

Unique forbidden β -decays at zero momentum transfer

Chien-Yeah Seng, Ayala Glick-Magid, Vincenzo Cirigliano

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**CIPANP
2025**



U.S. DEPARTMENT OF
ENERGY

Office of
Science

Nuclear Theory for New Physics Topical Collaboration

 Introduction: β -decays for exotic weak interactions

 **Forbidden** β -decays (tree level)

 Experimental Surge

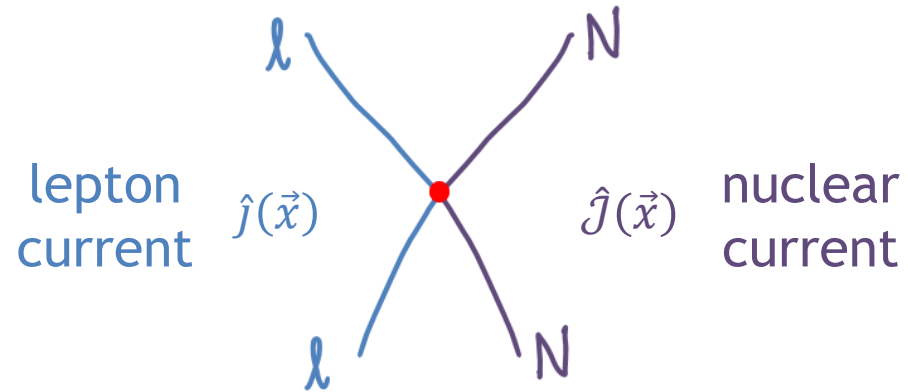
 Radiative Effects (loop level)

 $\Rightarrow$ ~~Forbidden~~ (and other opportunities)

 Summary & Outlook

Weak interaction

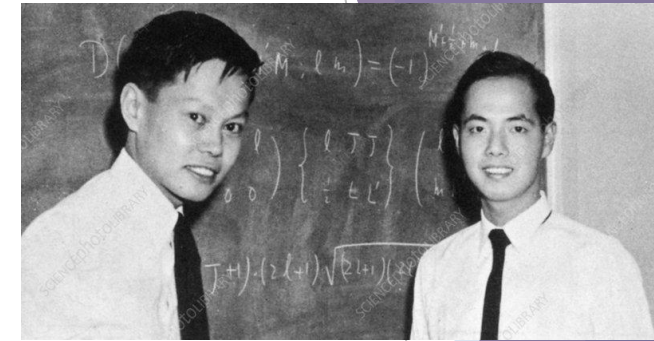
Low energy reaction of
leptons with nucleons



$$\hat{\mathcal{H}}_W \sim \mathbf{C} \hat{j}(\vec{x}) \cdot \hat{J}(\vec{x})$$

A-priori: {

- Scalar (C_S)
- PseudoScalar (C_P)
- Vector (C_V)**
- Axial vector (C_A)**
- Tensor (C_T)



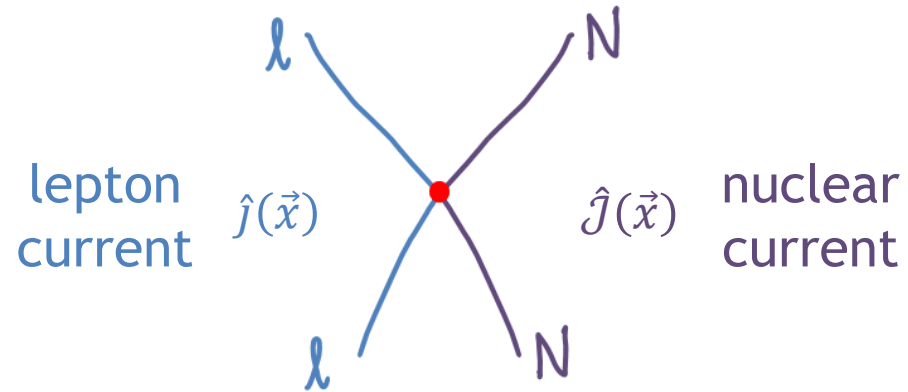
Theory: C.N. Yang and T.D. Lee (Nobel 1957)



Experiment: C.S. Wu:
Parity violation in *nuclear β -decays*
 $\Rightarrow$ Weak SM structure: “**V** – **A**”

