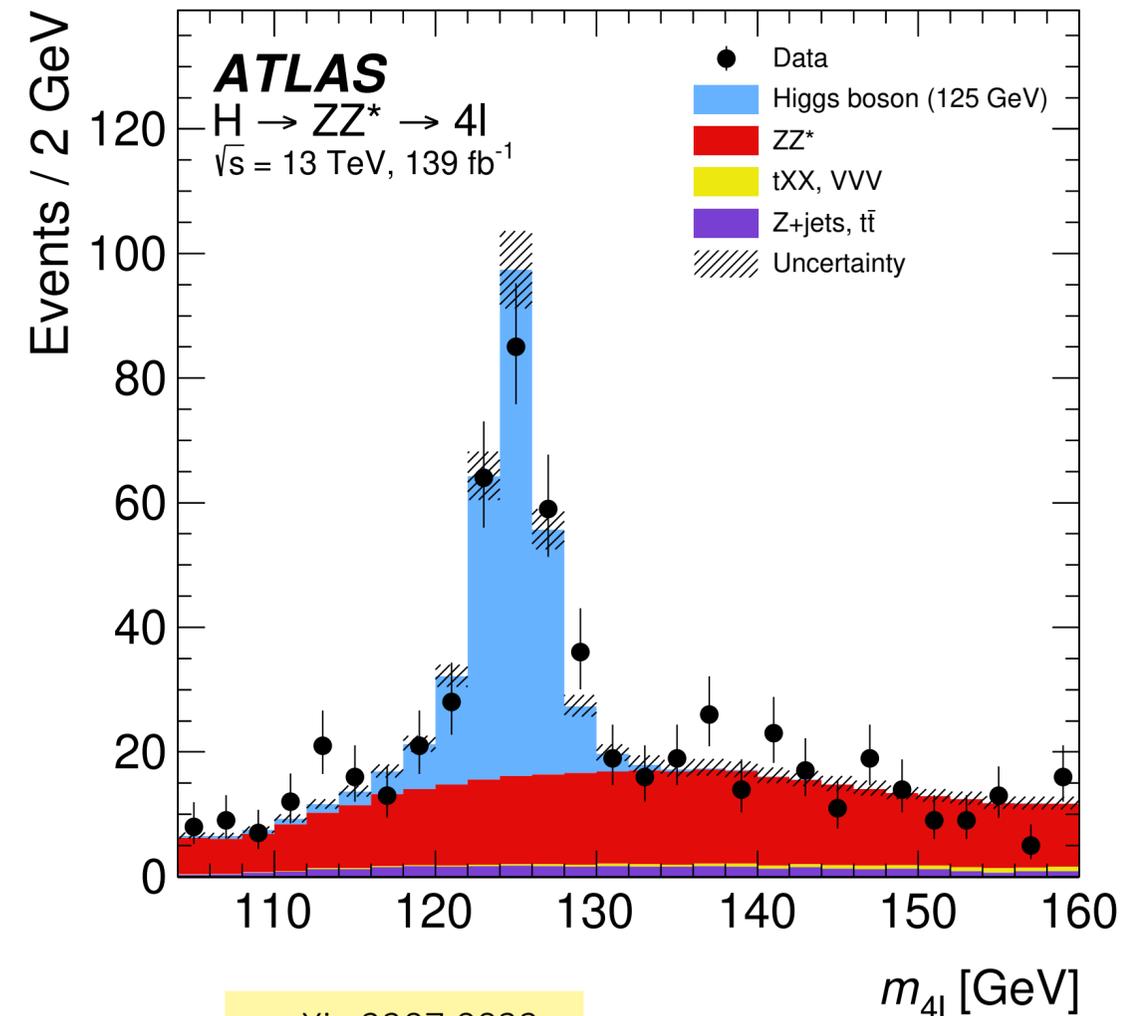
# Evolution of mass measurement:

Measurement of the Higgs Boson mass has come a long way since the 2012 Higgs discovery!

- ✓ More statistics (30x)
- ✓ Better understanding of the dominant backgrounds.
- ✓ Significant reduction in systematic uncertainties.
- ✓ Improved analysis techniques.



Physics Letters B 716 1 (2012) 1-29



arXiv:2207.0032



# Latest $H \rightarrow ZZ^* \rightarrow 4\ell$ mass result:

arXiv:2207.0032

In this channel: larger dataset + improved experimental techniques + more precise lepton calibration.  
Improvements since  $36 \text{ fb}^{-1}$  analysis:

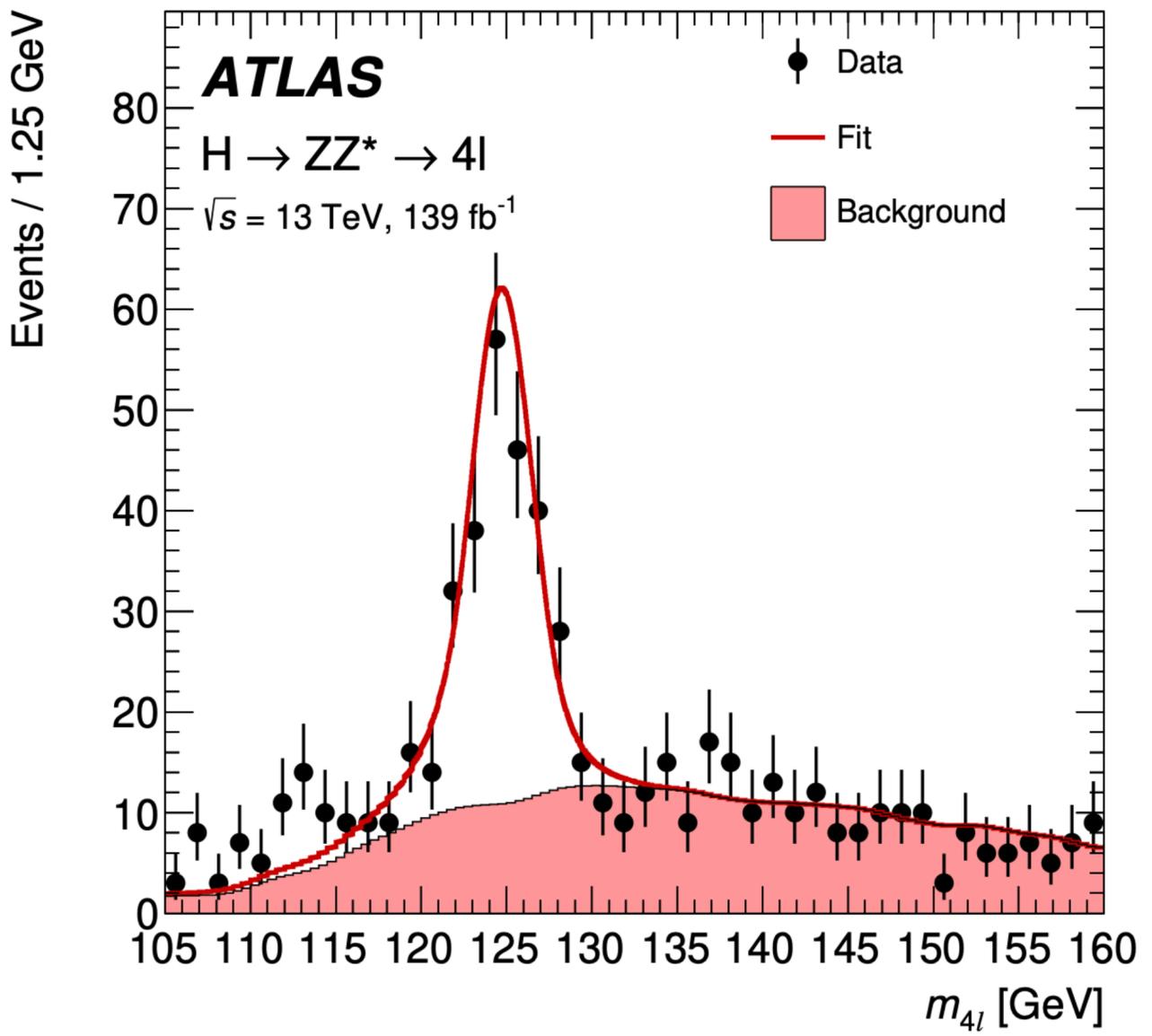
- ⇒ Factor 2 reduction in statistical uncertainty
- ⇒ ~ 20% reduction on systematic uncertainty relative to previous Run 2 publication

Size of leading systematics:

Systematic Uncertainty	Contribution [MeV]
Muon momentum scale	±28
Electron energy scale	±19
Signal-process theory	±14

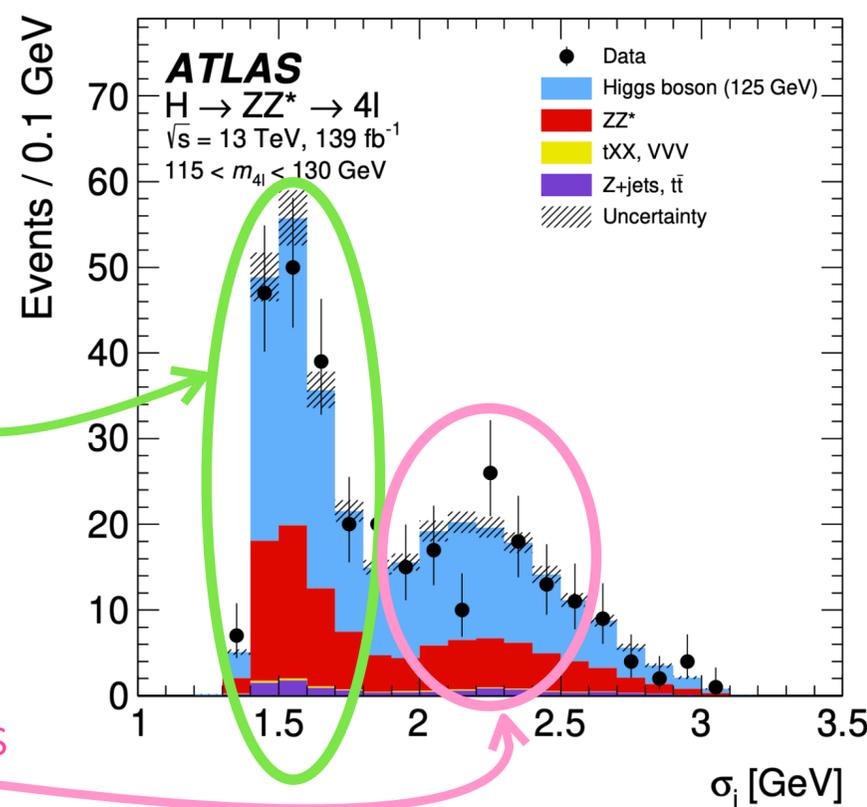
Analysis improvements:

1. New deep neural network (DNN) for S/B discrimination
2. Improved  $p_T^\mu$  calibration.
3. Per event,  $\sigma_i$  resolution estimate!



$\sigma_i \sim 1.5 \text{ GeV}$  for  $4\mu$  and  $2\mu 2e$  events

$\sigma_i \sim 2.1 \text{ GeV}$  for  $2e 2\mu$  and  $4e$  events



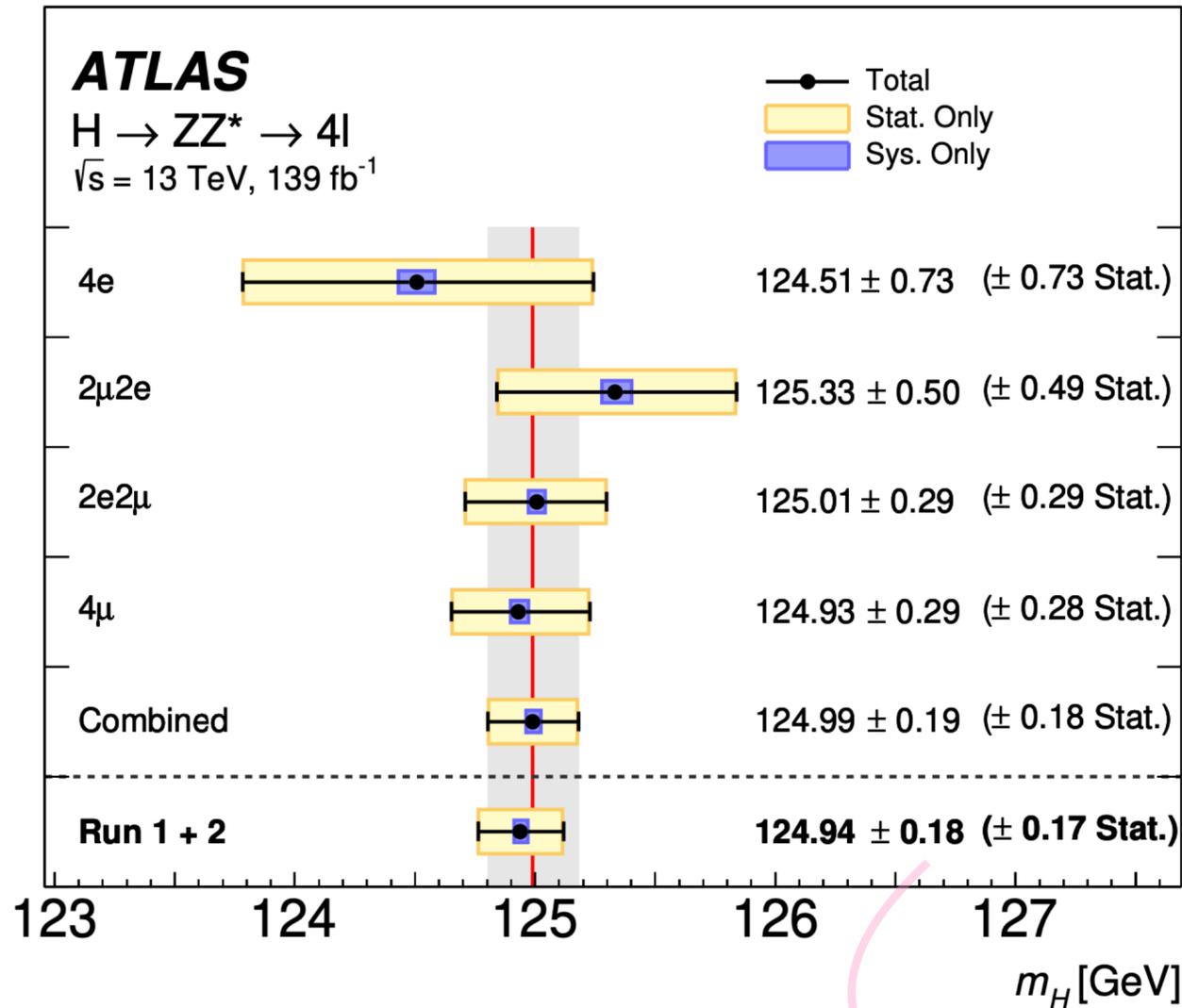
# Latest $H \rightarrow ZZ^* \rightarrow 4\ell$ mass result:

arXiv:2207.0032

Simultaneous fit for all decay channels.

*Strongly* stat error limited

~ 88% improvement compared to  $m_H^{H \rightarrow ZZ}$ , Run 1



$H \rightarrow ZZ^* \rightarrow 4\ell$  full Run 2 dataset ( $139 \text{ fb}^{-1}$ ):

$$m_H^{ZZ} = 124.99 \pm 0.18 \text{ (stat)} \pm 0.04 \text{ (sys) GeV}$$

$H \rightarrow ZZ^* \rightarrow 4\ell$  combination of Run 1 and Run 2 datasets:

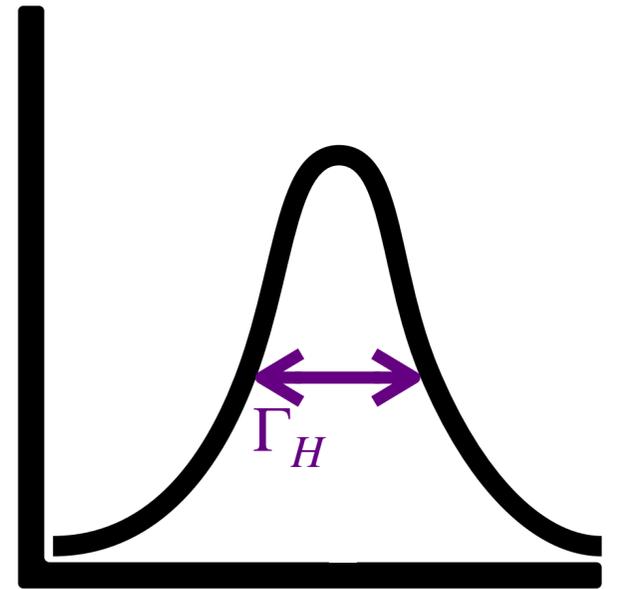
$$m_H^{ZZ} = 124.94 \pm 0.17 \text{ (stat)} \pm 0.03 \text{ (sys) GeV}$$

This  $m_H$  result is one of **the most precise** measurements of the LHC scientific program!

