

Renormalization Of Beta Decay At 3+ Loops

PRESENTED BY

PETER VANDER GRIEND

ON BEHALF OF

KAUSHIK BORAH, RICHARD J. HILL, RYAN PLESTID

CIPANP | MADISON WISCONSIN | JUNE 2025

ARXIV:2309.15929 ,
ARXIV:2309.07343 ,
ARXIV:2402.13307 ,
ARXIV:2402.14769 .

Caltech

Neutrino Theory Network

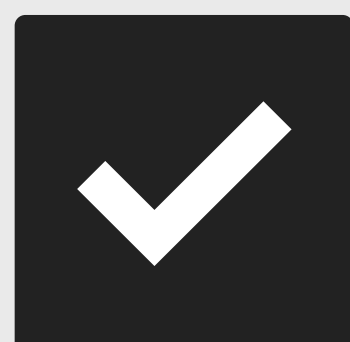




PART 1

EFT & β DECAY

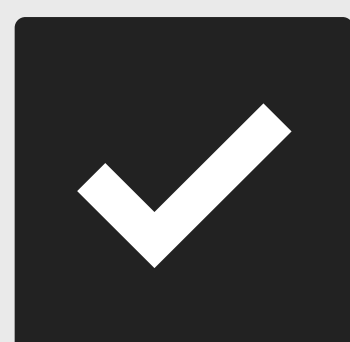
- Motivation & relevance for **fundamental physics**.
- Necessary **precision**, and requisite **loop orders**.



PART 2

FERMI FUNC.

- **Point-like** EFT of nuclei and leptons.
- The **Fermi function** from loops.



PART 3

RAD. CORR.

- Structure of **radiative corrections** from EFT.
- Renormalization group **resummation of logarithms**.

Quark Mixing In The SM

FUNDAMENTAL
CONSTANTS
OF NATURE

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$



CKM \equiv CABIBBO-KOBAYASHI-MASKAWA



Quark Mixing In The SM

BETA DECAY

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$



CKM \equiv CABIBBO-KOBAYASHI-MASKAWA



CKM Unitarity

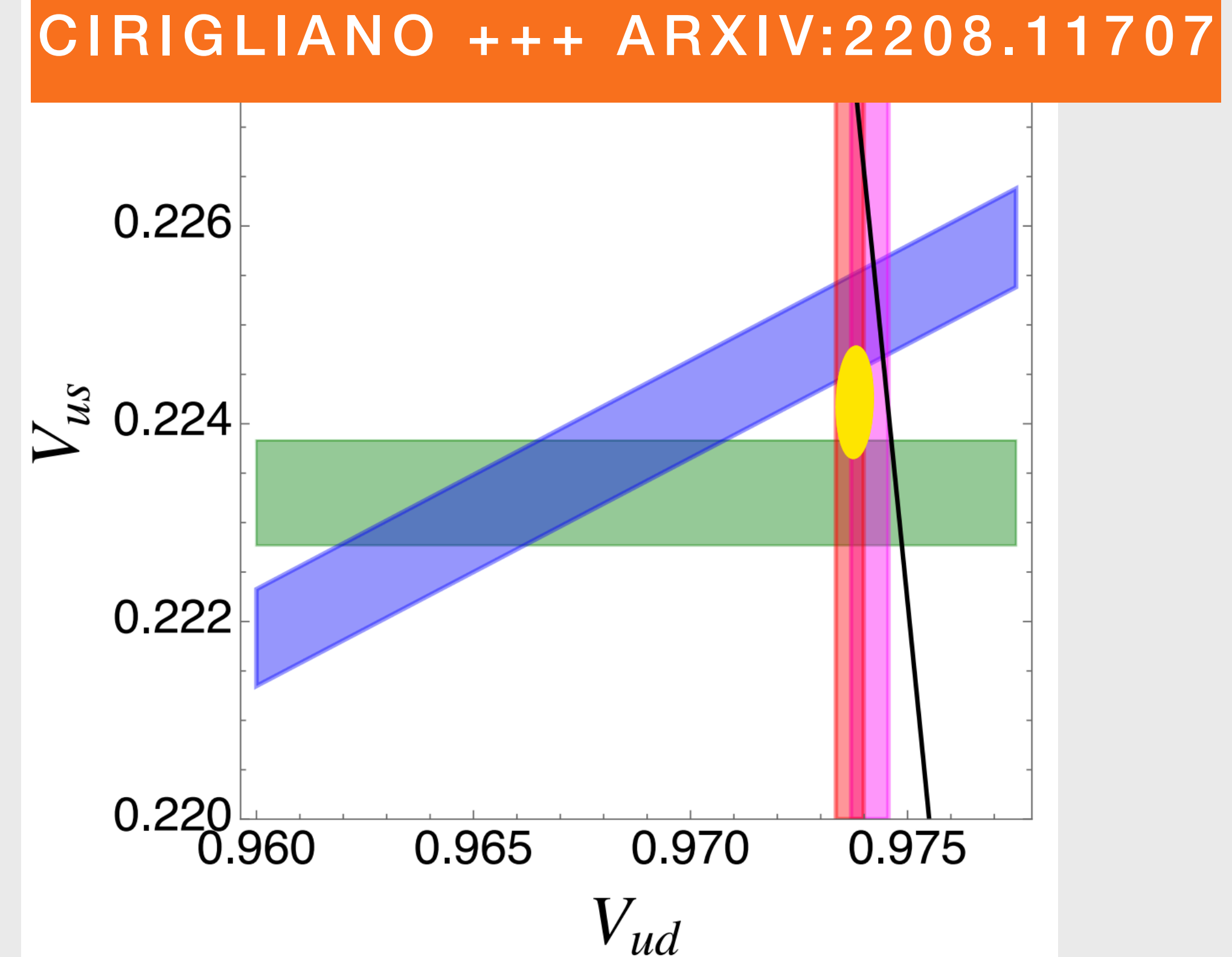
FIRST ROW UNITARITY

$$|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 1$$

IN WOLFENSTEIN NOTATION

$$1 - \lambda_{ud}^2 + \lambda_{us}^2 + O(\lambda^6) = 1$$

- Tension in first-row CKM unitarity at $\sim 10^{-3}$ level.
- **If** theory is under control: strong hints of new physics!

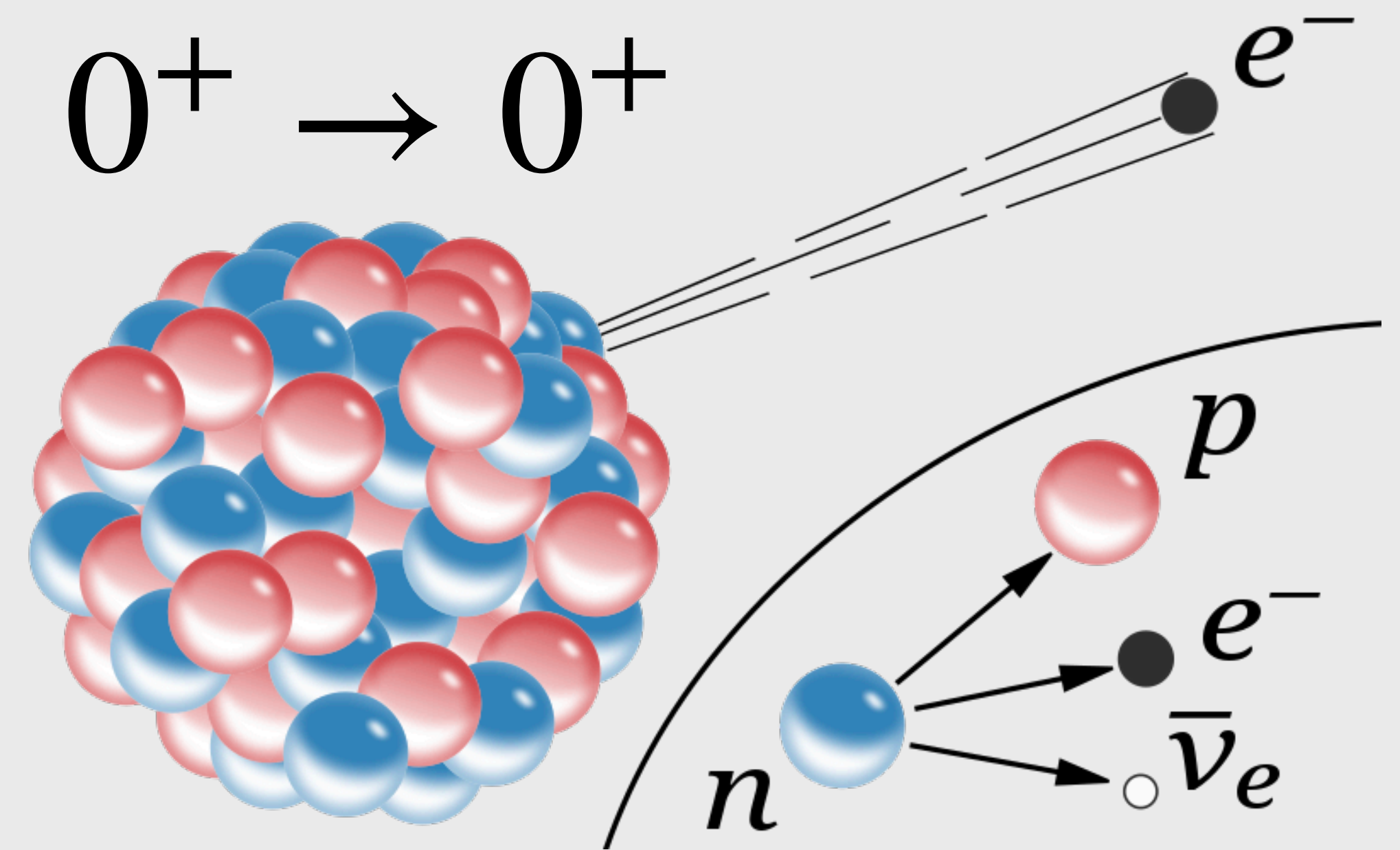


Calculate Matrix Element To High Order

- At this level of accuracy
- Need control over corrections in low-energy theory **at least** at