Weak interaction

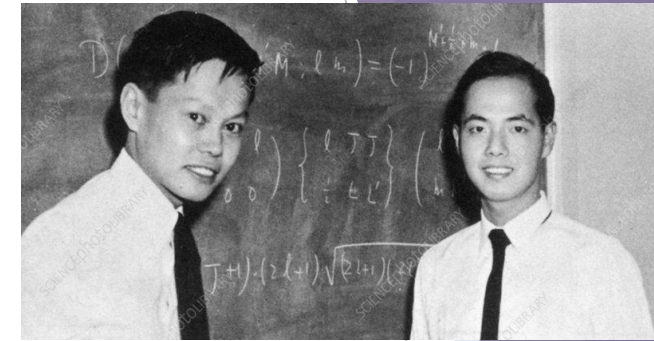
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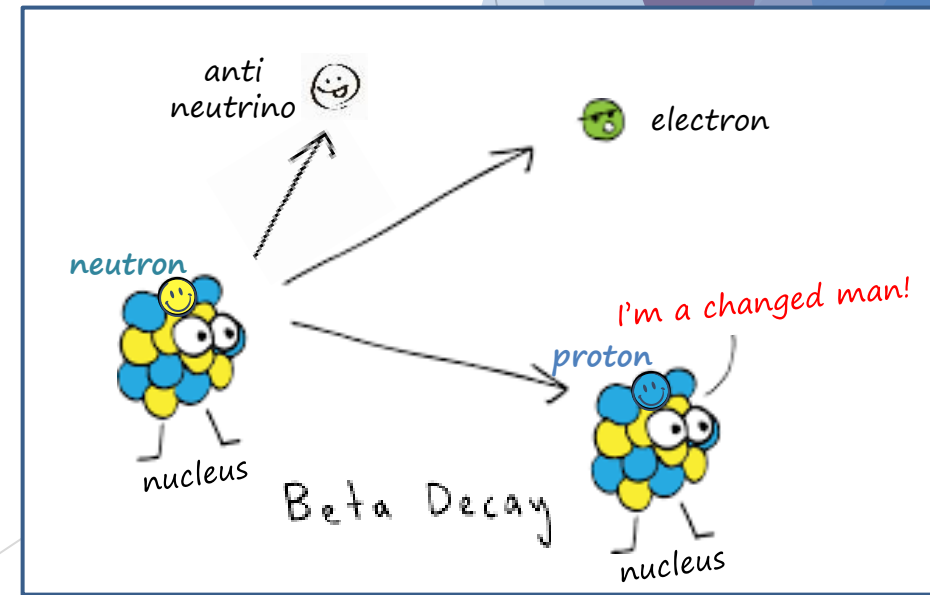
The SM is incomplete

$\gg$ Ongoing searches for C_S, C_P, C_T
in precision *nuclear β -decay* experiments

Nuclear β -decay

Low momentum transfer: $q \sim 0 - 10 \text{ MeV}/c$

Beta decay, Khan Academy, cdn.kastatic.org/ka-perseus-images/8d978444f15f9bbc3bcadb0549816bc7c264b977.svg



Nuclear β -decay

Low momentum transfer: $q \sim 0 - 10 \text{ MeV}/c$

angular
momentum parity

Transitions $J^{\Delta\pi}$:

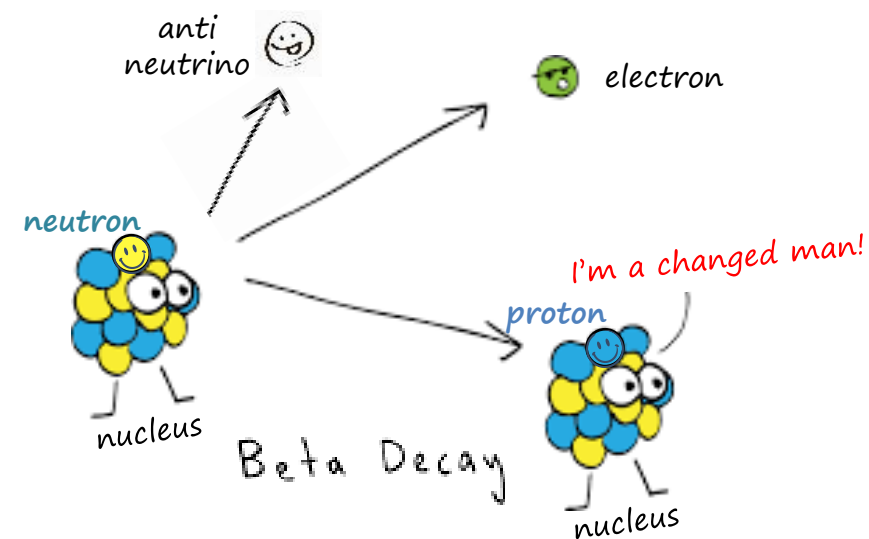
“Allowed”
(when $q \rightarrow 0$)