0.14% uncertainty on  $m_H$ !

WIDTH

2.



Contains information about the interactions of the Higgs with all other fundamental particles.

Since width value depends on Higgs couplings to SM particles-> very important to measure it precisely!

**If theory  $\neq$  measurement  $\Rightarrow$  new particles that also couple to the Higgs? New physics?**

The SM predicts the total Higgs width to be:

$$\Gamma_H^{SM} = 4.07 \text{ MeV}$$

**Direct measurement:**

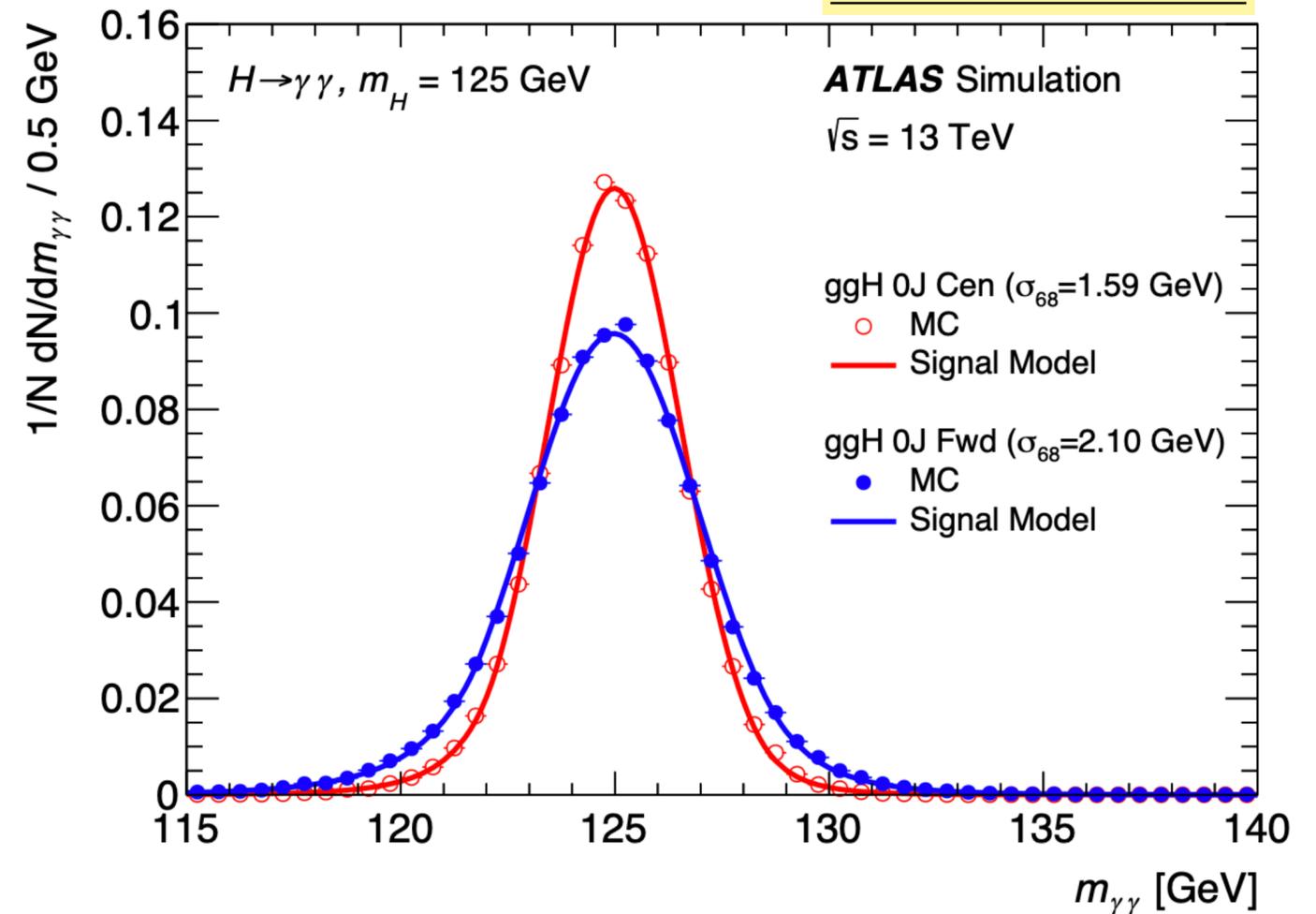
Technique: Scan over invariant mass profile of the Higgs decay products, in peak on-shell region.

Feasibility: Possible for other particles with larger width - the Z boson ( $\Gamma_Z^{SM} = 2.09 \text{ GeV}$ ).

Limitation: Restricted by *experimental resolution*

$$\Gamma_H < 1 - 2 \text{ GeV.}$$

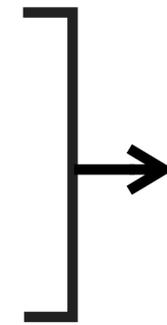
ATLAS-CONF-2020-026



Across  $m_H$  invariant mass distribution can define two regions of interest.

**On-shell:** region close to Higgs resonant peak, defined by  $s \sim m_H^2$

**Off-shell:** region far from peak of distribution, defined as  $s \gg m_H^2$



**Indirect measurement techniques:**

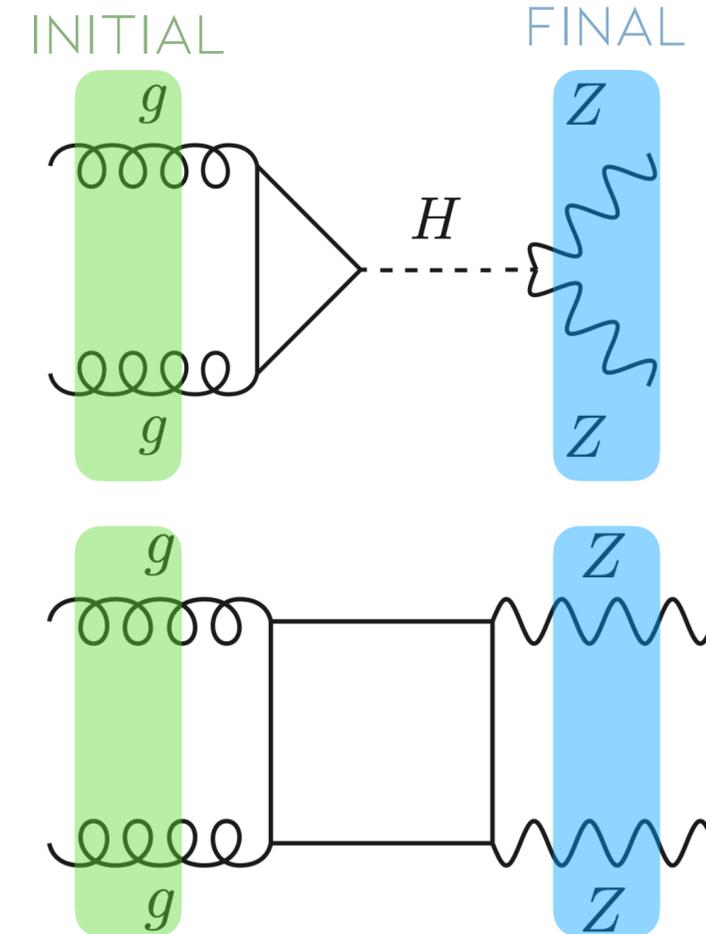
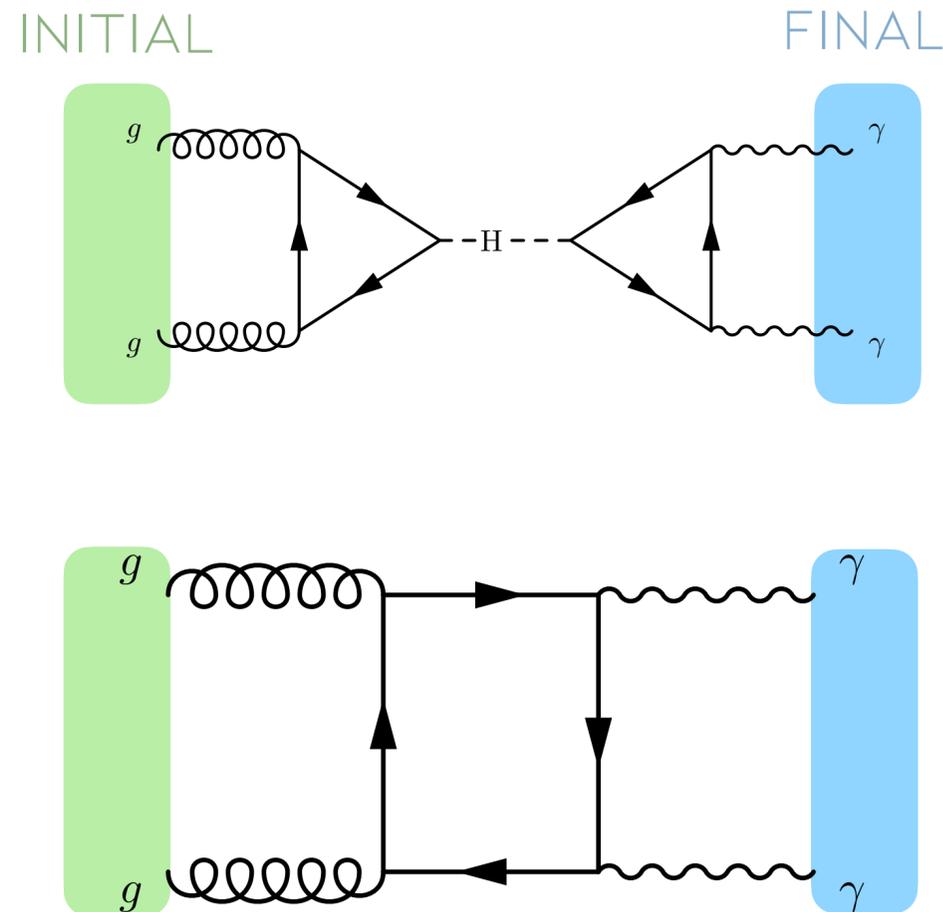
1. Interferometry (in on-shell region)
2. Off-shell method

Interference occurs between processes with the same initial and final state particles.

Interference impacts both the total cross-section, and the invariant mass line-shape.

SIGNAL

BACKGROUND



**Off-shell region:** [[Phys. Rev. D 90, 053003](#)],  
 [[arXiv:1307.4935v3](#)], [[Phys. Rev. D 89, 053011](#)]:

Combine on-shell and off-shell information.

- Particularly interesting in HVV (V = W or Z) decays - strongly enhanced at high energies => far from resonant mH peak.
- Off-shell Higgs production enhanced at ~ 350 GeV because of top-quark loops in ggF.
- ~15% of Higgs cross-section is for masses > 2 × m<sub>Z</sub> due to increased phase space of on-shell Z production.

Cross-section given the Higgs propagator:

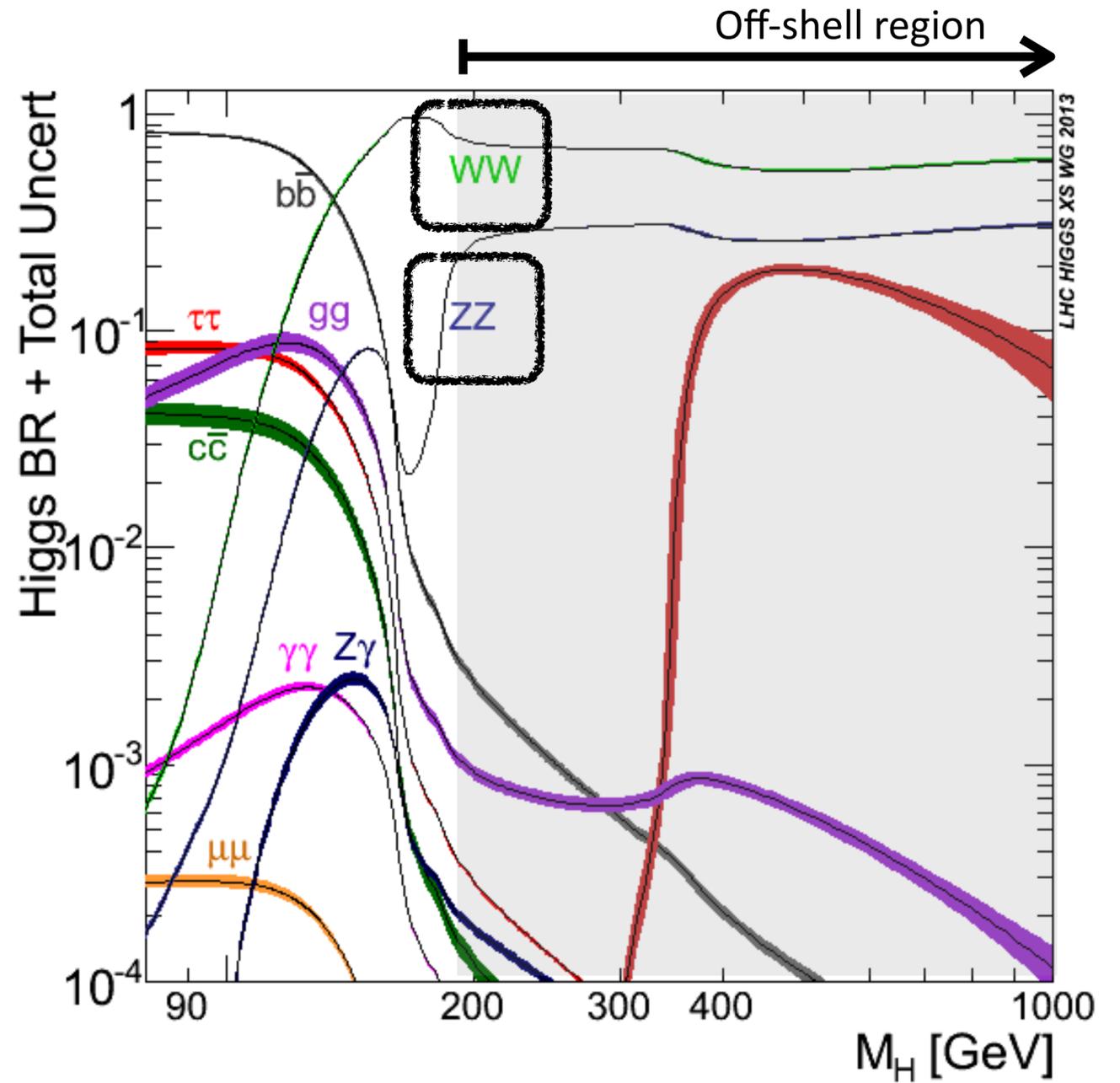
$$\sigma \sim \frac{g_{\text{prod}}^2 \times g_{\text{decay}}^2}{(s - m_H^2)^2 + m_H^2 \Gamma_H^2}$$

