

Exploiting stellar explosion induced by the QCD phase transition in large-scale neutrino detectors

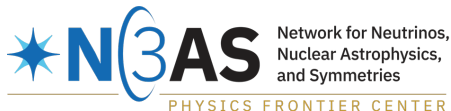
Anna M. Suliga

arXiv: 2208.14469

with T. Pitik, D. Heimsath, and B. Balantekin



CIPANP Orlando, FL,
September 1, 2022

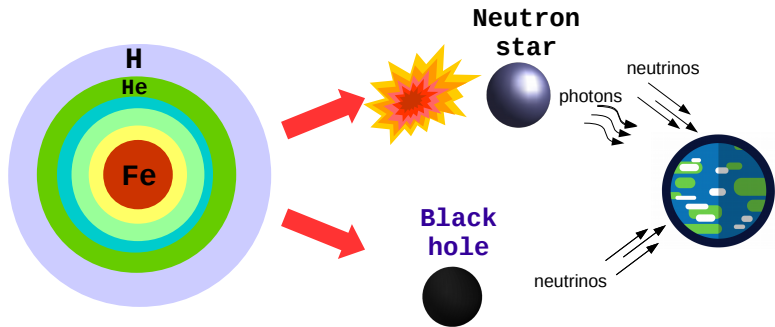


Introduction and Motivation

Why are neutrinos important for a core-collapse supernova?

Neutrinos:

- $\sim 10^{58}$ of them emitted from a single core collapse
- only they (+ GW) can reveal the deep interior conditions
- only they (+ GW) are emitted from the collapse to a black hole



Why core-collapse supernovae are good physics probes?

Advantages

- extreme physical conditions not accessible on Earth: very high densities, long baselines etc.
- within our reach to detect (SK, JUNO, XENON, PandaX...)

What can we learn with a variety of detectors?

- explosion mechanism [Bethe & Wilson \(1985\)](#), [Fischer et al. \(2011\)](#)...
- yields of heavy elements [Woosley et al. \(1994\)](#), [Surman & McLaughlin \(2003\)](#)...
- compact object formation [Warren et al. \(2019\)](#), [Li, Beacom et al. \(2020\)](#)...
- neutrino mixing [Balantekin & Fuller \(2013\)](#), [Tamborra & Shalgar \(2020\)](#)...
- non-standard physics [Suliga et al. \(2019\)](#), (2020) [Suliga & Tamborra \(2020\)](#) .2/13

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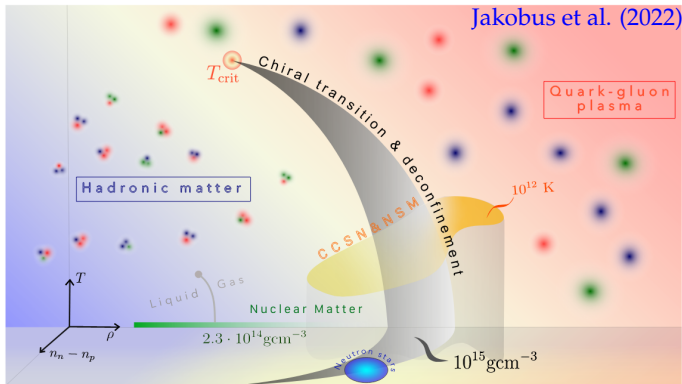
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QCD phase diagram



- Does the protocompact star contain non-leptonic degrees of freedom other than neutrons and protons?
- How to identify the presence of quark matter in astrophysical objects?

Where the quark matter can appear in astrophysical objects?