- Fermi (0^+)
- Gamow-Teller (1^+)

“Forbidden”
(vanish for $q \rightarrow 0$)

- All the rest ($J^{\Delta\pi}$)

Beta decay, Khan Academy, cdn.kastatic.org/ka-perseus-images/8d978444f15f9bbc3bcadb0549816bc7c264b977.svg



Allowed decays

$$d\Gamma \propto |\langle \psi_f | \hat{H}_W | \psi_i \rangle|^2 \propto 1 + \mathbf{b} \frac{m_e}{E_e} + \mathbf{a} \vec{\beta} \cdot \hat{\nu}$$

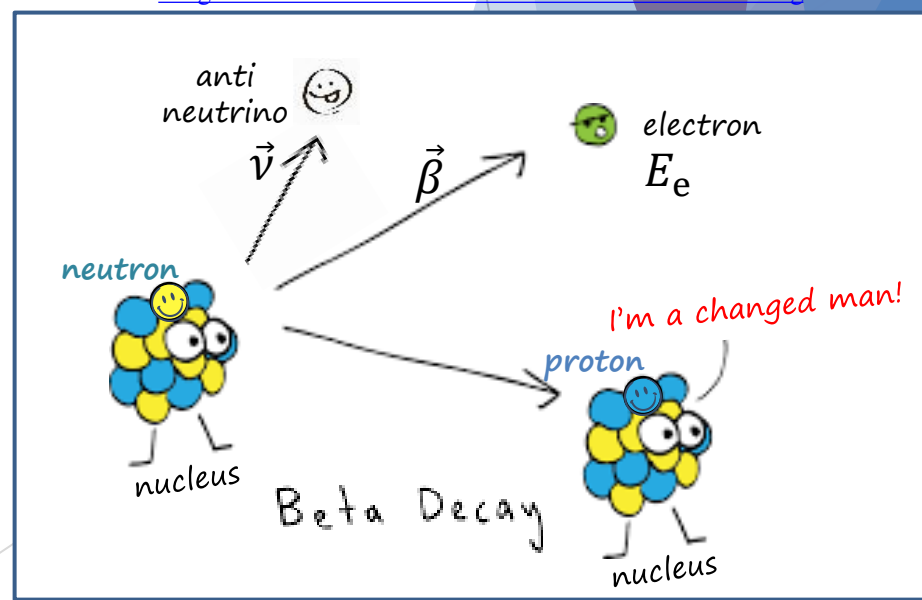
- 1) Energy spectrum: Fierz term $\mathbf{b} = 0 \pm \frac{C_T^+}{C_A}$

► Linear in C_T^+

- 2) Angular correlation: $\mathbf{a} = -\frac{1}{3} \left(1 - \frac{|C_T^+|^2 + |C_T^-|^2}{4|C_A|^2} \right)$

► Sensitive to right-handed neutrinos C_T^-

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Unique forbidden decays

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AGM & Gazit, PRD 2023

- 1) Energy spectrum: Fierz term $\mathbf{b} = 0 \overset{\text{SM}}{\pm} \frac{C_T^+}{C_A}$

► Linear in C_T^+

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- 3) Angular squared correlation: $\mathbf{a}_2 = \overset{\text{SM}}{\frac{1}{5}} \left(1 - \overset{\text{BSM}}{\frac{|C_T^+|^2 + |C_T^-|^2}{4|C_A|^2}} \right)$

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**4:30 PM @ Nuclear Structure for Neutrino & Astro:
Novel Simplification of tensor interactions**

*On the technique behind this result &
applications to β -decay, neutrinos, DM & $\mu \rightarrow e$*

