

# Interpreting Flavor and Electroweak Puzzles

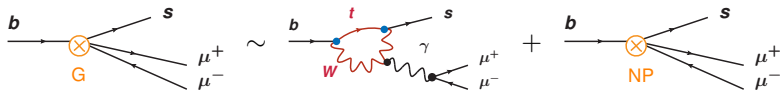
Wolfgang Altmannshofer  
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14th Conference on the Intersections of Particle and Nuclear Physics  
Lake Buena Vista, August 29 - September 4, 2022

# Basic Idea Behind Indirect Probes of New Physics

Example: Rare  $B$  decays



$$G \sim \frac{1}{16\pi^2} \frac{g^4}{m_W^2} \frac{m_t^2}{m_W^2} V_{tb} V_{ts}^* + \frac{C_{NP}}{\Lambda_{NP}^2}$$

measure  
precisely

calculate precisely  
the SM contribution

get information on  
NP coupling and scale

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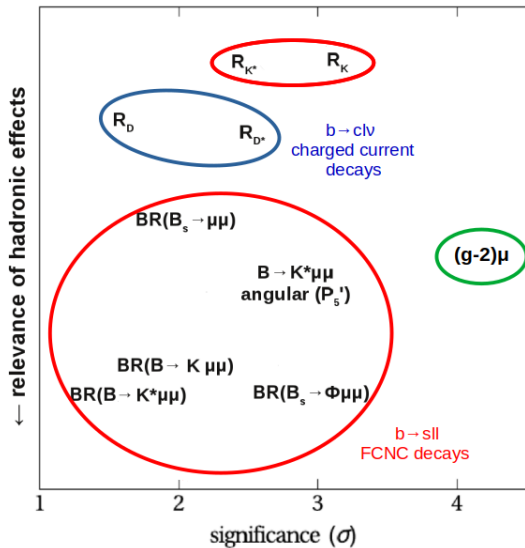
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get information on  
NP coupling and scale

Anomalies at low energies can establish a new scale in particle physics  
 $\Rightarrow$  “no-loose theorems”, “guaranteed” discoveries at colliders, ...

(at least in principle)

# Anomalies and Puzzles in 2022



see talks by:

Asutosh Kotwal

Martin Fertl

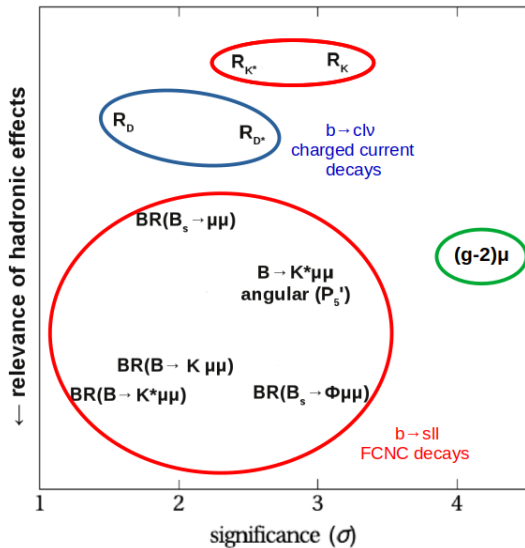
Soeren Prell

Carla Marin Benito (Fr)

(plot inspired by

Zoltan Ligeti)

# Anomalies and Puzzles in 2022



$m_W$

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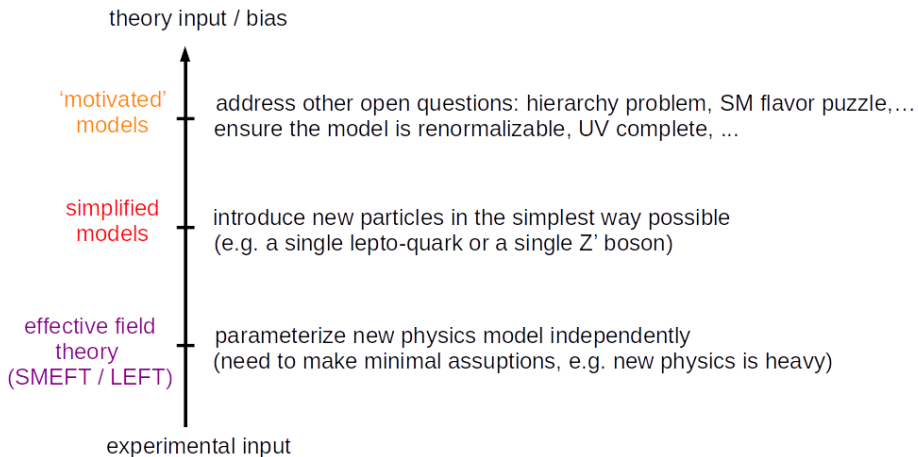
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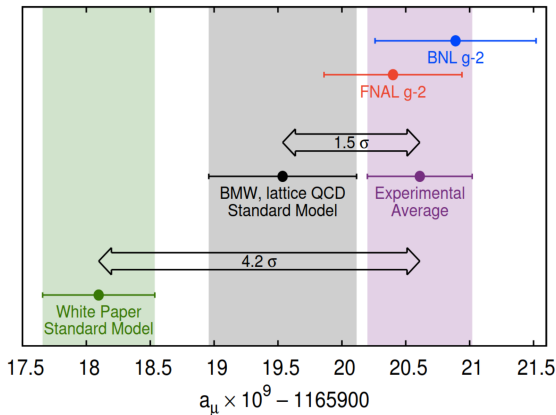
# Bottom-Up Approach to the Anomalies



(inspired by Marco Nardecchia)

# Implications of the muon $g-2$

# Anomalous Magnetic Moment of the Muon



**4.2  $\sigma$  discrepancy between the experimental average (Fermilab g-2, 2104.03281) and the SM consensus (Aoyama et al. 2006.04822)**

