

**N3AS – iTHEMS Meeting on
Quantum Information Science in Multimessenger Astrophysics
June 16–18, 2024, RIKEN, Japan**

Structure learning ansatz for Lipkin model

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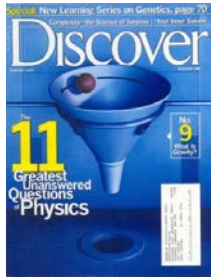
June 16, 2024



東京大学
THE UNIVERSITY OF TOKYO

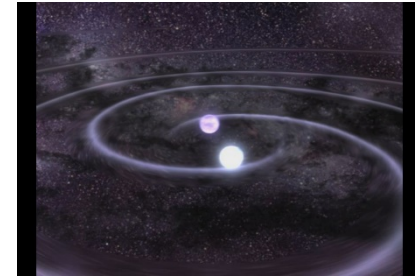
r -process nucleosynthesis & nuclear β decays

The 11 greatest unanswered questions of physics

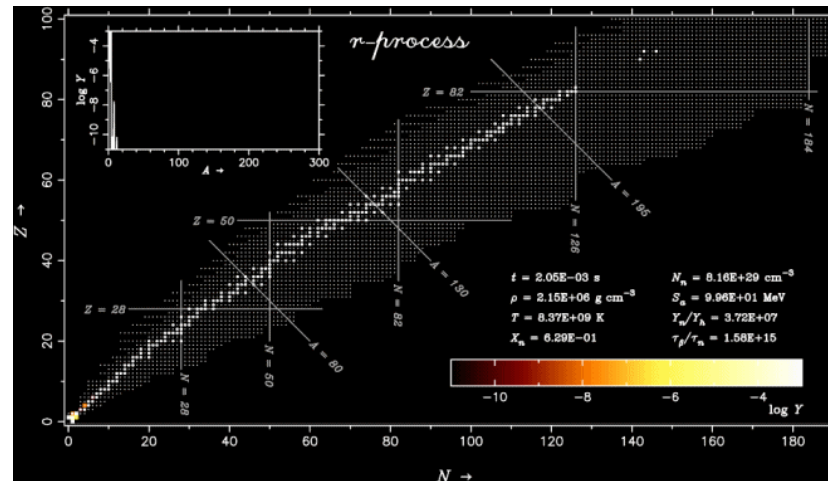


Question 3

How were the heavy elements from iron to uranium made?



Rapid neutron-capture process (r -process)

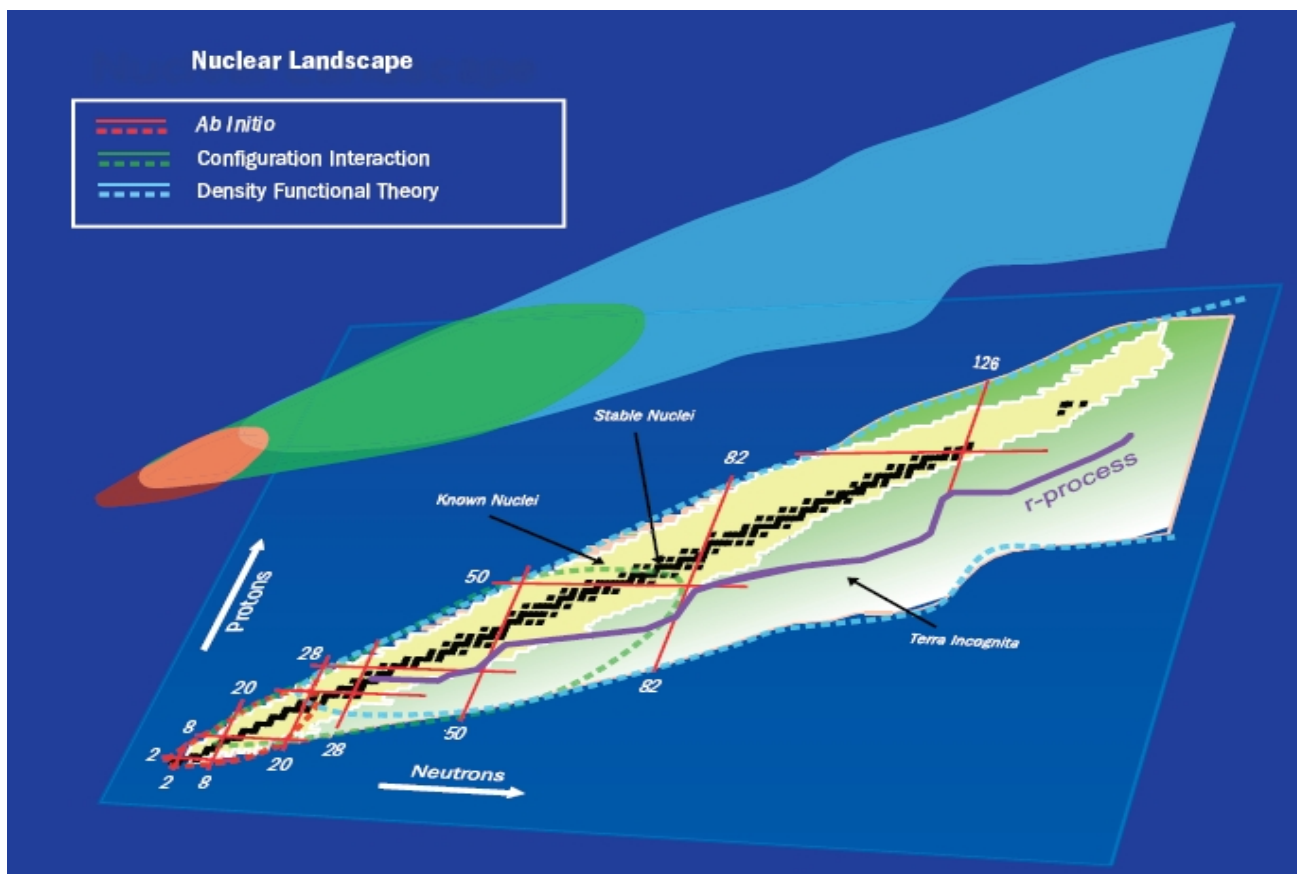


Courtesy of S. Wanajo

- Nuclear masses \rightarrow **path** of r -process
- β -decay rates \rightarrow **timescale** of r -process
- β -delayed n -emission branchings \rightarrow **final abundance pattern** of r -process

Key exp. @ RIKEN

State-of-the-art nuclear methodologies

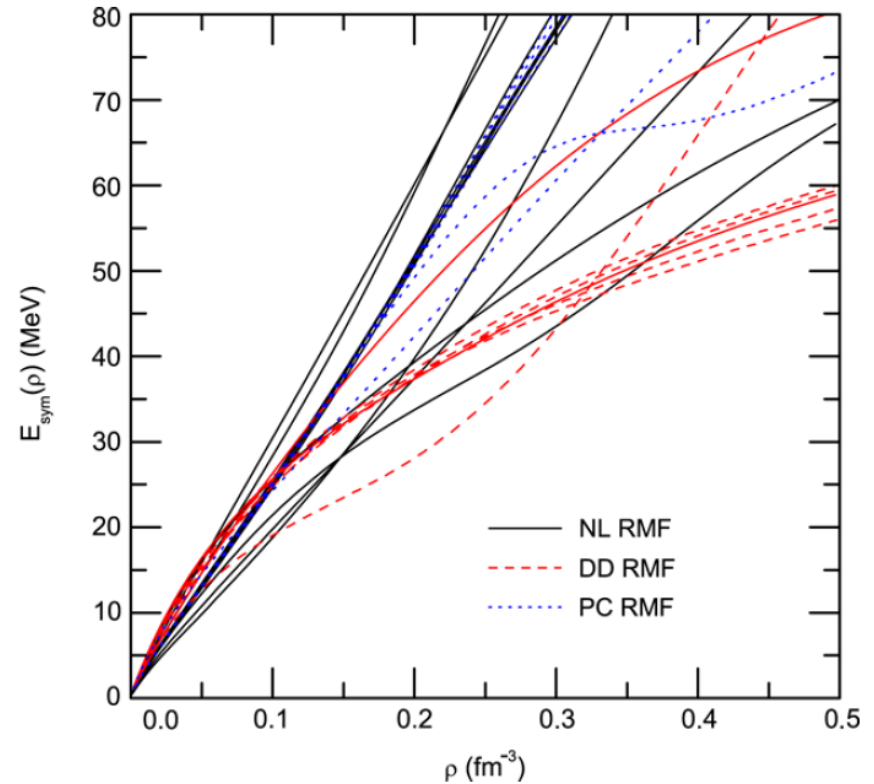
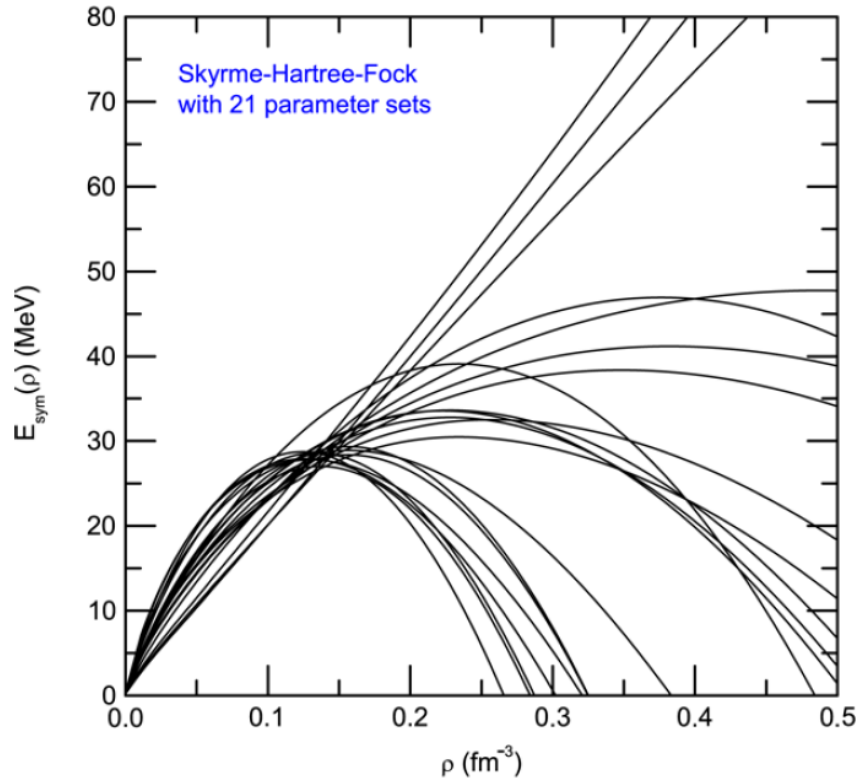