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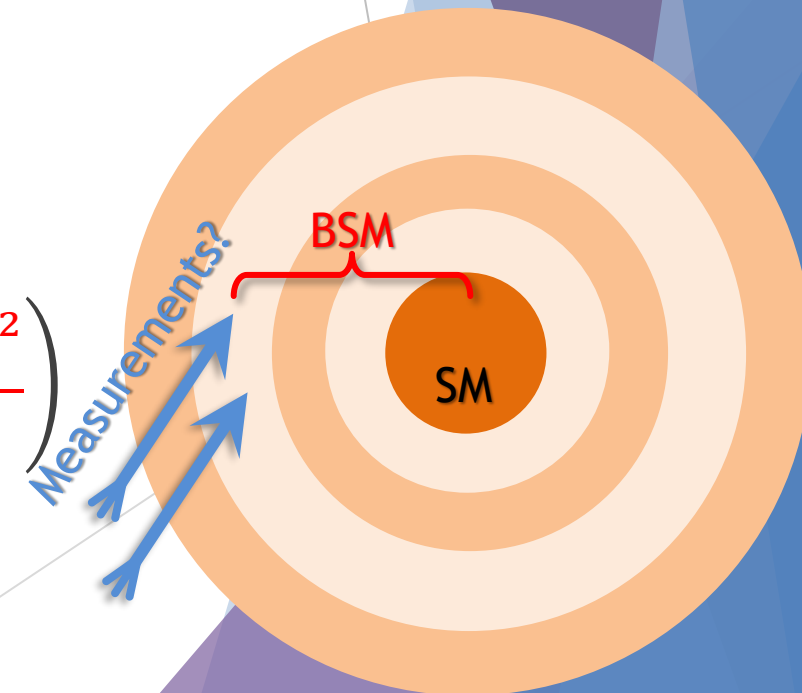
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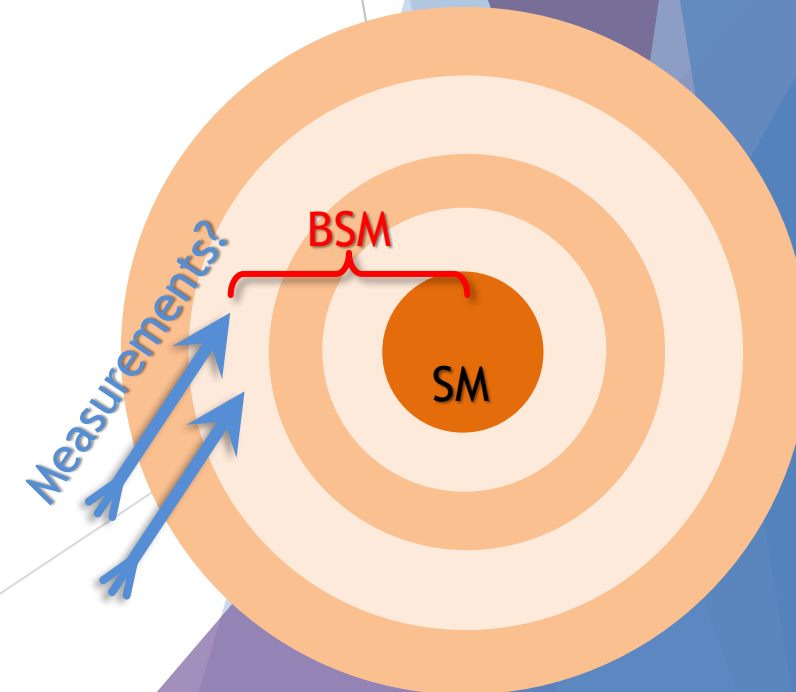
[AGM & Gazit, PRD 2023](#)

> Experiments:

Searches for deviations from
the SM “V-A” structure

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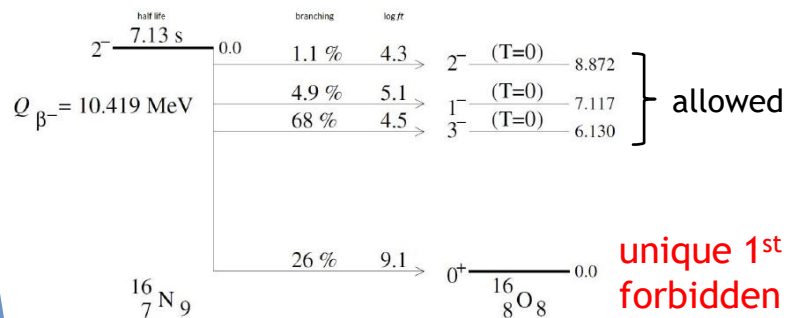
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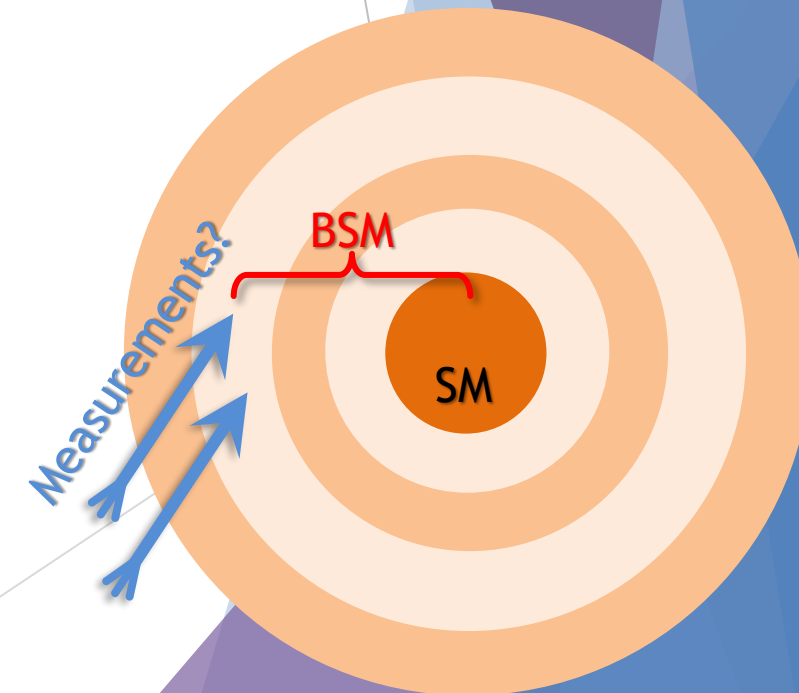
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¹⁶N: Large energy separation between forbidden and allowed branches



[Ohayon, Chocron, Hirsh, AGM, et al., Hyp.Int.2018](#)