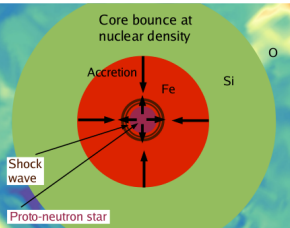
- quark matter in accreting neutron stars
Lin et al. (2006), Abdikamalov et al. (2008), Espino, Paschalidis (2021), ...
- in protoneutron stars after the CCSN explosion
Pons et al. (2001), Keranen et al. (2004)
- in protocompact stars during early postbounce phase
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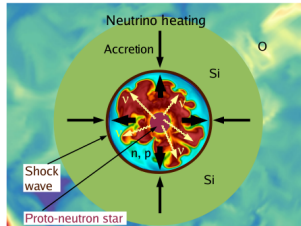
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Different phases of core-collapse supernova explosion

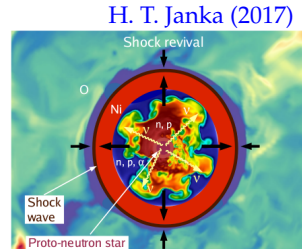
- Infall phase, ν_e burst ~ 40 ms



- Accretion phase, ~ 100 ms



- Cooling phase, ~ 10 s

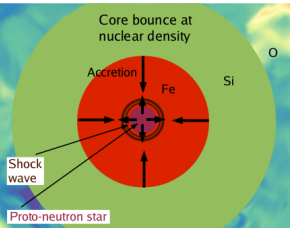


What drives the supernova explosions?

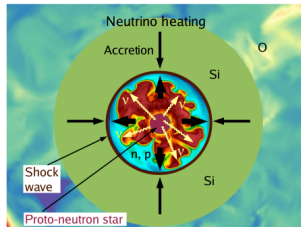
- neutrino heating Colgate & White (1966), Bethe & Wilson (1985)
- magneto-rotational mechanism LeBlanc and Wilson (1970), Takiwaki et al. (2009)
- particles beyond the Standard Model Fuller et al. (2008), Suliga et al. (2020) ...
- phase transition to quark matter Sagert et al. (2008)...

Different phases of core-collapse supernova explosion

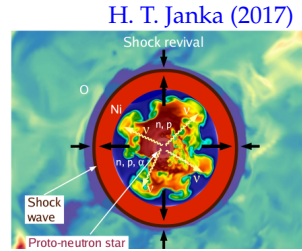
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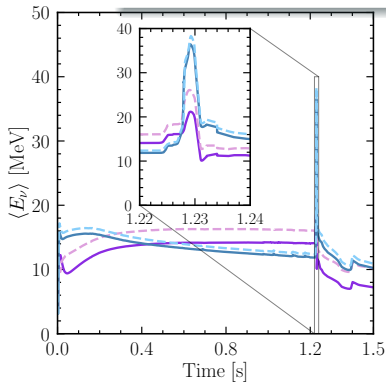
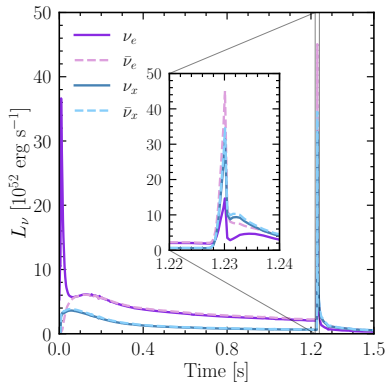
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Numerical modeling

Neutrino Emission Properties from the QHPT CCSN

1D SN model Fischer et al. (2017)

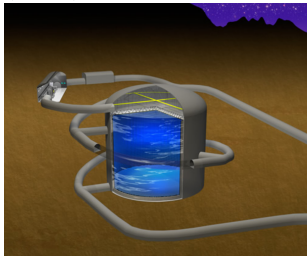


- second sharp neutrino burts dominated by $\bar{\nu}_e$
- non-exploding models can explode

Astrophysical neutrino detection

Large scale neutrino detectors

Hyper-Kamiokande (2027)



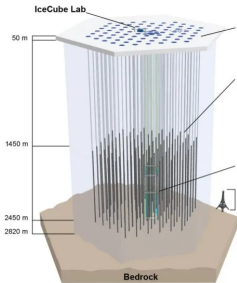
fiducial volume

217 kton

main detection channel



Ice-Cube Observatory DUNE (2030)