<http://www.unedf.org/>

- Density functional theory (DFT) aims at understanding both ground-state and excited-state properties of thousands of nuclei in a consistent and predictive way.

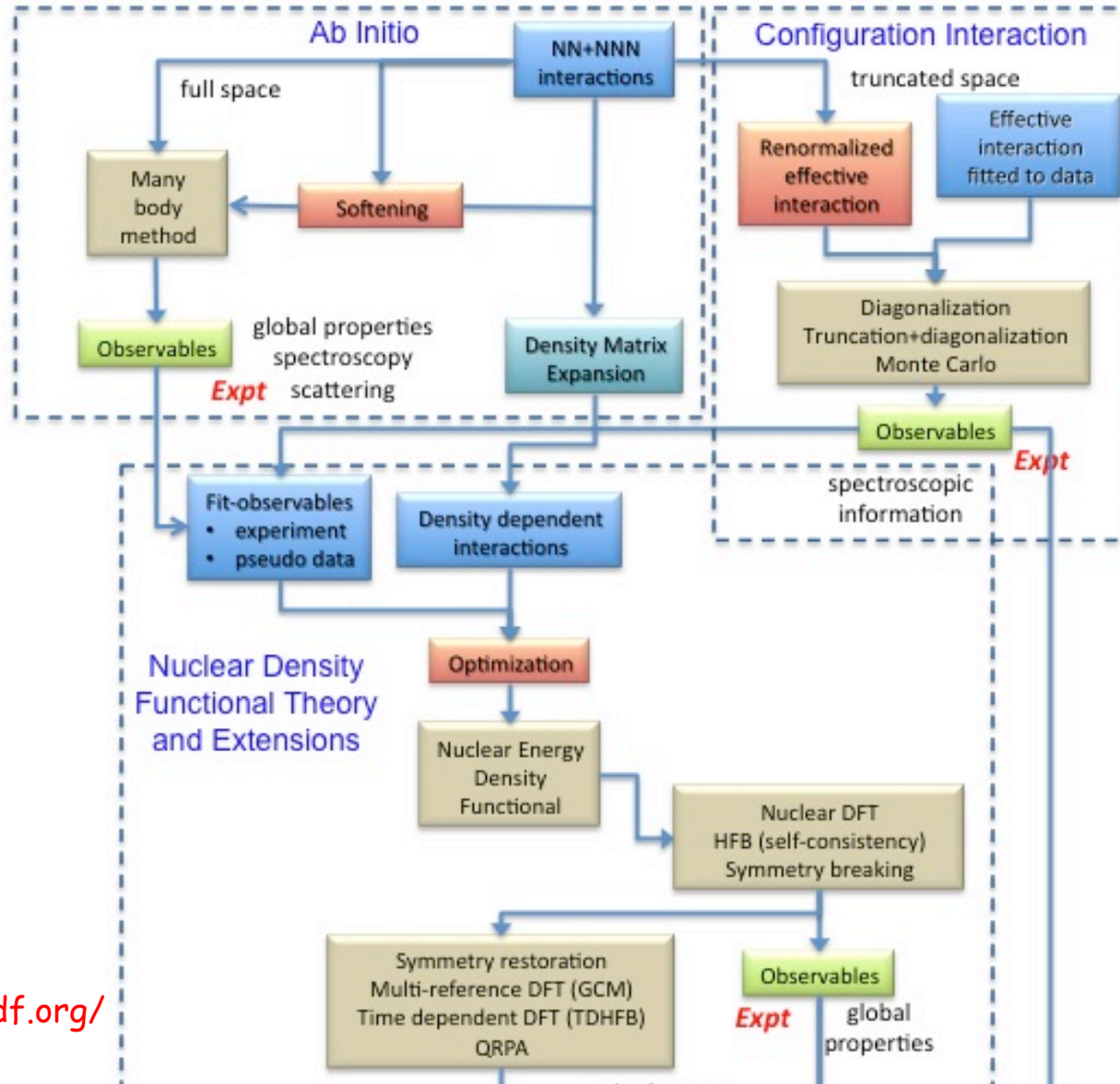
Phenomenological energy density functionals

▣ Predictions of phenomenological EDF (~10 years ago)



Possible strategies (I)

UNEDF



A recent review

Progress in Particle and Nuclear Physics 109 (2019) 103713



Contents lists available at [ScienceDirect](https://www.sciencedirect.com)

Progress in Particle and Nuclear Physics

journal homepage: www.elsevier.com/locate/ppnp



Review

Towards an *ab initio* covariant density functional theory for nuclear structure

Shihang Shen^{a,b,c}, Haozhao Liang^{d,e}, Wen Hui Long^{f,g}, Jie Meng^{a,h,i,*},
Peter Ring^{a,j}

53 pages

with 26 figs and 436 references

A dream for new-generation DFT

quantum-field-theory
oriented DFT

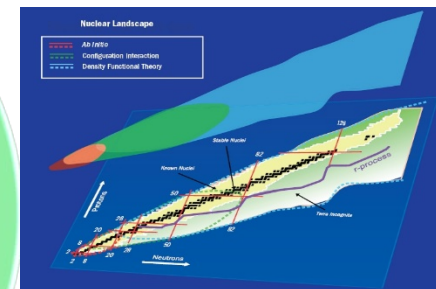
EDF from
effective action
 $F_{\text{HK}}[\rho] \sim \Gamma[\rho]/\beta$
(Legendre transform)

non-perturbative
nature by
**renormalization
group**
 $\partial_k \Gamma_k[\rho] = \text{Tr}\{\dots\}$
(flow eq.)

theoretical
uncertainties
from **EFT**
 $\Gamma^{(2)}, \Gamma^{(3)}, \Gamma^{(4)} \dots$
(power counting)

Interdisciplinary:
(lattice) QCD
hadron
cold atom
condensed matter
quantum chemistry

.....



