

Neural quantum state tomography

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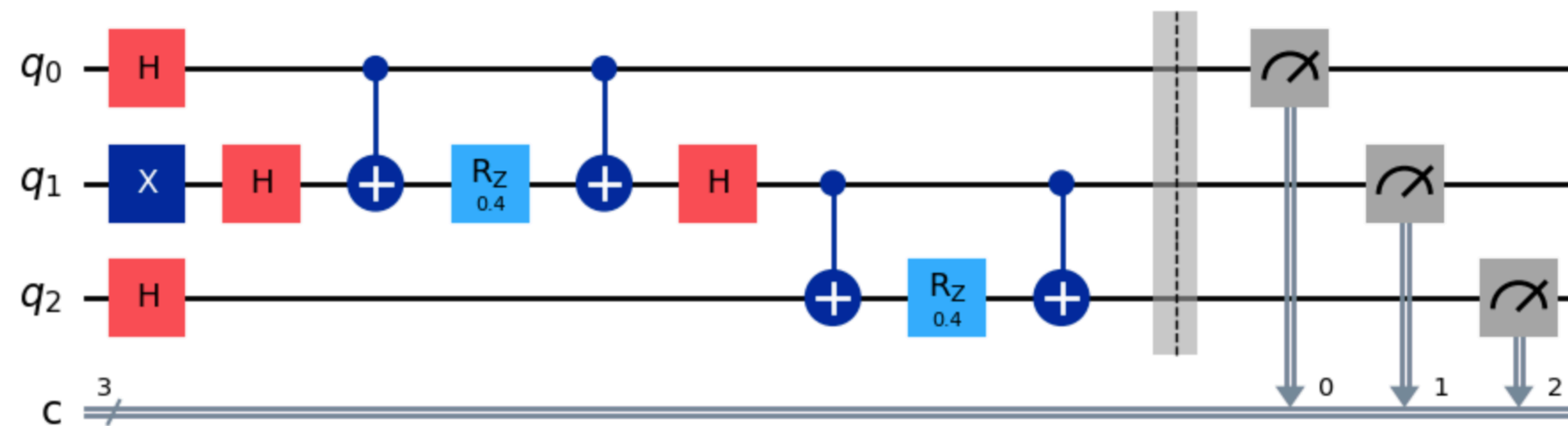
[arXiv:2405.06864](https://arxiv.org/abs/2405.06864)



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Quantum states in quantum computing

Typical quantum circuit



Quantum computing is powerful
because it does massive parallel computing

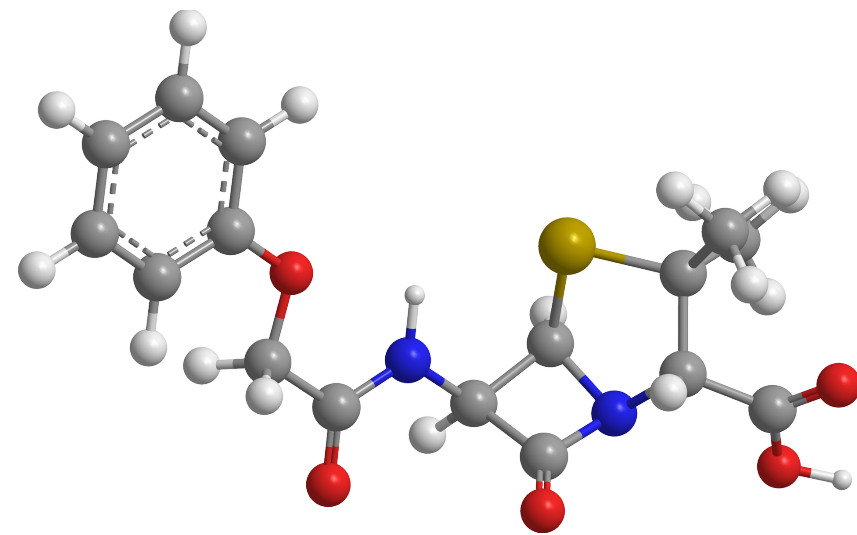
Measurement:
Exponential information loss!

$$|\psi\rangle = \sum_{k=1}^{2^n} a_k |k\rangle$$

$|k\rangle$

Challenge of quantum measurements

Chemistry



$$H = \sum f_{ij} a_i^\dagger a_j + \sum f_{ijkl} a_i^\dagger a_j^\dagger a_k a_l$$

