Quantum computati valence two-neutron s towards solving full nuclear

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Ref: SY, T. Sato, T. Ogata, T. Naito, M. Kimura, Physl

Motivations for Quantum Computing in NP(Nuclear Physics) 1

- strongly interacting protons and neutrons (let me consider only "nucleons" dof)
- nuclear force is highly nontrivial

three-nucleon force

first principle understanding by Lattice QCD c.f. Ishii, Aoki, Hatsuda, PRL 99, 022001 (2007).

still have very large uncertainty for nuclear structure/reaction(dynamics)

- divergent of the dimension: 10^{16} ! 10^{19} !!, 10^{23} !!! most of nuclei cannot be solved in exact manner \rightarrow zoo of nuclear many-body methods \rightarrow need to prepare for "Quantum Computing era" for NP

Difference from other quantum many-body systems

Quantum chemistry:

Nuclear physics:

"99 > % of energy of a molecule in equilibrium

is explained within Hartree-Fock level"

(i.e., single Slater determinant)

rest 1 % is called correlation energy:

Møller ‒ Plesset (MP a.k.a MBPT)

Coupled Cluster Single and Double (CCSD)

CCSD + Triple (CCSDT)

Full Configuration interaction (Full-CI) accurate but heavy Interaction is highly non-perturbative

uncertain (many channels, three-nucleon force,..)

⁵⁶Ni under modern Nuclear Force (Chiral EFT)

 $HF = -302.716$ MeV

 $HF + MP2 + MP3 = -473.089$ MeV $(MP2 = -152.533, MP3 = -17.716)$

How dare people say perturbation theory !!

c.f. Energy (Exp.) = -483.996 MeV

- 1. initialize the qubits |00...>, Hartree-Fock, etc.
- 2. ansatz: Operating Unitary gates \rightarrow prepare trial wave functions
- 3. Operating Unitary gates (Hamiltonian) \rightarrow measurement \rightarrow energy
- 4. Optimize the parameters classically to minimize the energy

Popular choice for the ansatz: the u </u>nitary c </u>oupled c luster (UCC) or its variants</u>

schematics for valence two-neutron systems $_{\{n, 1, i, iz, tz\}}$ 4

blue: valence space, pale orange: inert core

valence two-neutron systems: p shell example 15

ground state $(J=0)$ can be described only by time-reversal pairs $(jz = m \& -m)$

 $|110000 \rangle \rightarrow |100 \rangle$

 $|000110\rangle \rightarrow |001\rangle$

|001001

original Hamiltonian can be rewritten by pair creation/annihilation operators

$$
H = \sum_{i} \varepsilon_{i} \hat{a}_{i}^{\dagger} \hat{a}_{i} + \frac{1}{4} \sum_{ijkl} V_{ijkl} \hat{a}_{i}^{\dagger} \hat{a}_{j}^{\dagger} \hat{a}_{l} \hat{a}_{k},
$$
\n
$$
A_{\tilde{i}} = c_{\tilde{i}} c_{\tilde{i}},
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A_{\tilde{i}} = c_{\tilde{i}} c_{\tilde{i}},
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N_{\tilde{i}} = c_{\tilde{i}}^{\dagger} c_{\tilde{i}},
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N_{\tilde{i}} = \tilde{a}_{\tilde{i}}^{\dagger} c_{\tilde{i}},
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$$
M_{\tilde{i}} = \sum_{i} \bar{h}_{i} A_{i}^{\dagger} A_{i} + \sum_{i \leq j} \bar{V}_{ij} A_{i}^{\dagger} A_{j},
$$

expressing the "pair" as a single qubit, # of qubits needed is reduced to half

ansatz: circuit for trial wave function

You need only singly occupied configurations \rightarrow ansatz can be expressed by (C-)Ry & CNOTs

 $42Ca$ (10 qubits)