(see, however, the lattice results 2002.12347, 2206.06582, 2206.15084, 2207.04765)

$$\Delta a_\mu = (251 \pm 59) \times 10^{-11}$$



# Model Independent Analysis and New Physics Scale

The **leading effective operator** that modifies the anomalous magnetic moment of the muon and that respects  $SU(2)_L \times U(1)_Y$

$$\mathcal{L}_{\text{eff}} = \frac{C}{\Lambda_{\text{NP}}^2} H(\bar{\mu}\sigma_{\alpha\beta}\mu)F^{\alpha\beta} \quad \Rightarrow \quad \Delta a_\mu \simeq \frac{4m_\mu v C}{e\sqrt{2}\Lambda_{\text{NP}}^2}$$

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strong coupling  $\frac{1}{\Lambda_{\text{NP}}^2} H(\bar{\mu}\sigma_{\alpha\beta}\mu) F^{\alpha\beta}$   $\Lambda_{\text{NP}} \simeq 290 \text{ TeV}$

weak coupling  $\frac{e}{16\pi^2} \frac{1}{\Lambda_{\text{NP}}^2} H(\bar{\mu}\sigma_{\alpha\beta}\mu) F^{\alpha\beta}$   $\Lambda_{\text{NP}} \simeq 14 \text{ TeV}$

weak coupling + MFV  $\frac{ey_\mu}{16\pi^2} \frac{1}{\Lambda_{\text{NP}}^2} H(\bar{\mu}\sigma_{\alpha\beta}\mu) F^{\alpha\beta}$   $\Lambda_{\text{NP}} \simeq 280 \text{ GeV}$

(MFV = Minimal Flavor Violation)

# New Physics Models for $(g - 2)_\mu$

- ▶ In the strongly coupled case, the new physics scale could be extremely high, outside the reach of current and future colliders. (However, I am not aware of any actual model)
- ▶ Most explanations of  $(g - 2)_\mu$  predict:

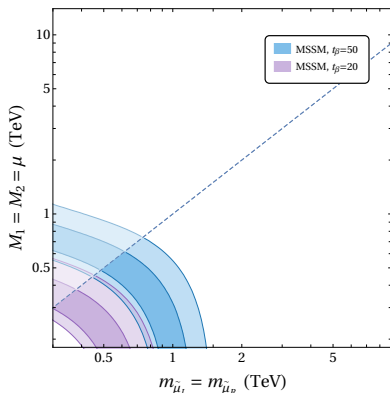
new physics not far above the electroweak scale  
("heavy new physics": SUSY, leptoquarks,  $Z'$ , ...)

or

new physics considerably below the electroweak scale  
("light new physics": dark photons, axions, light  $Z'$ , ...)

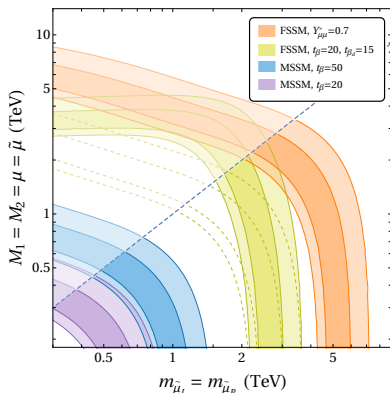
# Heavy New Physics Example: SUSY

- ▶ It is very well known that the **MSSM** can give sizeable contributions to  $(g-2)_\mu$  via  $\tan\beta$  enhanced slepton chargino/neutralino loops  
Athron et al. 2104.03691 + many others  
(apologies for the omission)
- ▶ S sleptons, charginos, neutralinos need to be pretty light
- ▶ **Compressed spectra** to avoid existing LHC constraints
- ▶ Good discovery prospects at the high luminosity LHC and  $e^+e^-$  colliders (ILC, CLIC)



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- ▶ In **non-minimal SUSY scenarios**, sleptons, charginos, neutralinos can be significantly heavier

WA, Gadam, Gori, Hamer 2104.08293

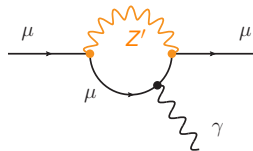
# Light New Physics Example: $L_\mu - L_\tau$

New gauge bosons are well known candidates to explain  $(g - 2)_\mu$

(e.g. Greljo et al. 2203.13731)

Dark photons have been ruled out for quite a while

Gauged  $L_\mu - L_\tau$  is one of the least constrained options



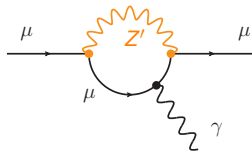
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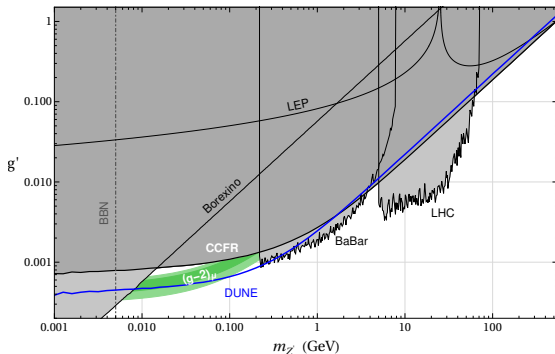
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WA, Gori, Pospelov, Yavin, 1406.2332;

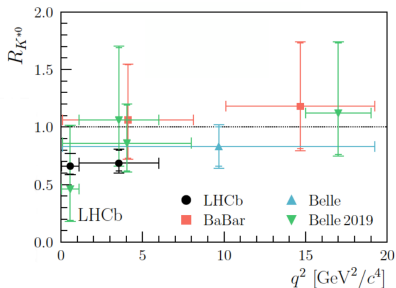
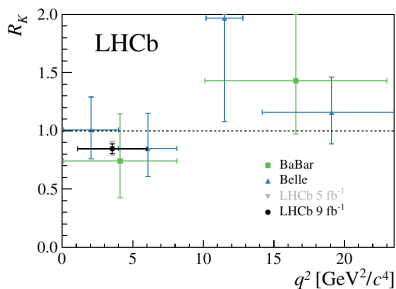
WA, Gori, Martin-Albo, Sousa, Wallbank 1902.06765



Implications of the  
 $b \rightarrow sll$  Anomalies  
( $R_K$ ,  $R_{K^*}$  and Friends)