$$\rightarrow \sum_k \tilde{f}_k P_k \quad (P_k : \text{Pauli operators})$$

Energy $E = \langle H \rangle = \sum_k \tilde{f}_k \langle P_k \rangle$

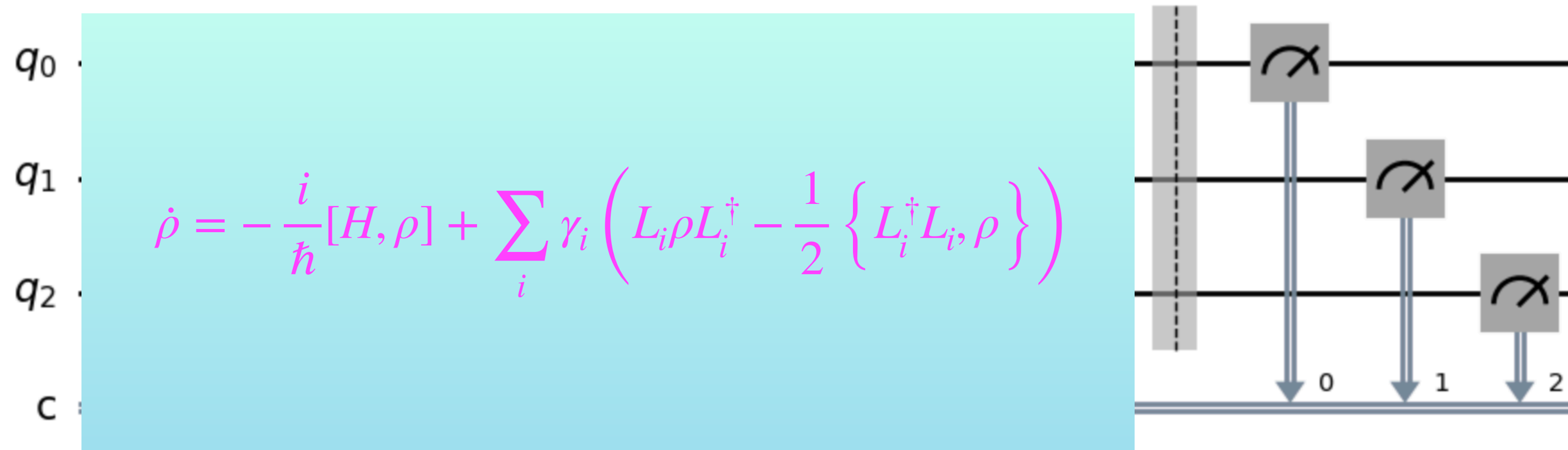
Hydrogen atom

$$P_k = [\text{'IIII'}, \text{'IIIZ'}, \text{'IIZI'}, \text{'IIZZ'}, \text{'IZII'}, \text{'IZIZ'}, \text{'ZIII'}, \text{'ZIIZ'}, \text{'IZZI'}, \text{'ZIZI'}, \text{'ZZII'}, \\ \text{'YYYY'}, \text{'XXYY'}, \text{'YYXX'}, \text{'XXXX'}]$$

n -qubit: $|P| \sim \mathcal{O}(n^4)$,

of measurements= N : Sampling error $\sim 1/\sqrt{N}$

Unknown quantum evolution



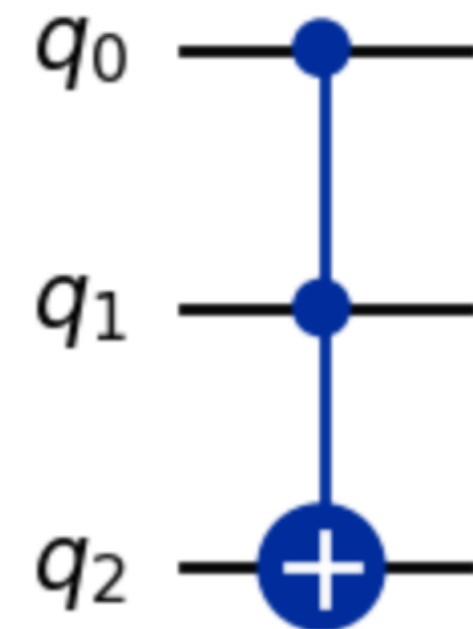
Is the final state entangled or separable?

For a Hermitian operator A $\begin{cases} \text{Entangled state} \rightarrow \text{Tr}(A\rho) < 0 \\ \text{Tr}(A\rho) \geq 0 \rightarrow \text{Separable state} \end{cases}$

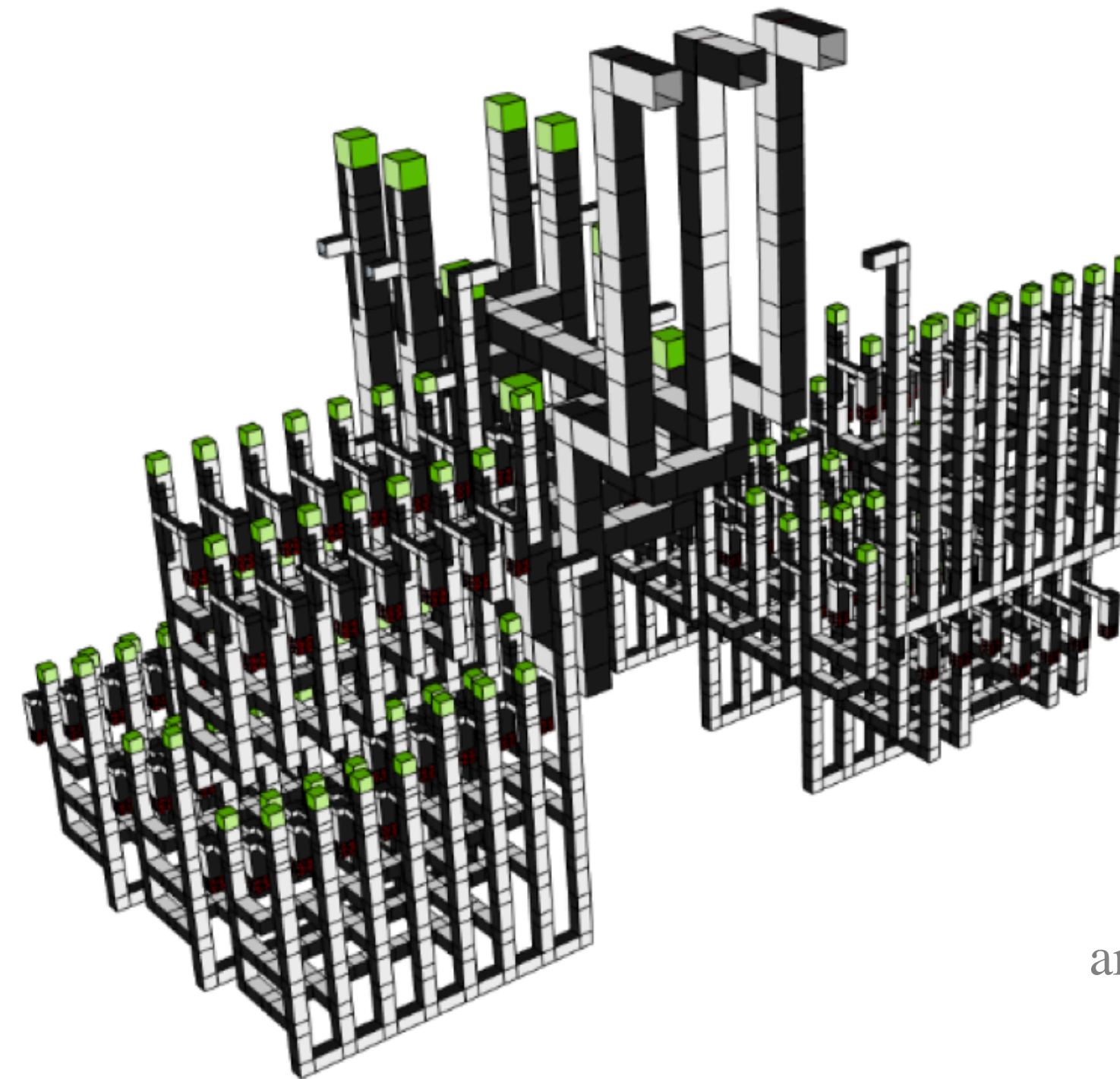
However $|\text{Tr}(A\rho)|$ can be very small: **Sensitive to sampling error**