A. Pérez

Measurement of energy

$$
H^{\text{pw}} = \sum_{i} \bar{h}_i A_i^{\dagger} A_i + \sum_{i \leq j} \bar{V}_{ij} A_i^{\dagger} A_j, \quad \hat{H}_{\text{qubit}}^{\text{pw}} = \sum_{i} \frac{\bar{h}_i + \bar{V}_{ii}}{2} (I)
$$

1st term: Pauli-Z measurement of the ansatz (occup

2nd term:

method A: measure expectation value of all Pau

method B: computational basis sampling (QunaS

M. Kohda, et al. \rightarrow variational, only two measurements (ansatz δ

Results: FTQC(noise-free) simulator

Results: NISQ (noisy) simulator

9

*R R R R P R <i>R <i>B R R B B B B and <i>B arral <i>essilient results**esgmmmm**esgmmms**a**esgmmn**a**esgmmn**a**a**d**a**d**a**d**a**d**a**d**a**d**d**d**d**d**d**d**d*

¹⁸O (6 qubits)case:

rearrangement of the circuit to the equivalent but shallow circuit *c* alculating the relative weights of jj-coupling orbitals *|*0⟩⁴ *Ry*(2*θ*4) *• •*

 α connectivity of qubits in the hard \sqrt{R} α β α β β α β β γ β γ β γ

*|*0⟩¹ *Ry*(2*θ*2) *• • •*

 \vert ()

Results on real hardware: ibm_brisbane (127

Summary

SY, T. Sato, T. Ogata, T. Naito, M.

 \triangleright We explored quantum algorithms for NP

- ansatz for valence two neutrons \rightarrow much f

 $-$ accurate: < 0.1% (⁶He & ¹⁸O), ~1%(⁴²Ca)

- importance of taking care of symmetries, c

- parameter optimization: sequential & deriva

> TODO (a lot, but some are already in pro

extensions to proton-neutron systems, m towards driplines \sim input for astrophysic

not only NISQ algorithms, towards (early

Collaborators & Acknowledgement

Collaborators:

UTokyo: Takeshi Sato, Takumi Ogata

RIKEN iTHEMS: Tomoya Naito

RIKEN Nishina center: Masaaki Kimura

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Logo Mark / Symbol Mark + Logo Type

Typical situation in NP

Full CI: current limitation of the

¹²C (e.g. Hoyle state) 7 major shell excitation ~ 10^{19} dim. \rightarrow valence CI : T.Otsuka

frontiers of neutron-rich nuclei: Ni \sim 10

 \rightarrow CI variants: Monte Carlo Shell Model (MCSM) Quasi-vacua Shell Model (QVSM)

> T. Toga N. Shin

occupations of pair configurations

Note: the parameters are fixed to the ones giving exact results for noisy simulator and hardware results, post-selection is done

computational basis sampling so-called computational basis sampling technique, which two factors in the above equation, *|*h*m|* i*|* ² and *[|]*h*n[|]* ⁱ*[|]* computational basis sampling

*|*0i⁵ : *Ry*(2✓5) *•*

w. Kohda, R. Imai, K. Kanno, K. Mitarai, W. Mizukami, and Y. O. Nakagawa,
M. Kohda, R. Imai, K. Kanno, K. Mitarai, W. Mizukami, and Y. O. Nakagawa, m. Konaa, K. imai, K. Kanno, K. Mitarai, W. Mizukami, and T. O. Nakagawa,
Quantum expectation-value estimation by computational basis sampling, Phys. Rev. Res. 4, 033173 (2022).

$$
\langle \psi | H | \psi \rangle = \sum_{m,n=0}^{N} \langle \psi | m \rangle \langle m | H | n \rangle \langle n | \psi \rangle
$$

\n
$$
= \sum_{m,n=0}^{N} |\langle m | \psi \rangle|^2 |\langle n | \psi \rangle|^2 \frac{\langle m | H | n \rangle}{\langle m | \psi \rangle \langle \psi | n \rangle}. \tag{7}
$$

\n
$$
\langle \psi | X_j \otimes X_k | \psi \rangle = \sqrt{\sigma_j^2 \sigma_k^2} \operatorname{sgn} [\sigma_j \sigma_k], \tag{8}
$$