# Evidence for Lepton Flavor Universality Violation



$$R_{K^{(*)}} = \frac{BR(B \rightarrow K^{(*)} \mu \mu)}{BR(B \rightarrow K^{(*)} e e)} \stackrel{\text{SM}}{\simeq} 1$$

$$R_{K^+}^{[1,6]} = 0.846_{-0.039-0.012}^{+0.042+0.013} \quad (3.1\sigma)$$

$$R_{K^{*0}}^{[0.045,1.1]} = 0.66_{-0.07}^{+0.11} \pm 0.03 \quad (\sim 2.5\sigma)$$

$$R_{K^{*0}}^{[1.1,6]} = 0.69_{-0.07}^{+0.11} \pm 0.05 \quad (\sim 2.5\sigma)$$

$$R_{K_S}^{[1.1,6]} = 0.66_{-0.14-0.04}^{+0.20+0.02} \quad (\sim 1.5\sigma)$$

$$R_{K^{*+}}^{[0.045,6]} = 0.70_{-0.13-0.04}^{+0.18+0.03} \quad (\sim 1.5\sigma)$$

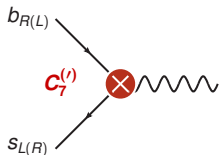
$$R_{\rho K}^{[0.1,6]} = 0.86_{-0.11}^{+0.14} \pm 0.05 \quad (\sim 1\sigma)$$

LHCb 2103.11769, LHCb 1705.05802, 1912.08139, 2110.09501; also Belle 1904.02440, 1908.01848

# Model Independent Analysis

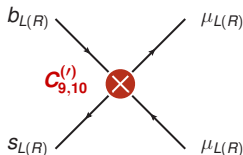
$$\mathcal{H}_{\text{eff}}^{b \rightarrow s} = -\frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \frac{e^2}{16\pi^2} \sum_i (C_i \mathcal{O}_i + C'_i \mathcal{O}'_i)$$

magnetic dipole operators



$$C_7^{(i)} (\bar{s} \sigma_{\mu\nu} P_{R(L)} b) F^{\mu\nu}$$

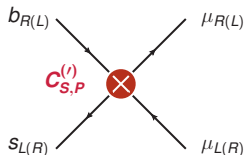
semileptonic operators



$$C_9^{(i)} (\bar{s} \gamma_\mu P_{L(R)} b) (\bar{\mu} \gamma^\mu \mu)$$

$$C_{10}^{(i)} (\bar{s} \gamma_\mu P_{L(R)} b) (\bar{\mu} \gamma^\mu \gamma_5 \mu)$$

scalar operators

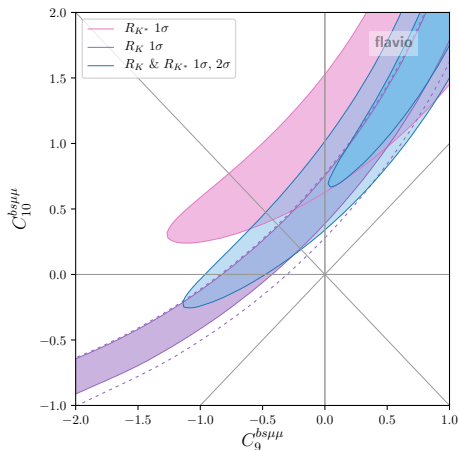


$$C_S^{(i)} (\bar{s} P_{R(L)} b) (\bar{\mu} P_{L(R)} \mu)$$

neglecting tensor operators and additional scalar operators

(they are dimension 8 in SMEFT: Alonso, Grinstein, Martin Camalich 1407.7044)

# Global Fits of Rare $b \rightarrow sl\ell$ Decays



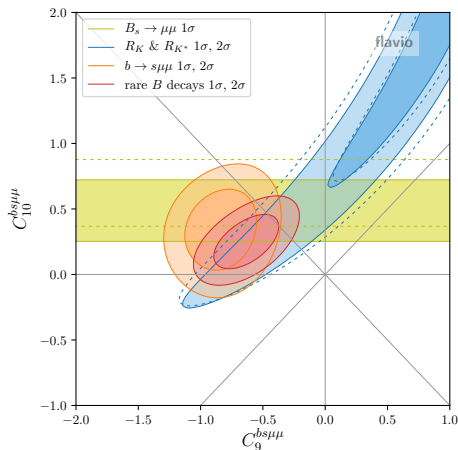
$$C_9^{bs\mu\mu}(\bar{s}\gamma_\alpha P_L b)(\bar{\mu}\gamma^\alpha \mu)$$

$$C_{10}^{bs\mu\mu}(\bar{s}\gamma_\alpha P_L b)(\bar{\mu}\gamma^\alpha \gamma_5 \mu)$$

● LFU ratios

WA, Stangl 2103.13370 (other recent fits: Geng et al.  
 2103.12738; Cornella et al. 2103.16558; Alguero et al.  
 2104.08921; Hurth et al. 2104.10058; Gubernari et al.  
 2206.03797)

# Global Fits of Rare $b \rightarrow s\ell\ell$ Decays



$$C_9^{bs\mu\mu}(\bar{s}\gamma_\alpha P_L b)(\bar{\mu}\gamma^\alpha \mu)$$

$$C_{10}^{bs\mu\mu}(\bar{s}\gamma_\alpha P_L b)(\bar{\mu}\gamma^\alpha \gamma_5 \mu)$$

- LFU ratios
- $B_s \rightarrow \mu^+ \mu^-$  branching ratio  
(with latest CMS update probably compatible with SM-like  $C_{10}$ )
- $b \rightarrow s\mu\mu$  observables
- overall remarkable consistency

WA, Stangl 2103.13370 (other recent fits: Geng et al. 2103.12738; Cornella et al. 2103.16558; Alguero et al. 2104.08921; Hurth et al. 2104.10058; Gubernari et al. 2206.03797)