cf. <http://www.unedf.org/>



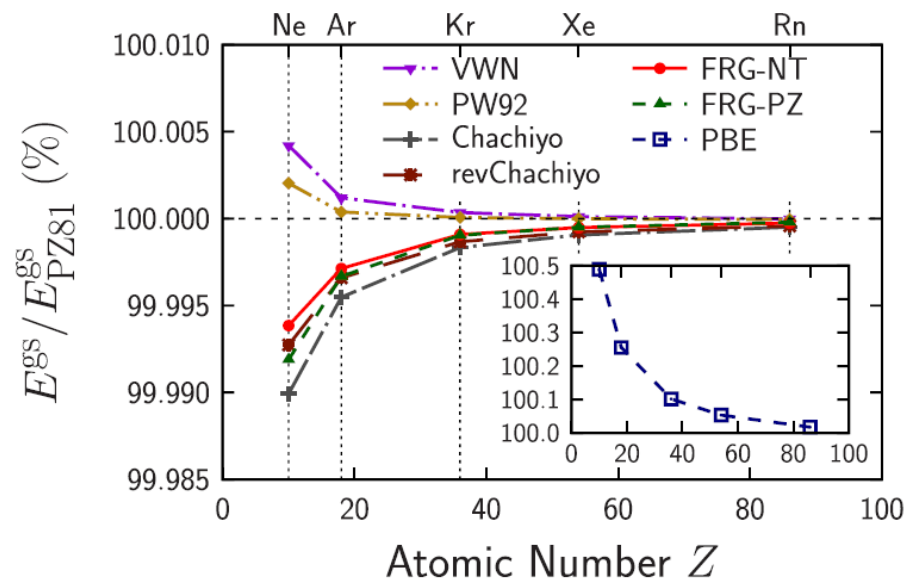
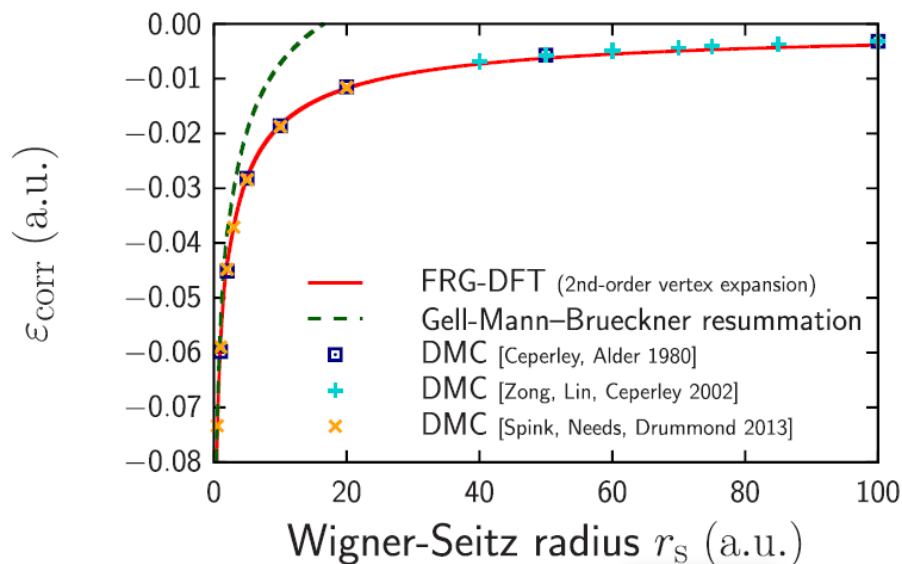
IUPAP Young Scientist Prize
@ INPC2016, Australia

also cf.
Schwenk & Polonyi, arXiv:0403011 [nucl-th]
Kutzelnigg, JMS 768, 163 (2006)
Drut, Furnstahl, Platter, PPNP 64, 120 (2010)
Braun, JPG 39, 033001 (2012)
Metzner et al., RMP 84, 299 (2012)
Drews & Weise, PPNP 93, 69 (2017)

.....

Progress in 3D electron systems

- ◆ Correlation energy of 3D electron gas
- ◆ Ground-state energies of Ne, Ar, Kr, Xe, and Rn atoms



Yokota & Naito, *Phys. Rev. Research* **3**, L012015 (2021).

$$F_{\text{HK}}[\rho] = T_0[\rho] + E_{\text{H}}[\rho] + E_{\text{x}}^{\text{LDA}}[\rho] + E_{\text{c}}^{\text{LDA}}[\rho] + E_{\text{xc}}^{\text{GGA}}[\rho, |\nabla\rho|] + E_{\text{xc}}^{\text{Meta-GGA}}[\rho, |\nabla\rho|, \nabla^2\rho, \tau] + \dots$$

ongoing

A pioneering work: QC for atomic nuclei

PHYSICAL REVIEW LETTERS **120**, 210501 (2018)

Editors' Suggestion

Featured in Physics

Cloud Quantum Computing of an Atomic Nucleus

E. F. Dumitrescu,¹ A. J. McCaskey,² G. Hagen,^{3,4} G. R. Jansen,^{5,3} T. D. Morris,^{4,3} T. Papenbrock,^{4,3,*}
R. C. Pooser,^{1,4} D. J. Dean,³ and P. Lougovski^{1,†}

¹*Computational Sciences and Engineering Division, Oak Ridge National Laboratory, Oak Ridge, Tennessee 37831, USA*

²*Computer Science and Mathematics Division, Oak Ridge National Laboratory, Oak Ridge, Tennessee 37831, USA*

³*Physics Division, Oak Ridge National Laboratory, Oak Ridge, Tennessee 37831, USA*

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⁵*National Center for Computational Sciences, Oak Ridge National Laboratory, Oak Ridge, Tennessee 37831, USA*

Physics

by Stefano Gandolfi*

VIEWPOINT

Cloud Quantum Computing Tackles Simple Nucleus

Researchers perform a quantum computation of the binding energy of the deuteron using a web connection to remote quantum devices.

Model setup and quantum programming

Dumitrescu et al., *PRL* **120**, 210501 (2018)

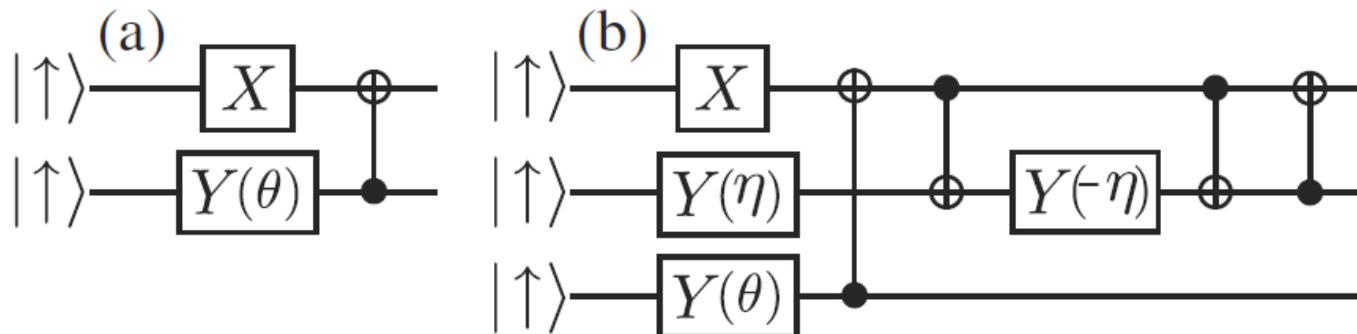
➤ Variational wave function

$$U(\theta) \equiv e^{\theta(a_0^\dagger a_1 - a_1^\dagger a_0)} = e^{i(\theta/2)(X_0 Y_1 - X_1 Y_0)},$$

$$U(\eta, \theta) \equiv e^{\eta(a_0^\dagger a_1 - a_1^\dagger a_0) + \theta(a_0^\dagger a_2 - a_2^\dagger a_0)} \\ \approx e^{i(\eta/2)(X_0 Y_1 - X_1 Y_0)} e^{i(\theta/2)(X_0 Z_1 Y_2 - X_2 Z_1 Y_0)}.$$