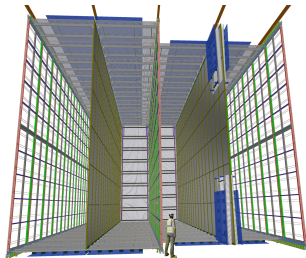
fiducial volume

3500 kton

main detection channel



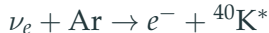
DUNE (2030)



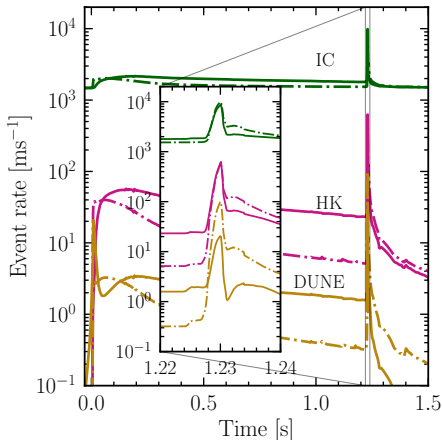
fiducial volume

40 kton

main detection channel



Neutrino Event Rates



Impact of neutrino conversions

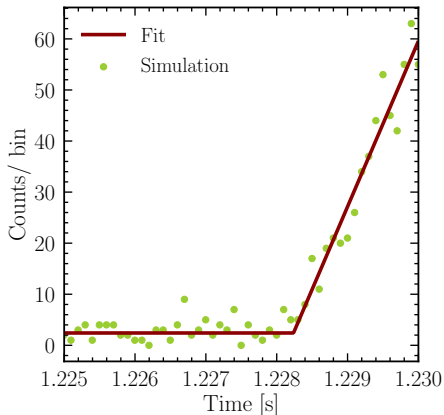
- Event rate in the antineutrino detectors comparable for both conversion scenarios
- Event rate in the neutrino detector larger for the full conversion case

$$R(t) = N_t \int_{E_\nu^{\min}}^{\infty} dE_\nu \int_{E_{\text{th}}}^{E_{\max}} dE \varepsilon \sigma_i(E, E_\nu) F_{\nu\beta}(E_\nu, t)$$

Main Results

Timing the Neutrino Signal

HK: No conversion



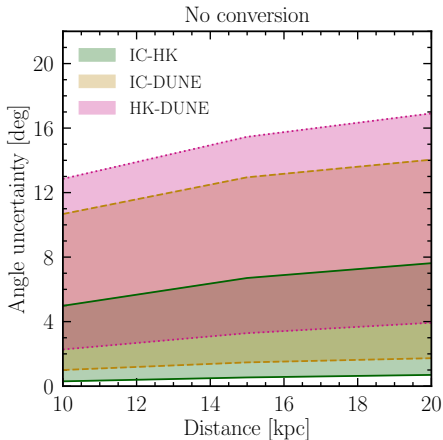
$$R_{\text{exp}} = \begin{cases} R_*, & \text{if } t < t_0 \\ R_* + a(t - t_0), & \text{otherwise} \end{cases},$$

Detectors	No conversion	Full conversion
	B_{ij} [ms]	
IC-HK	-0.32 ± 0.10	-0.32 ± 0.10
IC-DUNE	-0.11 ± 0.48	-0.27 ± 0.20
HK-DUNE	0.22 ± 0.50	0.05 ± 0.22
	$\delta(\theta_{ij})$ (min, max) [deg]	
IC-HK	(0.30, 5.00)	(0.29, 4.90)
IC-DUNE	(1.00, 10.67)	(0.41, 6.90)
HK-DUNE	(2.27, 12.85)	(1.00, 8.54)
	95% C.L. upper limit on m_ν [eV]	
IC	$0.16^{+0.03}_{-0.04}$	$0.21^{+0.05}_{-0.05}$
HK	$0.22^{+0.05}_{-0.06}$	$0.30^{+0.07}_{-0.09}$
DUNE	$0.80^{+0.21}_{-0.29}$	$0.58^{+0.14}_{-0.19}$