$O(Z^2\alpha^3)$ i.e. 3+ loops

- **Enhancements from charge of nucleus are important.**



- Precision goal: 100 ppm

Historical Approach

$$\mathcal{F}t \equiv ft(1 + \delta'_R)(1 + \delta_{\text{NS}} - \delta_C) = \frac{K}{2G_V^2(1 + \Delta_R^V)},$$

- The “ ft ” value includes the Fermi function (Dirac w.f.) .
- Δ_R^V is a short-distance correction.
- δ_{NS} and δ_C are related to nuclear structure.
- δ'_R computed in the “independent particle model”.

THIS TALK: A MODEL-INDEPENDENT EFT APPROACH
TO BOTH THE FERMI FUNCTION AND δ'_R

Main Idea

- At $\sim \text{MeV}$ energies nuclei appear point-like.
- Long-distance QED corrections can be compute in an EFT.
- At "leading-power" only sensitive to charge of nucleus.

DEGREES OF FREEDOM


1. Heavy nuclei A and B .
 2. Electrons.
 3. Photons.
- Just like HQET but nuclei are the "heavy quarks".
 - Gauge field is photon.

Main Idea

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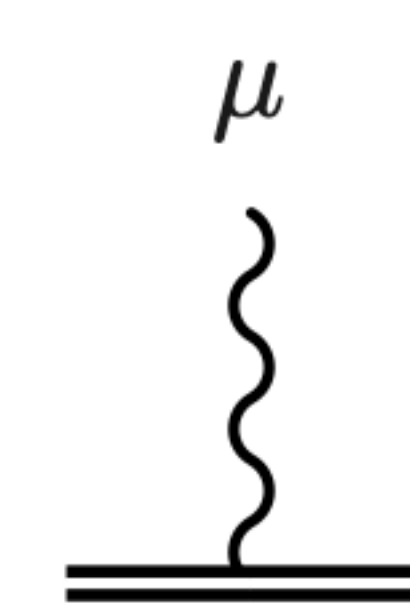
SIMPLE FEYNMAN RULES