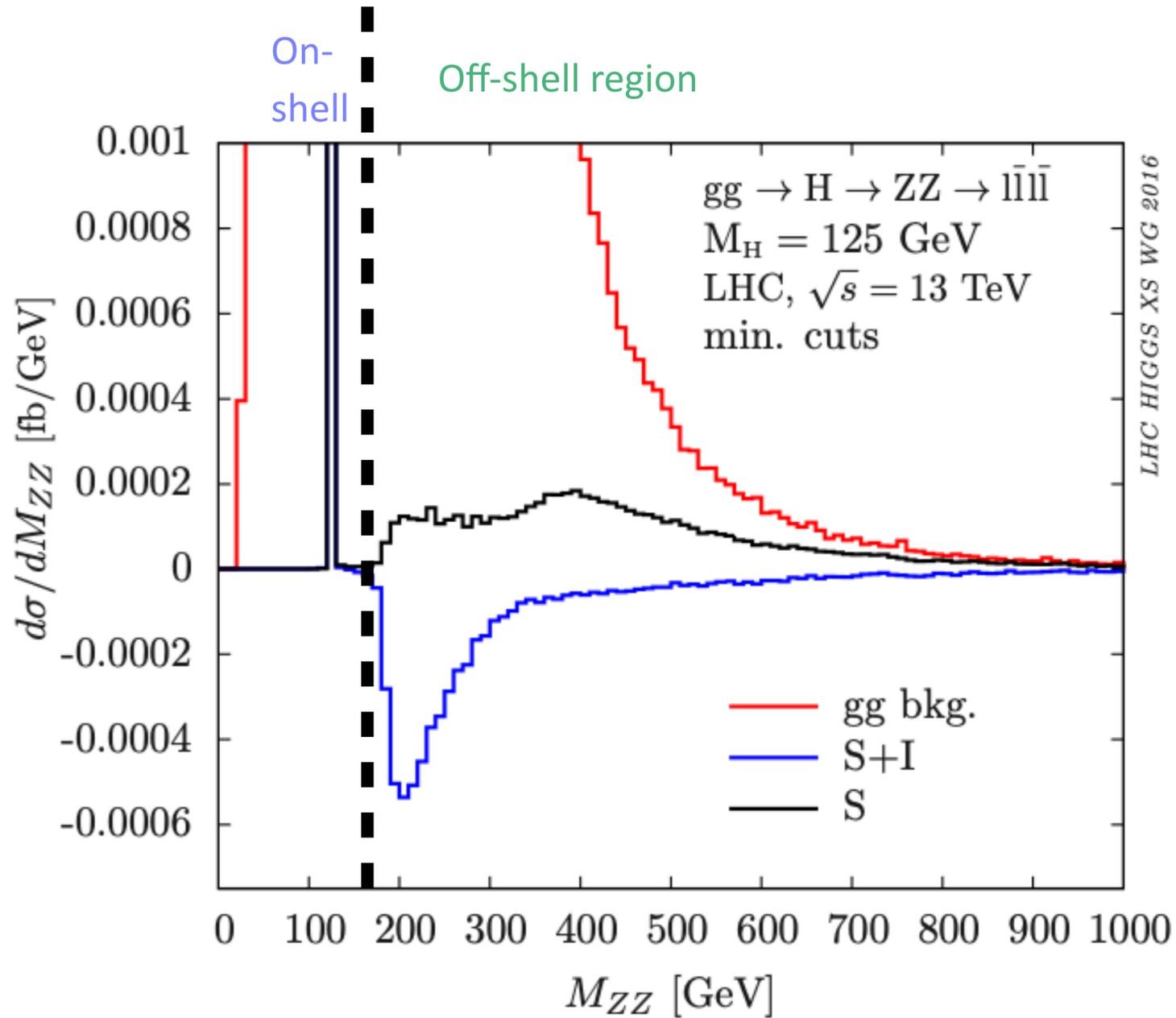
**On-shell** limit:  $s \sim m_H^2$

$$\sigma_{\text{on-shell}} \sim \frac{g_{\text{prod}}^2 \times g_{\text{decay}}^2}{\Gamma_H}$$

**Off-shell** limit  $s \gg m_H^2$ :

$$\sigma_{\text{off-shell}} \sim g_{\text{prod}}^2 \times g_{\text{decay}}^2$$





Interference effects = sizeable and destructive.

Determination of Higgs width,  $\Gamma_H$  :

take ratio between off-shell and on-shell event yields:

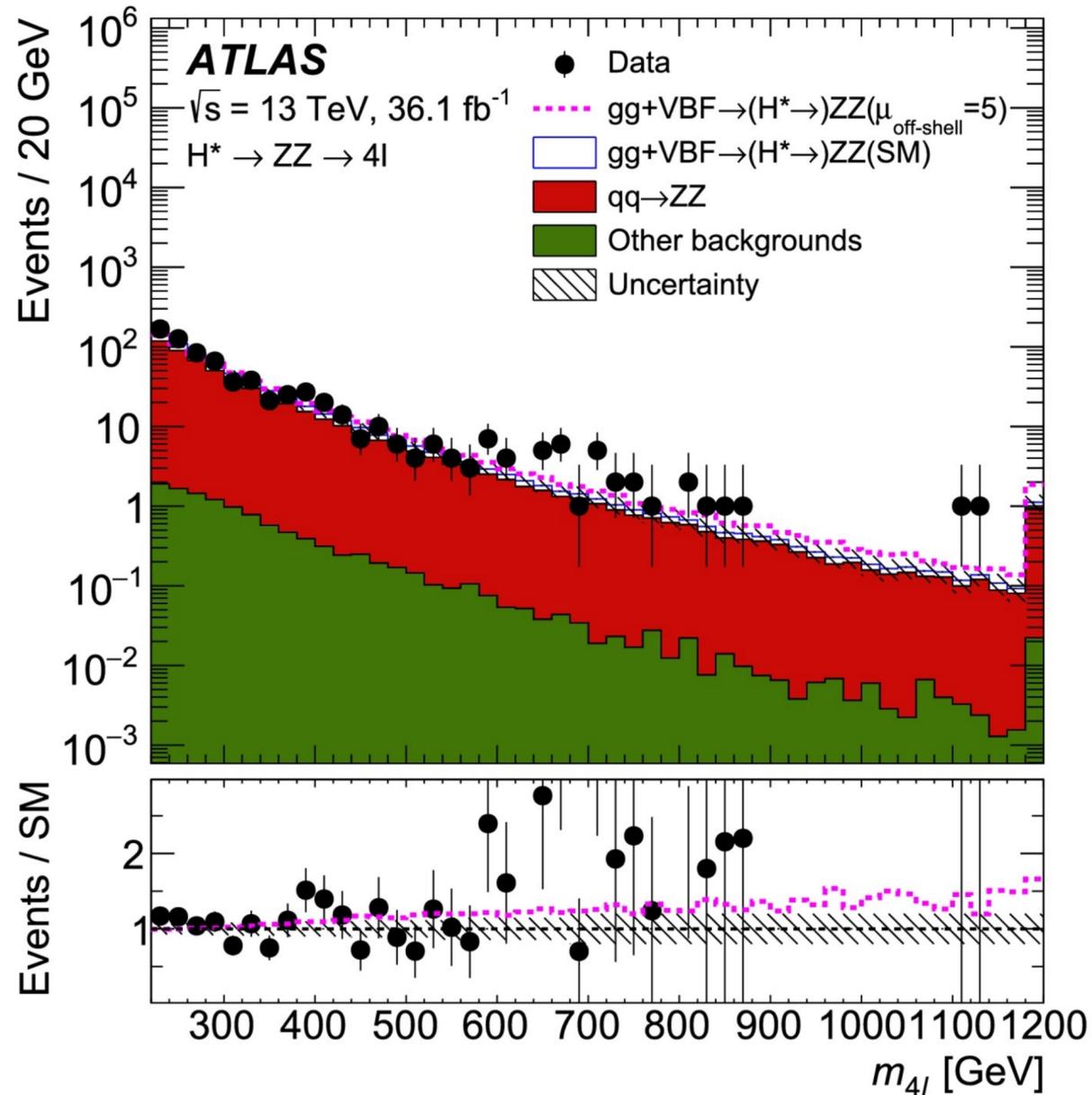
$$\frac{\mu_{\text{off-shell}}}{\mu_{\text{on-shell}}} = \frac{\Gamma_H}{\Gamma_H^{SM}}$$

**WARNING:** Method relies on the couplings cancellation.

- Assumes: on-shell and off-shell couplings are the same, despite very different energies.
- Theoretical model dependence not present in on-shell only methods.
- Several BSM effects could de-correlate on- and off-shell couplings, e.g. additional Higgs states [[PhysRevD.92.075038](https://arxiv.org/abs/1508.07503)].

Latest results use  $36 \text{ fb}^{-1}$  of data.

Measurements use  $H \rightarrow ZZ^* \rightarrow 4\ell$  and  $H \rightarrow ZZ^* \rightarrow 2\ell 2\nu$  decays.



Discriminating variables:

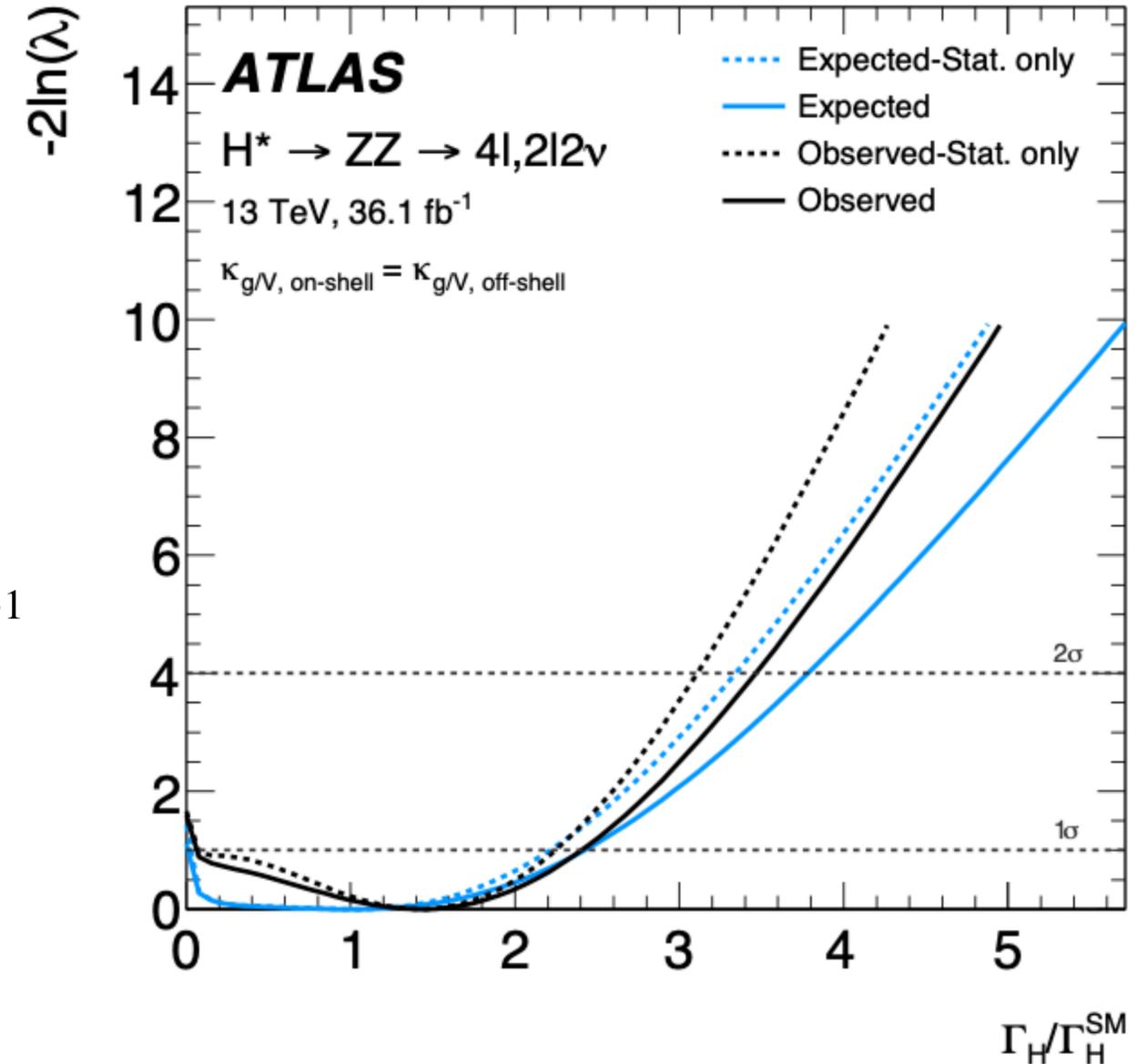
$4\ell$  channel:  $m_{4\ell}$  and  $D_{\text{ME}}$

$2\ell 2\nu$  channel:  $m_T^{ZZ}$

On-shell rates from  $36 \text{ fb}^{-1}$  analysis:

$$\mu_{\text{on-shell}} = 1.28^{0.21}_{0.19}$$

[[arXiv:1712.02304](https://arxiv.org/abs/1712.02304)]



Observed results for the Higgs boson total width:

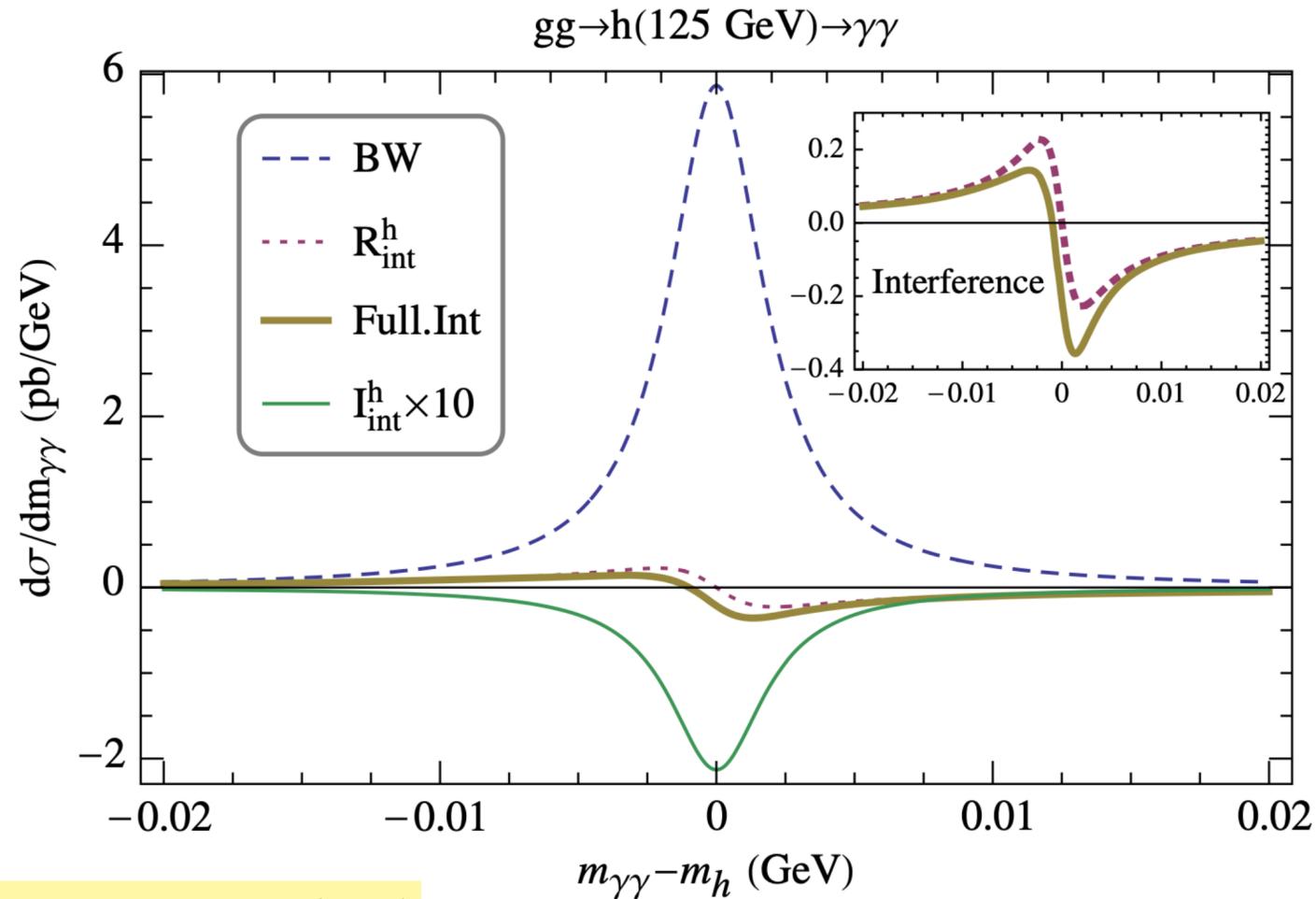
$$\Gamma_H < 14.4 \text{ MeV @ 95\% CL}$$