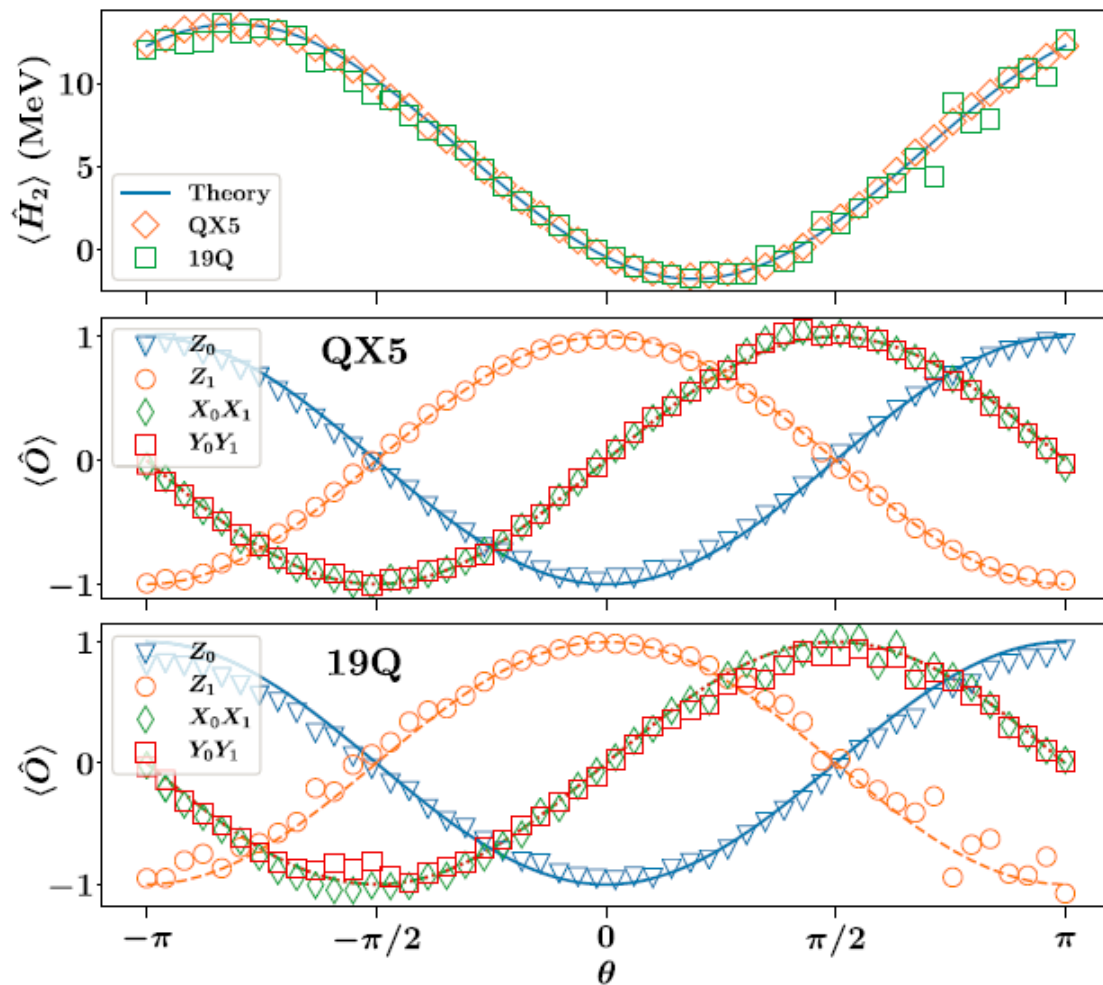
➤ Computing architectures

https://en.wikipedia.org/wiki/Quantum_logic_gate
<https://algassert.com/quirk>



- **QX5** and **19Q** chips: with a single qubit **connected to up to three neighbors**
- It works here ← **only requires up to two connections** for each qubit

Results



$$H_2 = 5.906\,709I + 0.218\,291Z_0 - 6.125Z_1 - 2.143\,304(X_0X_1 + Y_0Y_1),$$

$$U(\theta) \equiv e^{\theta(a_0^\dagger a_1 - a_1^\dagger a_0)} = e^{i(\theta/2)(X_0Y_1 - X_1Y_0)},$$

$$E_2^{\text{QX5}} = -1.80 \pm 0.05 \text{ MeV}$$

$$E_2^{19\text{Q}} = -1.72 \pm 0.03 \text{ MeV}$$

thus

$$E_2 = -1.74 \pm 0.03 \text{ MeV}$$

■ Experimentally determined energies for H_2

Model setup and main results

Dumitrescu et al., *PRL* **120**, 210501 (2018)

- **Deuteron Hamiltonian** (discrete variable representation in HO basis)

$$H_N = \sum_{n,n'=0}^{N-1} \langle n'|(T+V)|n\rangle a_{n'}^\dagger a_n.$$

where $\langle n'|T|n\rangle = \frac{\hbar\omega}{2} \left[(2n+3/2)\delta_n^{n'} - \sqrt{n(n+1/2)}\delta_n^{n'+1} \right.$

$$\left. - \sqrt{(n+1)(n+3/2)}\delta_n^{n'-1} \right],$$

$$\langle n'|V|n\rangle = V_0\delta_n^0\delta_n^{n'}.$$

$$V_0 = -5.68658111 \text{ MeV}$$

$$\hbar\omega = 7 \text{ MeV}$$

- **Results**

<i>E</i> from exact diagonalization				
<i>N</i>	<i>E_N</i>	$O(e^{-2kL})$	$O(kLe^{-4kL})$	$O(e^{-4kL})$
2	-1.749	-2.39	-2.19	
3	-2.046	-2.33	-2.20	-2.21
<i>E</i> from quantum computing				
<i>N</i>	<i>E_N</i>	$O(e^{-2kL})$	$O(kLe^{-4kL})$	$O(e^{-4kL})$
2	-1.74(3)	-2.38(4)	-2.18(3)	
3	-2.08(3)	-2.35(2)	-2.21(3)	-2.28(3)

$$E_{\text{exact}} = -2.22 \text{ MeV}$$

Lipkin model

➤ Lipkin Hamiltonian

$$H = \frac{1}{2}\varepsilon \sum_{p\sigma} \sigma a_{p,\sigma}^\dagger a_{p,\sigma} + \frac{1}{2}V \sum_{pp'\sigma} a_{p,\sigma}^\dagger a_{p',\sigma}^\dagger a_{p',-\sigma} a_{p,-\sigma} + \frac{1}{2}W \sum_{pp'\sigma} a_{p,\sigma}^\dagger a_{p',-\sigma}^\dagger a_{p',\sigma} a_{p,-\sigma},$$

**VALIDITY OF MANY-BODY APPROXIMATION METHODS
FOR A SOLVABLE MODEL**

(I). Exact Solutions and Perturbation Theory

**VALIDITY OF MANY-BODY APPROXIMATION METHODS
FOR A SOLVABLE MODEL**

(II). Linearization Procedures

**VALIDITY OF MANY-BODY APPROXIMATION METHODS
FOR A SOLVABLE MODEL**

(III). Diagram Summations

**VALIDITY OF MANY-BODY APPROXIMATION METHODS
FOR A SOLVABLE MODEL**

(IV). The Deformed Hartree-Fock Solution

D. AGASSI and H. J. LIPKIN

The Weizmann Institute of Science, Rehovoth, Israel

and

N. MESHKOV

Catholic University of America, Washington, D.C. †

Lipkin model

➤ Quasi-spin formulation

$$J_+ = \sum_p a_{p,+1}^\dagger a_{p,-1}, \quad J_- = \sum_p a_{p,-1}^\dagger a_{p,+1}, \quad J_z = \frac{1}{2} \sum_{p\sigma} \sigma a_{p,\sigma}^\dagger a_{p,\sigma}.$$

➤ Hamiltonian

$$H = \varepsilon J_z + \frac{1}{2} V (J_+^2 + J_-^2) + \frac{1}{2} W (J_+ J_- + J_- J_+).$$

➤ Exact solutions ($N = 2, 3, 4, 6, 8$ with $W = 0$)

- for $N = 2$:

$$\frac{E}{\varepsilon} = 0, \pm \left[1 + \left(\frac{V}{\varepsilon} \right)^2 \right]^{\frac{1}{2}},$$

- for $N = 3$:

$$\frac{E}{\varepsilon} = \pm \left\{ \frac{1}{2} \pm \left[1 + 3 \left(\frac{V}{\varepsilon} \right)^2 \right]^{\frac{1}{2}} \right\},$$

- for $N = 4$:

$$\frac{E}{\varepsilon} = 0, \pm 2 \left[1 + 3 \left(\frac{V}{\varepsilon} \right)^2 \right]^{\frac{1}{2}},$$

$$\frac{E}{\varepsilon} = \pm \left[1 + 9 \left(\frac{V}{\varepsilon} \right)^2 \right]^{\frac{1}{2}},$$

Lipkin model

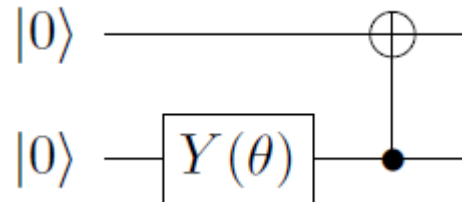
- Qubit representation of Lipkin Hamiltonian ($W = 0$)

$$H = -\frac{1}{2}\varepsilon \sum_{i=1}^N Z_i + \frac{1}{4}V \sum_{i,j=1}^N (X_i X_j - Y_i Y_j).$$

- Trial wave functions ($N = 2$)

$$|\psi\rangle = U(\theta)|00\rangle = \cos\frac{\theta}{2}|00\rangle + \sin\frac{\theta}{2}|11\rangle.$$

- Quantum circuit ($N = 2$)



- Number of parameters, $O(2^N)$, is needed for a complete expression of the trial wave functions.