COULOMB-FIELD



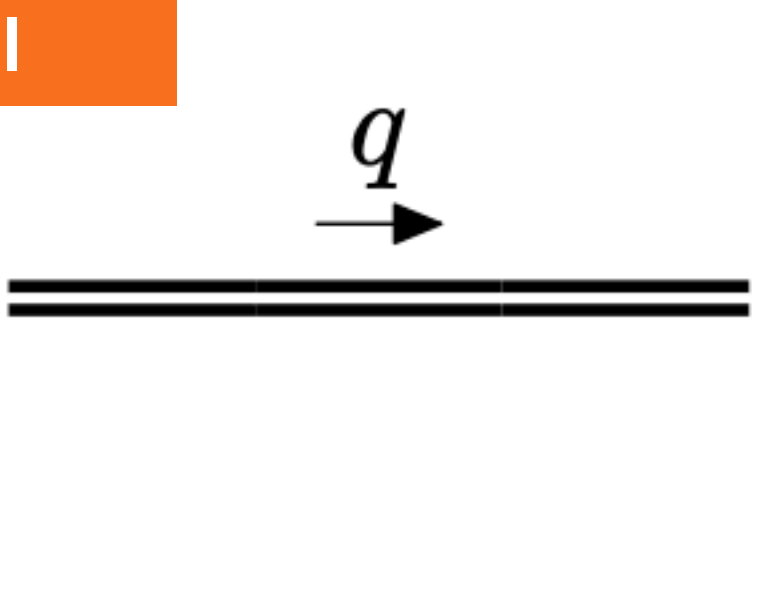
$$= iZe \delta_0^\mu 2\pi \delta(q^0)$$

NUCLEI



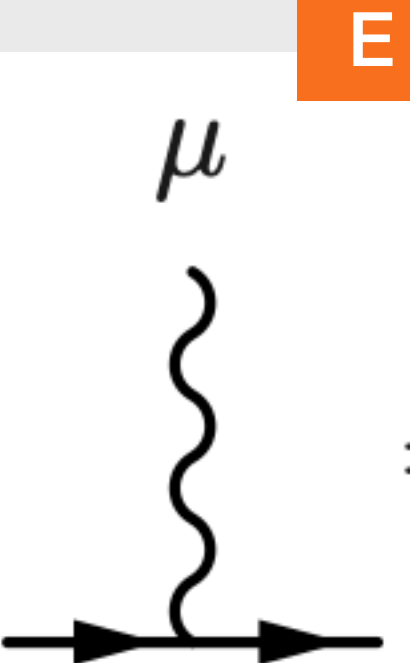
$$= ie \delta_0^\mu$$

NUCLEI



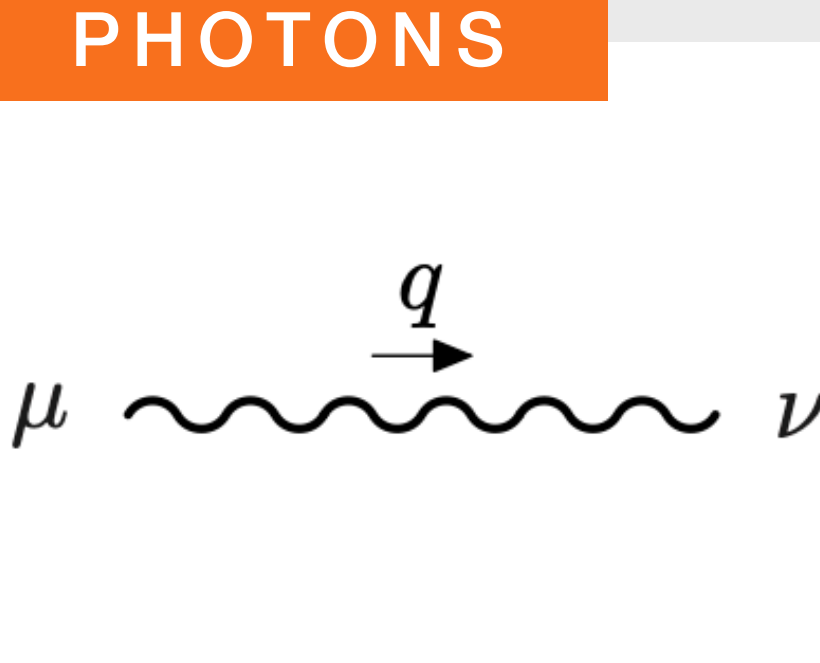
$$= \frac{i}{q^0 + i0}$$

ELECTRONS



$$= -ie \gamma^\mu$$

PHOTONS



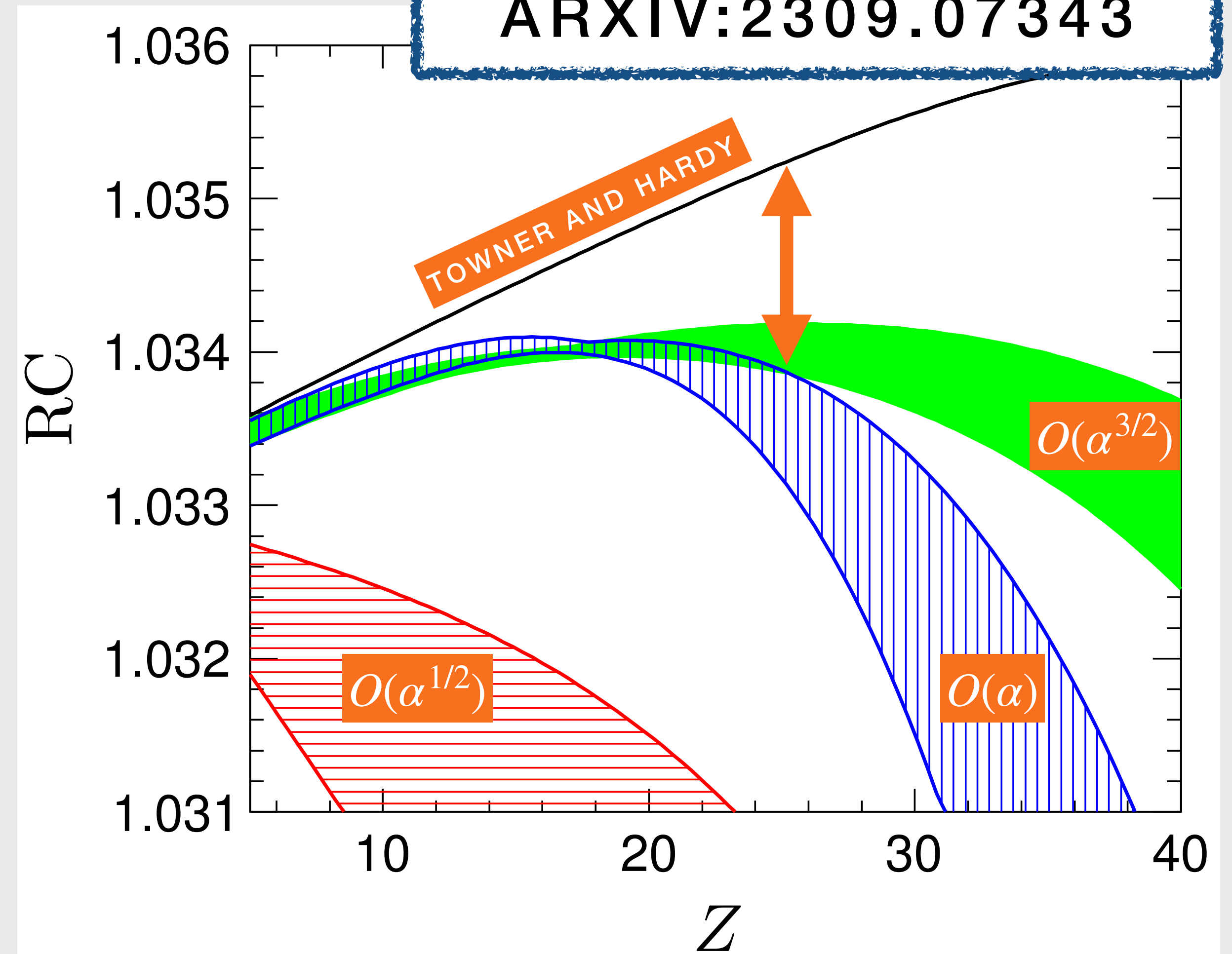
$$= \frac{-i}{q^2 - \lambda^2 + i0}$$

Impact For Flavour Physics

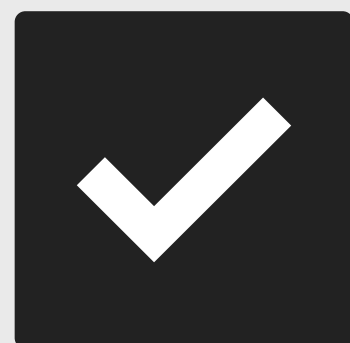
SHIFTING δ_3

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transition	$(\Delta a) \times Z^2 \alpha^3 \log(\Lambda/m)$
$^{14}\text{O} \rightarrow ^{14}\text{N}$	-1.1×10^{-4}
$^{26m}\text{Al} \rightarrow ^{26}\text{Mg}$	-3.2×10^{-4}
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COUNTING $Z \sim \log \sim 1/\sqrt{\alpha}$



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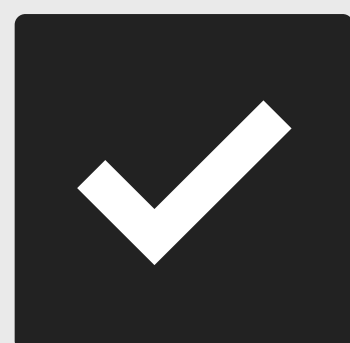
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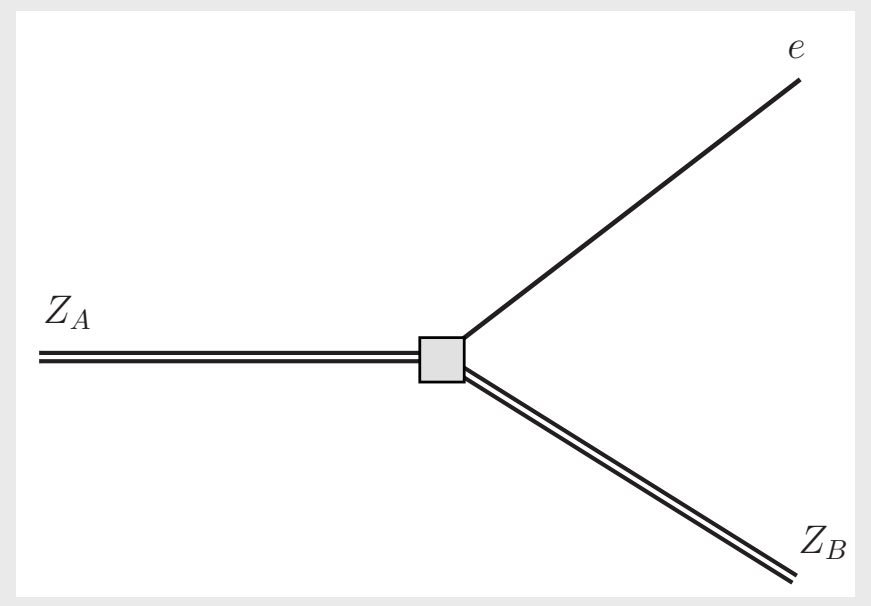
PART 3

RAD. CORR.

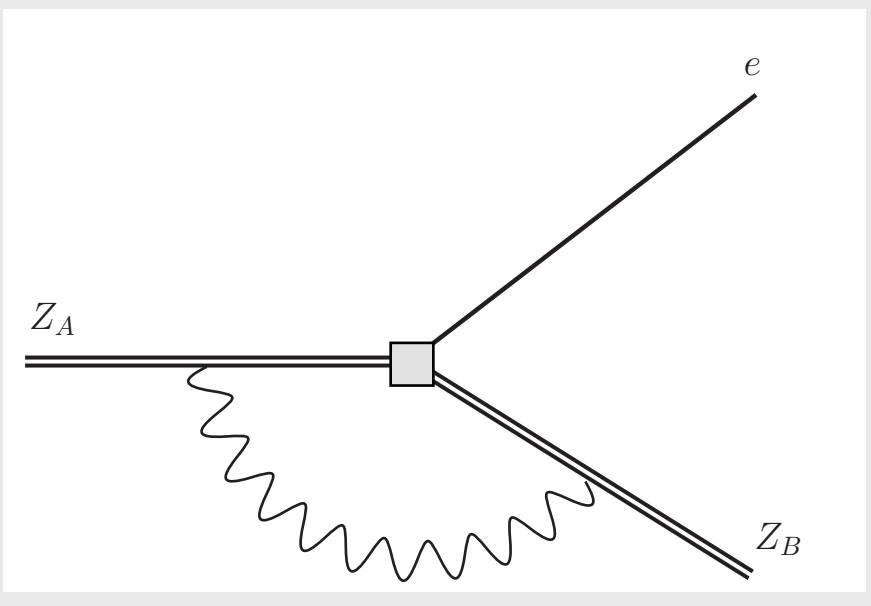
- Structure of **radiative corrections** from EFT.
- Renormalization group **resummation of logarithms**.

We Just Compute Diagrams!