Fault tolerant quantum computing

Significant overhead



NISQ

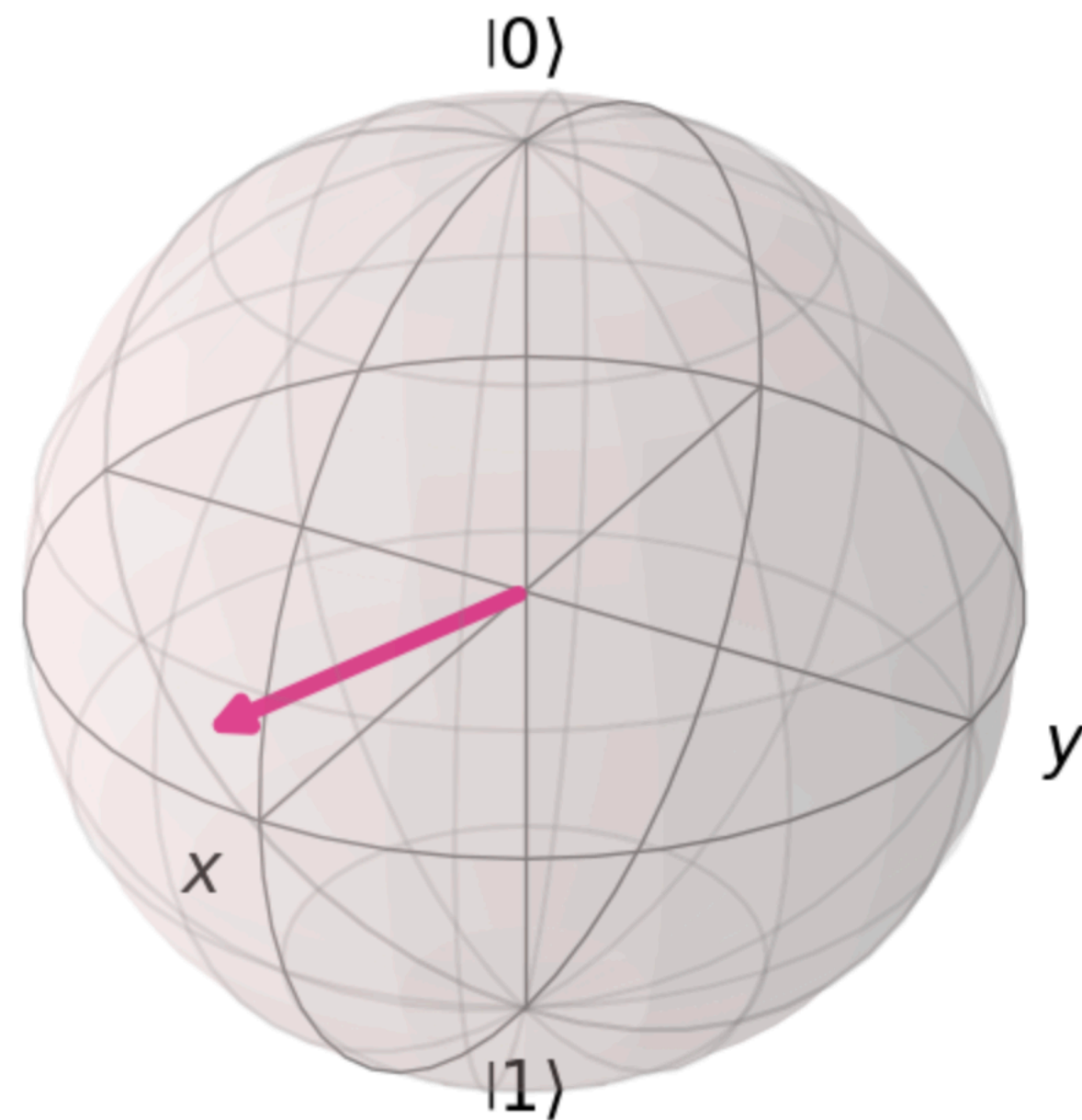


FTQC

arXiv:2011.04149

Reducing measurement complexity is useful in FTQC too!