UCC and structure learning ansatz

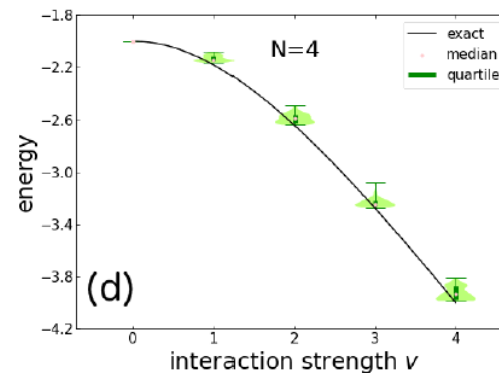
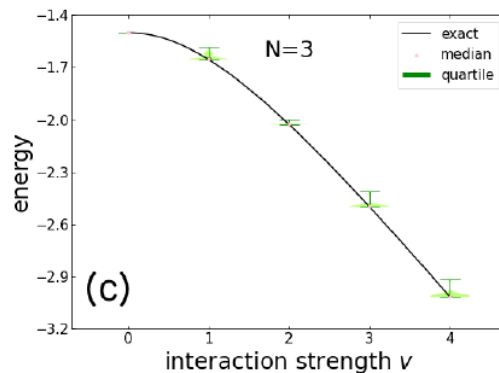
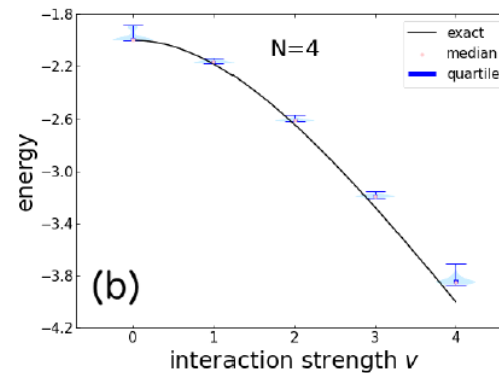
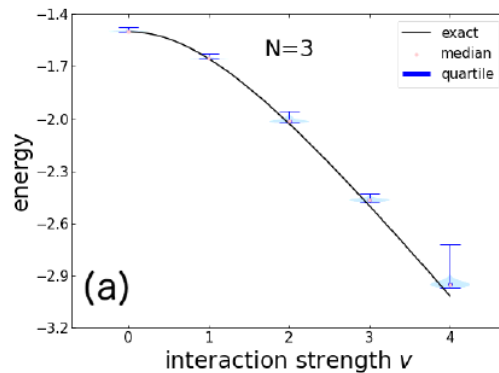
Chinese Physics C Vol. 46, No. 2 (2022) 024106

Quantum computing for the Lipkin model with unitary coupled cluster and structure learning ansatz*

Asahi Chikaoka(近岡旭)^{1,2} Haozhao Liang(梁豪兆)^{1,2†}

¹Department of Physics, Graduate School of Science, The University of Tokyo, Tokyo 113-0033, Japan

²RIKEN Nishina Center, Wako 351-0198, Japan



UCC ansatz

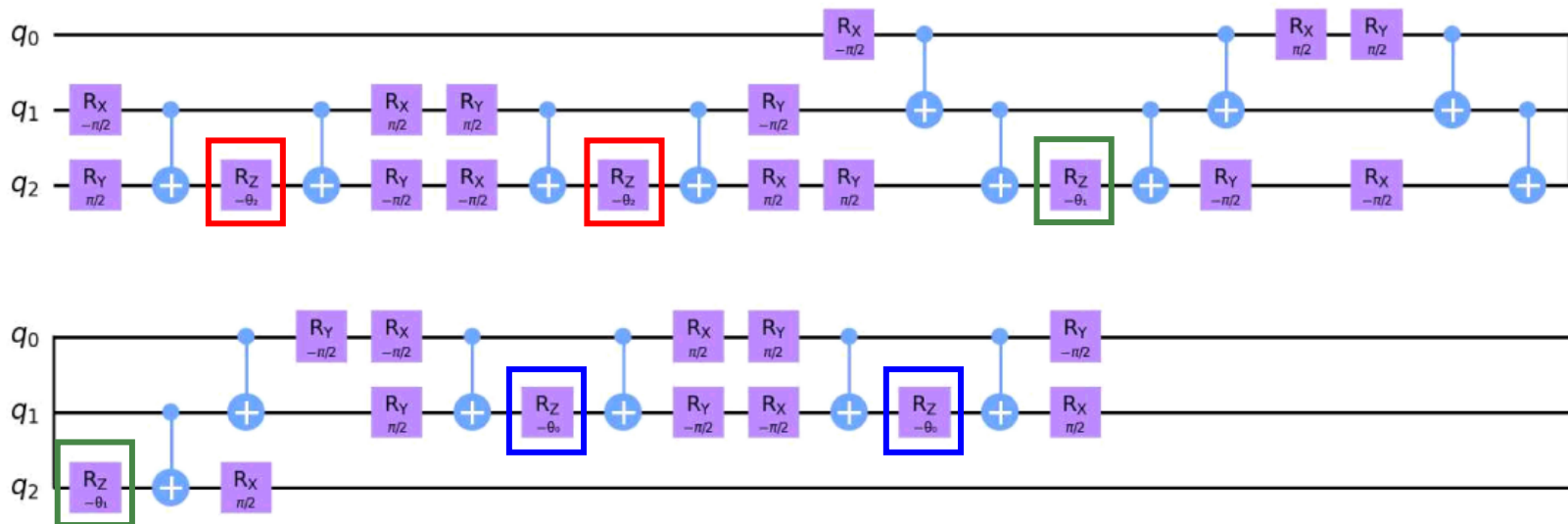
➤ Trial wave functions

Chikaoka and HZL, *Chin. Phys. C* **46**, 024106 (2022)

$$U(\theta) \equiv \exp \left[\sum_{ij} \theta_{ij} (a_i^\dagger a_j^\dagger - a_j a_i) \right]$$

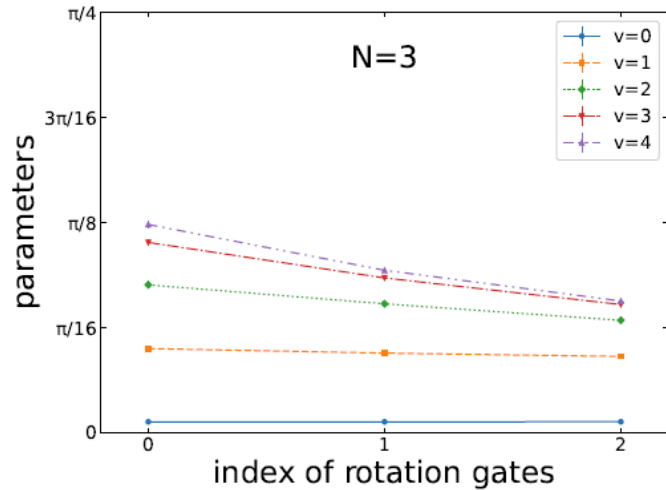
$$\mapsto \exp \left\{ i \sum_{ij} \theta_{ij} \frac{(-1)^{j-i-1}}{2} \left[X_i \left(\prod_{k=i+1}^{j-1} Z_k \right) Y_j \right. \right. \\ \left. \left. + Y_i \left(\prod_{k=i+1}^{j-1} Z_k \right) X_j \right] \right\},$$

➤ Quantum circuit ($N=3$)

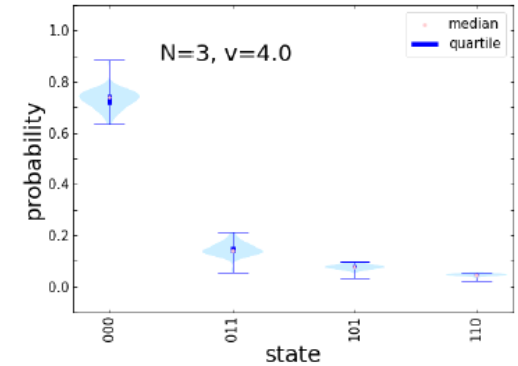
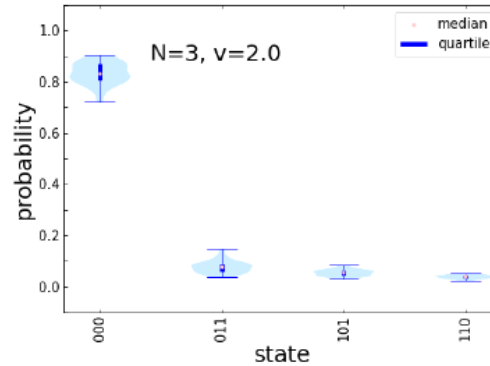


UCC ansatz (N=3)