**Interferometry - on-shell region:** [[Phys. Rev. Lett. 111, 111802](#)], [[Phys. Rev. Lett. 119, 181801](#)], [[Phys. Rev. D 86, 073016](#)]

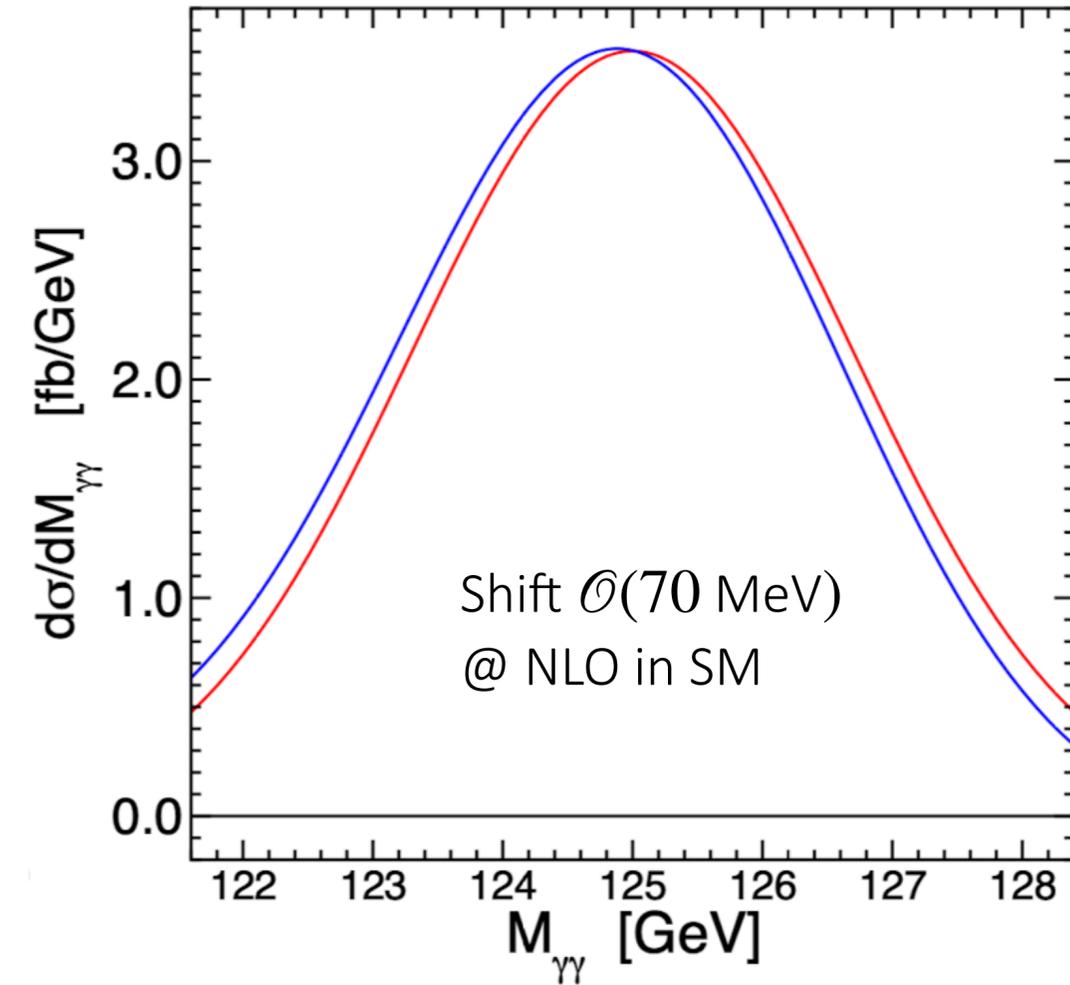
1. Apparent shift in mass distribution (due to antisymmetric shape effects)
2. Overall reduction in the total yield  $\mathcal{O}(2\%)$ .

- Measured mass shift is proportional to the Higgs boson width:  $\Delta m_H \propto \sqrt{\Gamma_H / \Gamma_H^{SM}}$
- No constraints made on the couplings.

arXiv:1208.1533



Phys. Rev. Lett. 119, 181801 (2017)



With the HL-LHC dataset, and subsequent theoretical improvements, expect bounds on the order of  $\Gamma_H \sim \mathcal{O}(10) \times \Gamma_H^{SM}$  [[CERN-LPCC-2018-04](#)].

3.

CP + SPIN



# The Higgs boson spin and CP quantum numbers:

The SM predicts the Higgs to be spin 0 (only SM particle), and have positive parity (CP even):  $J^{CP} = 0^{++}$

## Measuring the spin:

Run I: Evidence for spin 0 nature of the Higgs.

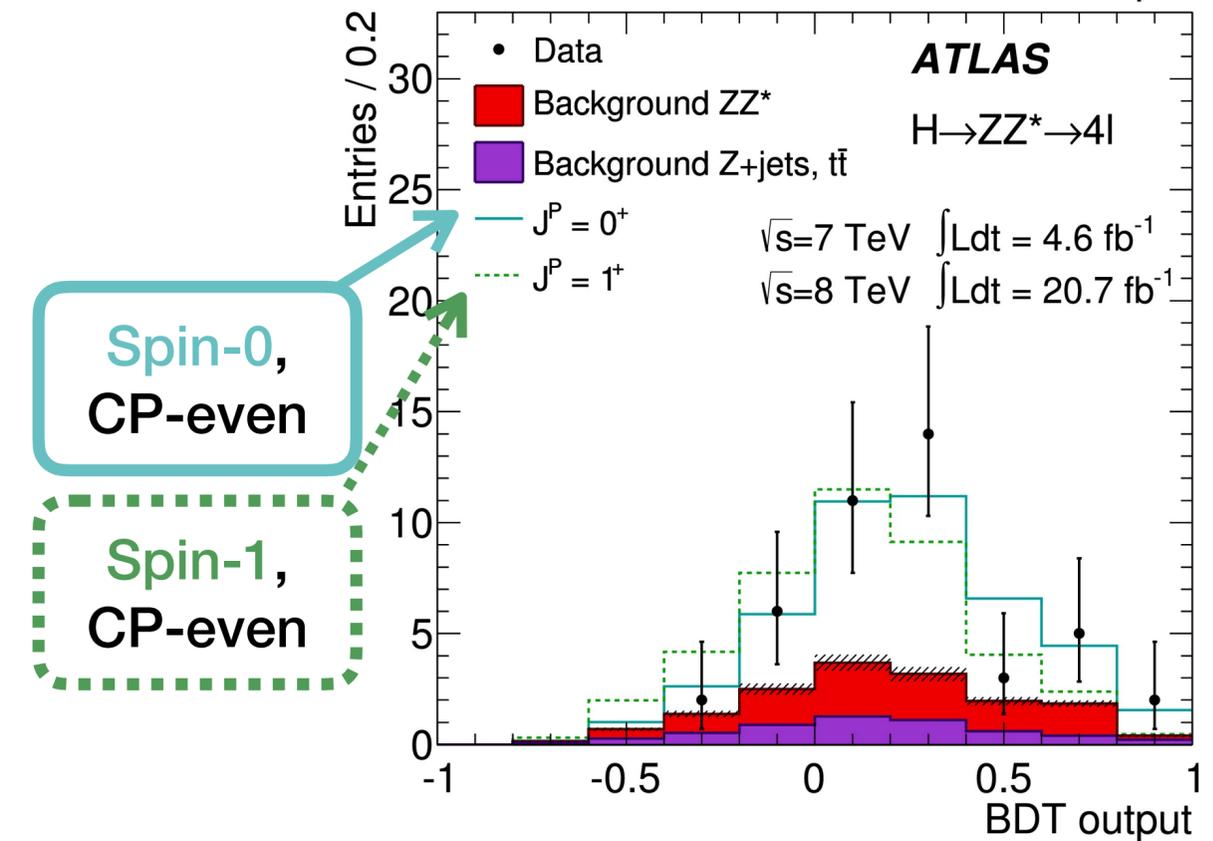
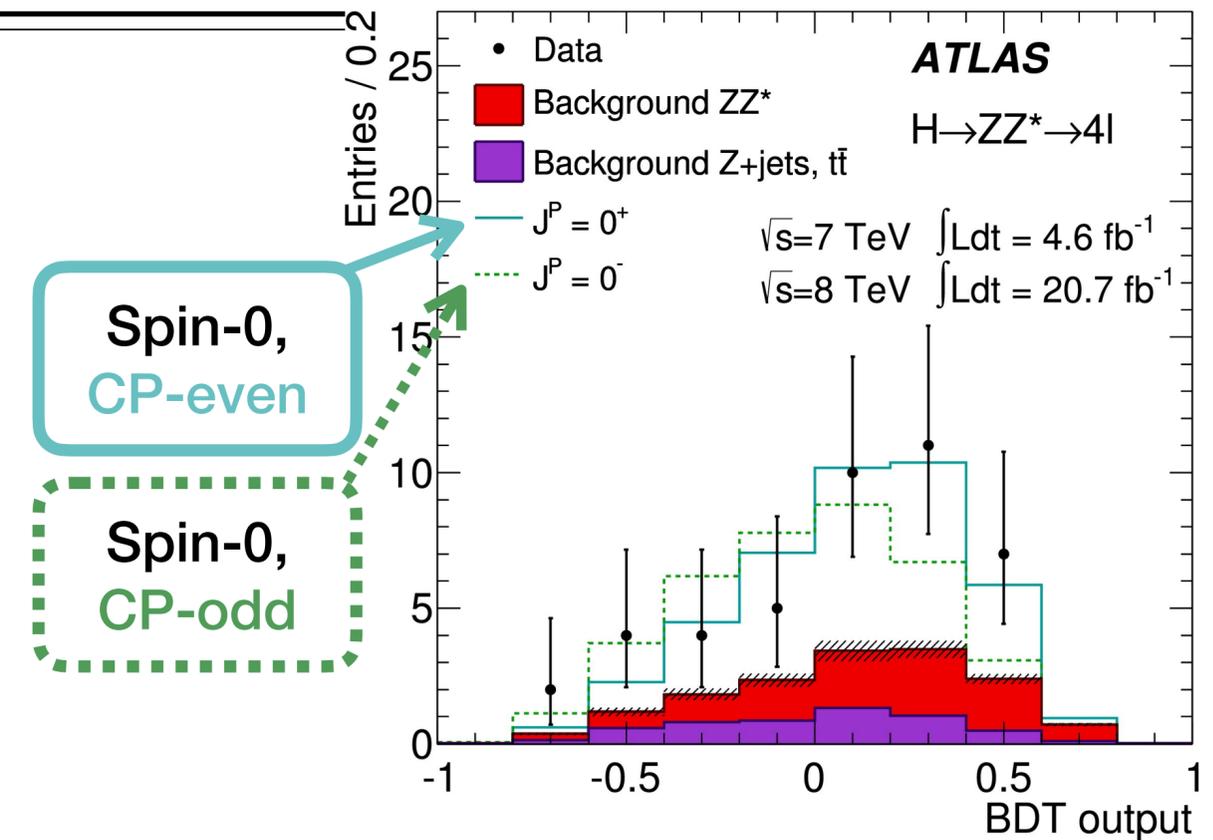
Spin 1 and 2 hypotheses excluded at  $> 99.9\%$  CL using  $H \rightarrow \gamma\gamma$ ,  $H \rightarrow WW^*$ ,  $H \rightarrow ZZ^*$ .

## Higgs purely CP-even? Possibility of CP-violation in Higgs sector?

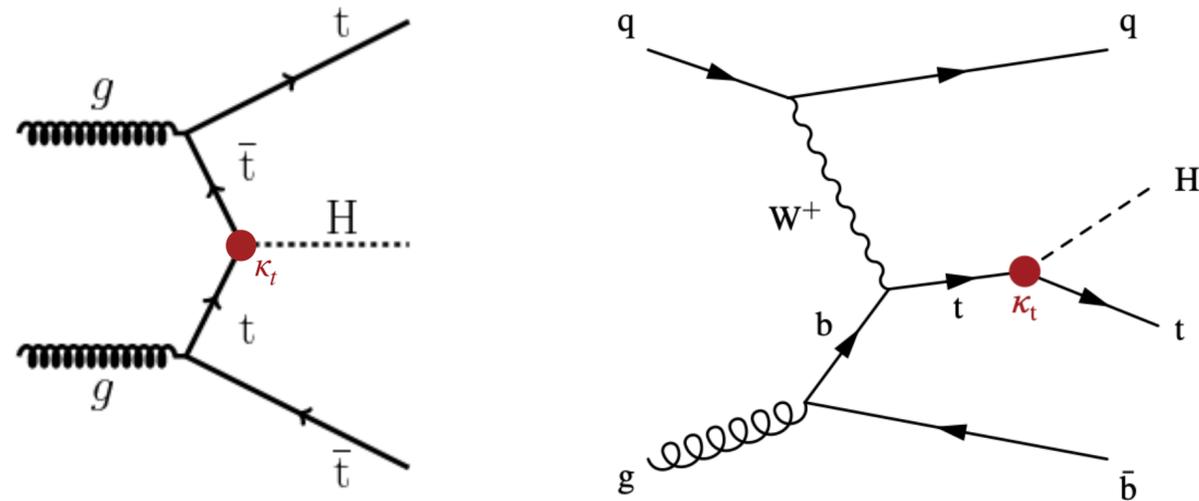
Presence of a CP-odd (pseudo-scalar) admixture has not yet been excluded.