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AGM & Gazit, PRD 2023

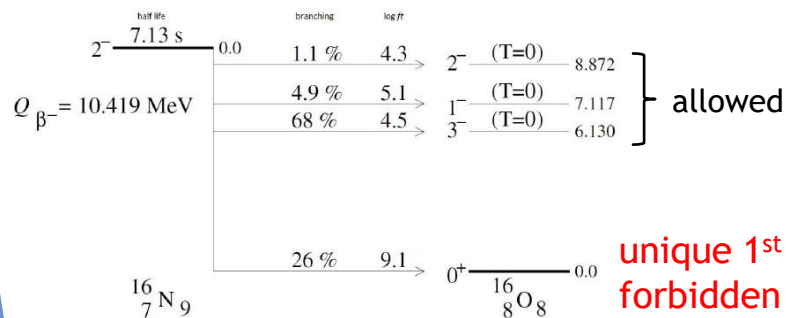
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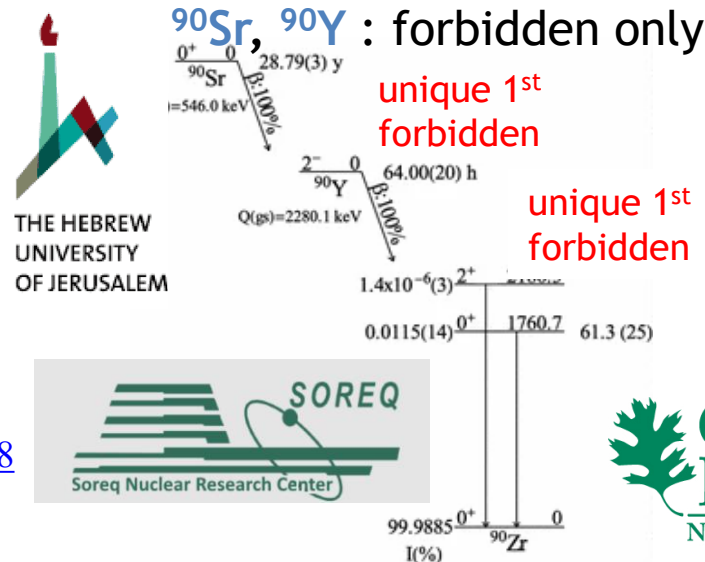
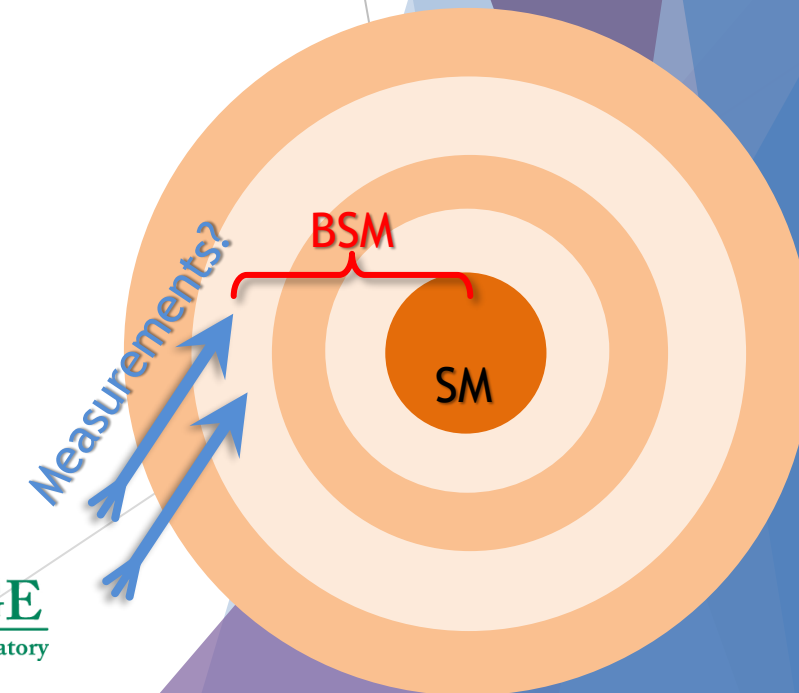


Fig.: Morozov *et al.* J.Rad.Nuc.Chem.2010



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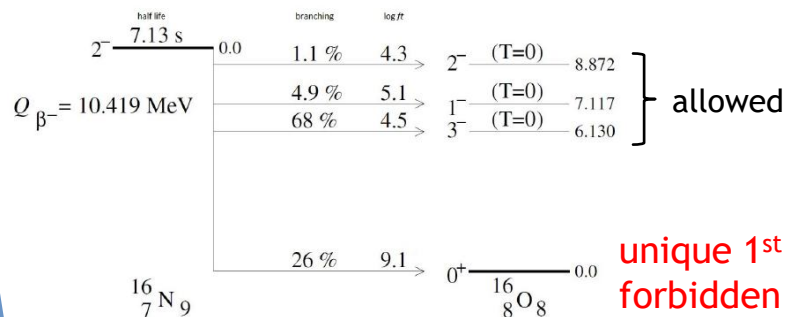
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⇒ Development of **MTAS @ ORNL**
Modular Total Absorption Spectrometer
to *distinguish between decay transitions*

[Shuai, Rasco, et al., PRD 2022](#)

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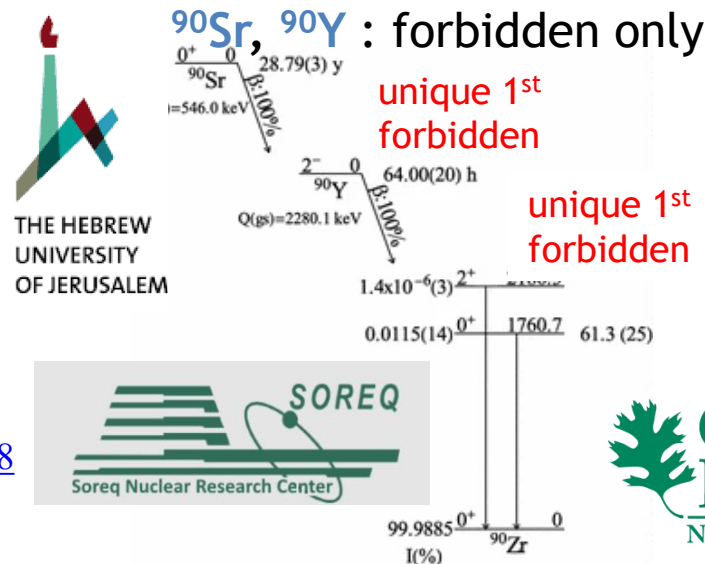