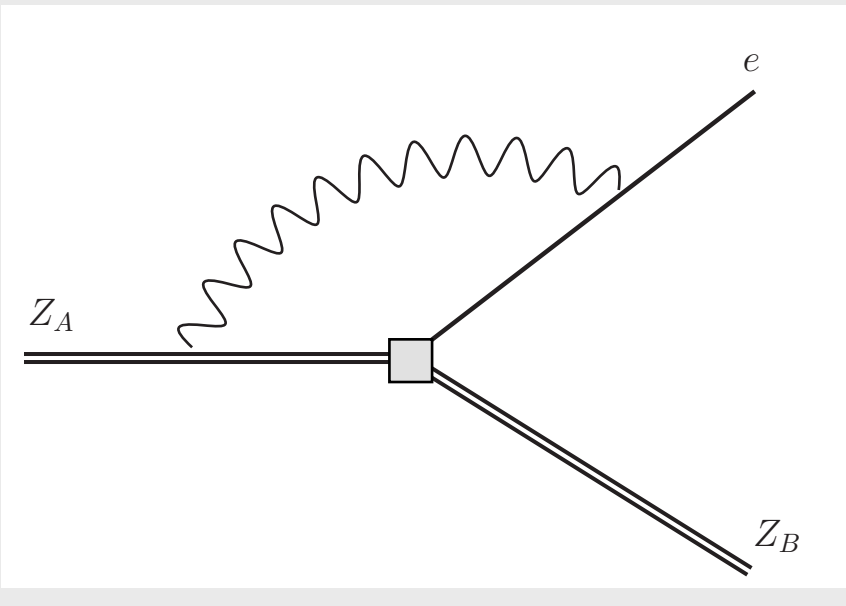
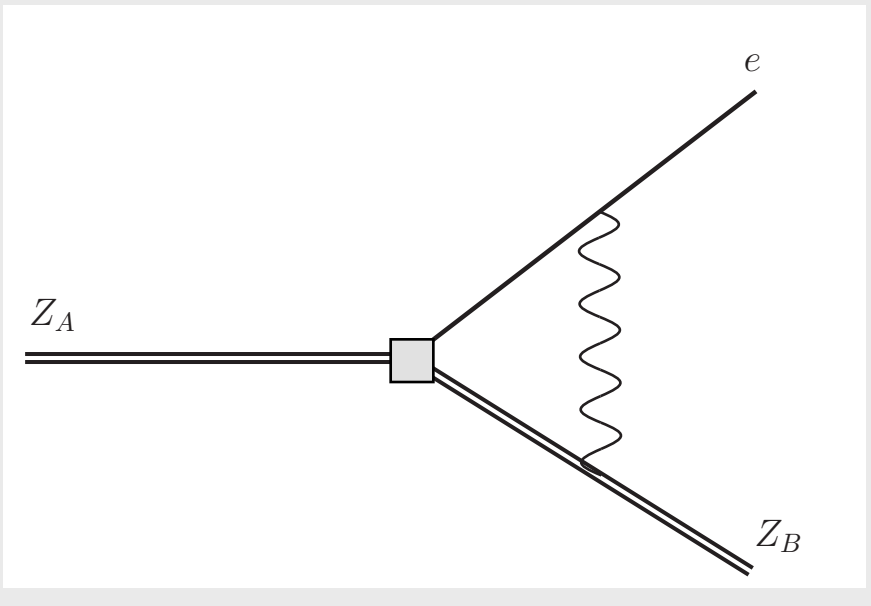
WAVEFUNCTION RENORMALIZATION NOT SHOWN



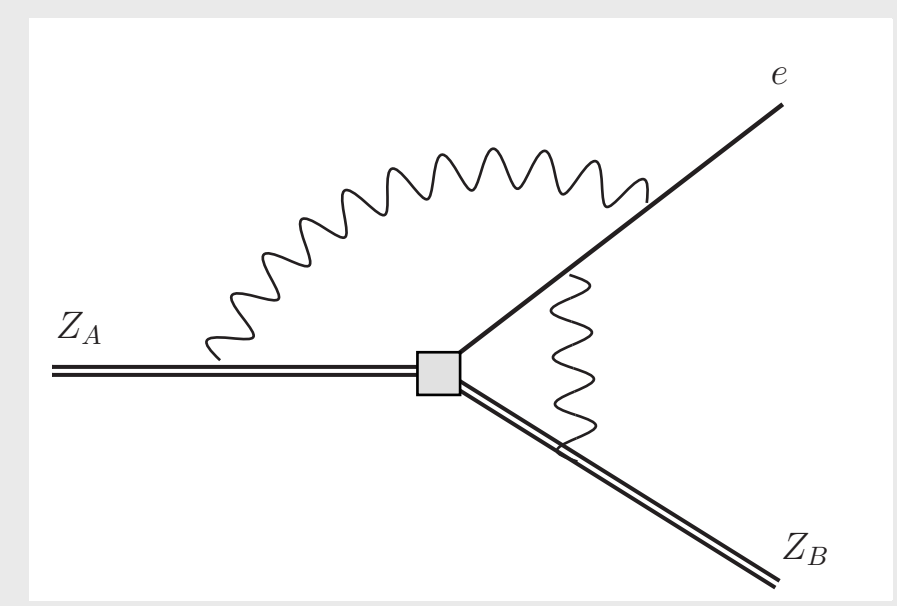
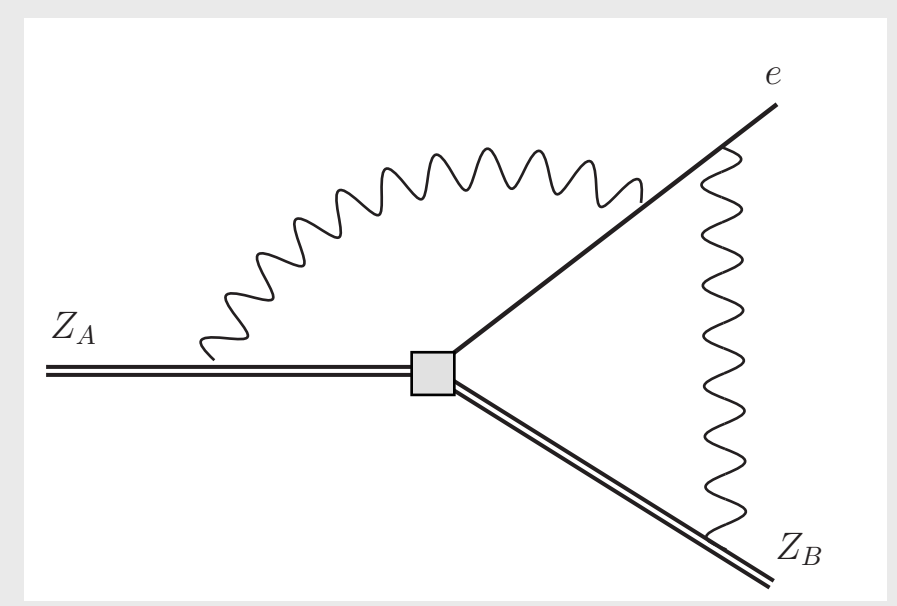
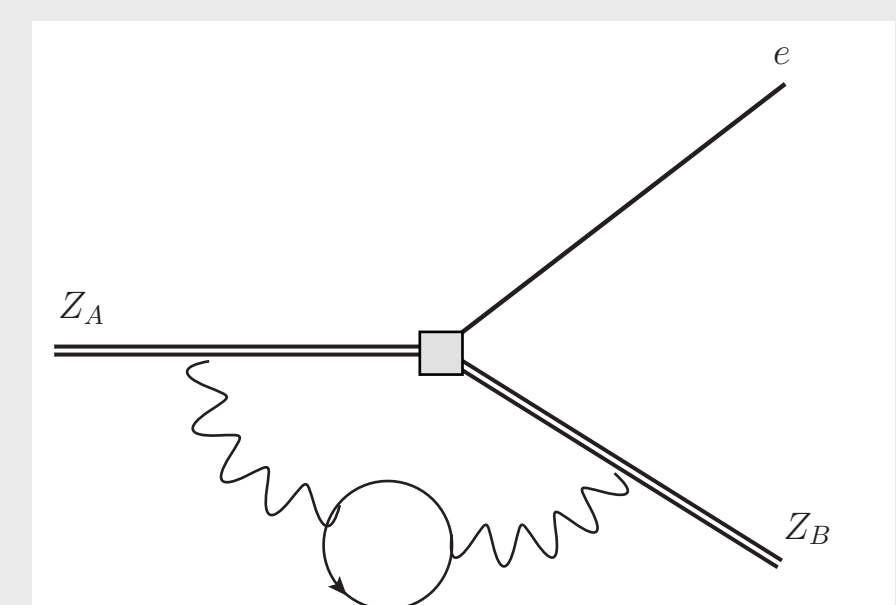
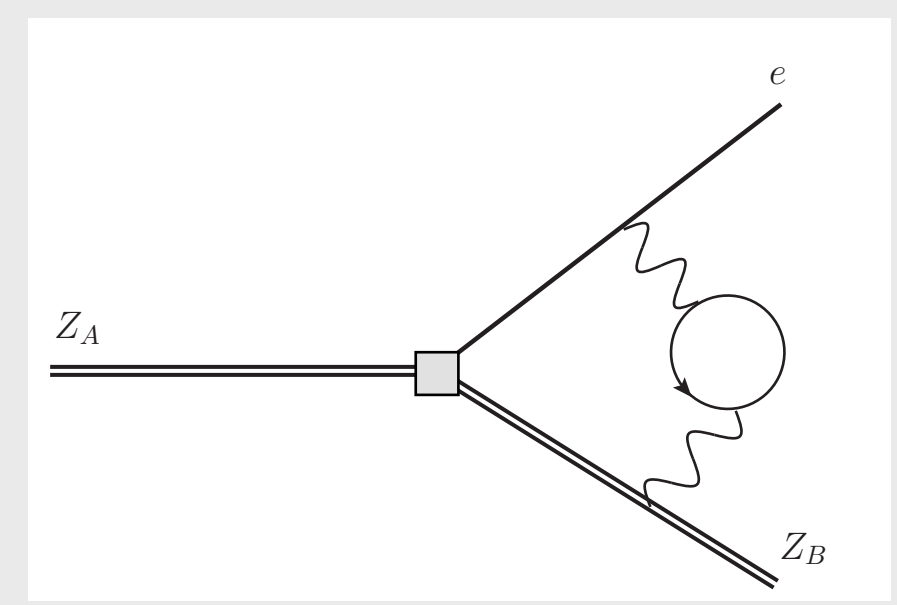
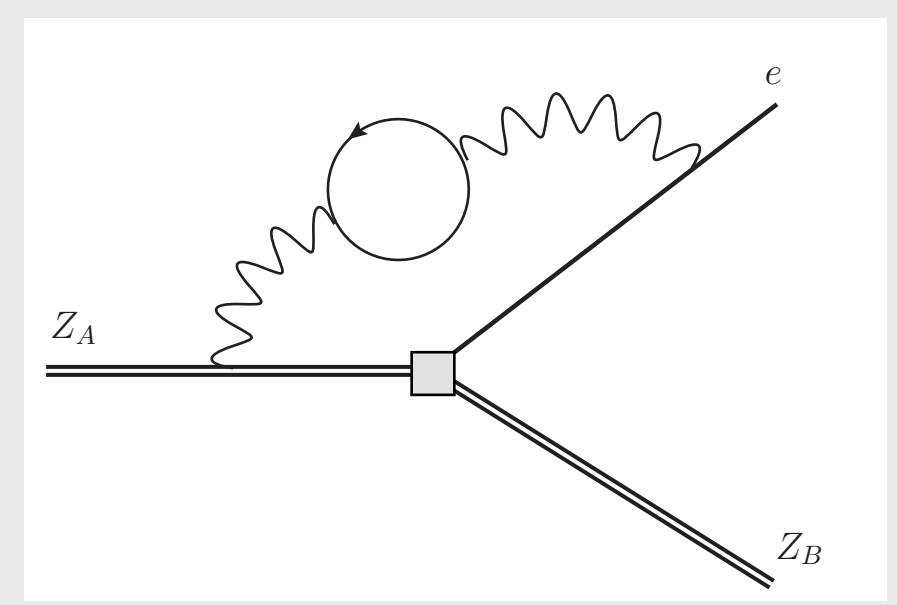
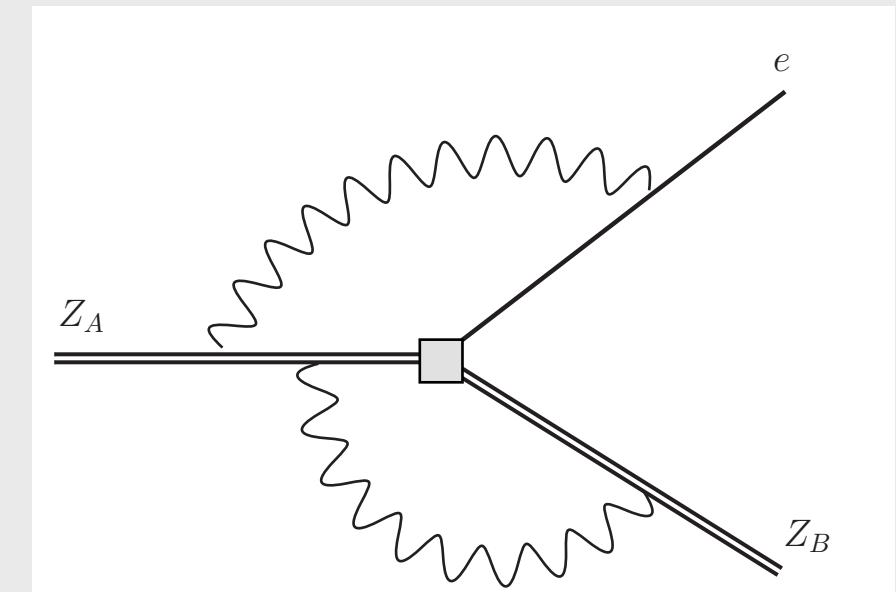
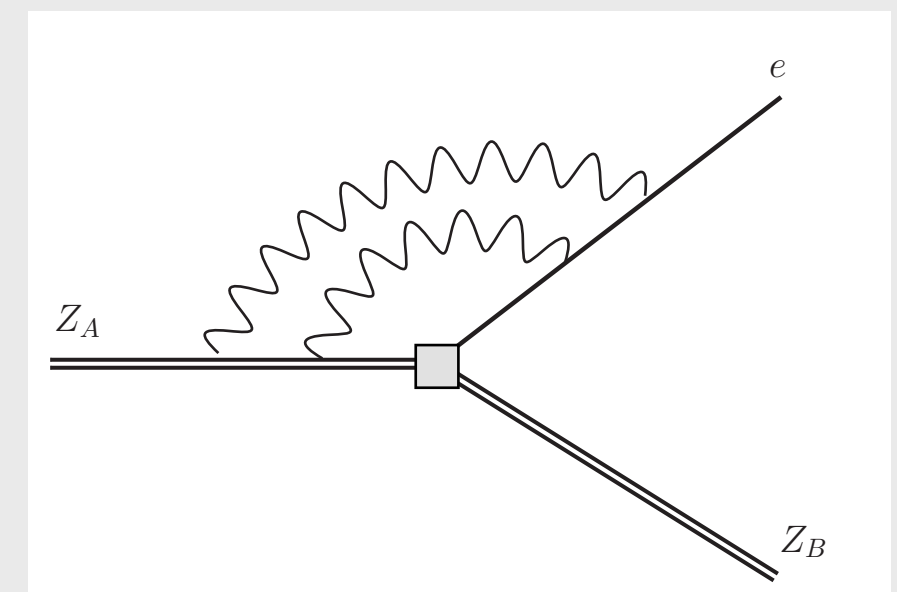
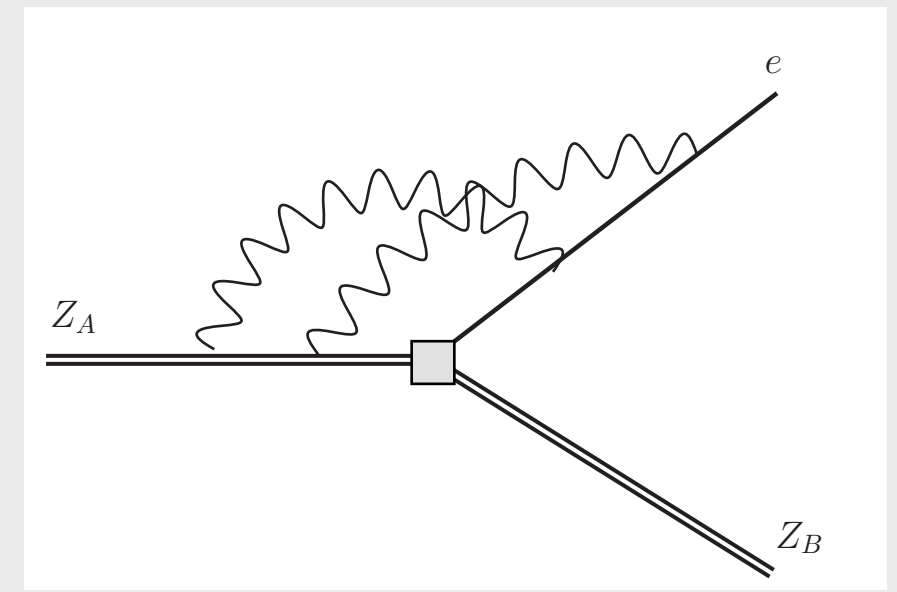
TREE-LEVEL