State Tomography



$$\rho = \frac{1}{2} \left(I + a_x X + a_y Y + a_z Z \right)$$

$$a_x = \text{tr}(\rho X) = \langle X \rangle$$

$$a_y = \text{tr}(\rho Y) = \langle Y \rangle$$

$$a_z = \text{tr}(\rho Z) = \langle Z \rangle$$

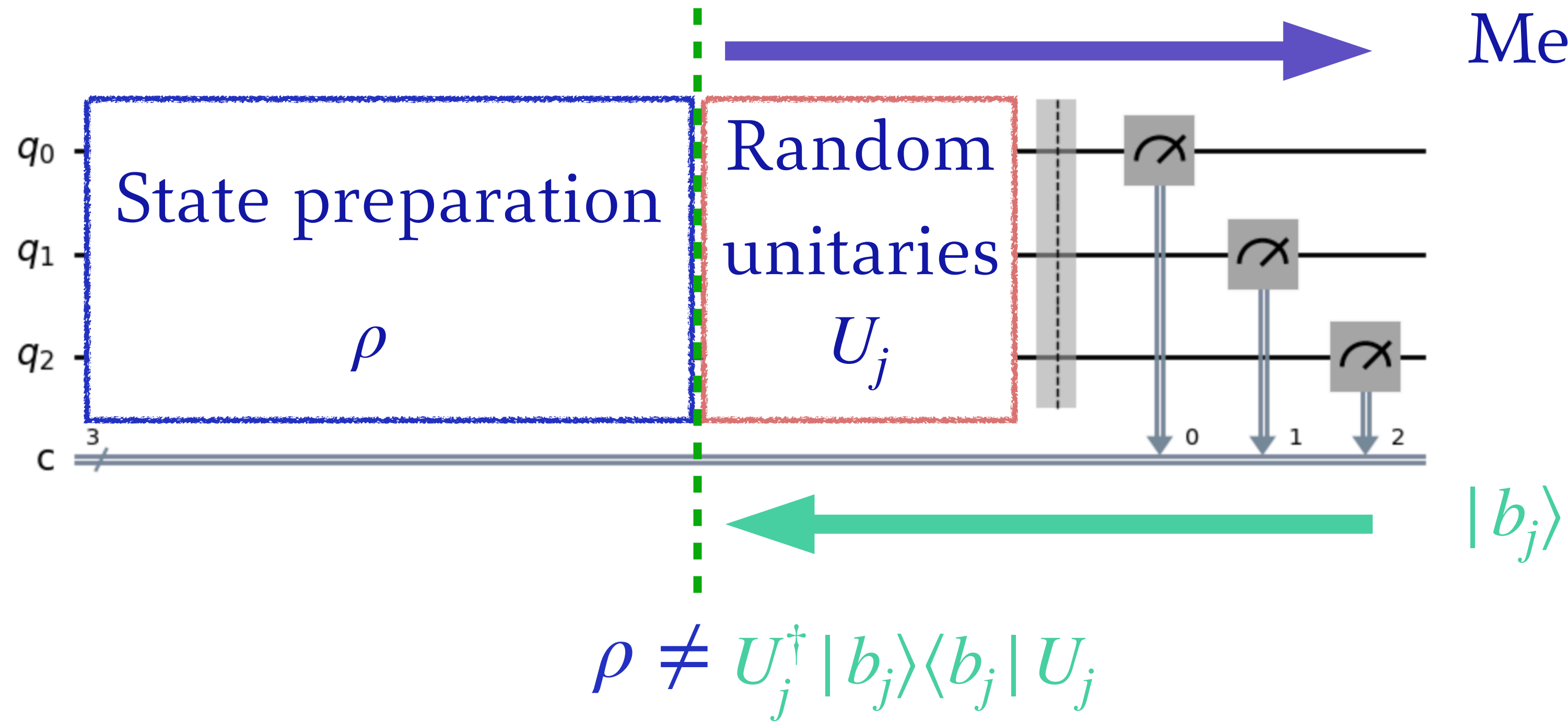
Multi-qubit case

Measure all possible Paulis (4^n)

$\langle IXIZ \rangle, \langle YYZX \rangle, \dots$

Classical Shadow Tomography

[Huang-Kueng-Preskill '20]



$$\rho \neq U_j^\dagger |b_j\rangle\langle b_j| U_j$$

$$\mathcal{M}(\rho) = \mathbb{E} \left[U_j^\dagger |b_j\rangle\langle b_j| U_j \right]$$

$$\rho = \mathbb{E} \left[\mathcal{M}^{-1} \left(U_j^\dagger |b_j\rangle\langle b_j| U_j \right) \right]$$

Classical Shadow Tomography

$$\hat{\rho}_i = \mathcal{M}^{-1} \left[U_i^\dagger |b_i\rangle\langle b_i| U_i \right]$$

$$\mathbb{E}(\hat{\rho}) = \rho$$

$$\langle O_a \rangle = \frac{1}{K} \sum_{i=1}^K \text{tr} (\hat{\rho}_i O_a)$$

If we want to estimate M observables with accuracy ϵ ,

Sampling complexity : $K = \mathcal{O}(\log(M) \max_i \|O_i\|_{\text{shadow}}^2 / \epsilon^2)$

Exponentially efficient

Classical Shadow Tomography

However, each $\hat{\rho}$ is not necessarily physical

For instance, if we measure $|0\rangle \Rightarrow \hat{\rho} = \begin{pmatrix} 2 & 0 \\ 0 & -1 \end{pmatrix}$

Suppose we want to compute the entanglement entropy

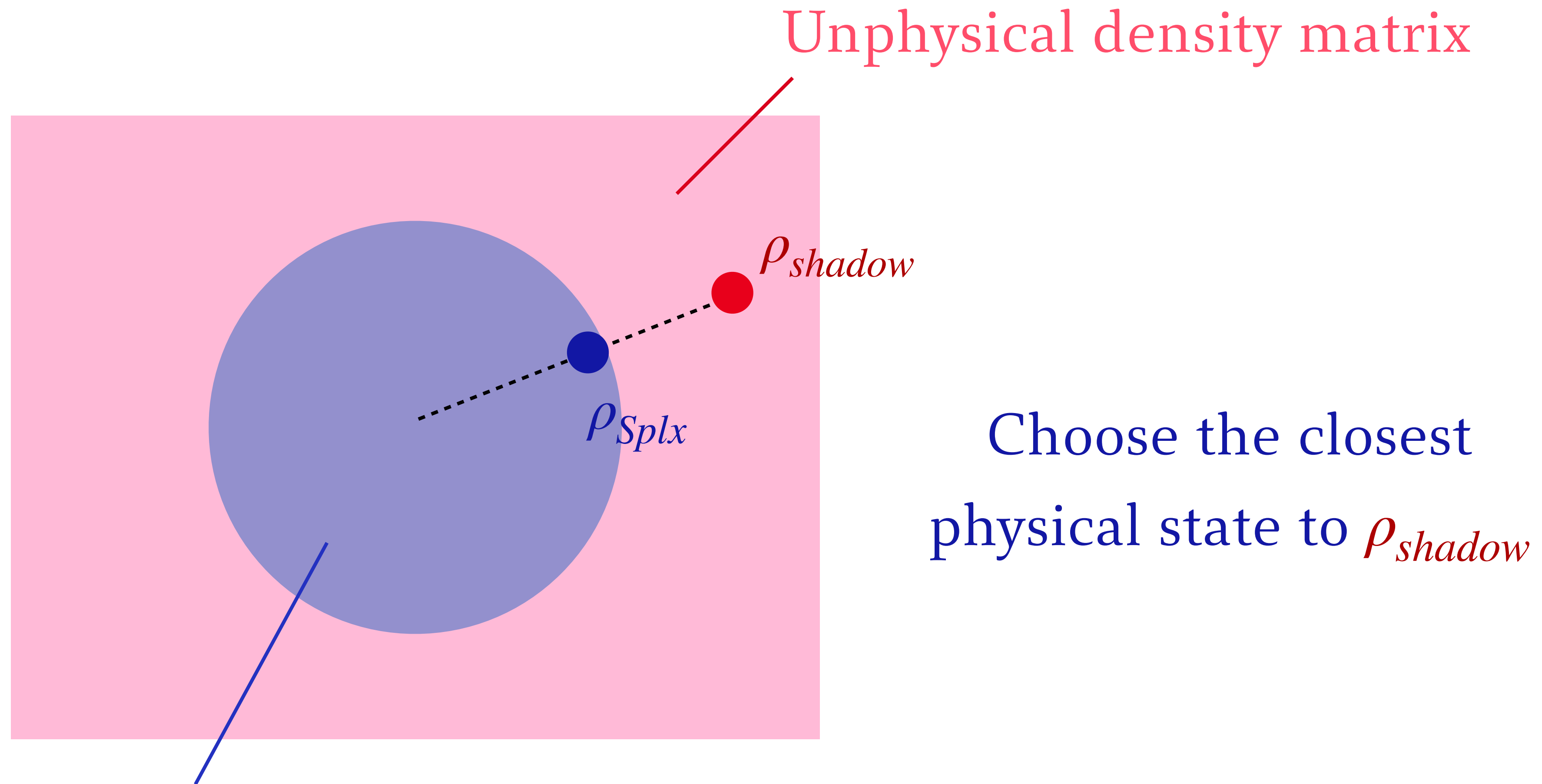
( This is just for intuition, sampling number = 1)

