# Quantum computation of valence two-neutron systems towards solving full nuclear Hamiltonians

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Ref: SY, T. Sato, T. Ogata, T. Naito, M. Kimura, PhysRevC.109.064305 (2024)

Motivations for Quantum Computing in NP(Nuclear Physics)<sup>1</sup>

- strongly interacting protons and neutrons (let me consider only "nucleons" dof)
- nuclear force is highly nontrivial



three-nucleon force

first principle understanding by Lattice QCD c.f. Ishii, Aoki, Hatsuda, PRL 99, 022001 (2007).

still have very large uncertainty for nuclear structure/reaction(dynamics)

divergent of the dimension: 10<sup>16</sup>! 10<sup>19</sup>!!, 10<sup>23</sup>!!!
 most of nuclei cannot be solved in exact manner
 → zoo of nuclear many-body methods
 → need to prepare for "Quantum Computing era" for NP

# Difference from other quantum many-body systems

Quantum chemistry:



Nuclear physics:



"99 > % of energy of a molecule in equilibrium

is explained within Hartree-Fock level"

(i.e., single Slater determinant)

rest 1 % is called **correlation energy**:

<u>M</u>øller – <u>P</u>lesset (MP a.k.a MBPT)

 $\underline{C}$ oupled  $\underline{C}$ luster  $\underline{S}$ ingle and  $\underline{D}$ ouble (CCSD)

 $CCSD + \underline{T}riple (CCSDT)$ 

Full Configuration interaction (Full-CI) accurate but heavy Interaction is highly non-perturbative

uncertain (many channels, three-nucleon force,..)

<sup>56</sup>Ni under modern Nuclear Force (Chiral EFT)

HF = - 302.716 MeV

HF + MP2 + MP3 = -473.089 MeV (MP2 = -152.533, MP3 = -17.716)

How dare people say perturbation theory !!

c.f. Energy (Exp.) = -483.996 MeV



- initialize the qubits
   I00...>, Hartree-Fock, etc.
- 2. ansatz:
   Operating Unitary gates
   → prepare trial wave functions
- 3. Operating Unitary gates (Hamiltonian)
   →measurement → energy
- 4. Optimize the parameters classically to minimize the energy

Popular choice for the ansatz: the <u>u</u>nitary <u>c</u>oupled <u>c</u>luster (UCC) or its variants

#### schematics for valence two-neutron systems

{n, l, j, jz, tz}

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blue: valence space, pale orange: inert core

#### valence two-neutron systems: p shell example

ground state (J=0) can be described only by time-reversal pairs (jz = m & -m)



original Hamiltonian can be rewritten by pair creation/annihilation operators

$$H = \sum_{i} \varepsilon_{i} \hat{a}_{i}^{\dagger} \hat{a}_{i} + \frac{1}{4} \sum_{ijkl} V_{ijkl} \hat{a}_{i}^{\dagger} \hat{a}_{j}^{\dagger} \hat{a}_{l} \hat{a}_{k}, \qquad A_{\tilde{i}}^{\dagger} = c_{\tilde{i}}^{\dagger} c_{\tilde{i}}^{\dagger}, \qquad \left[A_{\tilde{i}}, A_{\tilde{j}}^{\dagger}\right] = \delta_{\tilde{i}\tilde{j}} (1 - N_{\tilde{i}}), \\ A_{\tilde{i}} = c_{\tilde{i}} c_{i}, \qquad \left[N_{\tilde{i}}, A_{\tilde{j}}^{\dagger}\right] = 2\delta_{\tilde{i}\tilde{j}} A_{\tilde{j}}^{\dagger}. \qquad H^{\text{pw}} = \sum_{i} \bar{h}_{i} A_{i}^{\dagger} A_{i} + \sum_{i \leq j} \bar{V}_{ij} A_{i}^{\dagger} A_{j},$$

expressing the "pair" as a single qubit, # of qubits needed is reduced to half

$$\begin{bmatrix} |110000\rangle \rightarrow |100\rangle \\ |001001\rangle \rightarrow |010\rangle \\ |000110\rangle \rightarrow |001\rangle \end{bmatrix}$$

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## ansatz: circuit for trial wave function



A. Pérez-Obiol et al., <u>Sci. Rep. 13</u>, <u>12291</u> (2023)

## Measurement of energy

$$H^{\mathrm{pw}} = \sum_{i} \bar{h}_{i} A_{i}^{\dagger} A_{i} + \sum_{i \leq j} \bar{V}_{ij} A_{i}^{\dagger} A_{j}, \quad \hat{H}^{\mathrm{pw}}_{\mathrm{qubit}} = \underbrace{\sum_{i} \frac{\bar{h}_{i} + \bar{V}_{ii}}{2} \left(I - Z_{i}\right)}_{i} + \underbrace{\frac{1}{2} \sum_{i < j} \bar{V}_{ij} \left(X_{i} X_{j} + Y_{i} Y_{j}\right)}_{i < j}$$

1<sup>st</sup> term: Pauli-Z measurement of the ansatz (occupation numbers)

2<sup>nd</sup> term:

method A: measure expectation value of all Pauli spins (XX, YY)

method B: computational basis sampling (QunaSys&Osaka U. group)

→ variational, only two measurements (ansatz & ansatz + H-gates) M. Kohda, et al., Phys. Rev. Res. 4, 033173 (2022).