Measured CP-odd contribution = new physics

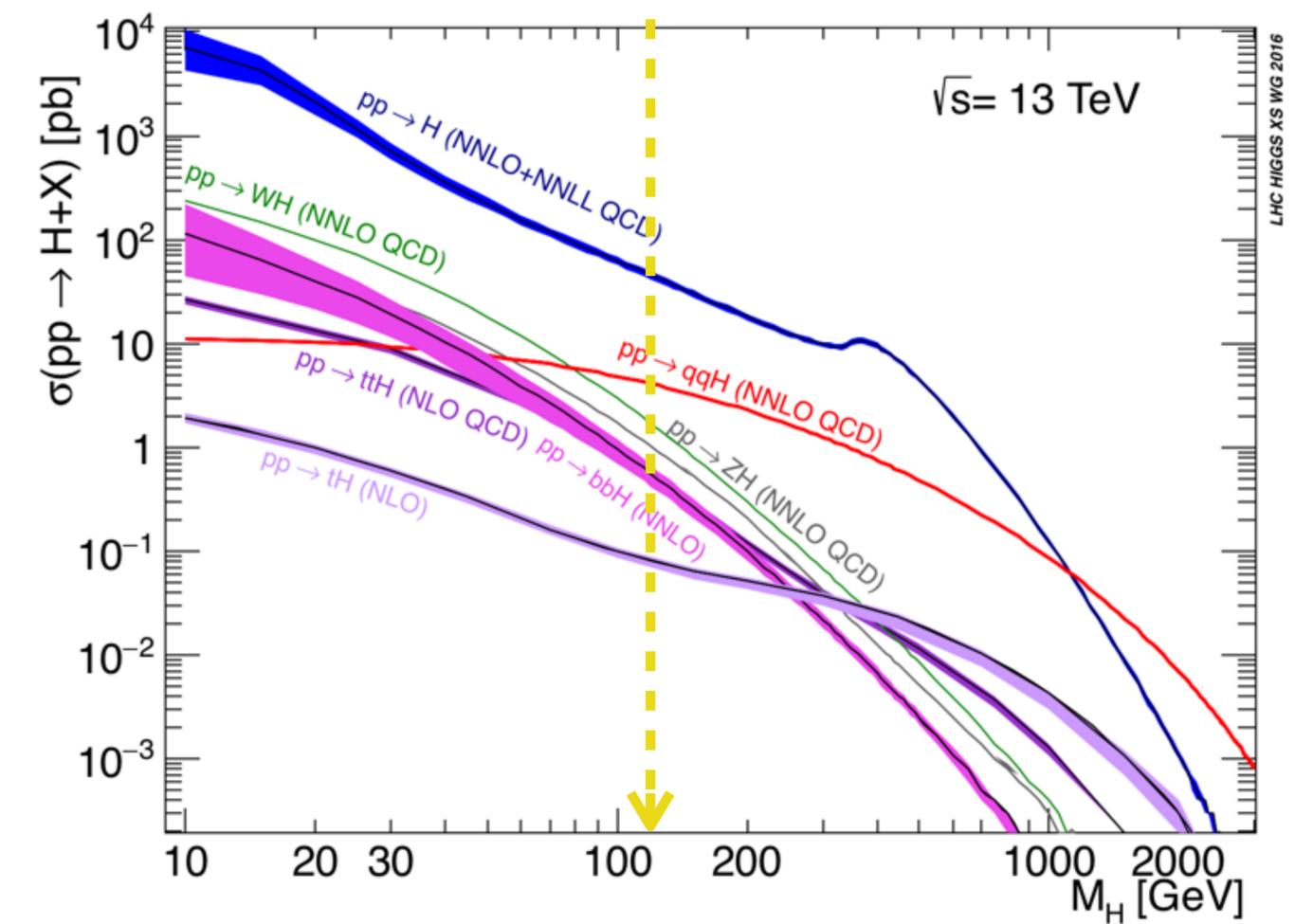
$\Rightarrow$  contribute to the observed baryon asymmetry of the universe.



CP properties of the top-Higgs Yukawa coupling studied directly using Higgs Boson production in association with top quarks, e.g. top quark pair,  $t\bar{t}H$  or single top  $tH$  processes.



$m_H = 125$  GeV



$$\mathcal{L}_{eff} = -\frac{m_t}{v} \{ \bar{\psi}_t \kappa_t [ \cos(\alpha) + i \sin(\alpha) \gamma_5 ] \psi_t \} H$$

$\kappa_t (> 0)$  is the top Yukawa coupling parameter (=1) for SM

$\alpha$  is the CP mixing angle  
SM: Pure CP-even,  $\alpha=0^\circ$

Pure CP-odd  
 $\alpha = \pm 90^\circ$

Any other  $\alpha$  value  $\rightarrow$  new source of CP violation in SM!

Latest ATLAS results, using  $139 \text{ fb}^{-1}$  collected at 13 TeV consider  $H \rightarrow \gamma\gamma$  and  $H \rightarrow b\bar{b}$ .

Measured rates:

$$\sigma_{t\bar{t}H} \times B_{\gamma\gamma} = 1.64^{+0.38}_{-0.36} \text{ (stat)}^{+0.17}_{-0.14} \text{ (sys) fb}$$

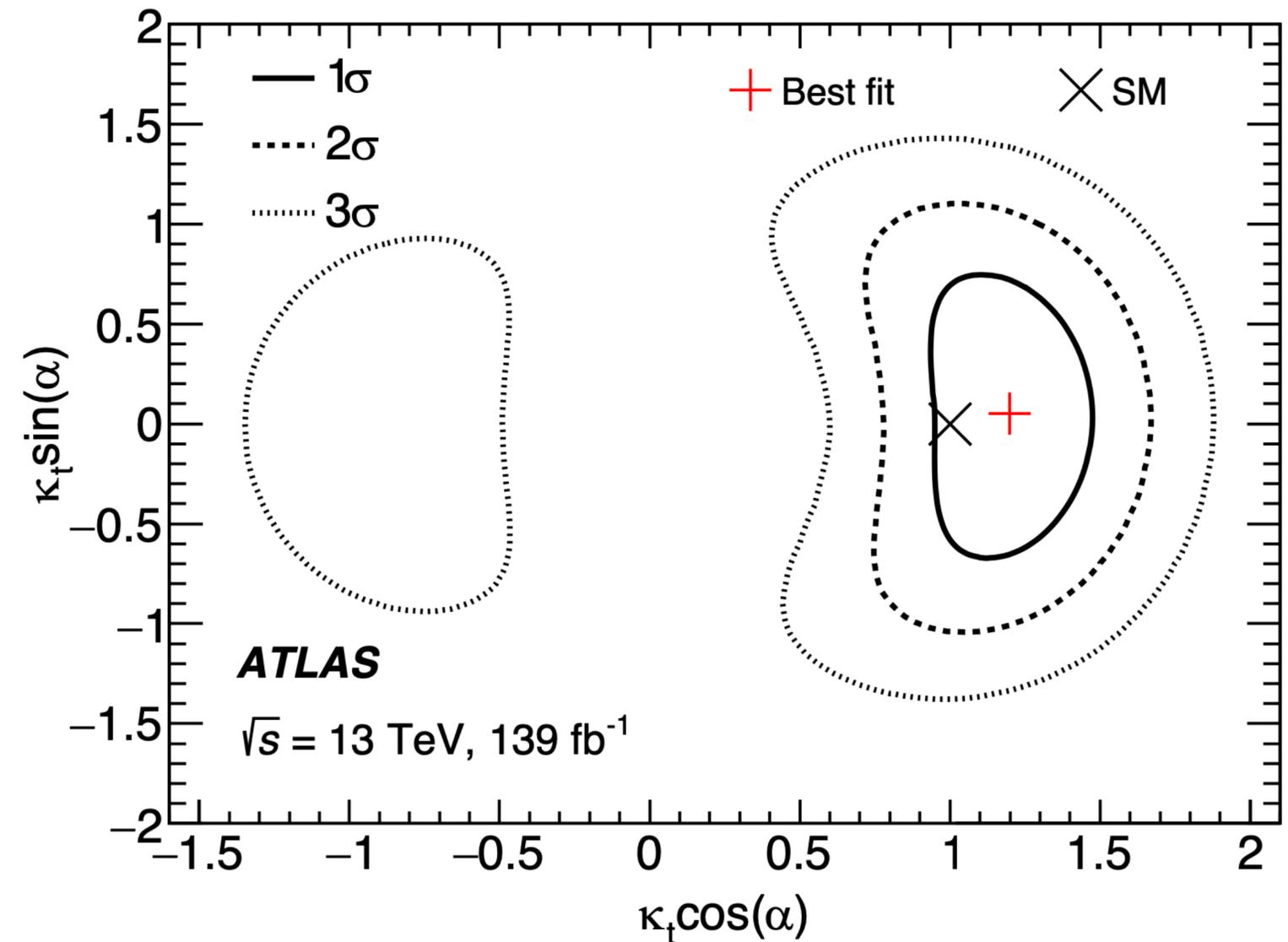
Assuming a CP-even coupling and constraining all non- $t\bar{t}H$  Higgs processes to SM predictions.

Measured rate for  $t\bar{t}H$ :

$$1.43^{+0.33}_{-0.31} \text{ (stat)}^{+0.21}_{-0.15} \text{ (sys)} \times SM_{\text{expectation}}$$

Pure CP-odd coupling excluded at  $3.9\sigma$   
 → *Severely* constrained!

Limit set on CP mixing angle:  
 $|\alpha| > 43^\circ$  excluded @ 95% CL



All measurements are consistent with SM expectations.

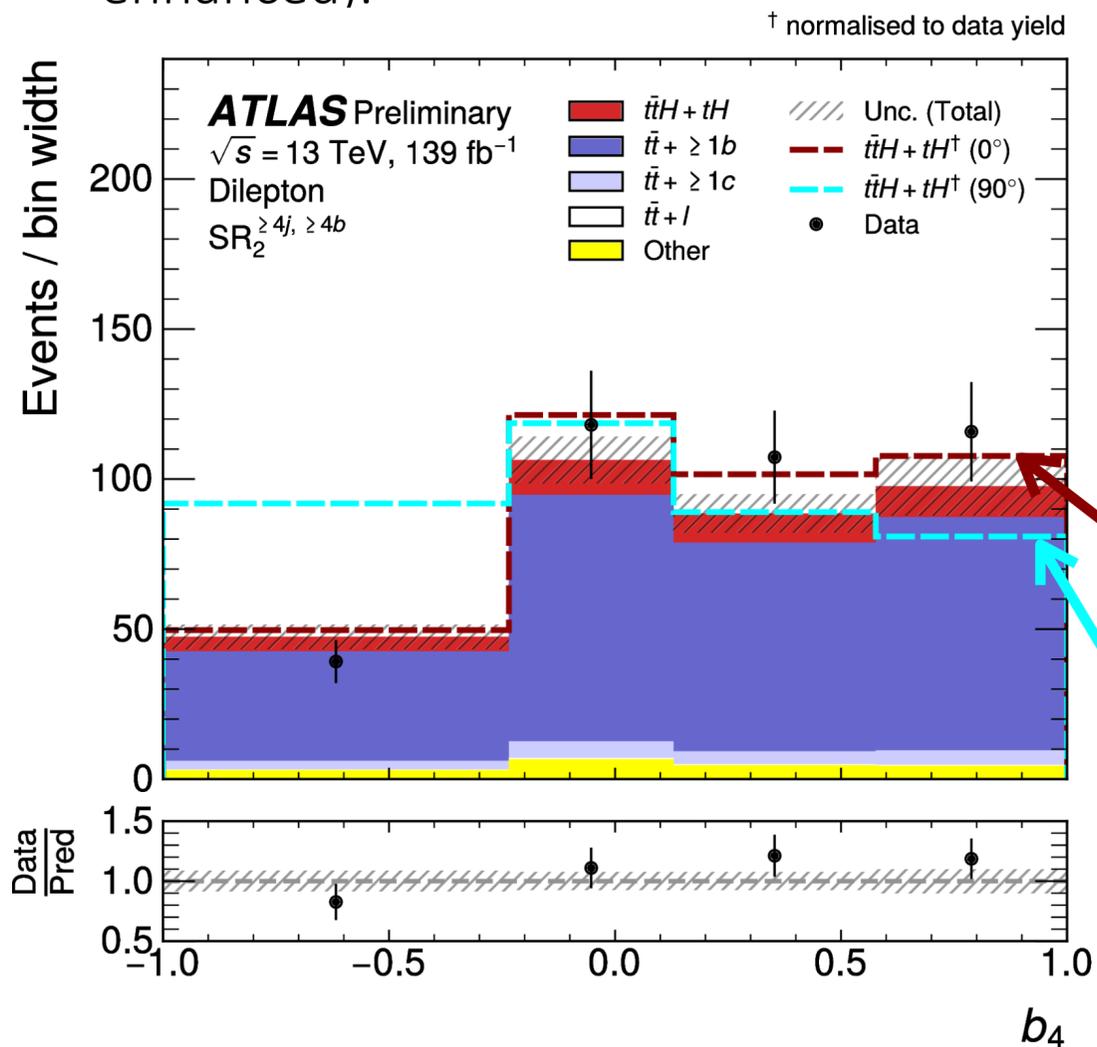
# Probing CP properties of the top-Higgs Yukawa in $H \rightarrow b\bar{b}$ :

ATLAS-CONF-2022-016

Channel:  $H \rightarrow b\bar{b}$  produced in association with top quarks ( $t\bar{t}H$  or  $tH$ ).  
 Target final states where  $\geq 1$  top decays semi-leptonically to  $e$  or  $\mu$ .

Analysis optimised for  $t\bar{t}H(\rightarrow b\bar{b})$ , not the  $tH$  signal  $\rightarrow$  small yield, due to destructive interference.

Higgs  $p_T$  sensitive to CP  $\rightarrow$  exploited by considering boosted region (tH enhanced).

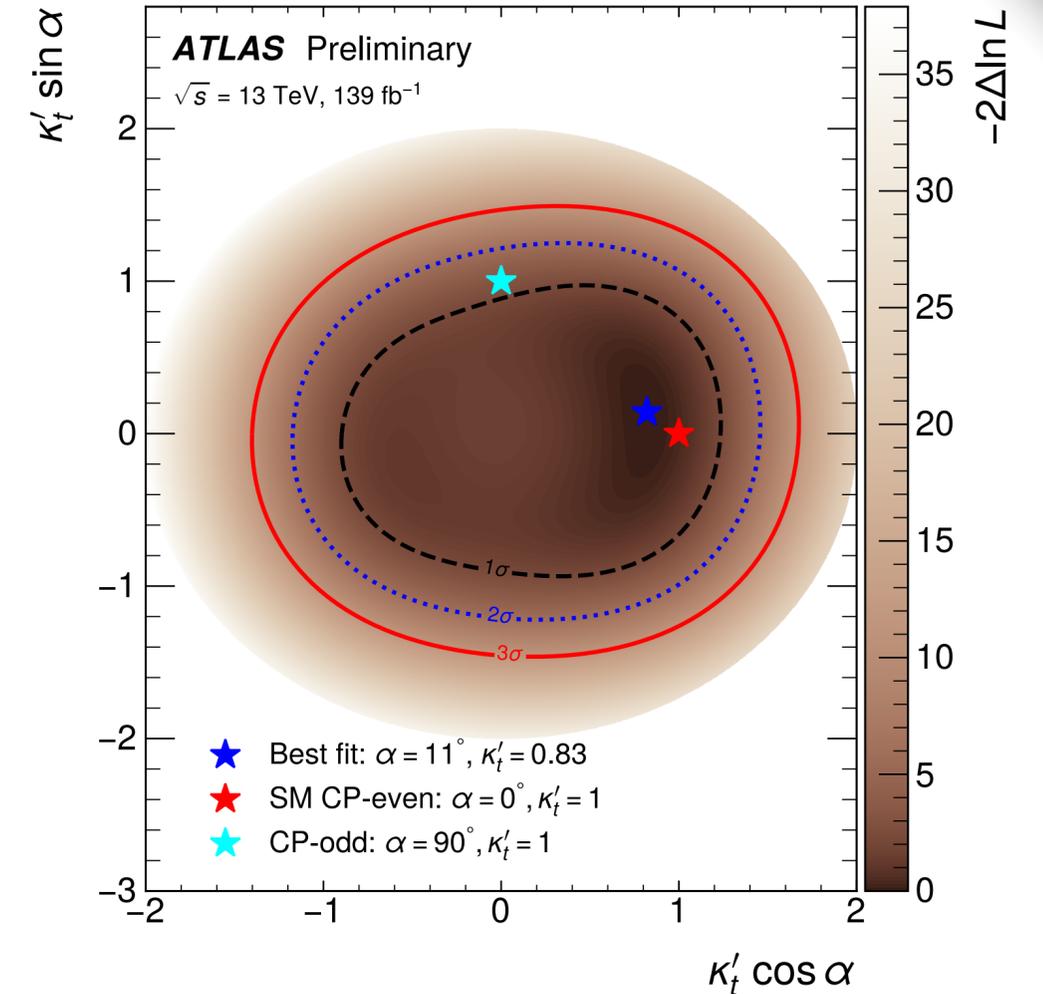


**CP sensitive observable:  $b_4$**   
 exploits enhanced  $t$  production close to beam pipe and in opposite directions, for CP-odd  $t\bar{t}H$  production.

$$b_4 = \frac{p_1^z p_2^z}{|\vec{p}_1| |\vec{p}_2|}$$

$\alpha = 0^\circ$ ,  
 CP-even

$\alpha = 90^\circ$ ,  
 CP-odd



Values of  $\alpha$  and  $\kappa_t$  determined simultaneously using a binned profile likelihood fit.