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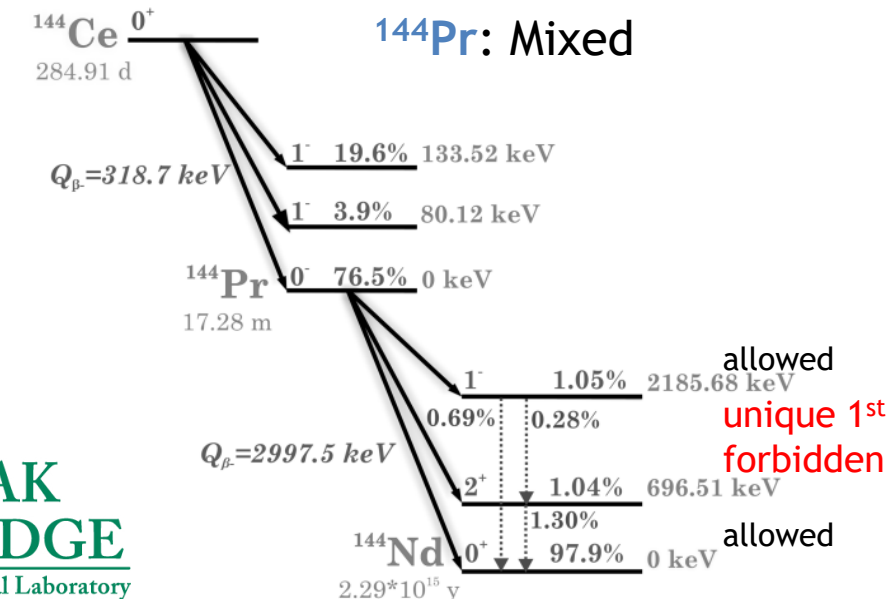


Fig.: [Journal of Physics: Conference Series 1390 \(2019\) 012117](#)

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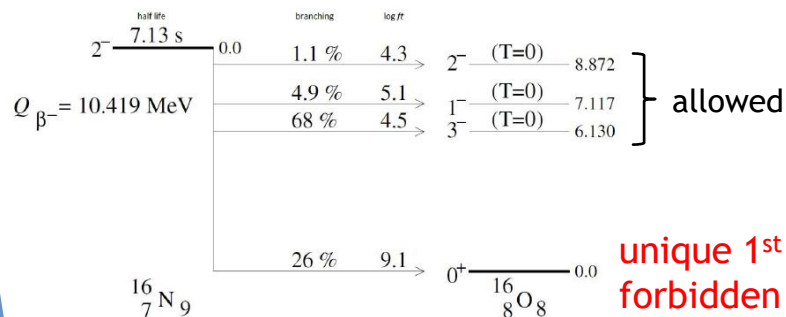
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>> More accurate theory is needed

¹⁶N:

forbidden and allowed branches



[Ohayon, Chocron, Hirsh, AGM, et al., Hyp.Int.2018](#)



⁹⁰Sr
28.79(3) y
 $Q = 546.0$ keV
B(100%)

unique 1st
forbidden

⁹⁰Y
64.00(20) h
 $Q(\text{gs}) = 2280.1$ keV
B(100%)

unique 1st
forbidden

$1.4 \times 10^{-6} (3) 2^+$
 $0.0115(14) 0^+$
1760.7 61.3 (25)

99.9885 0⁺
0 90Zr
I(%)

Fig.: [Morozov et al. J.Rad.Nuc.Chem.2010](#)

¹⁴⁴Ce 0⁺
284.91 d
 $Q_{\beta^-} = 318.7$ keV

14

Measurements

~~BSM~~

SM??