# The New Physics Scale

unitarity bound  $\frac{4\pi}{\Lambda_{\text{NP}}^2} (\bar{s}\gamma_\nu P_L b)(\bar{\mu}\gamma^\nu \mu)$   $\Lambda_{\text{NP}} \simeq 120 \text{ TeV} \times (C_9^{\text{NP}})^{-1/2}$

generic tree  $\frac{1}{\Lambda_{\text{NP}}^2} (\bar{s}\gamma_\nu P_L b)(\bar{\mu}\gamma^\nu \mu)$   $\Lambda_{\text{NP}} \simeq 35 \text{ TeV} \times (C_9^{\text{NP}})^{-1/2}$

MFV tree  $\frac{1}{\Lambda_{\text{NP}}^2} V_{tb} V_{ts}^* (\bar{s}\gamma_\nu P_L b)(\bar{\mu}\gamma^\nu \mu)$   $\Lambda_{\text{NP}} \simeq 7 \text{ TeV} \times (C_9^{\text{NP}})^{-1/2}$

generic loop  $\frac{1}{\Lambda_{\text{NP}}^2} \frac{1}{16\pi^2} (\bar{s}\gamma_\nu P_L b)(\bar{\mu}\gamma^\nu \mu)$   $\Lambda_{\text{NP}} \simeq 3 \text{ TeV} \times (C_9^{\text{NP}})^{-1/2}$

MFV loop  $\frac{1}{\Lambda_{\text{NP}}^2} \frac{1}{16\pi^2} V_{tb} V_{ts}^* (\bar{s}\gamma_\nu P_L b)(\bar{\mu}\gamma^\nu \mu)$   $\Lambda_{\text{NP}} \simeq 0.6 \text{ TeV} \times (C_9^{\text{NP}})^{-1/2}$

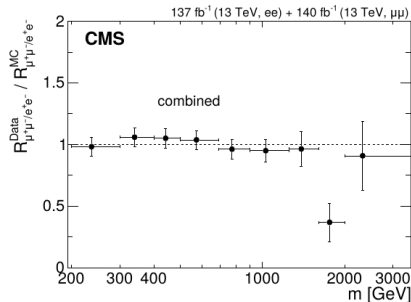
(MFV = Minimal Flavor Violation)

# Model Independent Approach at the LHC

If the new physics is not accessible directly at the LHC, high energy tails of di-lepton spectra are in principle still affected

(Greljo, Marzocca 1704.09015)

$$R = \frac{\sigma(pp \rightarrow \mu\mu)}{\sigma(pp \rightarrow ee)}$$



CMS 2103.02708 (also ATLAS 2105.13847)

$$C_9^{bs\mu\mu}(\bar{s}\gamma_\alpha P_L b)(\bar{\mu}\gamma^\alpha \mu)$$

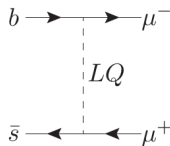
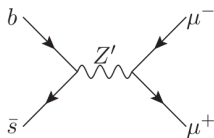
$$C_{10}^{bs\mu\mu}(\bar{s}\gamma_\alpha P_L b)(\bar{\mu}\gamma^\alpha \gamma_5 \mu)$$

- ▶ flavor changing operators are probed up to scales of few TeV
- ▶ **order of magnitude is missing** to probe the  $b \rightarrow sll$  anomalies
- would need a 100 TeV collider

# Simplified Models for $R_K$ and $R_{K^*}$

possible tree level explanations:

- ▶  $Z'$  Bosons
- ▶ Lepto-Quarks



upper bounds on flavor violating couplings from  $B_s$  mixing imply  
**upper bounds on the particle masses** (e.g. Di Luzio et al. 1909.11087)

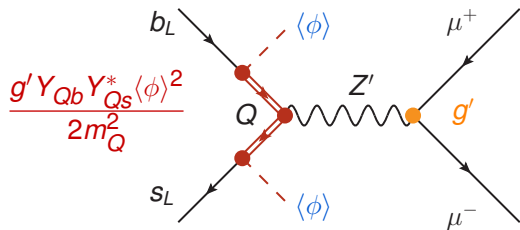
- ▶  $m_{Z'} \lesssim g_\mu \times 5\text{TeV}$
- ▶  $m_{LQ} \lesssim (30 - 60)\text{TeV}$  (depending on the lepto-quark representation)

→ a weakly coupled  $Z'$  might be in reach of the LHC

# My Favorite $Z'$ Model

$Z'$  based on gauged  $L_\mu - L_\tau$  (He, Joshi, Lew, Volkas PRD 43, 22-24)  
with effective flavor violating couplings to quarks

WA, Gori, Pospelov, Yavin 1403.1269; WA, Yavin 1508.07009



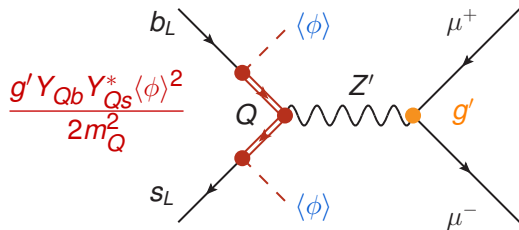
$Q$ : heavy vectorlike fermions with mass  $\sim 1 - 10$  TeV  
 $\phi$ : scalar that breaks  $L_\mu - L_\tau$



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predicted Lepton  
Universality Violation!