Pure CP-odd coupling excluded at  $1.2 \sigma$   
 Best fit values: CP-mixing angle,  $\alpha = 11^{+55}_{-77}$   
 Coupling strength  $\kappa_t = 0.83^{+0.30}_{-0.46}$

# Defining a CP sensitive observable in $H \rightarrow \tau\tau$ :

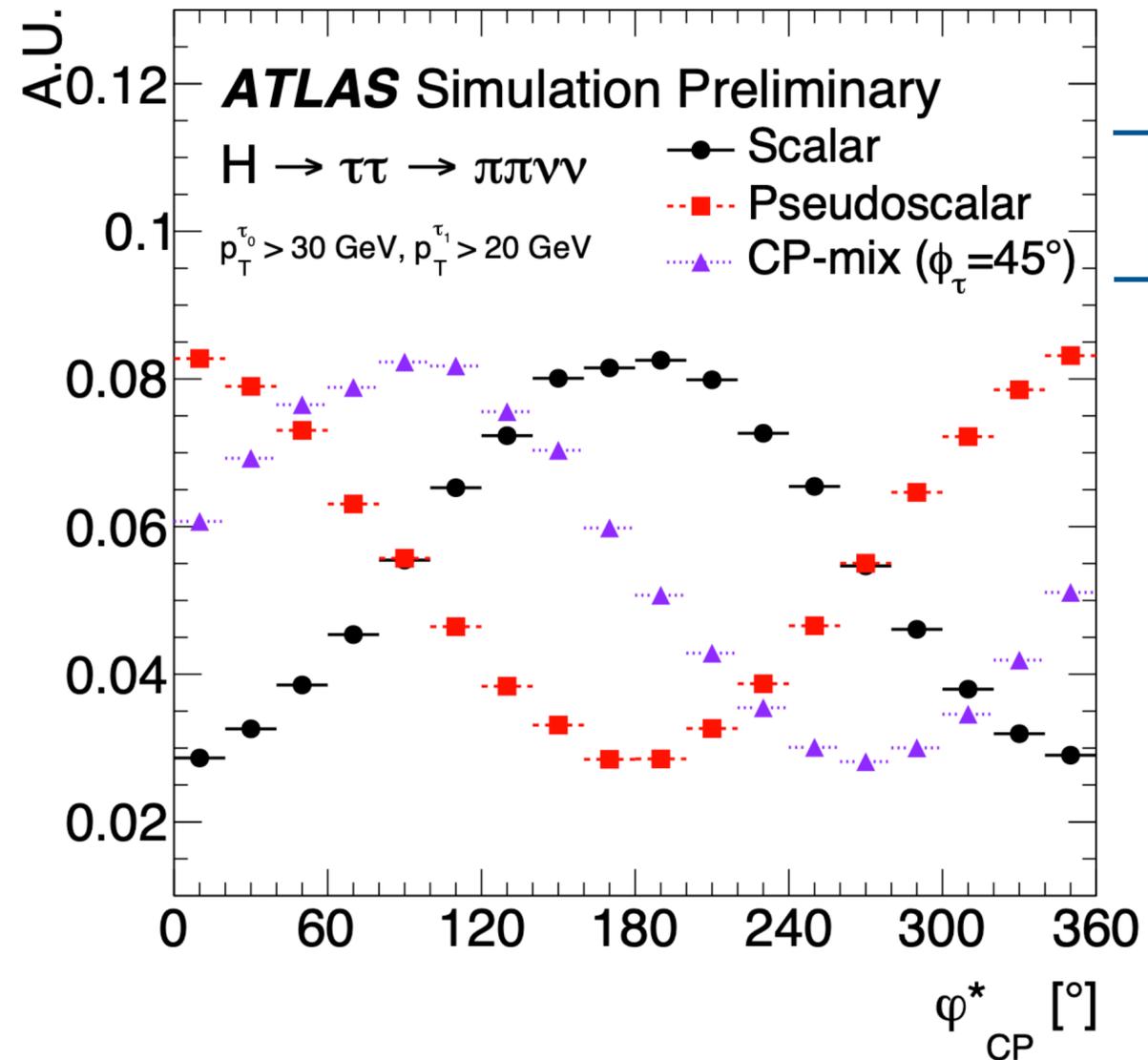
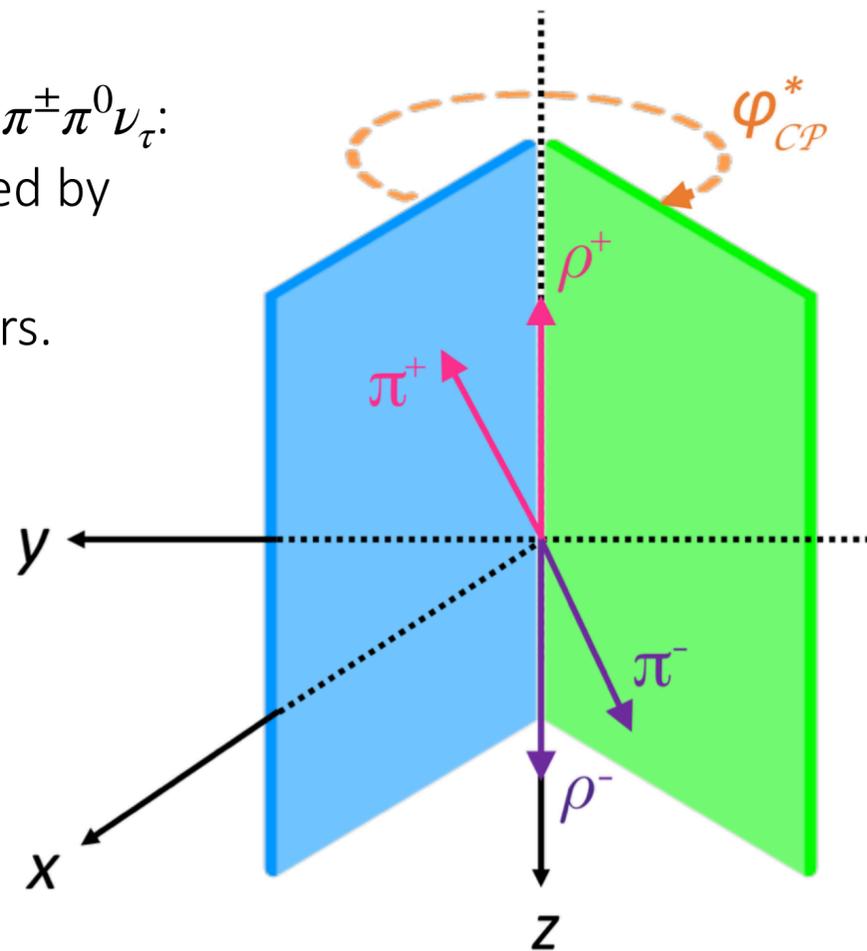
ATLAS-CONF-2022-032

Higgs decays to  $\tau$  lepton pairs allow for direct probe of the CP properties of the Higgs to  $\tau$  Yukawa coupling.  
 CP-mixing angle,  $\phi_\tau$  information encoded in correlations between transverse spin components of  $\tau$  leptons.

**CP sensitive observable: acoplanarity angle,  $\varphi_{CP}^*$**  between the decay planes.

⇒ Constructed based on visible decay products, where method is different for each decay channel.

For  $\tau^\pm \rightarrow \rho^\pm \nu_\tau \rightarrow \pi^\pm \pi^0 \nu_\tau$ :  
 decay plane formed by the pion four-momentum vectors.



Shape of  $\varphi_{CP}^*$  distribution under different Higgs CP hypotheses.

# Determination of the Higgs CP state in $H \rightarrow \tau\tau$ :

ATLAS measurement using  $139 \text{ fb}^{-1}$  of data, collected at 13 TeV.

ATLAS-CONF-2022-032

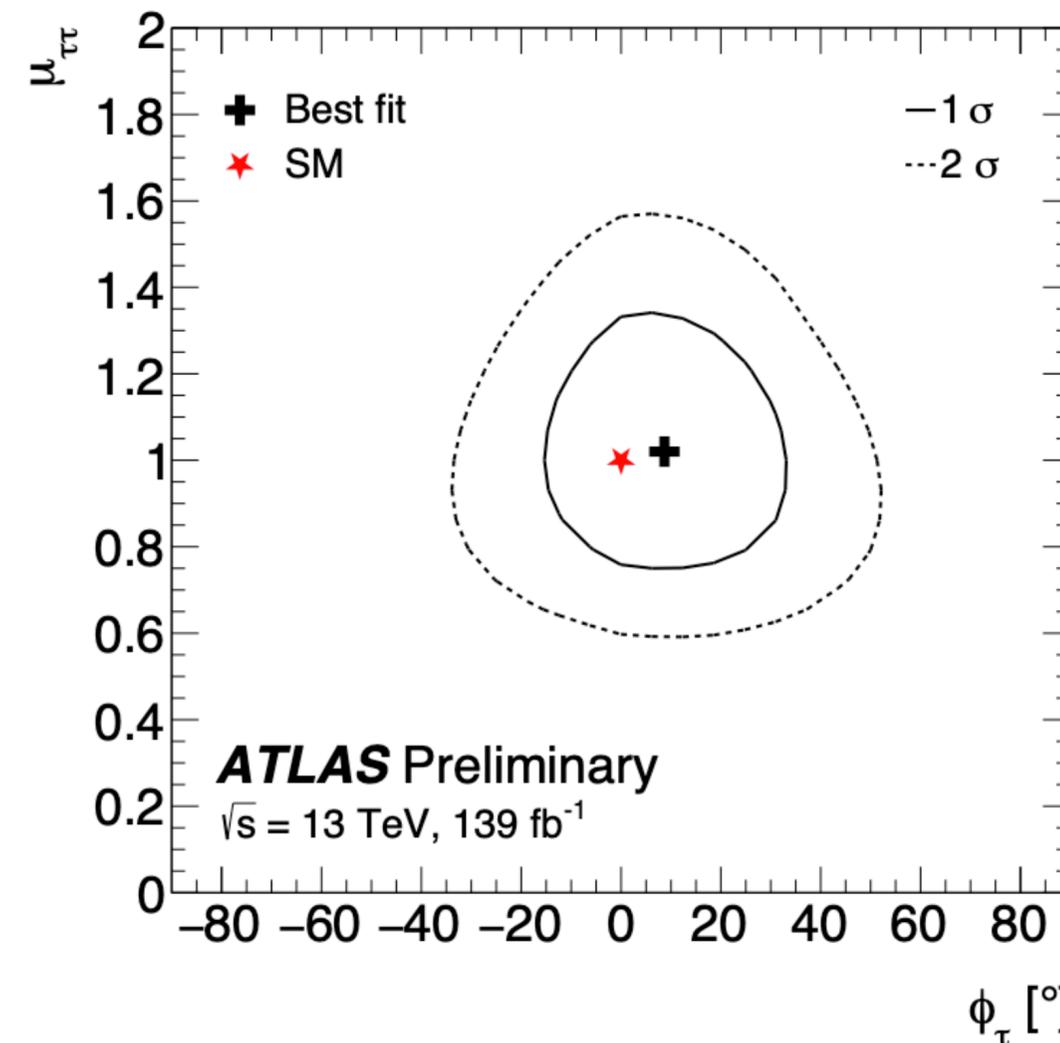
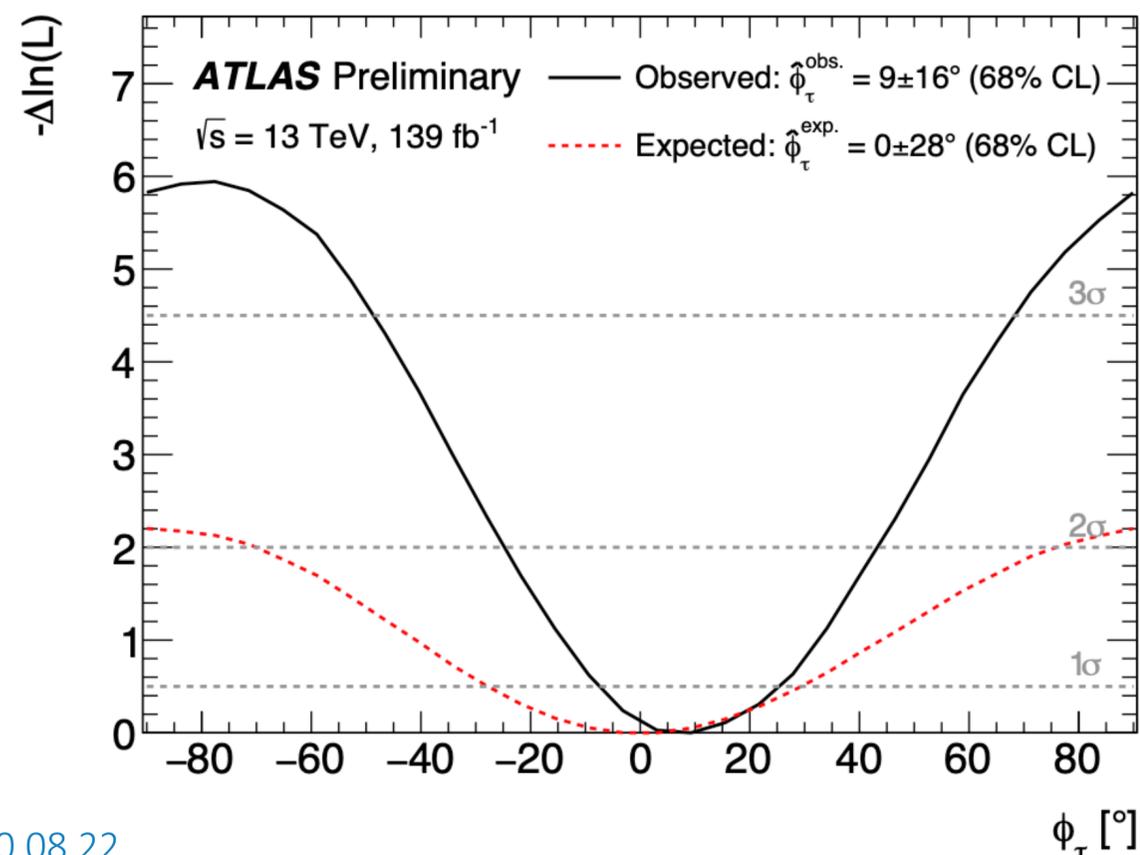
Observed value of  $\phi_\tau$  is:  $9 \pm 16^\circ$

Total uncertainty dominated by **statistical** uncertainties.

Leading systematic uncertainties:

- jets ( $4.3^\circ$  impact)
- theory uncertainties on  $H \rightarrow \tau\tau$  ( $1.5^\circ$ )

Uncertainties on the  $\tau$  decay reconstruction  $< 1^\circ$  on  $\phi_\tau$ .