ONE LOOP



TWO LOOP



Number Of Diagrams Grows Factorially

TREE-LEVEL

- 1 diagram.

ONE LOOP

- 3 diagrams.

TWO LOOP

- 21 diagrams.

THREE LOOP

- 144 diagrams.

- For the Fermi function we need 4+ loops.
- This is not feasible by brute force.

Solution: Make Use Of Simplified Feynman Rules

$$\begin{array}{c} q \\ \rightarrow \\ \hline \hline \end{array} = \frac{i}{q^0 + i0}$$

$$\begin{array}{c} \mu \\ \} \\ \hline \hline \end{array} = i(Z_A e) \delta_0^\mu$$

Number Of Diagrams Grows Factorially

TREE-LEVEL

- 1 diagram.

ONE LOOP

- 3 diagrams.

TWO LOOP

- 21 diagrams.

THREE LOOP

- 144 diagrams.

Reduce Number Of Diagrams
Avoid Difficult Integrals

Solution: Make Use Of Simplified Feynman Rules

$$\begin{array}{c} q \\ \rightarrow \\ \hline \hline \end{array} = \frac{i}{q^0 + i0}$$

$$\begin{array}{c} \mu \\ \} \\ \hline \hline \end{array} = i(Z_A e) \delta_0^\mu$$

Equivalent Feynman Rules

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TREE-LEVEL

- 1 diagram.

ONE LOOP

- 2 diagrams.

TWO LOOP


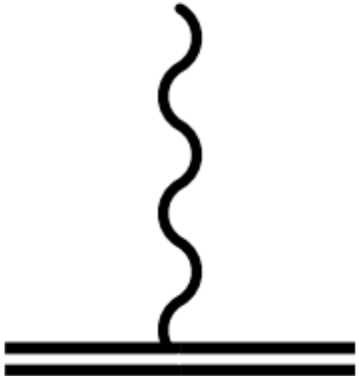
- 5 diagrams.