**Q:** heavy vectorlike fermions with mass  $\sim 1 - 10$  TeV  
 $\phi$ : scalar that breaks  $L_\mu - L_\tau$

# Probing the $L_\mu - L_\tau$ Parameter Space

WA, Gori, Martin-Albo, Sousa, Wallbank 1902.06765

Neutrino Tridentes

$B_s$  mixing

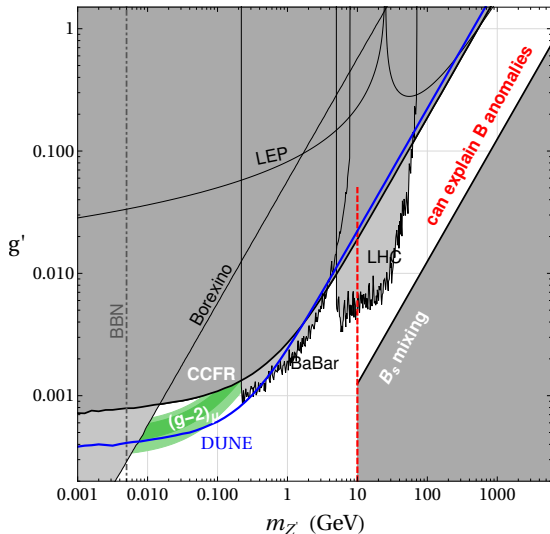
$(g-2)_\mu$

$\nu e$  scattering

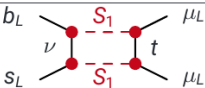
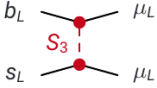
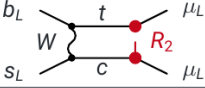

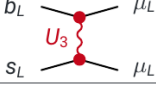
$Z \rightarrow \ell\ell$

$Z \rightarrow 4\mu$

$e^+e^- \rightarrow 4\mu$



# Simplified Leptoquark Models

Spin	$G_{\text{SM}}$	Name	Characteristic process	First time used for $b \rightarrow s\mu\mu$
0	$(\bar{3}, 1)_{1/3}$	$S_1$		Bauer, Neubert, arXiv:1511.01900
0	$(\bar{3}, 3)_{1/3}$	$S_3$		Hiller, Schmaltz, arXiv:1408.1627
0	$(3, 2)_{7/6}$	$R_2$		Bećirević, Sumensari, arXiv:1704.05835
1	$(3, 1)_{2/3}$	$U_1$		Barbieri et al., arXiv:1512.01560
1	$(3, 3)_{2/3}$	$U_3$		Fajfer, Košnik, arXiv:1511.06024

from talk by Peter Stangl LF(U)V workshop, Zurich, July 4

(the loop level leptoquarks struggle to accommodate the anomalies)

# Leptoquark Signatures at the LHC

e.g. Allanach, Gripaos, You 1710.06363, Hiller, Loose, Nisandzic 1801.09399

- Leptoquarks are **pair produced** through QCD interactions

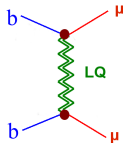
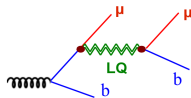
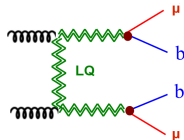
$$pp \rightarrow LQ LQ \rightarrow j(b)\mu^+ j(b)\mu^-$$

- Leptoquarks can be **singly produced** through their couplings to quarks/leptons

$$pp \rightarrow LQ \mu \rightarrow j(b)\mu^+ \mu^-$$

- Leptoquarks contribute to di-muon production

$$pp \rightarrow \mu^+ \mu^-$$



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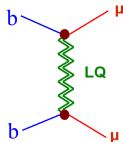
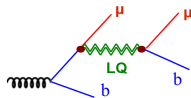
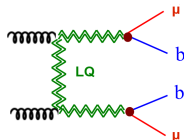
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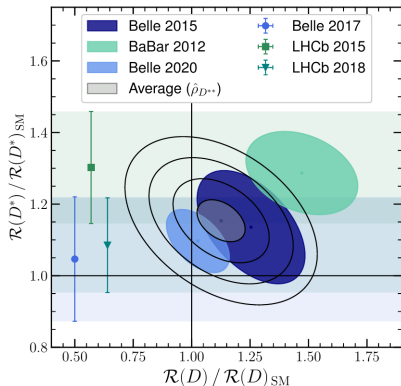
Also: excellent prospects to see these leptoquarks at a muon collider

Huang, Jana, Queiroz, Rodejohann 2103.01617, Asadi, Capdevilla, Cesarotti, Homiller 2104.05720

Implications of the  
 $b \rightarrow c\tau\nu$  Anomalies  
( $R_D, R_{D^*}$ )

# LFU in Charged Current Decays: $R_D$ and $R_{D^*}$

Bernlochner, Franco Sevilla, Robinson, 2101.08326



$$R_D = \frac{BR(B \rightarrow D\tau\nu)}{BR(B \rightarrow D\ell\nu)}$$

$$R_{D^*} = \frac{BR(B \rightarrow D^*\tau\nu)}{BR(B \rightarrow D^*\ell\nu)}$$

$$\begin{aligned} \ell = \mu, e & \quad (\text{BaBar/Belle}) \\ \ell = \mu & \quad (\text{LHCb}) \end{aligned}$$

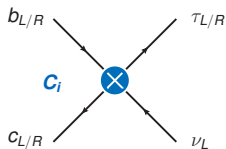
$$R_D^{\text{exp}}/R_D^{\text{SM}} = 1.13 \pm 0.10, \quad R_{D^*}^{\text{exp}}/R_{D^*}^{\text{SM}} = 1.15 \pm 0.06$$

**combined discrepancy with the SM:  $3.6\sigma$**

(the heavy flavor averaging group quotes  $3.1\sigma$ )

# Model Independent Analysis

$$\mathcal{H}_{\text{eff}} = \frac{4G_F}{\sqrt{2}} V_{cb} \mathcal{O}_{V_L} + \frac{1}{\Lambda^2} \sum_i C_i \mathcal{O}_i$$

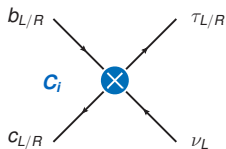


$\mathcal{O}_i =$  contact interactions  
with vector, scalar  
or tensor currents



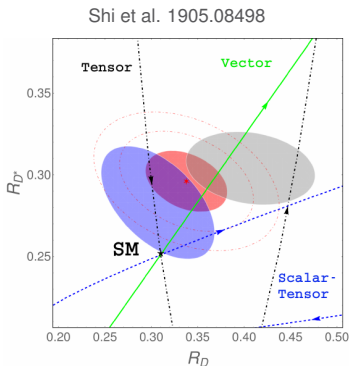
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$\mathcal{O}_i$  = contact interactions  
with vector, scalar  
or tensor currents

rescaling of the **SM vector operator** fits the data best  
combinations of operators  
are also possible