$$\text{“ } S = -\text{tr}(\rho \log \rho) \text{ ”}$$

$$= - (2 \log 2 + (-1) \log(-1))$$

Ill-defined

Solution 1 : Simplex projection



Unphysical density matrix

ρ_{shadow}

ρ_{Splx}

Choose the closest
physical state to ρ_{shadow}

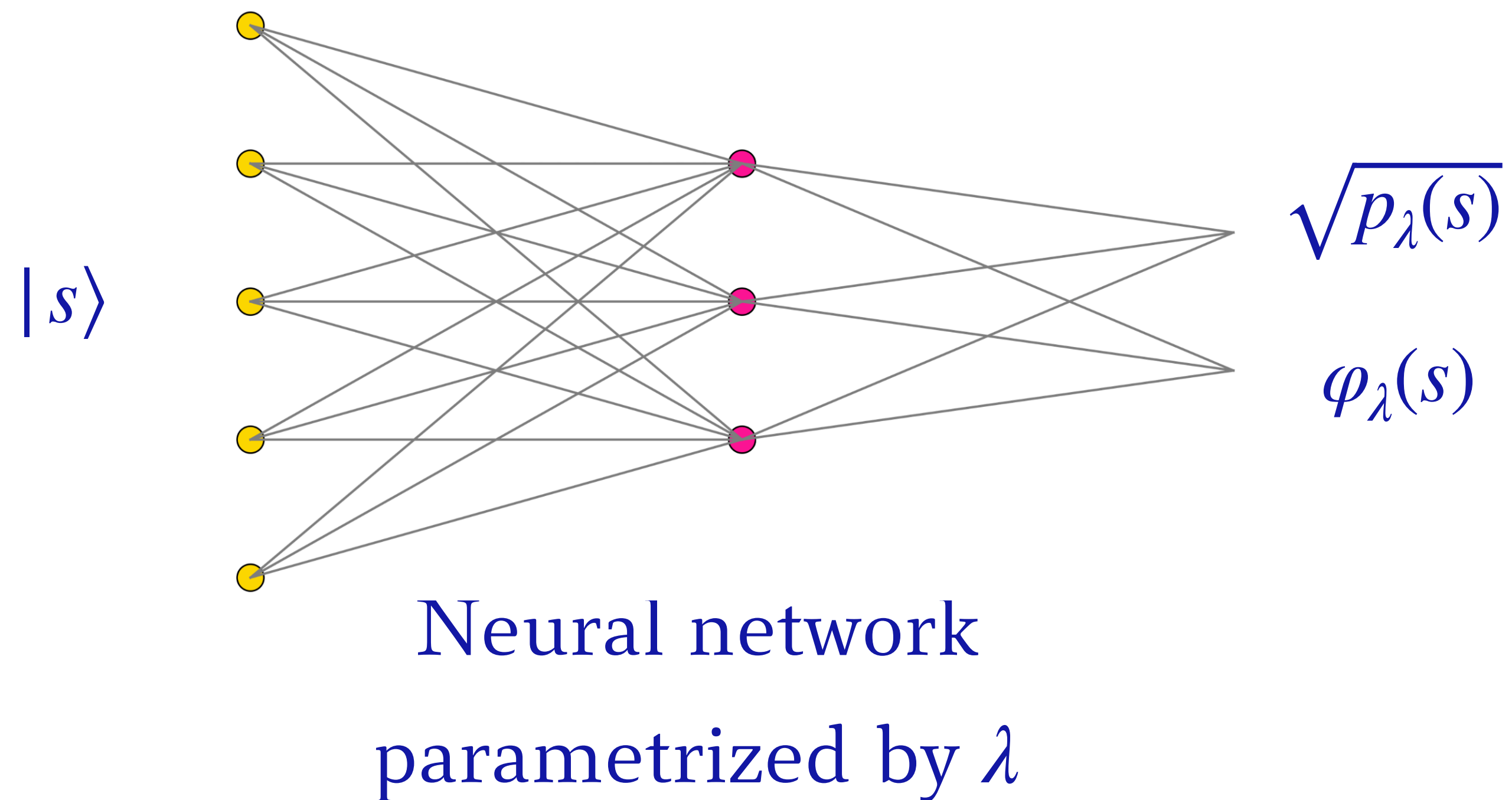
Physical density matrix (positive eigenvalues between 0 and 1)

Solution 2 : Neural network quantum state

[Torlai-Mazzola-Carrasquilla-Troyer-Melko-Carleo '18, Carrasquilla-Torlai-Melko-Aolita '19]