Results: FTQC(noise-free) simulator



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## Results: NISQ (noisy) simulator



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narrower circuits leads to noise-resilient results

#### <sup>18</sup>O (6 qubits)case:

rearrangement of the circuit to the equivalent but shallow circuit calculating the relative weights of jj-coupling orbitals



1()

connectivity of qubits in the handwage is another important factor

 $\mathbf{D}$ 



![](_page_11_Figure_1.jpeg)

![](_page_11_Picture_2.jpeg)

#### Results on real hardware: ibm\_brisbane (127 qubits)

![](_page_12_Figure_1.jpeg)

Run1-5: different runs

- w/ hatch (//): rearranged
- w/o hatch: original circuit

$$^{18}O \& ^{6}He \rightarrow < 0.1\%$$

$$^{42}Ca \rightarrow \sim 1\%$$

c.f. UCC-type results reported so far

<sup>6</sup> He	~	2%	PRC	105,	064308	(2022)
<sup>6</sup> Li	~	4%	PRC	106,	034325	(2022)
<sup>18</sup> O	~	3%	PRC	108,	064305	(2023)

![](_page_13_Picture_0.jpeg)

- > We explored quantum algorithms for NP on NISQ devices
  - ansatz for valence two neutrons  $\rightarrow$  much fewer CNOTs than UCC
  - accurate: < 0.1% (<sup>6</sup>He & <sup>18</sup>O), ~1%(<sup>42</sup>Ca)
  - importance of taking care of symmetries, connectivity, etc.
  - parameter optimization: sequential & derivative-free method works fine see appendix
- > TODO (a lot, but some are already in progress!)

extensions to proton-neutron systems, many-particle systems towards driplines  $\sim$  input for astrophysics

not only NISQ algorithms, towards (early-)FTQC

# Collaborators & Acknowledgement

Collaborators:

UTokyo: Takeshi Sato, Takumi Ogata

RIKEN iTHEMS: Tomoya Naito

RIKEN Nishina center: Masaaki Kimura

![](_page_14_Picture_5.jpeg)

![](_page_14_Picture_6.jpeg)

![](_page_14_Picture_7.jpeg)

Grant (only for me):

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![](_page_14_Picture_10.jpeg)

<u>Full CI</u>: current limitation of the Lanczos method  $\sim 10^{11}$  dim.

<sup>12</sup>C (e.g. Hoyle state) 7 major shell excitation ~ 10<sup>19</sup> dim. → No-core MCSM

<u>valence CI</u>:

frontiers of neutron-rich nuclei: Ni ~10<sup>16</sup> or much larger (10<sup>23</sup>!)

→ Cl variants: Monte Carlo Shell Model (MCSM) Quasi-vacua Shell Model (QVSM)

> <u>T. Togashi et al., PRL 121, 062501 (2018)</u> <u>N. Shimizu et al., PRC **103**, 014312 (2021)</u>

T.Otsuka et al., Nature Comm. 13:2234 (2022)

## occupations of pair configurations

![](_page_16_Figure_1.jpeg)

Note: the parameters are fixed to the ones giving exact results for noisy simulator and hardware results, post-selection is done

# computational basis sampling

M. Kohda, R. Imai, K. Kanno, K. Mitarai, W. Mizukami, and Y. O. Nakagawa, Quantum expectation-value estimation by computational basis sampling, Phys. Rev. Res. 4, 033173 (2022).

$$\langle \psi | H | \psi \rangle = \sum_{m,n=0}^{N} \langle \psi | m \rangle \langle m | H | n \rangle \langle n | \psi \rangle$$
$$= \sum_{m,n=0}^{N} |\langle m | \psi \rangle|^{2} |\langle n | \psi \rangle|^{2} \frac{\langle m | H | n \rangle}{\langle m | \psi \rangle \langle \psi | n \rangle}.$$
(7)
$$\langle \psi | X_{j} \otimes X_{k} | \psi \rangle = \sqrt{\sigma_{j}^{2} \sigma_{k}^{2}} \operatorname{sgn} [\sigma_{j} \sigma_{k}],$$
(8)
$$= \sqrt{\langle N_{j} \rangle \langle N_{k} \rangle} \operatorname{sgn} [\langle X_{j} X_{k} \rangle]$$
(9)

↑ ansatz plus H-gates on all qubits

↑ equiv. to Z measurement of the ansatz circuit

You just need 2 measurements, ansatz & ansatz + H-gates

# Derivative-free & Sequential optimization of the parameters

K. M. Nakanishi, K. Fujii, and S. Todo, Sequential minimal optimization for quantum-classical hybrid algorithms, Phys. Rev. Res. 2, 043158 (2020).

$$\mathcal{L}(\boldsymbol{\theta}) = \sum_{k} w_{k} \langle \varphi | U^{\dagger}(\boldsymbol{\theta}) \mathcal{H}_{k} U(\boldsymbol{\theta}) | \varphi \rangle, \qquad (B1) = \text{Energy expectation value}$$

Optimizing Ry-gate rotation angle

$$\mathcal{L}(\theta_1) = A\cos\left(\theta_1 - B\right) + C. \tag{B2}$$

Optimizing cRy-gate rotation angle

$$\mathcal{L}(\theta_j) = A\cos\theta_j + B\sin\theta_j + C\cos 2\theta_j + D\sin 2\theta_j + E.$$
(B3)

# Optimization of the circuit parameters

parameters so far are fixed to the ones giving exact results by hand

![](_page_19_Figure_2.jpeg)

K. M. Nakanishi, K. Fujii, and S. Todo:

Sequential minimal optimization for quantum-classical hybrid algorithms, Phys. Rev. Res. 2, 043158 (2020).