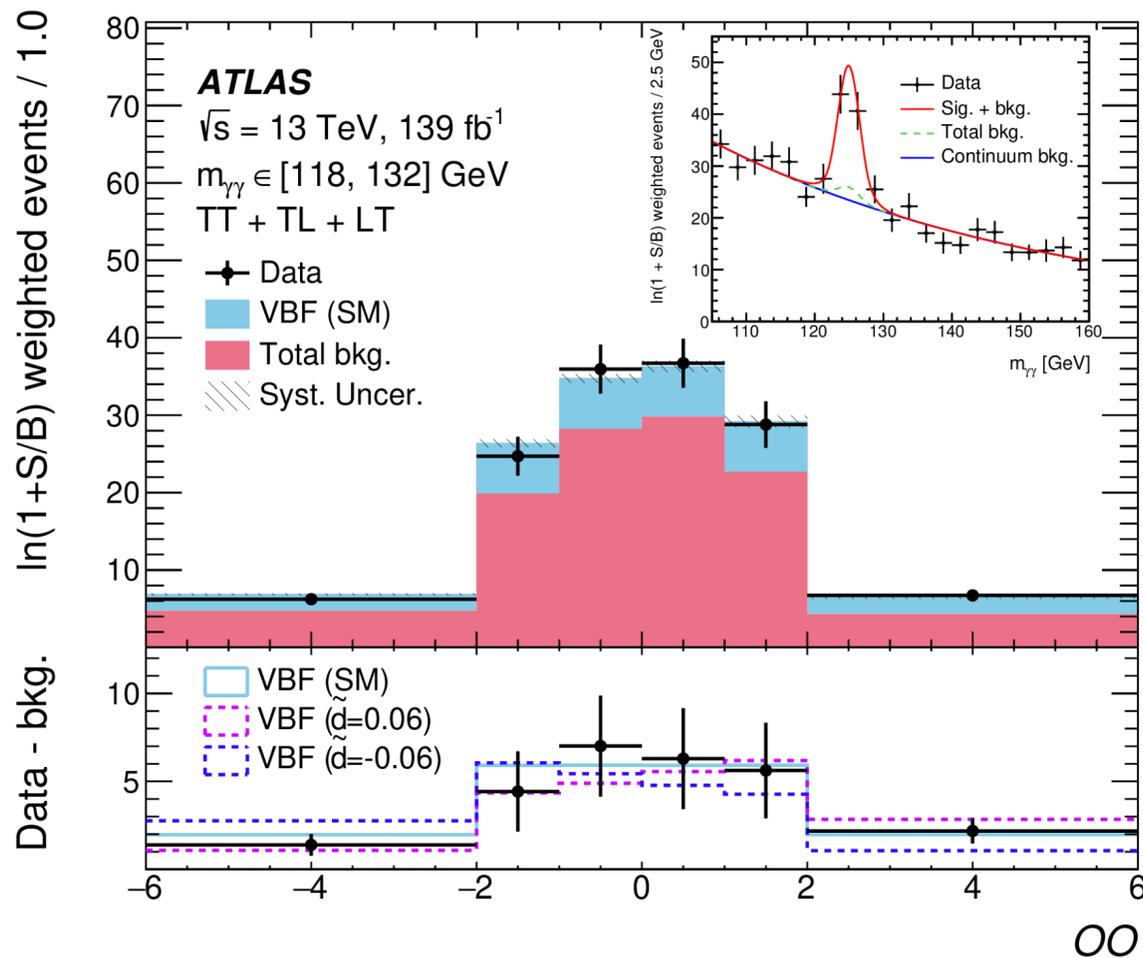
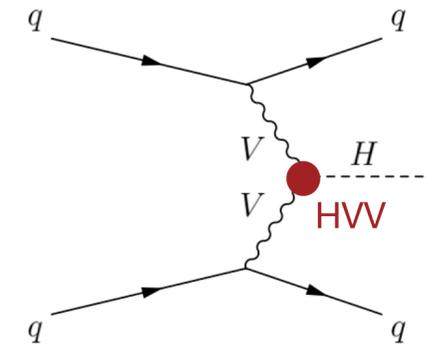
Pure CP-odd hypothesis is excluded at  $3.4\sigma$ !

**Compatible with the SM expectation within the measurement uncertainties.**

CP-odd component in  $HVV$  decays described by adding dim-6 operators to the SM Lagrangian, using an effective field theory (EFT) approach:

$$|\mathcal{M}|^2 = |\mathcal{M}_{\text{SM}}|^2 + \tilde{d} \cdot 2\text{Re}(\mathcal{M}_{\text{SM}}^* \mathcal{M}_{\text{CP-odd}}) + \tilde{d}^2 \cdot |\mathcal{M}_{\text{CP-odd}}|^2$$

Test CP invariance of the  $HVV$  coupling - Optimal Observable (OO) method:  $OO = 2 \cdot \text{Re}(\mathcal{M}_{\text{SM}}^* \cdot \mathcal{M}_{\text{CP-odd}}) / |\mathcal{M}_{\text{SM}}|^2$



SM: symmetric, and mean = 0.

Wilson coefficients multiplying CP-odd operators describe CP-odd couplings.

1. HISZ [[Phys. Rev. D 48, 2182](#)]: constrain  $\tilde{d}$
2. Warsaw [[arXiv:1008.4884v3](#)]: constrain  $c_{H\tilde{W}}$

All CP-even operator coefficients  $\Rightarrow$  zero

MOST STRINGENT!

	68% (exp.)	95% (exp.)	68% (obs.)	95% (obs.)
$\tilde{d}$ (inter. only)	[-0.027, 0.027]	[-0.055, 0.055]	[-0.011, 0.036]	[-0.032, 0.059]
$\tilde{d}$ (inter.+quad.)	[-0.028, 0.028]	[-0.061, 0.060]	[-0.010, 0.040]	[-0.034, 0.071]
$\tilde{d}$ from $H \rightarrow \tau\tau$	[-0.038, 0.036]	-	[-0.090, 0.035]	-
<b>Combined <math>\tilde{d}</math></b>	<b>[-0.022, 0.021]</b>	<b>[-0.046, 0.045]</b>	<b>[-0.012, 0.030]</b>	<b>[-0.034, 0.057]</b>
$c_{H\tilde{W}}$ (inter. only)	[-0.48, 0.48]	[-0.94, 0.94]	[-0.16, 0.64]	[-0.53, 1.02]
$c_{H\tilde{W}}$ (inter.+quad.)	[-0.48, 0.48]	[-0.95, 0.95]	[-0.15, 0.67]	[-0.55, 1.07]

**Results all compatible with SM!**

Successful and interesting 10 years of Higgs properties measurements.

1. Combined Run I+II Higgs mass measurement in  $H \rightarrow ZZ^* \rightarrow 4\ell$  with precision of **0.14%** !
2. Higgs width has been constrained to:  

$$\Gamma_H < 14.4 \text{ MeV @ 95\% CL}$$
3. Pure CP odd Higgs excluded at  $3.9\sigma$  ( $3.4\sigma$ ) with  $H \rightarrow \gamma\gamma$  ( $H \rightarrow \tau\tau$ ) decays. **CP mixing not yet ruled out.**

13.6 TeV collisions started in July 2022.  
Early Run3 analyses have begun.

Run3 + the HL-LHC will bring very interesting new results.



CERN

Full set of public ATLAS results:

<https://twiki.cern.ch/twiki/bin/view/AtlasPublic/HiggsPublicResults>

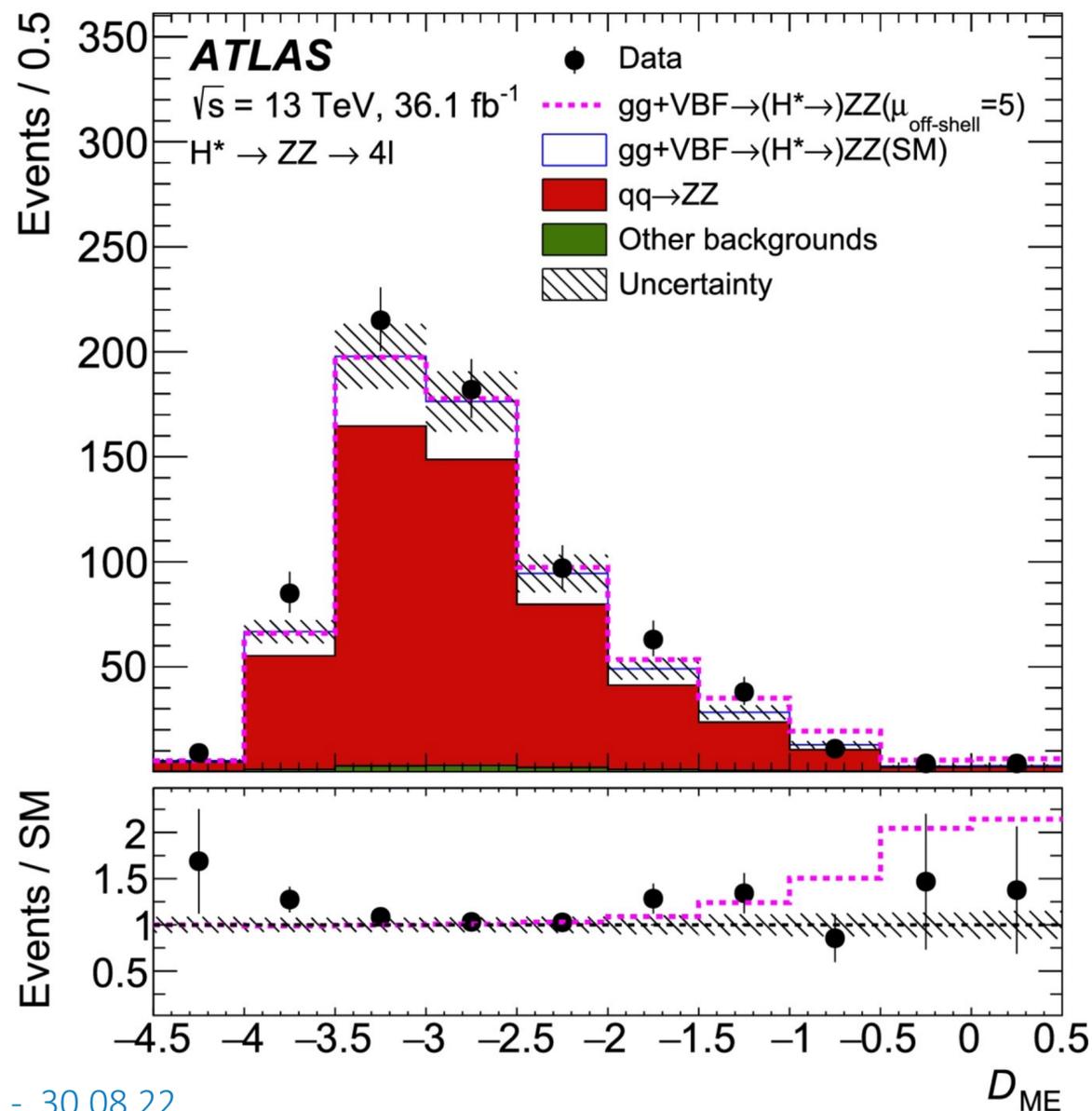
More questions or comments?

[ashley.ellen.mcdougall@cern.ch](mailto:ashley.ellen.mcdougall@cern.ch)

EXTRA MATERIAL

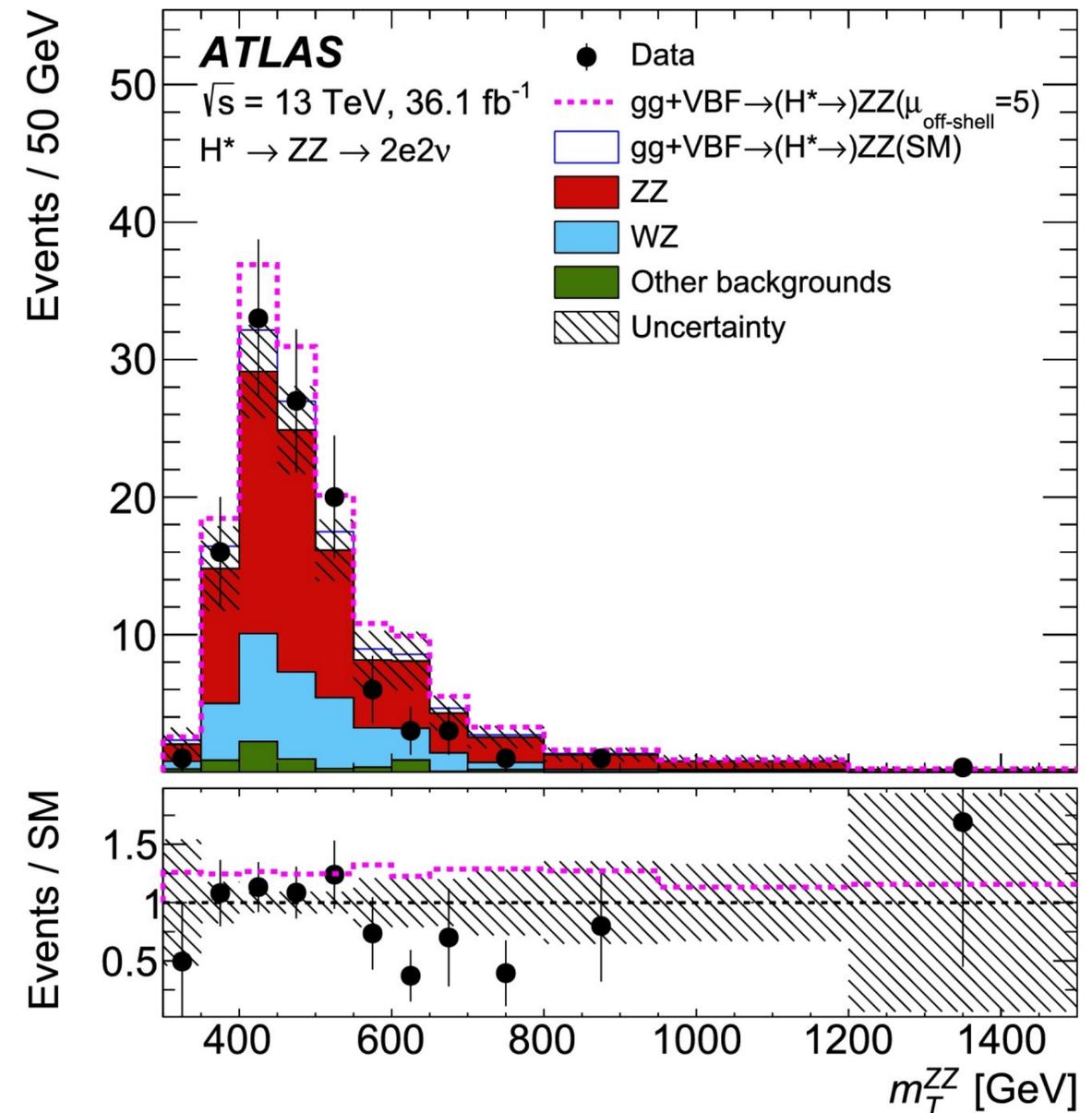
ME-based discriminant used in the 4l channel is defined as per XXXX:

$$D_{ME} = \log_{10} \left( \frac{P_H}{P_{gg+c} \cdot P_{q\bar{q}}} \right)$$



Observable of interest in  $2\ell 2\nu$  channel: transverse mass of the ZZ system,  $m_T^{ZZ}$ :

$$m_T^{ZZ} \equiv \sqrt{\left[ \sqrt{m_Z^2 + (p_T^{\ell\ell})^2} + \sqrt{m_Z^2 + (E_T^{\text{miss}})^2} \right]^2 - \left| \vec{p}_T^{\ell\ell} + \vec{E}_T^{\text{miss}} \right|^2}$$