$$\Delta t_{ij}^{\text{true}} = \frac{(\mathbf{r}_i - \mathbf{r}_j) \cdot \mathbf{n}}{c} = \frac{D_{ij} \cos \theta}{c}$$

$$\Delta t_{ij}^{\text{measured}} = \Delta t_{ij}^{\text{true}} + B_{ij}$$

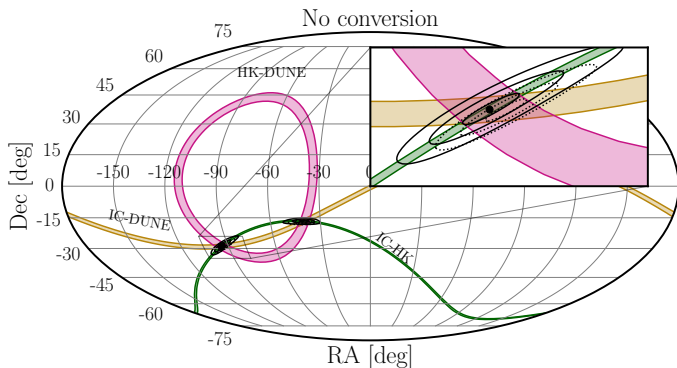
Determination of the uncertainty of the CCSN localization



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$$\delta(\theta_{ij}) \approx \begin{cases} \delta(\cos \theta_{ij}) / \sin \theta_{ij} & \text{if } \sin \theta_{ij} > \sqrt{\delta(\cos \theta_{ij})} \\ \sqrt{2\delta(\cos \theta_{ij})}, & \text{for } \theta_{ij} \ll \delta(\cos \theta_{ij}) \end{cases}$$

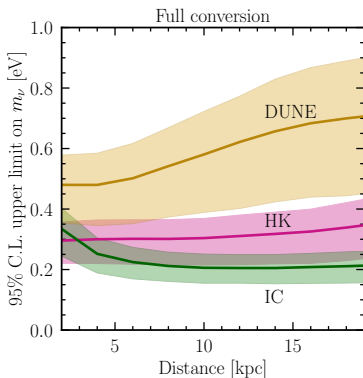
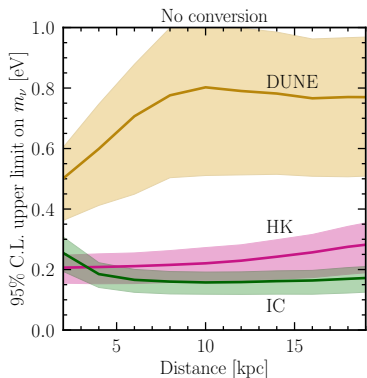
Determination of the CCSN localization



- improvement by 4.5-10 times compared to neutronization burst
- comparable results for black hole forming supernovae
- not far off from elastic scattering on electrons

Sensitivity to the Absolute Neutrino Mass

$$\Delta t \approx 5.15 \left(\frac{D}{10 \text{ kpc}} \right) \left(\frac{m_\nu}{1 \text{ eV}} \right)^2 \left(\frac{10 \text{ MeV}}{E_\nu} \right)^2 \text{ ms}$$



- up to $\sim 10\times$ improvement compared to neutronization burst
- more stringent limits than from the laboratory experiments (0.8 eV)

Summary and Conclusions

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- QCP phase transition in the collapsing star can:
 - produce second core bounce
 - result in release of a second sharp neutrino burst
 - lead to r -process elements production
- Detection of the phase transition induced neutrino burst:
 - indicates the QCD phase transition in supernova
 - improves the precision of the supernova triangulation
 - sets competitive limits on the neutrino mass