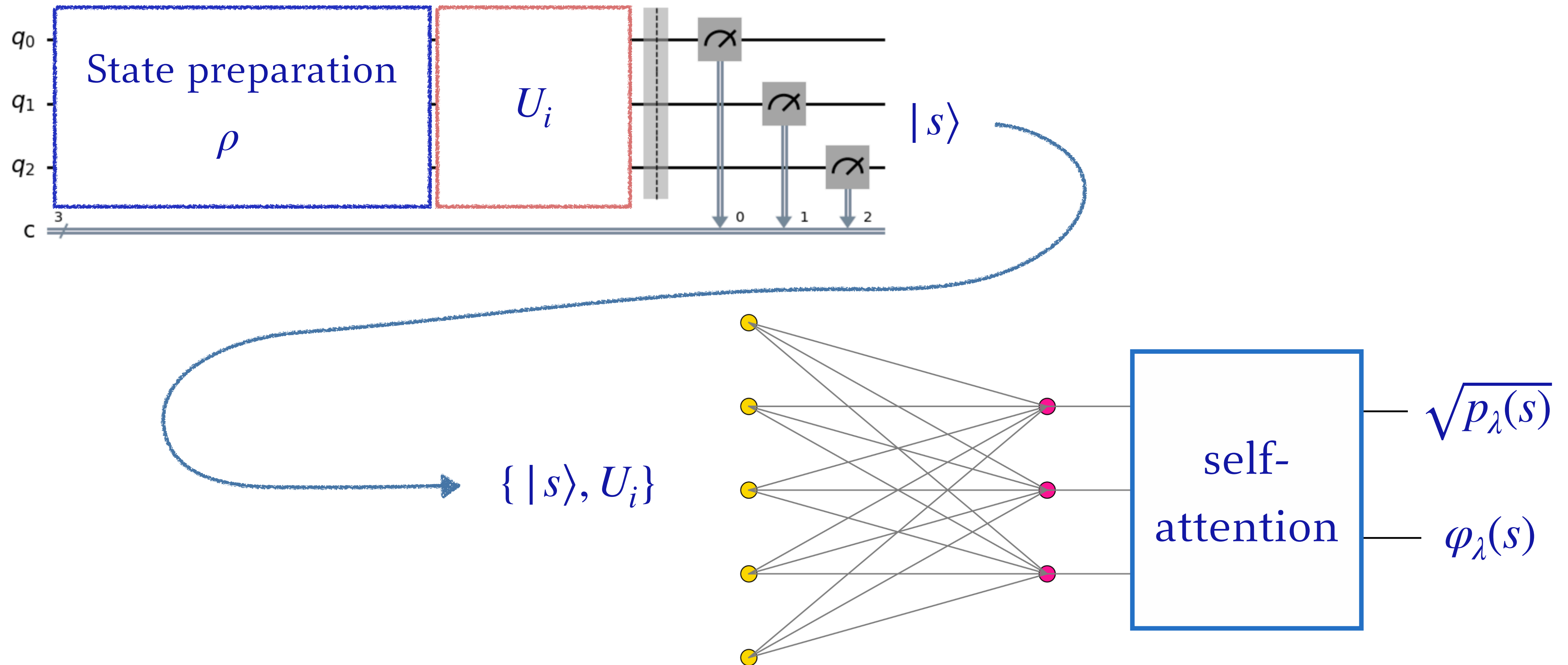
$$|\psi\rangle = \sum_s \sqrt{p_\lambda(s)} e^{i\varphi_\lambda(s)} \underline{|s\rangle}$$

Computational basis



Using neural network to reduce errors

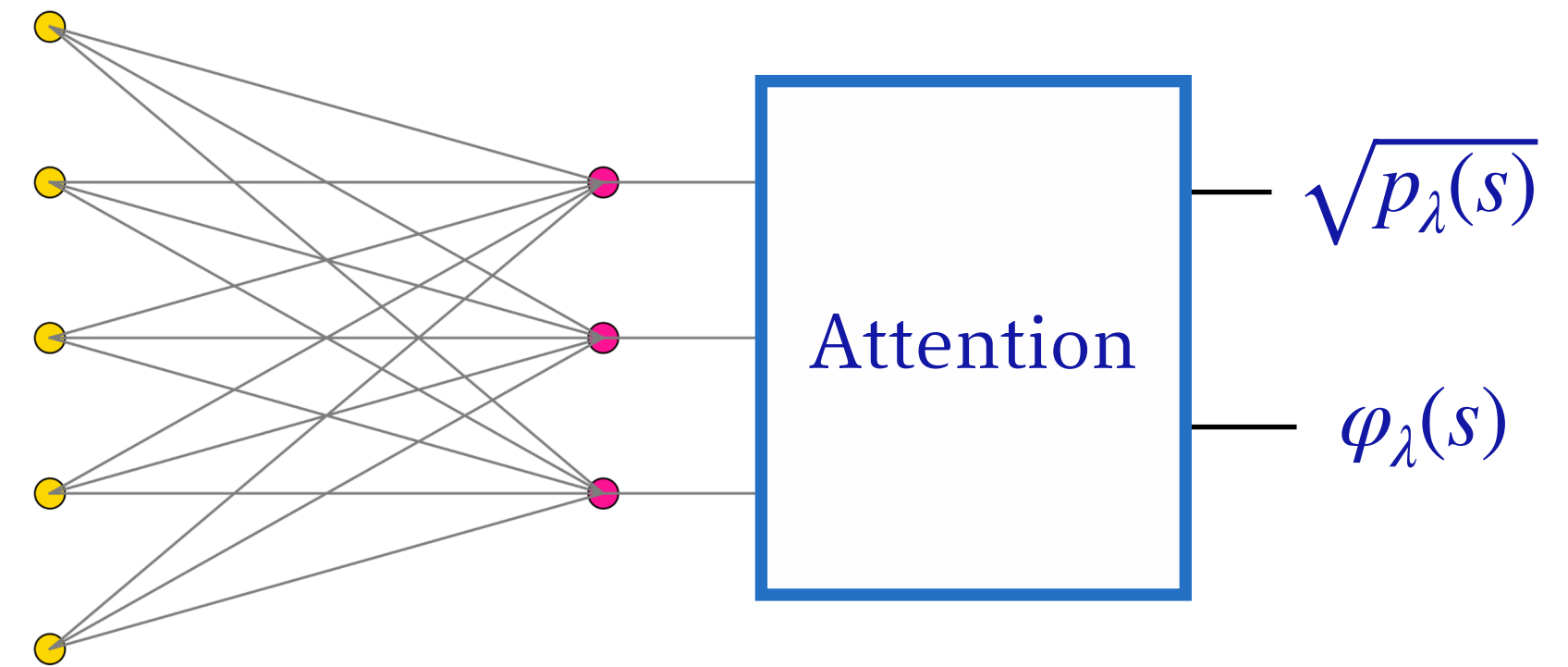
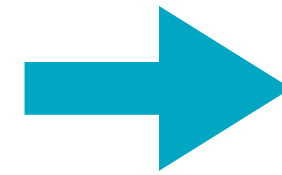
Neural shadow quantum state [Wei-Coish-Roangh-Muschik, '23]