THREE LOOP

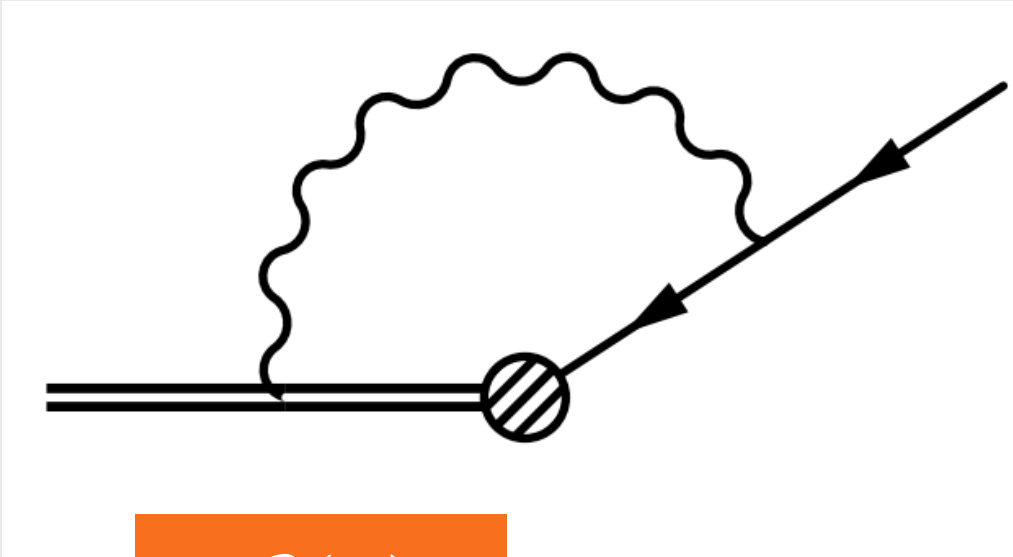
- 10 diagrams.

1 NUCLEUS WITH UNIT CHARGE

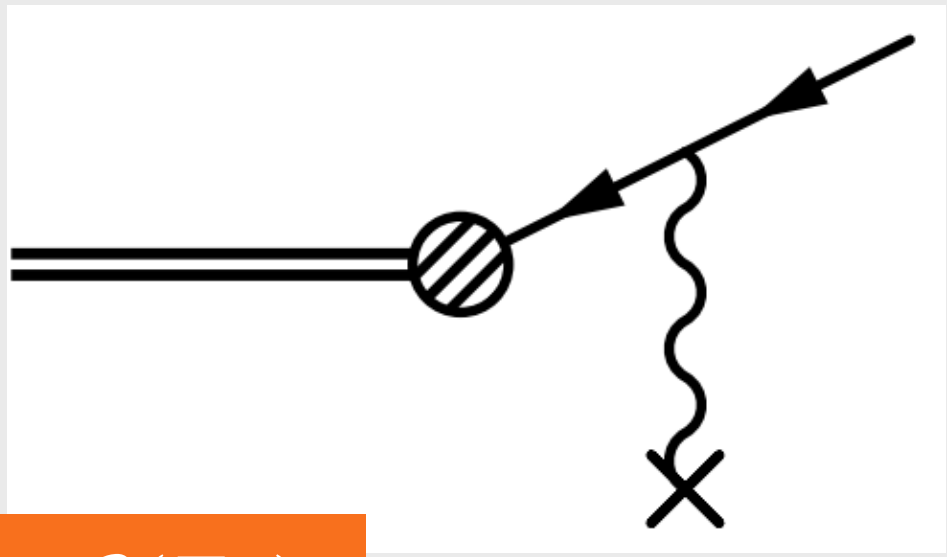
+ A BACKGROUND COULOMB FIELD

μ		μ
	$= iZe \delta_0^\mu 2\pi \delta(q^0)$	
		$= ie \delta_0^\mu$

ONE LOOP



$\mathcal{O}(\alpha)$

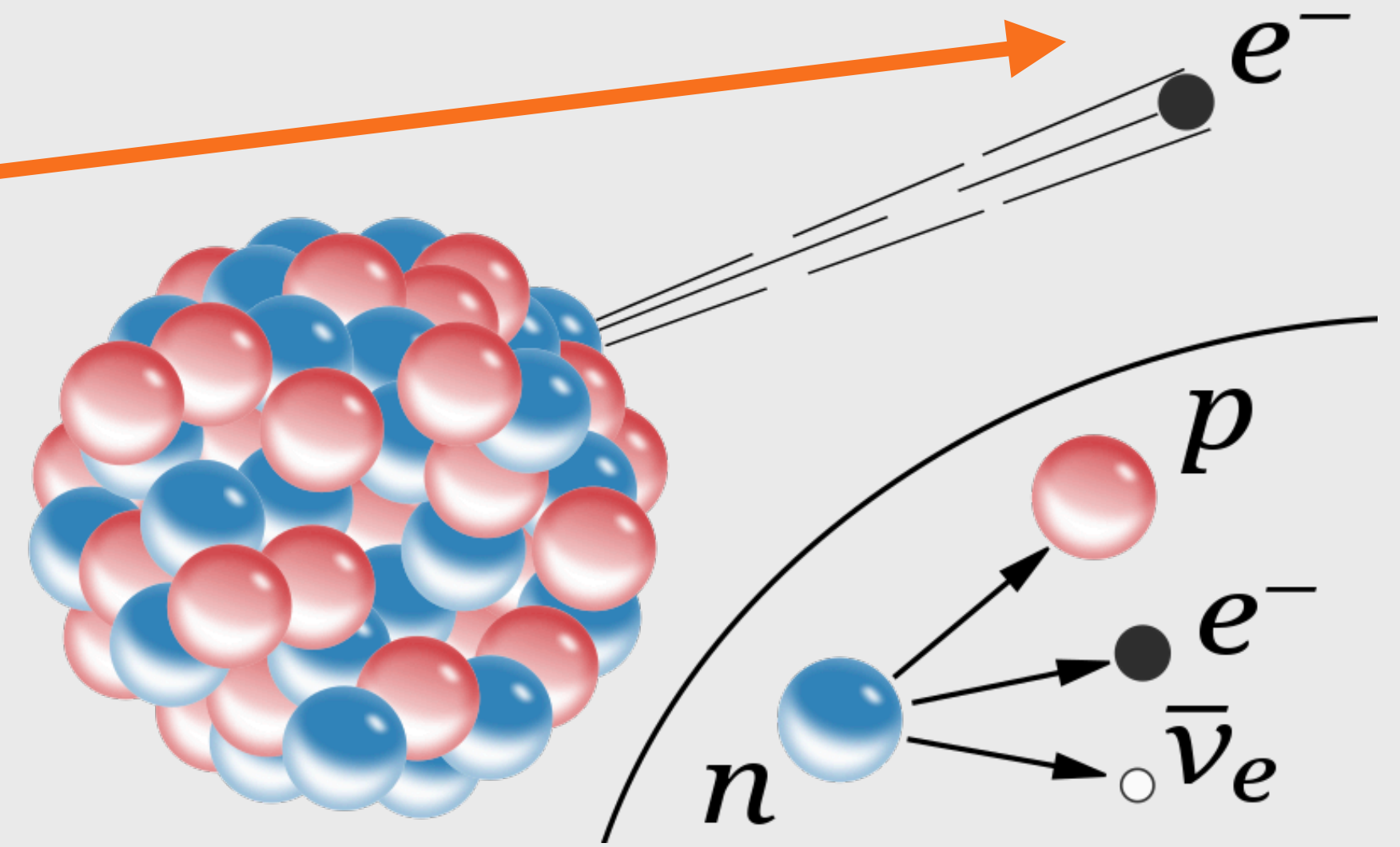


$\mathcal{O}(Z\alpha)$

Fermi Function

ATTRACTED TO NUCLEUS

- Coulomb effects are UV divergent.
- Largest effects are a series in $Z\alpha$
- Historically done with finite-distance regulator



$$\langle e^- | \bar{\psi}(\mathbf{x}) | 0 \rangle \sim \left(\frac{1}{|\mathbf{x}|} \right)^\nu$$

$$\nu = \sqrt{1 - Z^2 \alpha^2} - 1$$

Extraction Of Hard Matrix Element

$$\Psi(\mathbf{x}) = \mathcal{M}_S(\mu_S) \mathcal{M}_H(\mu_S, \mu_H) \tilde{\mathcal{M}}_{\mathbf{x}}(\mu_H, \mathbf{x})$$

KNOWN TO ALL ORDERS IN $Z\alpha$

NEW CALCULATION

EXTRACT WHAT WE WANT

$$\Psi(\mathbf{x})$$

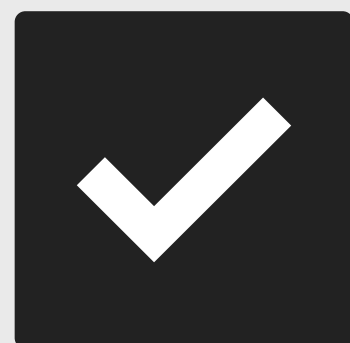
$$\mathcal{M}_H(\mu_S, \mu_H) = \frac{\Psi(\mathbf{x})}{\tilde{\mathcal{M}}_{\mathbf{x}}(\mu_H, \mathbf{x}) \mathcal{M}_S(\mu_S)}$$