\n
$$
= \sqrt{\langle N_j \rangle \langle N_k \rangle} \operatorname{sgn} [\langle X_j X_k \rangle] \tag{9}
$$

\n
$$
\uparrow \text{ ansatz plus H-gates on all qubits}
$$

can remove the configurations violating particle numbered particle numbered in $\frac{1}{2}$ conservative as α and the sequence as α and the sequence as α and the theorem is α and the theorem is α *ⁱ* ⌘ *|*h0 *···* 1*ⁱ ···* 0*|* i*|* is the projector onto the ↑ equiv. to Z measurement of the ansatz circuit

You just need 2 measurements, ansatz & ansatz + H-gates ated by measurements of the quantum circuits, and the quantum cir details on how to map these terms to map these terms to map these terms to map these terms to quantum circuits
The complex terms to map the complex terms to guarantee the complex terms to map the complex terms of the comp

Derivative-free & Sequential optimization of the parameters and the other with Hadamard gates added to every qubit. *^L*(✓) = ^X *w^k* h'*|U†* (✓) *HkU* (✓)*|*'i*,* (B1) U uti \mathcal{L}_{max} \mathcal{L}_{max} \mathcal{L}_{max} optimization \mathcal{L}_{max} α are a sequential optimization, \bigcap ity of optimization, but it provides entry α itive-tree & Seqi

K. M. Nakanishi, K. Fujii, and S. Todo, measurements, where **N** is the content of n , θ Phys. Rev. Res. 2. iSWAP gate, which requires auxiliary qubits and Sequential minimal optimization for quantum-classical hybrid algorithms,
Phys. Ray, Ras. 2, 043158 (2020) transformation corresponding to the circuit. *^L*(✓) = ^X *w^k* h'*|U†* (✓) *HkU* (✓)*|*'i*,* (B1) Phys. Rev. Res. 2, 043158 (2020). k. M ments of the sequence of the cone \mathcal{S} end the parties and the parties and the other with Phys \mathcal{R} . ous
Dhi because the former method requires only two measure-We followed the optimization method proposed in \mathcal{P}_c the optimization method proposed in \mathcal{P}_c f ev. Kes. Z, U431

expectation value of the Hamiltonian, in the Hamiltonian, is written in the Hamiltonian, in the Hamiltonian, i
The Hamiltonian, is written in the Hamiltonian, in the Hamiltonian, is written in the Hamiltonian, in the Hami

$$
\mathcal{L}\left(\bm{\theta}\right)=\sum_{k}w_{k}\left\langle \varphi\right|U^{\dagger}\left(\bm{\theta}\right)\mathcal{H}_{k}U\left(\bm{\theta}\right)\left|\varphi\right\rangle ,\qquad\left(\text{B1}\right)\text{ \ =Energy expectation value}
$$

quantum continuizi consider which quantum circuit is to be chosen among *Lotimizing Ruante rotation angle P* pullizing by gate rotation angle Optimizing Ry-gate rotation angle Optimizing ty-gate rotation angle

$$
\mathcal{L}(\theta_1) = A \cos (\theta_1 - B) + C. \tag{B2}
$$

appendix Big Display Display and the contraction of circuit parameters and contract parameters and contract pa Ω _{th} Ω _{th} Ω _{th} Ω of Ω in Ω of $\$ cry-yui α *CR* α ₋*ate* rotation angle ppurinzing cry-gale Foldtion angle Optimizing cRy-gate rotation angle

$$
\mathcal{L}(\theta_j) = A \cos \theta_j + B \sin \theta_j + C \cos 2\theta_j + D \sin 2\theta_j + E.
$$
\n(B3)

By measuring five di↵erent points, ✓*^j* , ✓*^j* + ⇡*/*5, . . . ,

Optimization of the circuit parameters

parameters so far are fixed to the ones giving exact

K. M. Nakanishi, K. Fujii, and S. Todo: Sequential minimal optimization for quantum-classical hybrid algorit