(also Murgui et al. 1904.09311, Asadi, Shih 1905.03311,  
Cheung et al. 2002.07272, ... )

# The New Physics Scale

unitarity bound  $\frac{4\pi}{\Lambda_{\text{NP}}^2} (\bar{c}\gamma_\nu P_L b)(\bar{\tau}\gamma^\nu P_L \nu)$   $\Lambda_{\text{NP}} \simeq 8.4 \text{ TeV}$

generic tree  $\frac{1}{\Lambda_{\text{NP}}^2} (\bar{c}\gamma_\nu P_L b)(\bar{\tau}\gamma^\nu P_L \nu)$   $\Lambda_{\text{NP}} \simeq 2.4 \text{ TeV}$

MFV tree  $\frac{1}{\Lambda_{\text{NP}}^2} V_{cb} (\bar{c}\gamma_\nu P_L b)(\bar{\tau}\gamma^\nu P_L \nu)$   $\Lambda_{\text{NP}} \simeq 0.5 \text{ TeV}$

(MFV = Minimal Flavor Violation)

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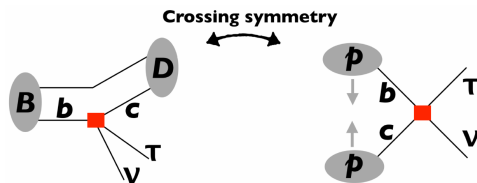
(MFV = Minimal Flavor Violation)

rather low scale  $\rightarrow$  model building is non-trivial

# Model Independent Approach at the LHC

Expect non-standard  
**mono-tau production**  
at the LHC

(possibly in association  
with b-jets)

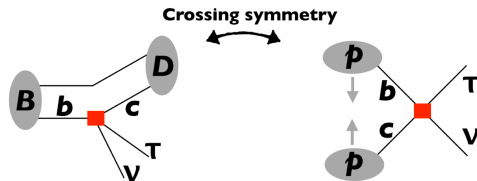


WA, Dev, Soni 1704.06659; Greljo et al. 1811.07920;  
Marzocca et al. 2008.07541; ...

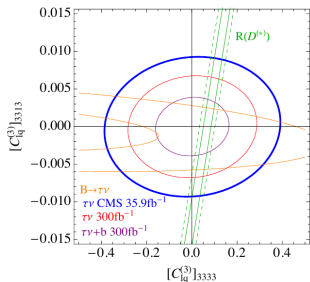
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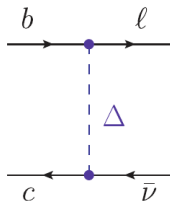
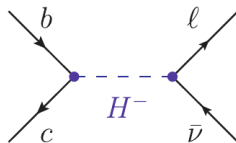
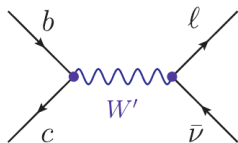
WA, Dev, Soni 1704.06659; Greljo et al. 1811.07920;  
 Marzocca et al. 2008.07541; ...



- ▶ Collider and low energy sensitivities are complementary
- ▶ High-luminosity LHC can probe relevant parts of parameter space

Need a tree level mediator: 3 options

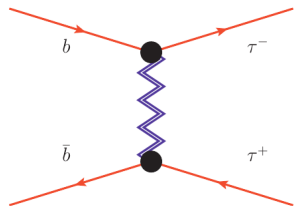
- 1)  $W'$  bosons excluded by direct searches
- 2) **Charged Higgs** bosons strongly constrained by  $B_c \rightarrow \tau \nu$  and  $B \rightarrow D^{(*)} \tau \nu$  kinematic distributions
- 3) **Leptoquarks** that couple dominantly to the 3rd generation can work.



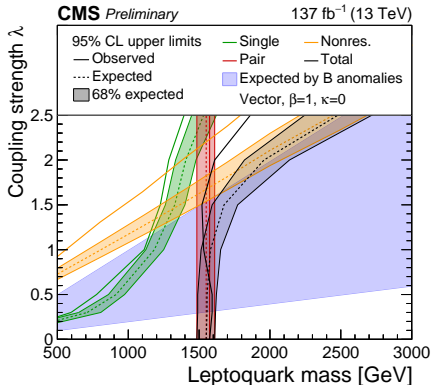
# Collider Signature of the Leptoquarks

- Robust collider signature of leptoquarks that explain  $R_D$  and  $R_{D^*}$ : non-standard di-tau production at high invariant mass

Faroughy et al. 1609.07138



CMS-PAS-EXO-19-016



# Combined Explanations of the B anomalies

- ▶  $U_1$  leptoquark can simultaneously explain  $R_{K^{(*)}}$  and  $R_{D^{(*)}}$  (recent studies: Cornella et al. 2103.16558; Angelescu et al. 2103.12504)
- ▶  $U_1$  could be the remnant of an extended gauge group: “4321 models”, (Pati-Salam)<sup>3</sup> models (Di Luzio et al. 1708.08450; Bordone et al. 1712.01368, ...)

Model	$R_{K^{(*)}}$	$R_{D^{(*)}}$
$S_3$ ( $\bar{\mathbf{3}}, \mathbf{3}, 1/3$ )	✓	✗
$S_1$ ( $\bar{\mathbf{3}}, \mathbf{1}, 1/3$ )	✗	✓
$R_2$ ( $\mathbf{3}, \mathbf{2}, 7/6$ )	✗	✓
$U_1$ ( $\mathbf{3}, \mathbf{1}, 2/3$ )	✓	✓
$U_3$ ( $\mathbf{3}, \mathbf{3}, 2/3$ )	✓	✗