Fermi-Function In The $\overline{\text{MS}}$ Scheme

$$\mathcal{M}_H(\mu_S, \mu_H) = e^{\frac{\pi\xi}{2} + i\xi \left(\log \frac{2p}{\mu_S} - \gamma_E \right) - i(\eta-1) \frac{\pi}{2}} \frac{2\Gamma(\eta - i\xi)}{\Gamma(2\eta + 1)}$$

$$\sqrt{\frac{\eta - i\xi}{1 - i\xi \frac{m}{E}}} \sqrt{\frac{E + \eta m}{E + m}} \sqrt{\frac{2\eta}{1 + \eta}} \left(\frac{2pe^{-\gamma_E}}{\mu_H} \right)^{\eta-1} \times \left[\frac{1 + M^*}{2} + \frac{1 - M^*}{2} \gamma^0 \right]$$

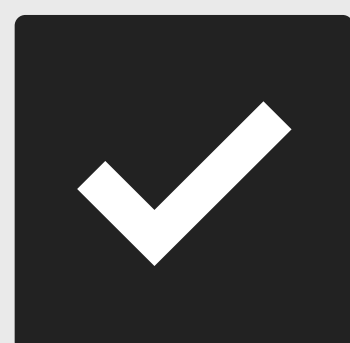
- $\eta = \sqrt{1 - Z^2\alpha^2}$
- $\xi = Z\alpha/\beta$
- $M = (E + m)(1 + i\xi m/E)/(E + \eta m)$



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Factorization Theorem

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- Amplitude depends on Wilson coefficient and matrix element.

$$d\Gamma \propto |C(\mu)|^2 |\mathcal{M}(\mu)|^2 + \mathcal{O}((pR)^2)$$

- Implies that all **short-distances** factorize from **long-distances**.

$$\mathcal{F}t \equiv ft(1 + \delta'_R)(1 + \delta_{\text{NS}} - \delta_C) = \frac{K}{2G_V^2(1 + \Delta_R^V)}$$

EFT Definition Of 'Outer' Corrections

$$\tilde{F}(Z, E) = \left[|\mathcal{M}(\mu)|^2 \right]_{\text{leading}-Z\alpha}$$

$$(1 + \tilde{\delta}_R) \equiv \frac{\langle |\mathcal{M}(\mu)|^2 \rangle}{\langle \tilde{F}(Z, E) \rangle}$$

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$(1 + \delta_{\text{NS}} - \delta_C)(1 + \Delta_R^V)$ ABSORBED IN WILSON COEFFICIENT

EFT Definition Of 'Outer' Corrections

$$(1 + \delta'_R) := \left[\frac{C(\mu_L)/C(\mu_H)}{\exp[(1 - \sqrt{1 - Z^2\alpha^2}) \log(\mu_H/\mu_L)]} \right]^2 \left(\frac{\int d\Pi \langle |\mathcal{M}_H|^2 \rangle}{\int d\Pi F(Z, E) \times \frac{4\eta}{(1+\eta)^2}} \right)_{\mu=\mu_L}$$

$$(1 + \tilde{\delta}_R) \equiv \frac{\langle |\mathcal{M}|^2(\mu) \rangle}{\langle \tilde{F}(Z, E) \rangle}$$

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EFT Definition Of Long-Distance Corrections

Resummation With RG + EFT

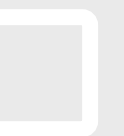
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$$d\Gamma \propto |C(\mu)|^2 |\mathcal{M}(\mu)|^2$$

No Large Logs

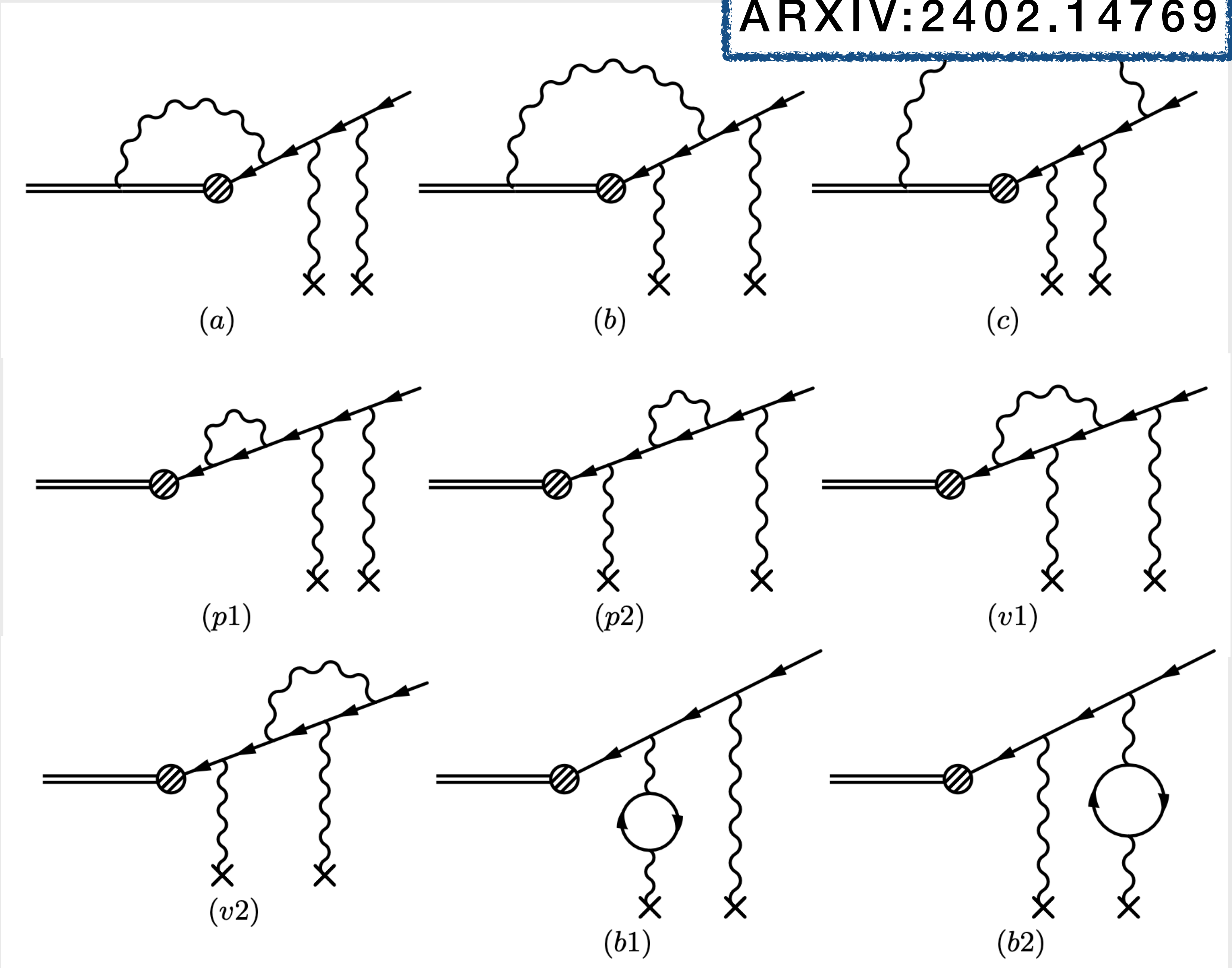
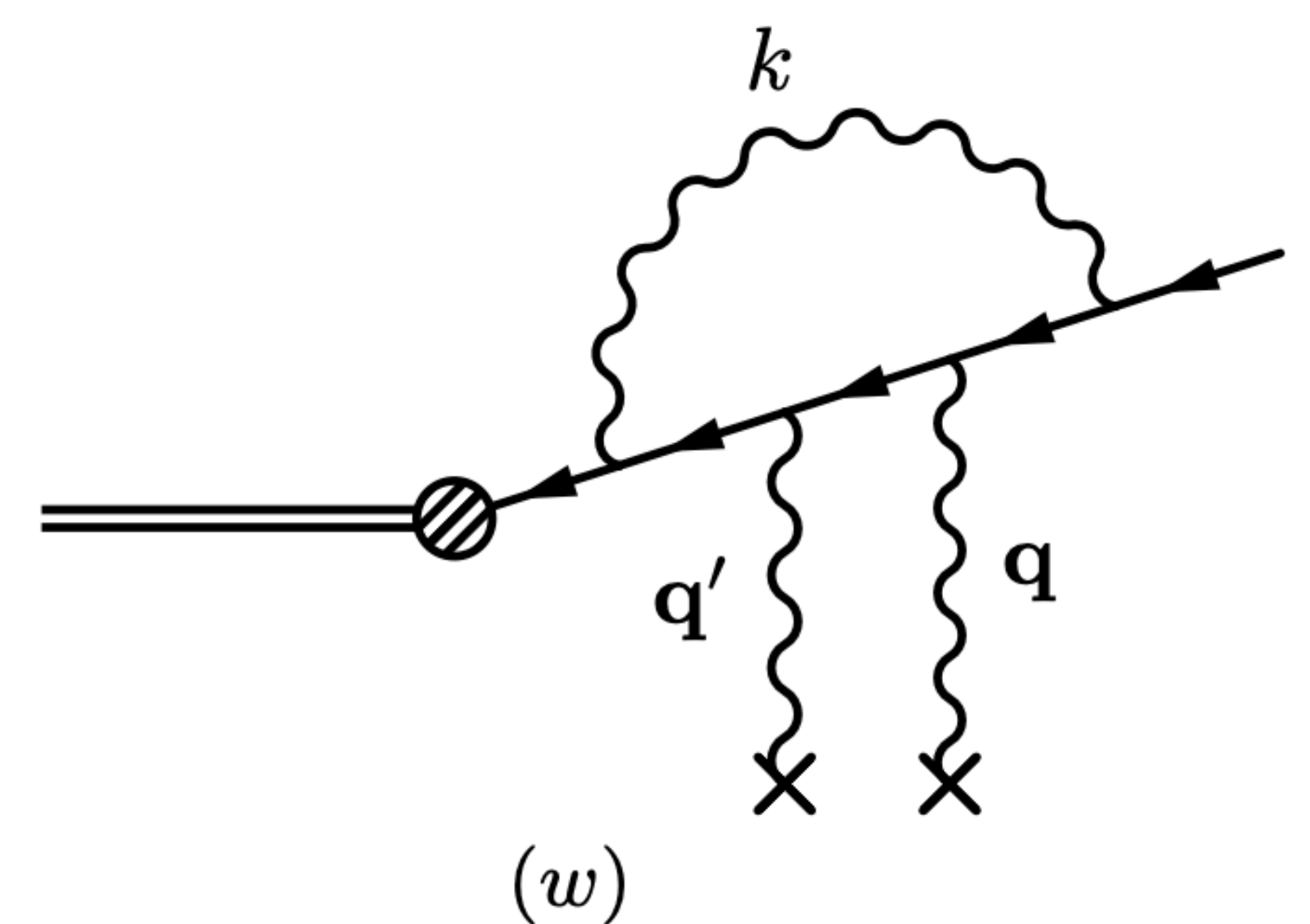
$$= C(\mu_H) \left[\frac{|C(\mu_L)|^2}{|C(\mu_H)|^2} \right] |\mathcal{M}(\mu_L)|^2$$