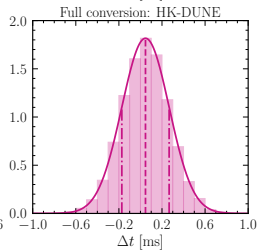
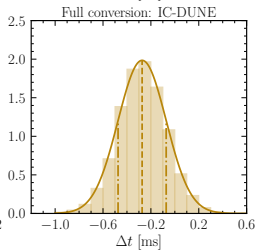
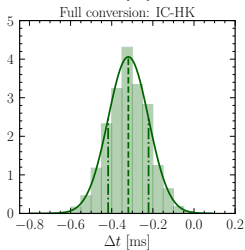
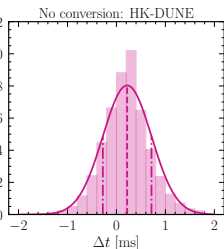
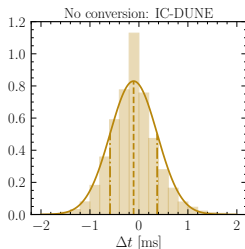
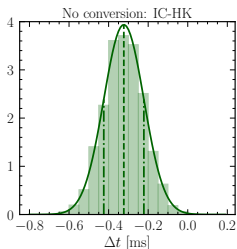
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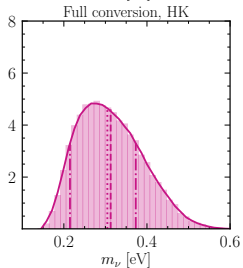
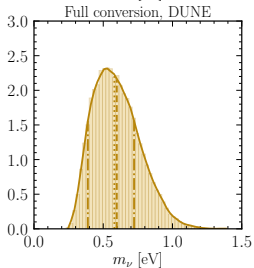
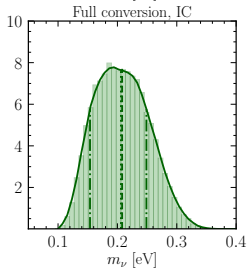
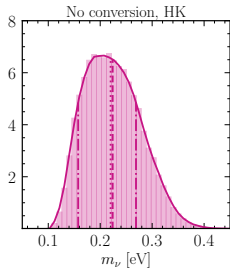
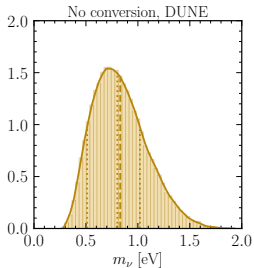
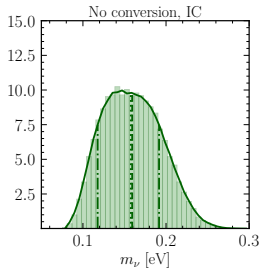
Thank you for the attention!

Backup

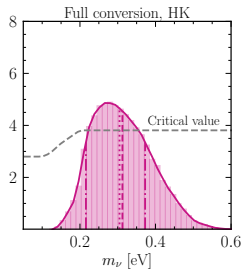
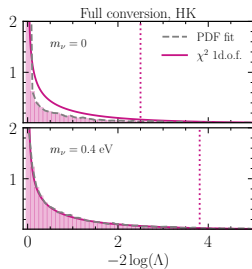
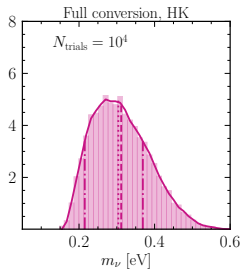
Histograms: Timing the neutrino signal