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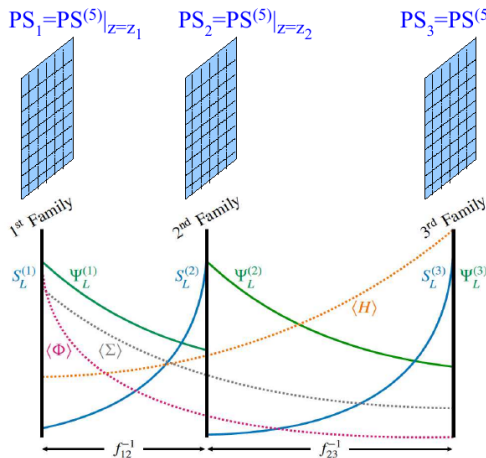
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$U_1$ ( $\mathbf{3}, \mathbf{1}, 2/3$ )	✓	✓
$U_3$ ( $\mathbf{3}, \mathbf{3}, 2/3$ )	✓	✗

- ▶ also attempts for simultaneous explanations in **RPV SUSY**  
Deshpande, He, 1608.04817; WA, Dev, Soni 1704.06659; Earl, Gregoire 1806.01343;  
Trifinopoulos 1807.01638; WA, Dev, Soni, Sui 2002.12910; Dev, Soni, Xu 2106.15647; ...

# Fleshed Out (Pati-Salam)<sup>3</sup> Model

Flavor anomalies from the  $U_1$  leptoquark of (Pati-Salam)<sup>3</sup>



Flavor  $\leftrightarrow$  special position (*topological defect*) in an extra (compact) space-like dimension

Dvali & Shifman, '00

Higgs and SU(4)-breaking fields with oppositely-peaked profiles, leading to the desired flavor pattern for masses & anomalies

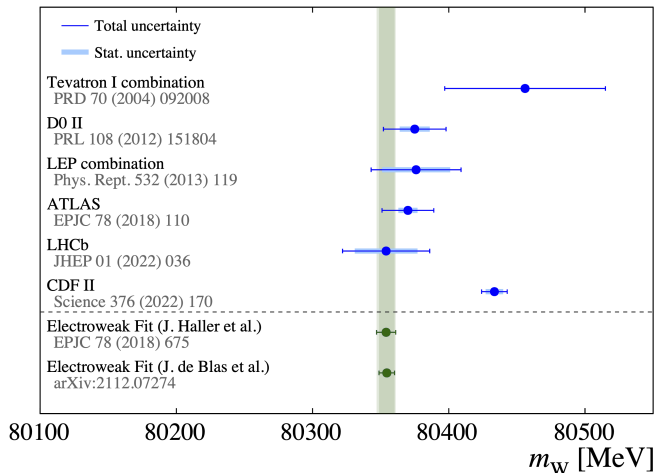
Bordone, Cornella, Fuentes-Martin, GI '17  
Fuentes-Martin, GI, Pages, Stefánek '20

Possible to implement anarchic neutrino masses via an inverse see-saw mechanism

(talk by Gino Isidori @ Beyond the Anomalies workshop, Durham 2021)

# Implications of the $W$ Mass

# W Mass Measurements



CDF measurement is  $7\sigma$  away from the SM prediction !?!

# The W Mass and New Physics

- ▶ The SM predicts a relation between the  $W$  mass, the  $Z$  mass, and weak mixing angle (precise relation is subject to higher order correction, choice of renormalization scheme, ...)

$$\frac{m_W^2}{m_Z^2 \cos^2 \theta_W} = \rho \simeq 1$$

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- ▶ changing the **weak mixing angle** (e.g. by modifying  $Z$  couplings)



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$$\frac{m_W^2}{m_Z^2 \cos^2 \theta_W} = \rho \simeq \frac{\sum_i (l_i(l_i + 1) - Y_i^2) v_i^2}{\sum_i 2 Y_i^2 v_i^2}$$

New physics could enter by:

- ▶ changing the **mass of the  $W$**  (e.g. mixing of  $W$  with a  $W'$ )
- ▶ changing the **mass of the  $Z$**  (e.g. mixing of  $Z$  with a  $Z'$ )
- ▶ changing the **weak mixing angle** (e.g. by modifying  $Z$  couplings)
- ▶ **changing the relation itself** (e.g. exotic Higgs sectors)

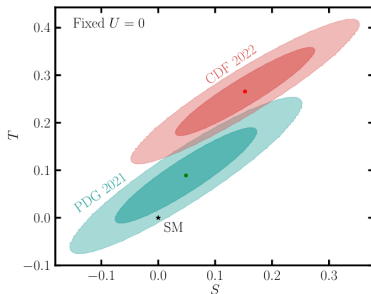
# Model Independent Approach

- ▶ New physics is often model independently described by **oblique corrections**

$$S \frac{gg'}{16\pi} \frac{1}{v^2} (H^\dagger \sigma^a H) W_{\mu\nu}^a B^{\mu\nu}$$

$$T \frac{e^2}{8\pi} \frac{1}{v^2} (H^\dagger \overleftrightarrow{D}_\mu H)^2$$

+ ...



- ▶ Need a generic **new physics scale of a few TeV**

Lu et al. 2204.03796

(also de Blas et al. 2204.04204; Strumia 2204.04191;

... + many others, apologies for the omission)

# Connection with the B Anomalies?

- ▶  $Z'$  models that can explain the  $b \rightarrow s\ell\ell$  anomalies might also explain shifts in the  $W$  mass
- ▶ Example 1: gauge a linear combination of hypercharge, baryon number, and individual lepton numbers

WA, Davighi, Nardecchia 1909.02021

$$X = a_Y Y - a_e(B/3 - L_e) - a_\mu(B/3 - L_\mu) - a_\tau(B/3 - L_\tau)$$

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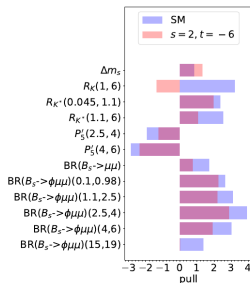
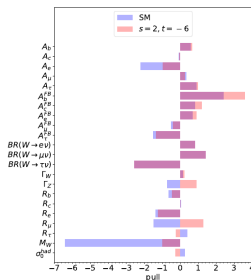
WA, Davighi, Nardecchia 1909.02021

$$X = a_Y Y - a_e (B/3 - L_e) - a_\mu (B/3 - L_\mu) - a_\tau (B/3 - L_\tau)$$

- ▶ Example 2: gauge a linear combination of third generation hypercharge, baryon number, and lepton number