Calculate With Renormalization Group



USE EIKONAL ALGEBRA TO REDUCE DIAGRAMS

$$\gamma_2^{(1)} = 16\pi^2 \left(6 - \frac{\pi^2}{3} \right)$$



MIXED EUCLIDEAN + LORENTZIAN INTEGRALS

New Result For Anomalous Dimension

Z^n \ Loops	1-loop	2-loop	3-loop	4-loop
Z^0	$\gamma_0^{(1)} = -3$	$\gamma_1^{(2)} = -16\zeta_2 + \frac{5}{2} + \frac{10}{3}n_e$	$\gamma_2^{(3)} = \text{GROZIN 2003}$	$\gamma_3^{(4)} = \text{GROZIN 2023}$
Z^1	$\gamma_0^{(0)} = 0$	$\gamma_1^{(1)} = \gamma_2^{(2)}$	$\gamma_2^{(2)} = \gamma_2^{(1)}$	$\gamma_3^{(3)} = \gamma_3^{(2)} - \gamma_3^{(0)}$
Z^2	—	$\gamma_1^{(0)} = -8\pi^2$	$\gamma_2^{(1)} = 16\pi^2 \left(6 - \frac{\pi^2}{3}\right)$	$\gamma_3^{(2)} = ?$
Z^3	—	—	$\gamma_2^{(0)} = 0$	$\gamma_3^{(1)} = 2\gamma_3^{(0)}$
Z^4	RESUMMATION COMPLETE THROUGH 3-LOOPS!			$\gamma_3^{(0)} = -32\pi^4$

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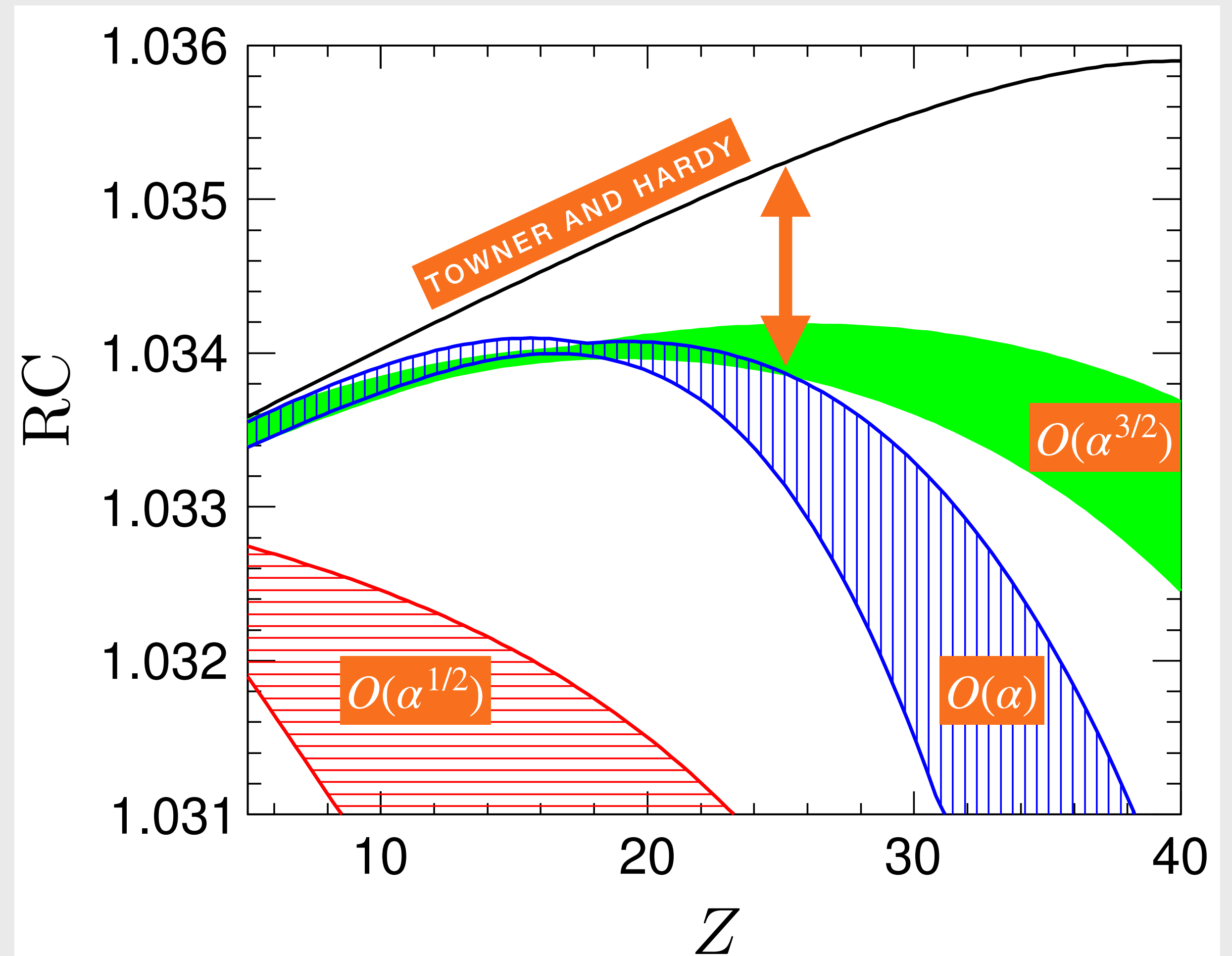
NEW INPUT



Impact For Flavour Physics

SHIFTING δ_3

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COUNTING $Z \sim \log \sim 1/\sqrt{\alpha}$

What's Next?

Order $Z\alpha^2$ Correction

- We will soon present two-loop finite corrections in the long-distance EFT.
- Supplies new model-independent correction relevant for interpreting superallowed beta decays.

Two-loop corrections for superallowed beta decays

PETER VANDER GRIEND^{1,2}, ZEHUA CAO¹, RICHARD J. HILL^{1,2}, AND RYAN PLESTID³

TO APPEAR LATER THIS YEAR

Power Corrections

- Theory so far worked out at leading power.
- Sub-leading power is straightforward, but matrix elements must be computed.
- Structure of factorization formulae will change.
- Many simple things to do here.

Matching Calculations

- Chiral-EFT and/or *ab initio* provides a microphysical framework to extract Wilson coefficient.
- Matching calculation still needs to be performed.

Summary

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- Long-distance QED corrections to beta decay are enhanced by nuclear charge.
- Compute in a point-like EFT with short-distance corrections absorbed into low-energy constants.
- First updates to outer radiative corrections in 40 years.
- Impacts CKM unitarity.