Mixed states [arXiv:2405.06864]

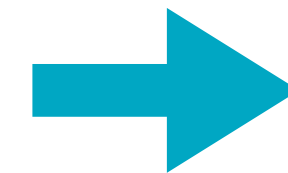
Pure state

$$|\psi\rangle$$



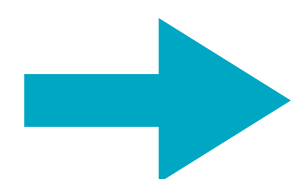
Mixed state

$$\rho = \begin{pmatrix} p_1|\psi_1\rangle\langle\psi_1| & 0 & 0 & \dots \\ 0 & p_2|\psi_2\rangle\langle\psi_2| & 0 & \dots \\ \vdots & \vdots & \ddots & \ddots \end{pmatrix}$$



Additional qubits (B) to purify the mixed state

$$|\Psi\rangle = \sqrt{p_1} |\psi_1\rangle_A |0\rangle_B + \sqrt{p_2} |\psi_2\rangle_A |1\rangle_B + \sqrt{p_3} |\psi_3\rangle_A |2\rangle_B + \dots$$

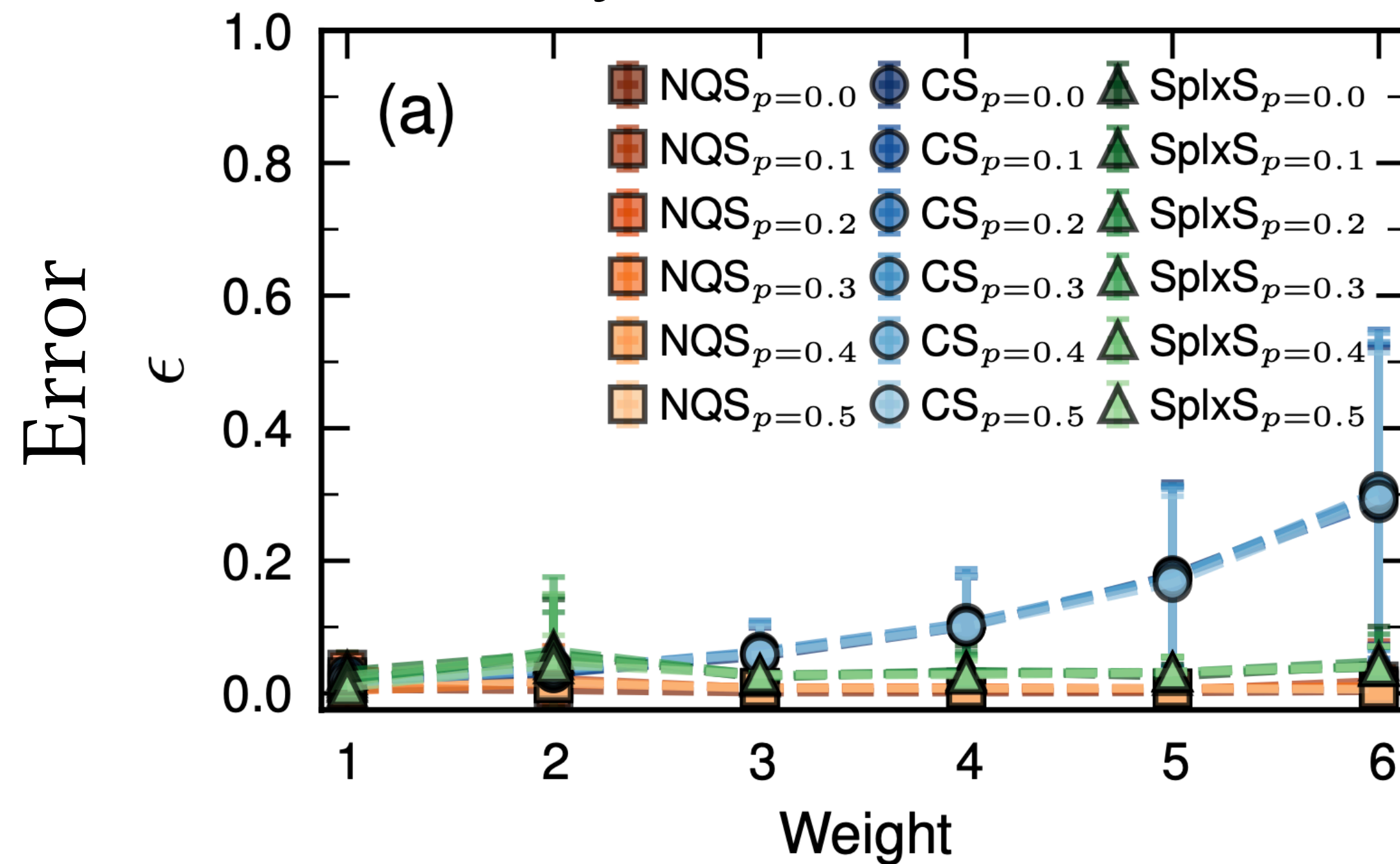


$$\rho = \text{tr}_B |\Psi\rangle\langle\Psi|$$

Numerical results