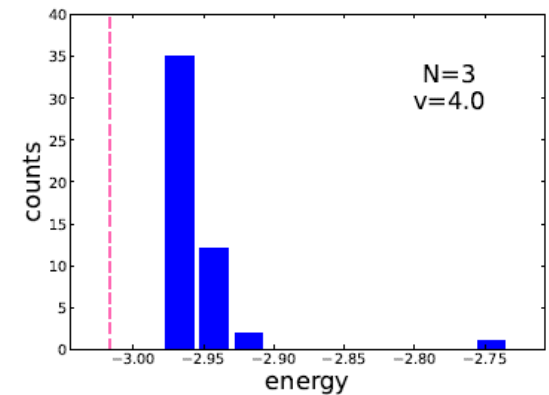
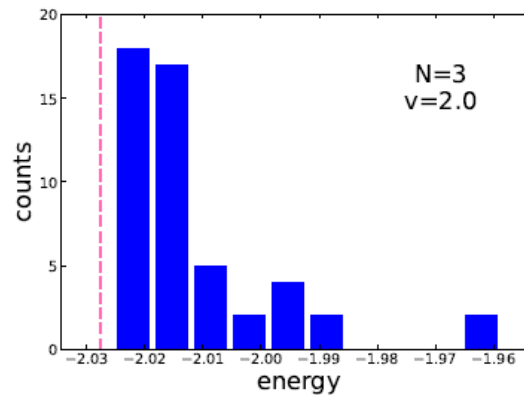
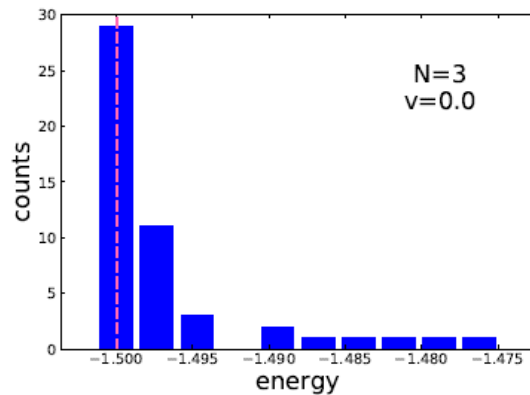
Parameters



State probabilities

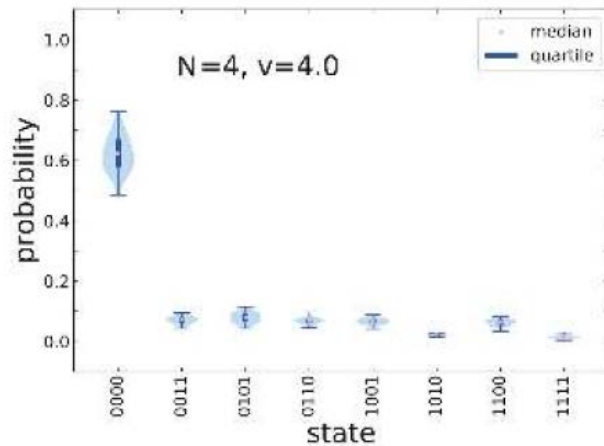
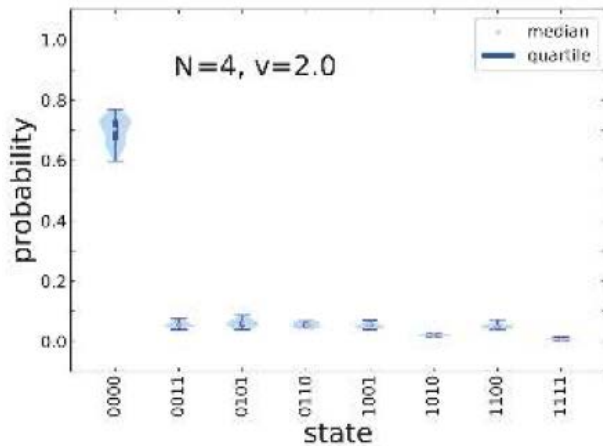


Ground-state energies

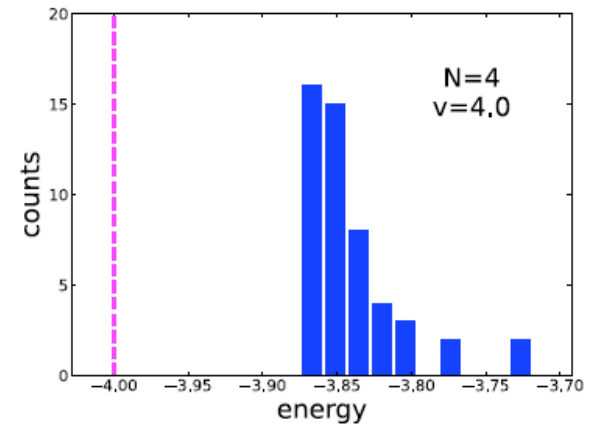
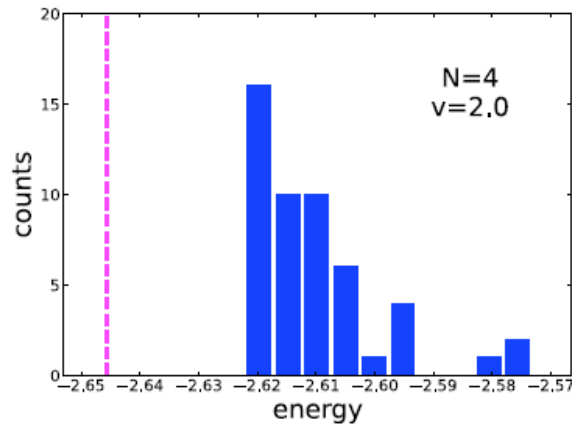
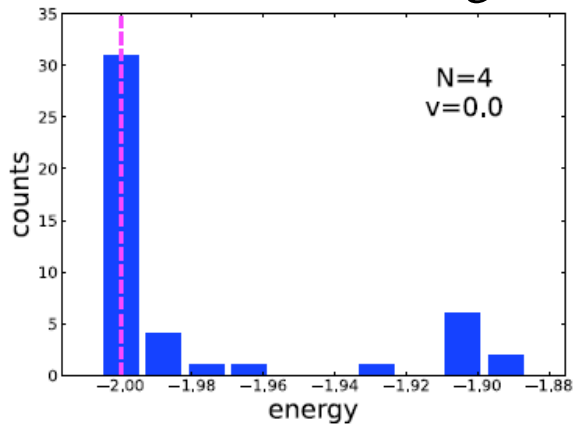


UCC ansatz (N=4)

➤ State probabilities



➤ Ground-state energies



Structure learning ansatz

➤ Trial wave functions

Chikaoka and HZL, *Chin. Phys. C* **46**, 024106 (2022)

Algorithm 1 Rotoselect

Input: Function calculating expectation values with respect to each quantum circuit U : $\langle M(U) \rangle$. Here, M represents a Hermitian operator. The quantum circuit U with the maximum value of the depth, D , is composed of rotation gates at the depth d , $U_d(\theta_d, H_d) = H_d(\theta_d)$ (e.g., $R_X(\theta_d) = \exp[-i\frac{\theta_d}{2}X]$), and the CNOTs. Here, θ_d is a parameter at the depth d and H_d is the element of the set of the rotation operators $\{I, R_X, R_Y, R_Z\}$. Axes of rotation gates, i.e. I, R_X, R_Y , or R_Z , are chosen in order to minimize the expectation value.

Output: Optimized quantum circuit U_{opt} . Here, U_{opt} is optimized with respect to θ_d and H_d .

Initialize $\theta_d \in (\pi, \pi]$ and $H_d \in \{I, R_X, R_Y, R_Z\}$ for $d = 1, \dots, D$ heuristically or at random. (In practice, initialize all $\theta_d = 0$ and all $H_d = I$.)

repeat

 for $d = 1, \dots, D$ do

 Compute $\theta_{d,P}^*$ for $P \in \{I, R_X, R_Y, R_Z\}$ using SMO method, where $\theta_{d,P}^*$ is the optimized parameter with the selected gate P .