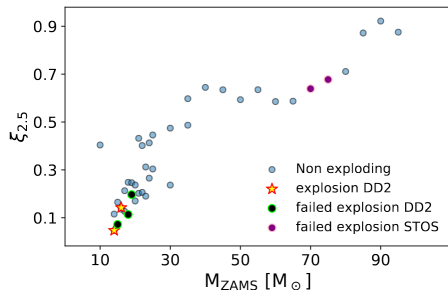
Histograms: neutrino mass limit



Relaxing Wilk's theorem approximation



The Role of the QCD Phase Transition in CCSNe

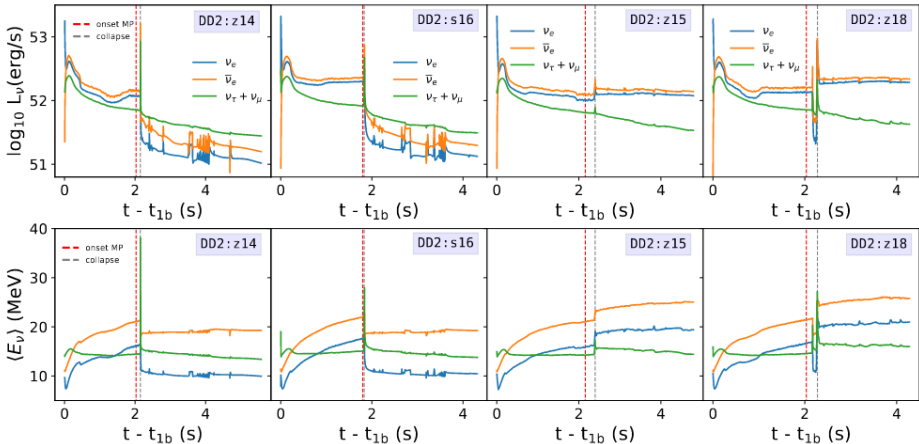


Compactness parameter

$$\xi_M = \frac{M/M_{\odot}}{r(M)/1000 \text{ km}}$$

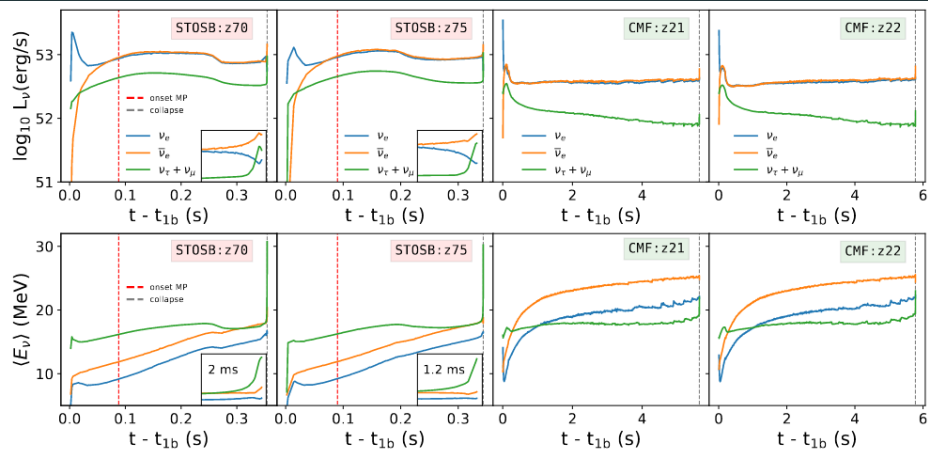
- Three equations of state: DD2F (1st order PT, Gibbs), STOS-B145 (1st order PT, Maxwellian), and CMF (smooth crossover)
- Successful explosions only for 2 models in DD2F
- Failed explosions in DD2F and STOS-B145

Neutrino signals: DD2F



- Low explosion energies $\sim 10^{50}$ erg
- Majority of models have second bounce 37/40
- Failed explosions only for zero metallicity

Neutrino signals: STOS-B145, CMF



- Relatively small increase in luminosity during 2nd bounce
- No models successfully explode
- No 2nd bounces in the CMF models

Quark deconfinement as a supernova explosion engine for massive blue supergiant stars