6 qubit GHZ state: $|00\dots 0\rangle + |11\dots 1\rangle$

Accuracy of Pauli observable

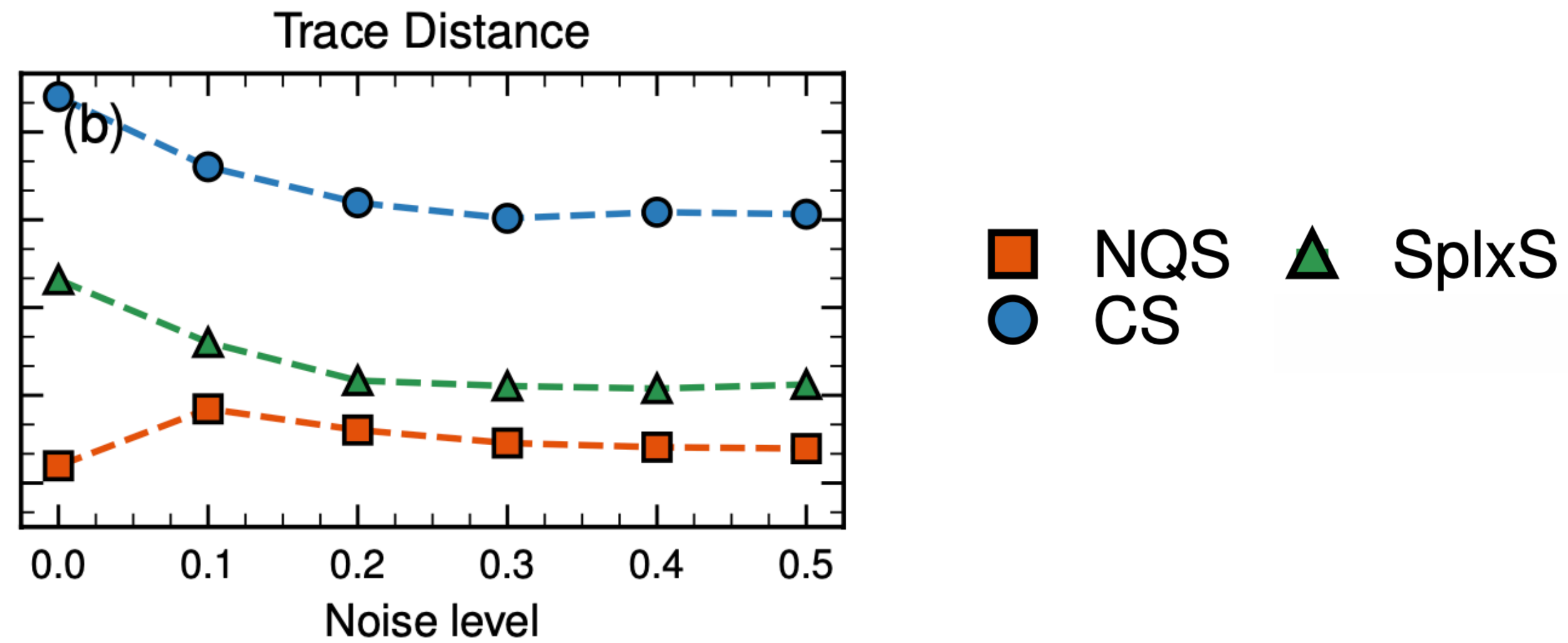


(Number of non-identity Paulis)

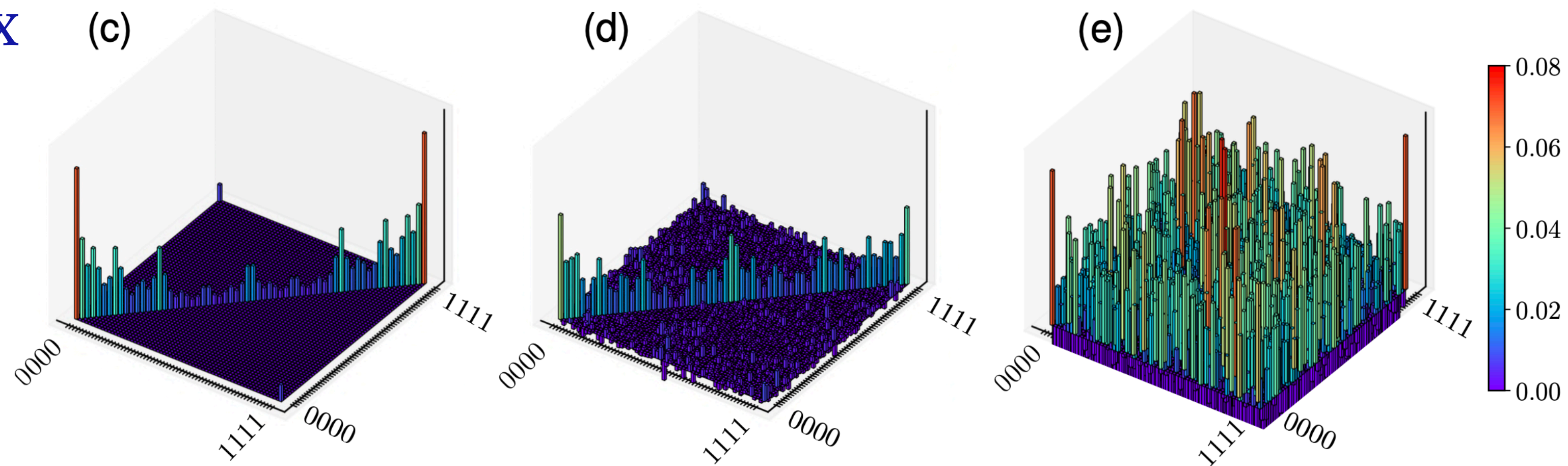
IXIII : weight 1

XXYZIY : weight 5

Numerical results



Density matrix



Theory

Neural shadow

Classical shadow

Conclusion

- Reducing sampling complexity is a critical challenge in quantum computing

Even without additional experimental information,
we can still improve the accuracy of estimating physical quantities
through post-processing

Neural shadow quantum states (autoregressive model) works well
for both pure and mixed states