$H_d \leftarrow \arg \min_P \langle M(U) \rangle |_{U_d(\theta_d, H_d) = U_d(\theta_{d,P}^*, P)}$

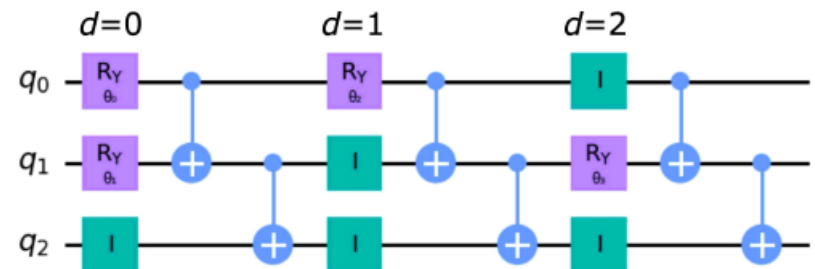
$\theta_d \leftarrow \theta_{d, H_d}^*$, where θ_{d, H_d}^* is the optimized parameter with the selected gate H_d

 end for

until stopping criterion is met

return optimized quantum circuit U_{opt}

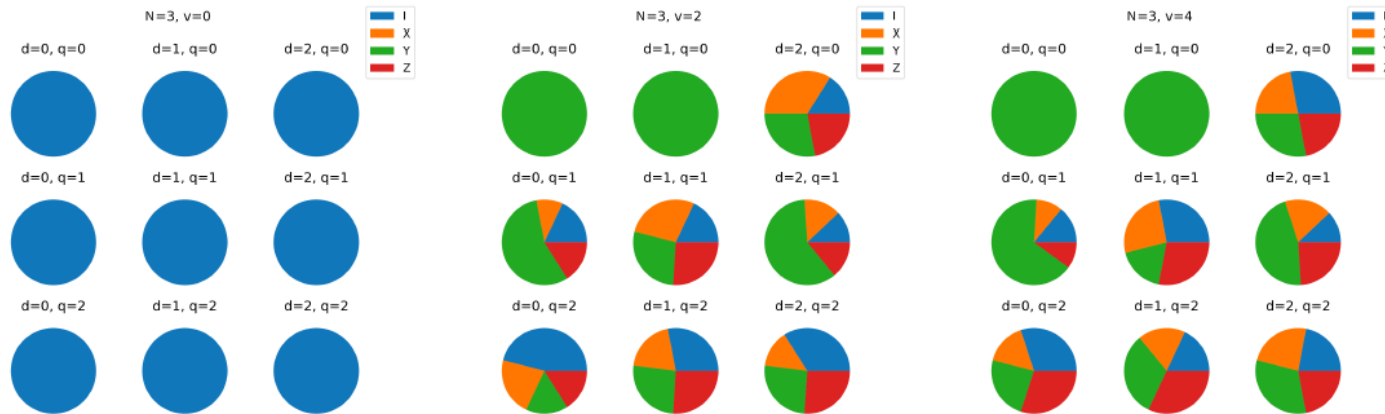
➤ Quantum circuit ($N = 3$) (an example)



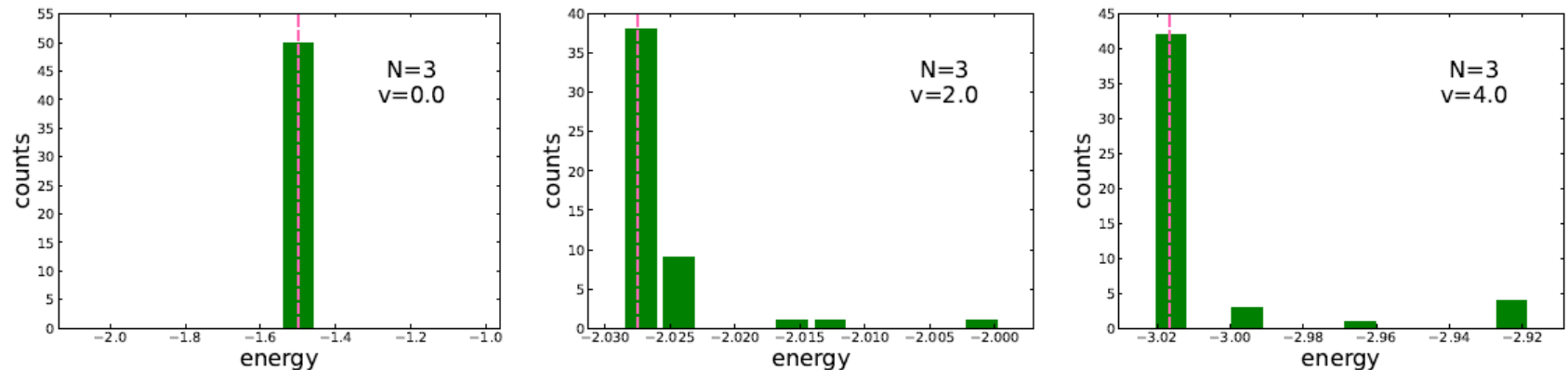
Structure learning ansatz (N=3)

Rotating axes

Chikaoka and HZL, *Chin. Phys. C* **46**, 024106 (2022)



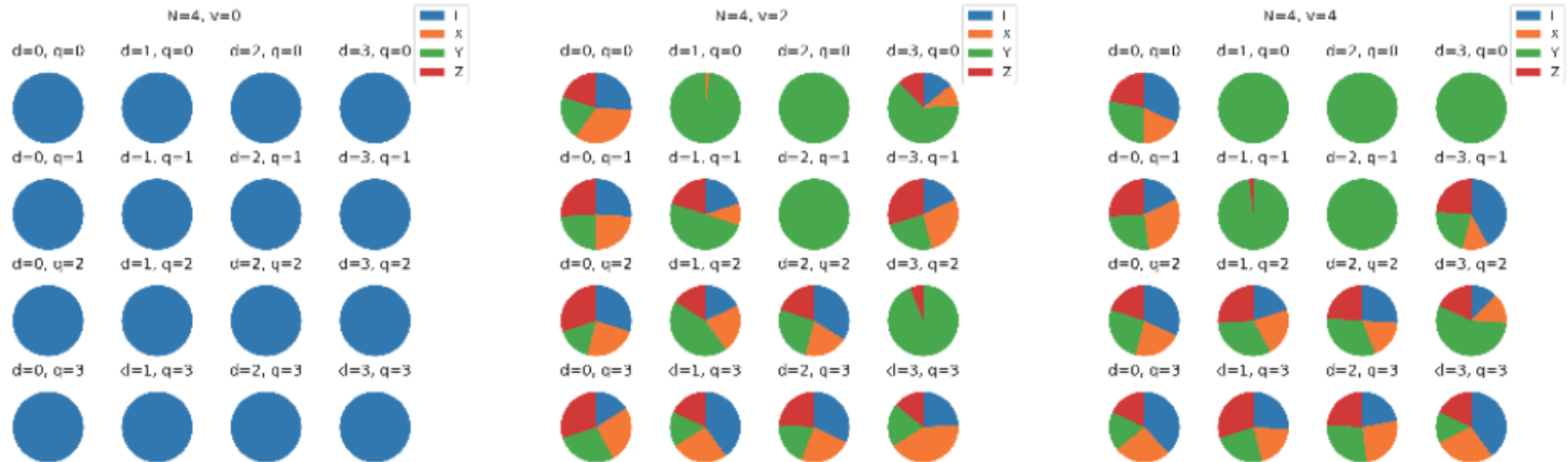
Ground-state energies



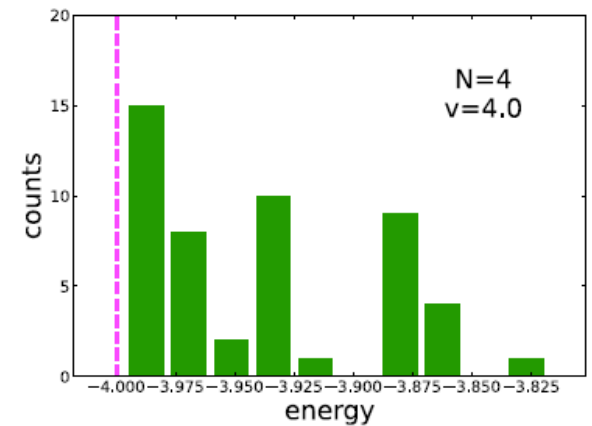
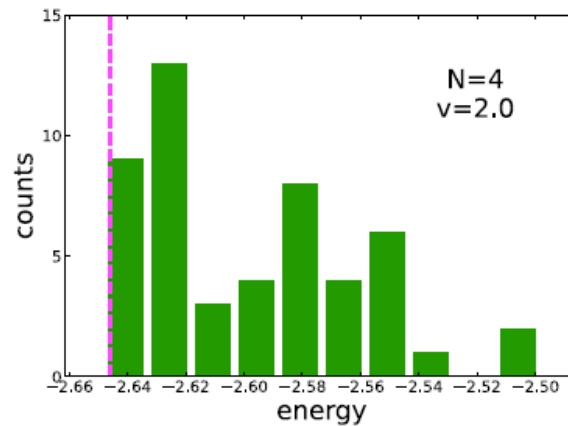
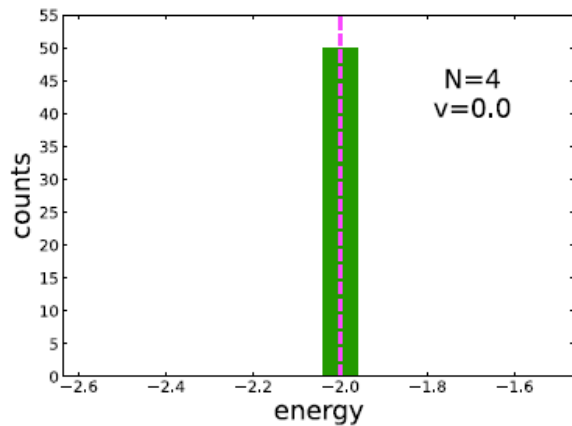
Structure learning ansatz (N=4)