Allanach, Davighi 2205.12252

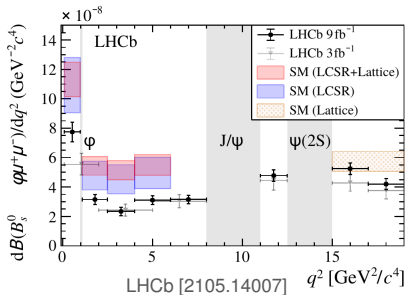
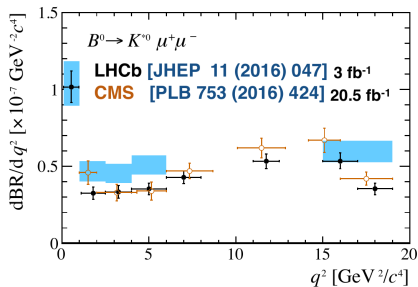
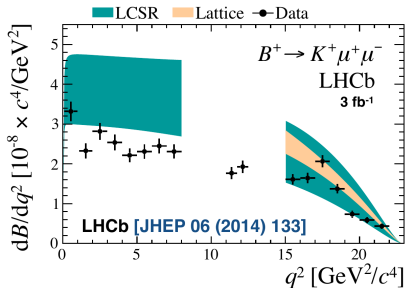
$$X = sY_3 + t(B_3 - L_3)$$



- ▶ Anomalies might be indirect signs of physics beyond the standard model.
  - ▶ Anomalies could establish a new mass scale in particle physics
- would have a transformative impact:  
motivate a large new physics model building effort  
and provide targets for direct searches at the LHC  
and future colliders

Back Up

# $b \rightarrow s\mu\mu$ Branching Ratios



Experimental results for

$$\text{BR}(B \rightarrow K\mu\mu)$$

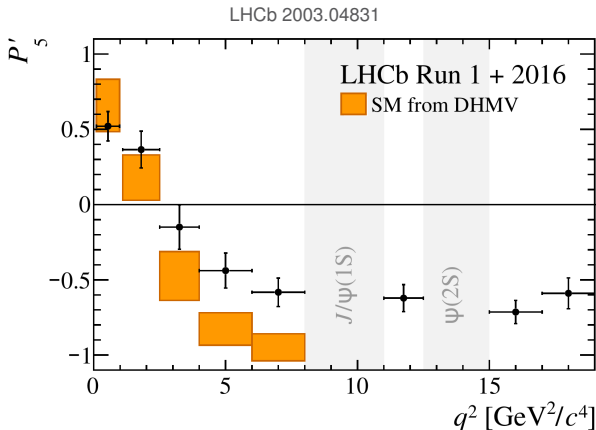
$$\text{BR}(B \rightarrow K^*\mu\mu)$$

$$\text{BR}(B_s \rightarrow \phi\mu\mu)$$

are consistently low  
across many  $q^2$  bins

# The $P'_5$ Anomaly

$P'_5 \sim$  a moment of the  $B \rightarrow K^* \mu^+ \mu^-$  angular distribution



**Anomaly persists** in the latest update of  $B^0 \rightarrow K^{*0} \mu^+ \mu^-$  with 2016 data.  
(Anomaly also seen in  $B^\pm \rightarrow K^{*\pm} \mu^+ \mu^-$  LHCb 2012.13241)



# Non-Standard $\mu^+\mu^- \rightarrow bs$ at a Muon Collider

$$\frac{d\sigma(\mu^+\mu^- \rightarrow b\bar{s})}{d\cos\theta} = \frac{3}{16}\sigma(\mu^+\mu^- \rightarrow bs) \left(1 + \cos^2\theta + \frac{8}{3}A_{\text{FB}}\cos\theta\right)$$

$$\frac{d\sigma(\mu^+\mu^- \rightarrow \bar{b}s)}{d\cos\theta} = \frac{3}{16}\sigma(\mu^+\mu^- \rightarrow bs) \left(1 + \cos^2\theta - \frac{8}{3}A_{\text{FB}}\cos\theta\right)$$

Total cross section **increases with the center of mass energy**

$$\sigma(\mu^+\mu^- \rightarrow bs) = \frac{G_F^2\alpha^2}{8\pi^3} |V_{tb}V_{ts}^*|^2 s \left(|C_9|^2 + |C_{10}|^2\right)$$

# Non-Standard $\mu^+ \mu^- \rightarrow bs$ at a Muon Collider

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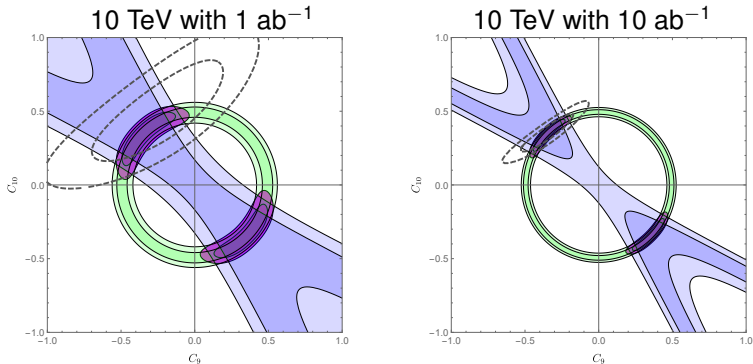
Forward backward asymmetry is sensitive to the **chirality structure**

$$A_{\text{FB}} = \frac{-3\text{Re}(C_9 C_{10}^*)}{2(|C_9|^2 + |C_{10}|^2)}$$

Need **charge tagging** to measure the forward backward asymmetry

# Sensitivity Projections

WA, Gadam, Profumo 2203.07495 and in preparation



- branching ratio (green) and forward backward asymmetry (blue) are highly complementary
- 10 TeV muon collider has better sensitivity than the current and projected rare B decay results (dashed)

(see also Huang et al. 2103.01617; Asadi et al. 2104.05720

Azatov et al. 2205.13552 for related studies)