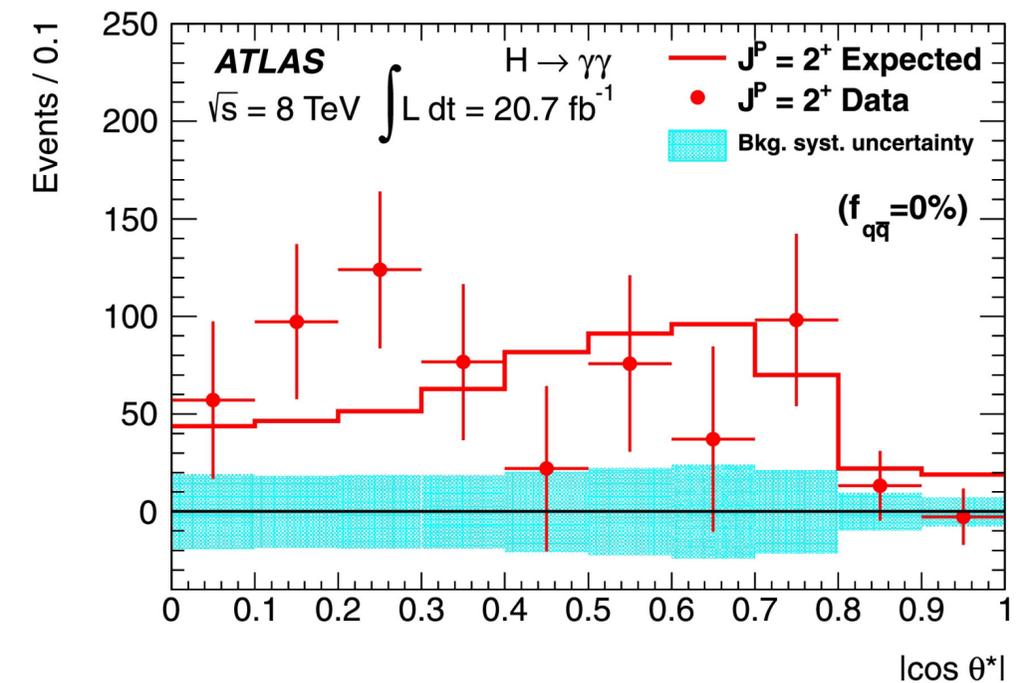
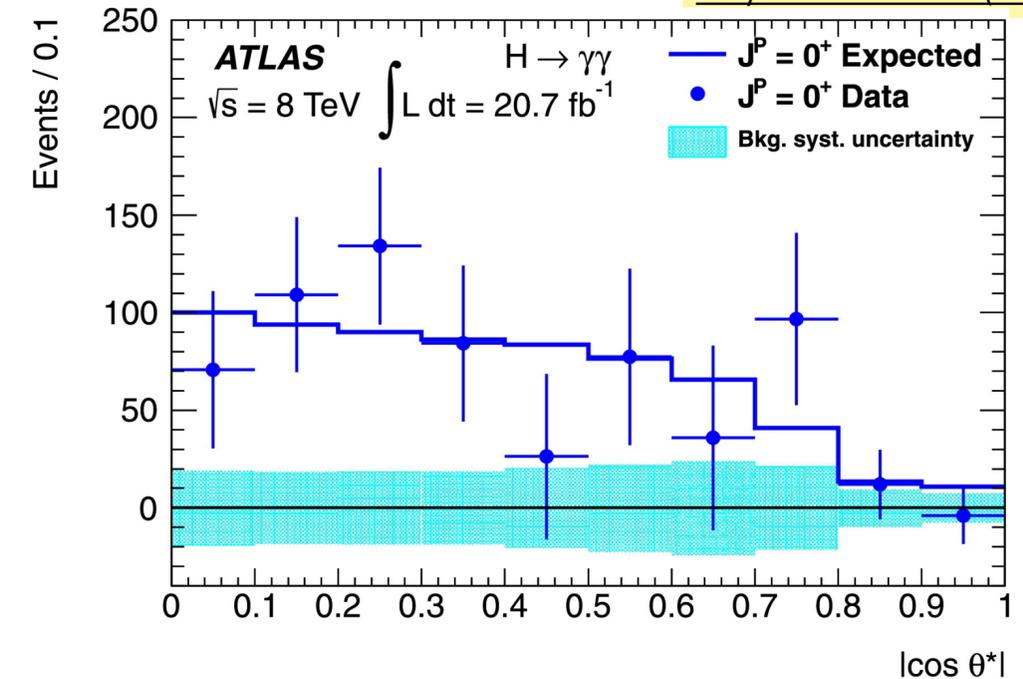
Spin is a property of the particle, where as Charge-Parity (CP) of the couplings.

$H \rightarrow \gamma\gamma$  analysis: sensitive observable,  $|\cos\theta^*|$  = polar angular distribution of the photons in the resonance rest frame, defined by:

$$|\cos\theta^*| = \frac{|\sinh(\Delta\eta^{\gamma\gamma})|}{\sqrt{1 + (p_T^{\gamma\gamma}/m_{\gamma\gamma})^2}} \frac{2p_T^{\gamma 1} p_T^{\gamma 2}}{m_{\gamma\gamma}^2},$$

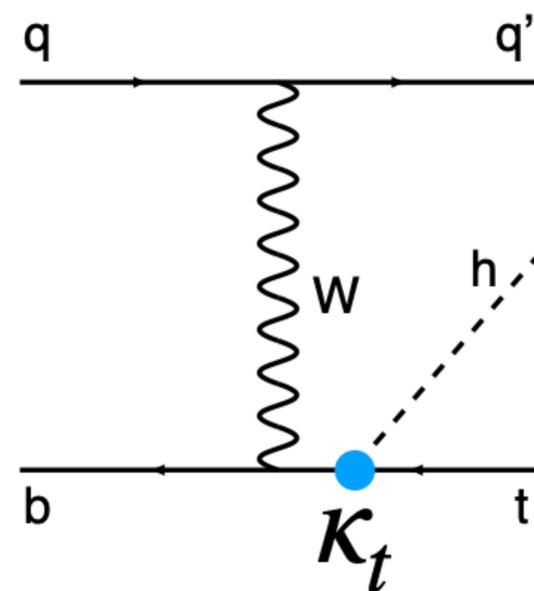
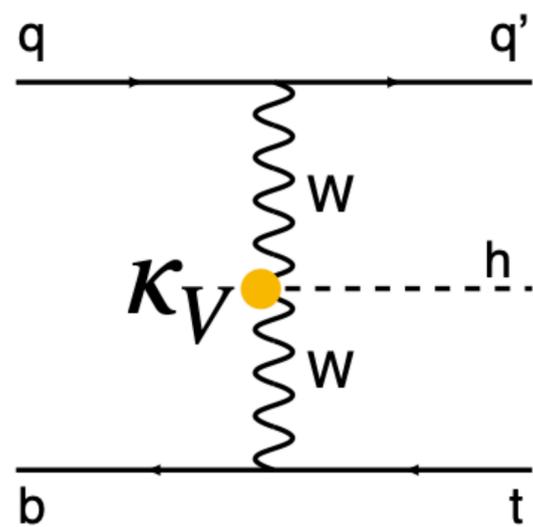
Spin information can be extracted from the distribution of the absolute value of the cosine of the polar angle  $\theta^*$  of the photons with respect to the z-axis of the Collins–Soper frame [ref].

Distributions show are background subtracted.  
 In this analysis, only J=0 and J=2 hypothesis were considered.  
 Normalised to fitted number of signal events.



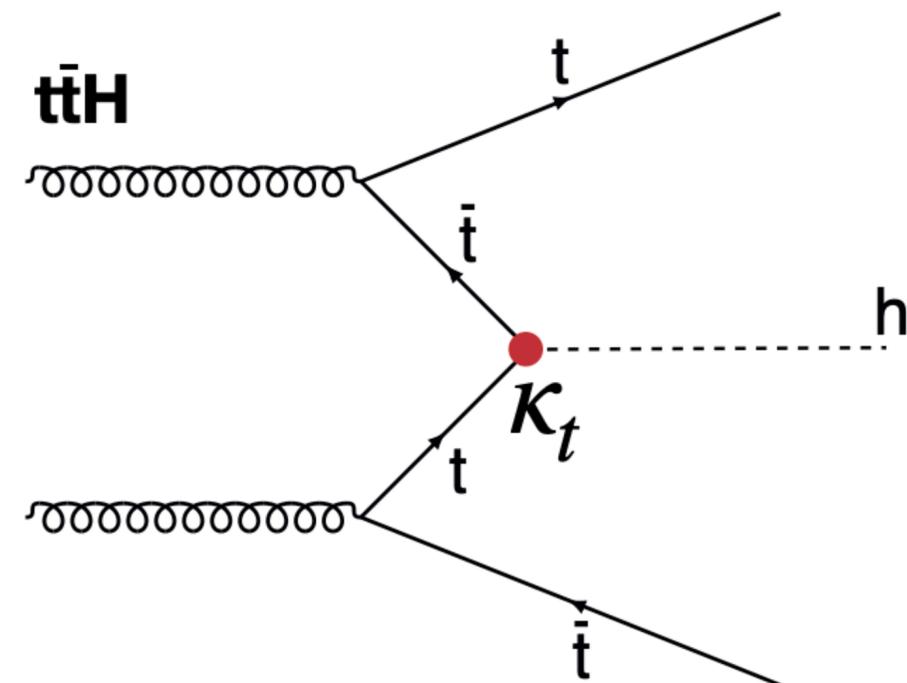
- $tH$  is sensitive to the sign of the top Yukawa coupling, as opposed to  $t\bar{t}H$
- Sensitivity comes from interference between possible diagrams.
- Use process to directly exclude  $\kappa_t$  in ATLAS.

**tH**



$$\sigma_{tH}/\sigma_{tH}^{SM} = 2.63\kappa_t^2 - 5.25\kappa_t\kappa_V + 3.58\kappa_V^2$$

**t $\bar{t}$ H**



$$\sigma_{t\bar{t}H}/\sigma_{t\bar{t}H}^{SM} = \kappa_t^2$$

$\phi_\tau$  information encoded in correlations between transverse spin components of  $\tau$  leptons (passed down to  $\tau$  decay products as consequence of parity violation in weak interactions).

Define **CP sensitive observable: acoplanarity angle,  $\varphi_{CP}^*$**  between the decay planes.

⇒ Constructed based on visible decay products, where method is different for each decay channel.

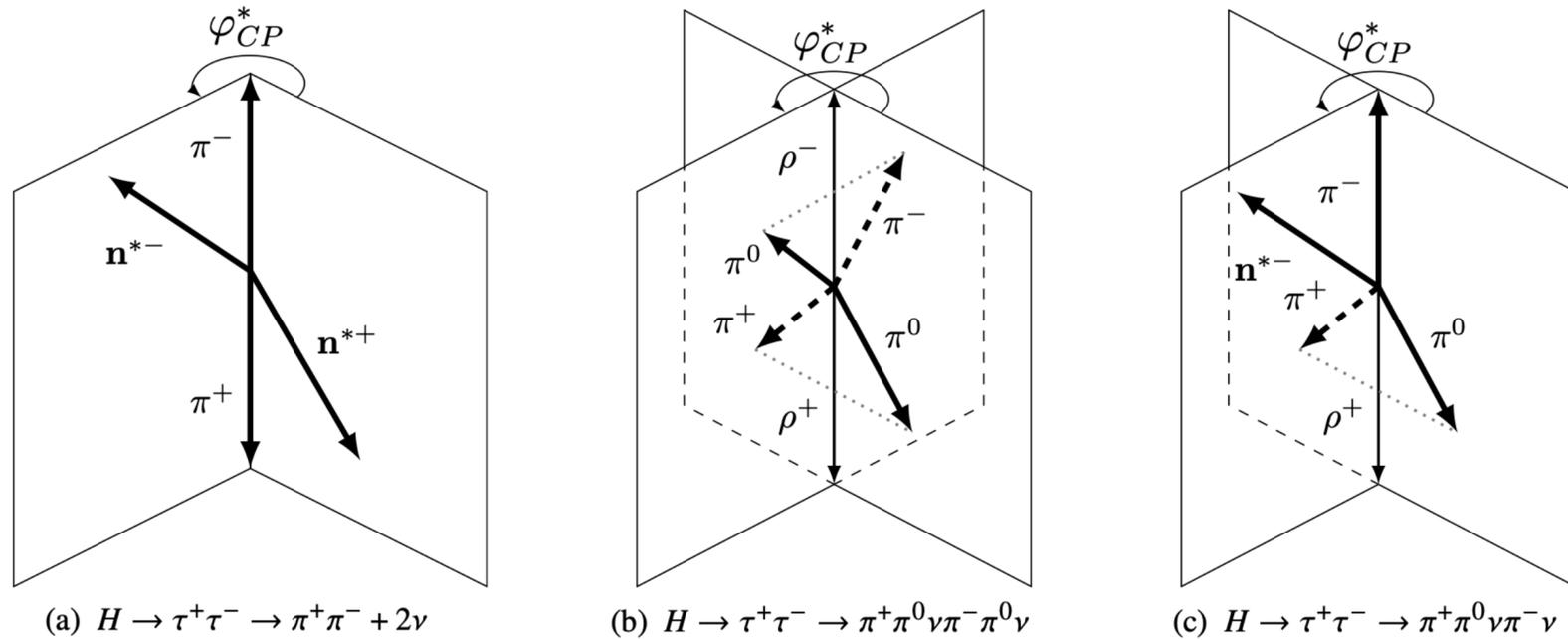


Figure 1: Illustration of the  $\tau$  lepton decay planes for constructing the  $\varphi_{CP}^*$  observable in (a)  $H \rightarrow \tau^+\tau^- \rightarrow \pi^+\pi^- + 2\nu$  decay using the impact parameter method, (b)  $H \rightarrow \tau^+\tau^- \rightarrow \pi^+\pi^0\nu\pi^-\pi^0\nu$  using the  $\rho$ -decay plane method, and (c)  $H \rightarrow \tau^+\tau^- \rightarrow \pi^+\pi^0\nu\pi^-\nu$  using the combined impact parameter and  $\rho$ -decay plane method. The decay planes are spanned by the spatial momentum vector of the charged decay particle of the  $\tau$  lepton ( $\pi^\pm$ ) and either its impact parameter  $\mathbf{n}^{*\pm}$  or the spatial momentum vector of the neutral decay particle of the  $\tau$  lepton ( $\pi^0$ ).

Impact parameter method:

$$\varphi^* = \arccos(\hat{\mathbf{n}}_{\perp}^{*+} \cdot \hat{\mathbf{n}}_{\perp}^{*-}) \quad \text{and} \quad O_{CP}^* = \hat{\mathbf{q}}^{*-} \cdot (\hat{\mathbf{n}}_{\perp}^{*+} \times \hat{\mathbf{n}}_{\perp}^{*-}),$$

$$\varphi_{CP}^* = \begin{cases} \varphi^* & \text{if } O_{CP}^* \geq 0 \\ 360^\circ - \varphi^* & \text{if } O_{CP}^* < 0, \end{cases}$$

$\rho$ -decay plane method:

$$\varphi^* = \arccos(\hat{\mathbf{q}}_{\perp}^{*0+} \cdot \hat{\mathbf{q}}_{\perp}^{*0-}) \quad \text{and} \quad O_{CP}^* = \hat{\mathbf{q}}^{*-} \cdot (\hat{\mathbf{q}}_{\perp}^{*0+} \times \hat{\mathbf{q}}_{\perp}^{*0-}),$$

$$\varphi^{*'} = \begin{cases} \varphi^* & \text{if } O_{CP}^* \geq 0 \\ 360^\circ - \varphi^* & \text{if } O_{CP}^* < 0. \end{cases} \quad \varphi_{CP}^* = \begin{cases} \varphi^{*'} & \text{if } y_+^\rho y_-^\rho \geq 0 \\ \varphi^{*'} + 180^\circ & \text{if } y_+^\rho y_-^\rho < 0. \end{cases}$$

$$y_{\pm}^\rho = \frac{E_{\pi^\pm} - E_{\pi^0}}{E_{\pi^\pm} + E_{\pi^0}}$$

Combined IP and  $\rho$  method:

$$\varphi_{CP}^* = \begin{cases} \varphi^{*'} & \text{if } y^\rho \geq 0 \\ \varphi^{*'} + 180^\circ & \text{if } y^\rho < 0, \end{cases}$$

# Studying CP properties of HVV couplings in VBF $H \rightarrow \gamma\gamma$ :

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	68% (exp.)	95% (exp.)	68% (obs.)	95% (obs.)
$\tilde{d}$ (inter. only)	[-0.027, 0.027]	[-0.055, 0.055]	[-0.011, 0.036]	[-0.032, 0.059]
$\tilde{d}$ (inter.+quad.)	[-0.028, 0.028]	[-0.061, 0.060]	[-0.010, 0.040]	[-0.034, 0.071]
$\tilde{d}$ from $H \rightarrow \tau\tau$	[-0.038, 0.036]	-	[-0.090, 0.035]	-
Combined $\tilde{d}$	[-0.022, 0.021]	[-0.046, 0.045]	[-0.012, 0.030]	[-0.034, 0.057]
$c_{H\tilde{W}}$ (inter. only)	[-0.48, 0.48]	[-0.94, 0.94]	[-0.16, 0.64]	[-0.53, 1.02]
$c_{H\tilde{W}}$ (inter.+quad.)	[-0.48, 0.48]	[-0.95, 0.95]	[-0.15, 0.67]	[-0.55, 1.07]

Results are compatible with the SM