Rotating axes

Chikaoka and HZL, *Chin. Phys. C* **46**, 024106 (2022)

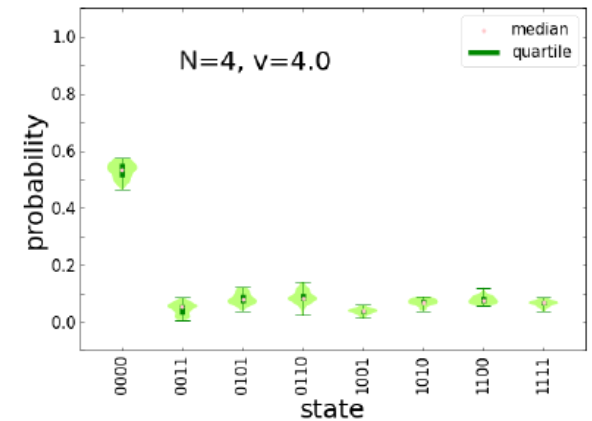
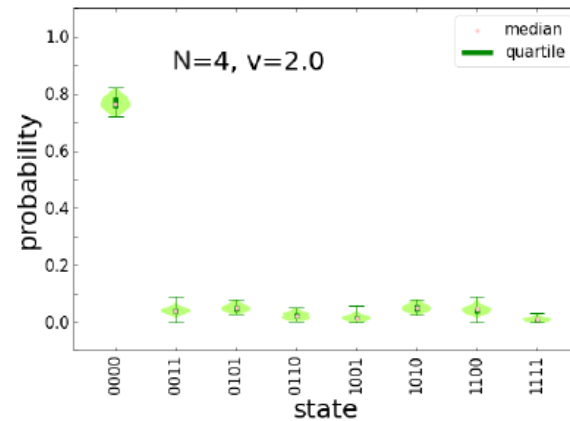
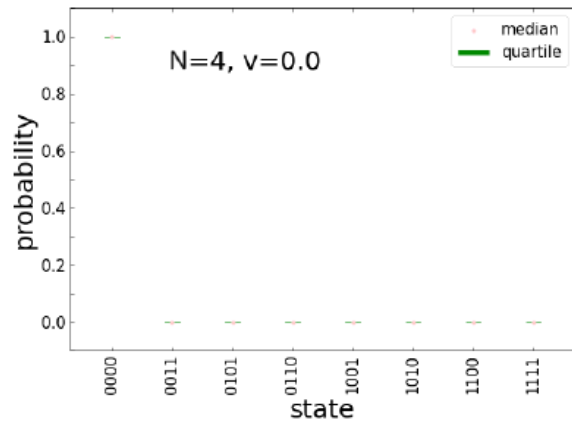
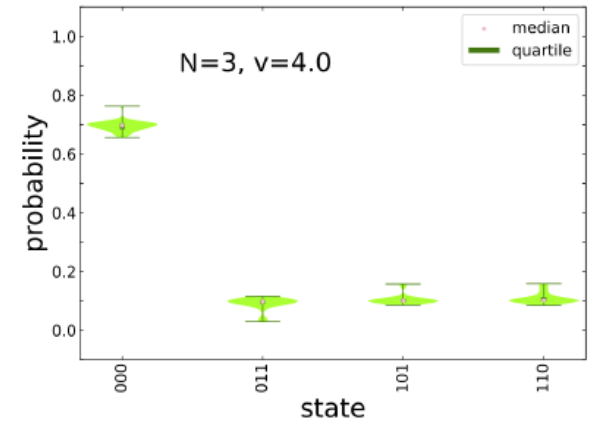
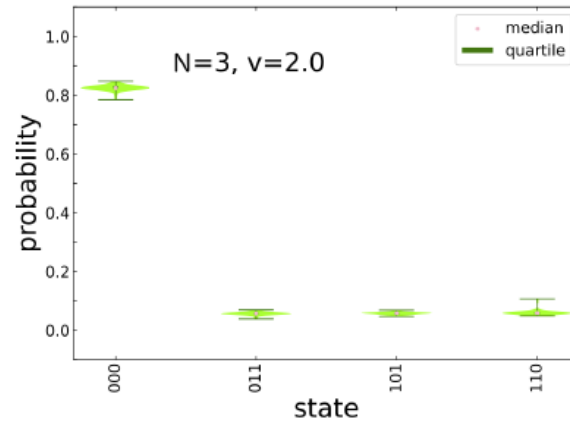
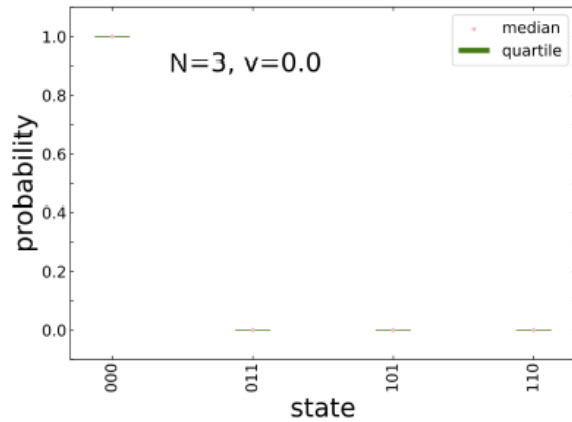


Ground-state energies



Structure learning ansatz

➤ State probabilities



Hybrid Quantum Annealing (HQA)

scientific reports

Irie, HZL, Doi, Gongyo, Hatsuda, *Sci. Rep.* **11**, 8426 (2021)

OPEN

Hybrid quantum annealing via molecular dynamics

Hirotsuka Irie^{1,2}, Haozhao Liang^{3,4}, Takumi Doi^{2,3}, Shinya Gongyo^{2,3} & Tetsuo Hatsuda²

➤ Concept of HQA

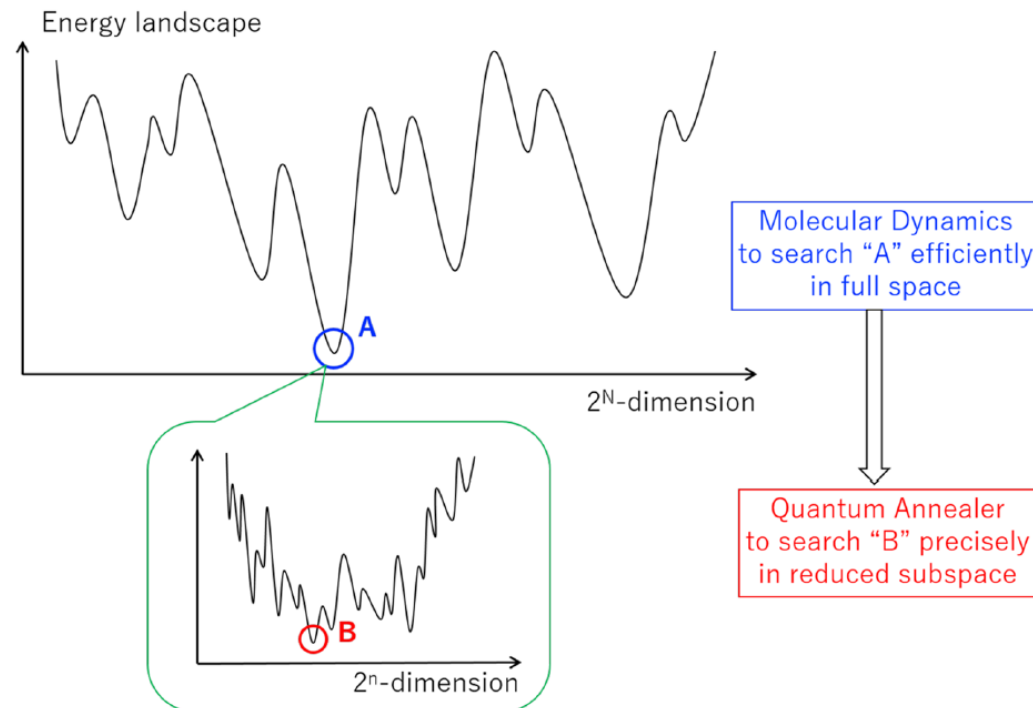


Figure 1. Concept of hybrid quantum annealing via molecular dynamics.

Quantum computing for nuclear physics

Lipkin model on a quantum computer

Michael J. Cervia, A. B. Balantekin, S. N. Coppersmith, Calvin W. Johnson, Peter J. Love, C. Poole, K. Robbins, and M. Saffman
Phys. Rev. C **104**, 024305 – Published 3 August 2021

Simulating excited states of the Lipkin model on a quantum computer

Manqoba Q. Hlatshwayo, Yinu Zhang, Herlik Wibowo, Ryan LaRose, Denis Lacroix, and Elena Litvinova
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