Naïve SM

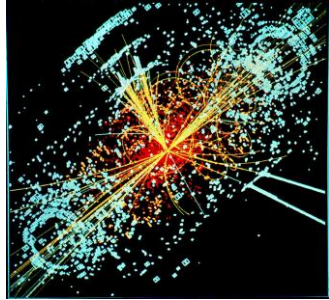


Fig.: [Journal of Physics: Communications Series 1390 \(2019\) 012117](#)

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High energy frontier



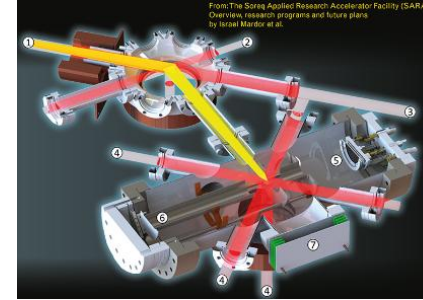
Lucas Taylor / CERN - <http://cdsweb.cern.ch/record/628469>
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LHC
TeV scale

Searches for BSM physics

C_T^+ (C_T^-) Tensor left (right)
coupling constants

Precision frontier

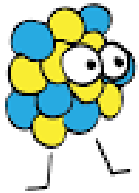


Mardor et al., [Eur. Phys. J. A 54, 91](#) (2018)

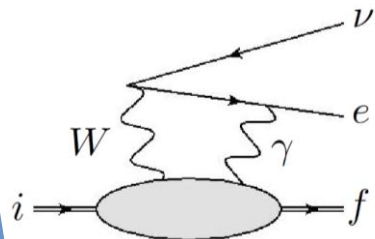
Nuclear phenomena
 10^{-3} precision level

new experiments

Requires Precision Theory – two challenges:



➤ *Nuclear-structure many-body problem*



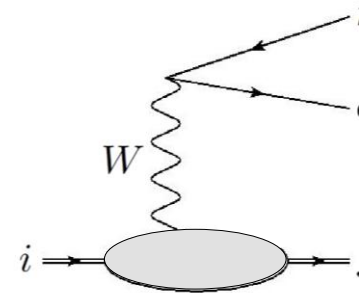
➤ *High-order radiative effects*

➡ Now

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On the technique behind forbidden's rate & applications to β -decay, neutrinos, DM & $\mu \rightarrow e$

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$$d\Gamma \propto |\langle 0^+ || \hat{H}_W || 2^- \rangle|^2 \propto 1 + \mathbf{b} \frac{m_e}{E_e} + \mathbf{a} \vec{\beta} \cdot \hat{\mathbf{v}} + \mathbf{a}_2 [\beta^2 - (\hat{\beta} \cdot \hat{\mathbf{v}})^2]$$



“Forbidden”
(vanish for $q \rightarrow 0$)

- All the rest ($J^{\Delta\pi}$)

angular
momentum

parity

$$J_i^{\pi_i} \rightarrow J_f^{\pi_f}$$

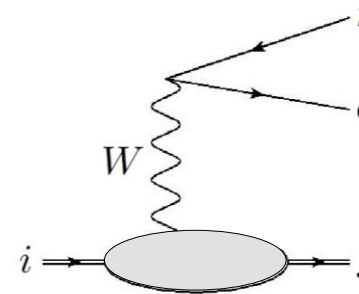
$$|J_i - J_f| \leq J \leq J_i + J_f$$

$$\Delta\pi = \pi_i \cdot \pi_f$$

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$$d\Gamma \propto |\langle 0^+ || \hat{H}_W || 2^- \rangle|^2 \propto \left\{ 1 + \mathbf{b} \frac{m_e}{E_e} + \mathbf{a} \vec{\beta} \cdot \hat{\nu} + \mathbf{a}_2 [\beta^2 - (\vec{\beta} \cdot \hat{\nu})^2] \right\} |\langle 0^+ || \hat{\mathcal{O}}_J || 2^- \rangle|^2$$

$$\propto q^{J-1}$$



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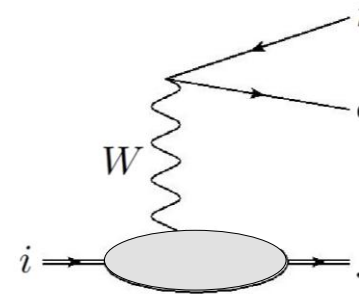
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$$\propto q^{J-1} \xrightarrow{q \rightarrow 0} 0$$



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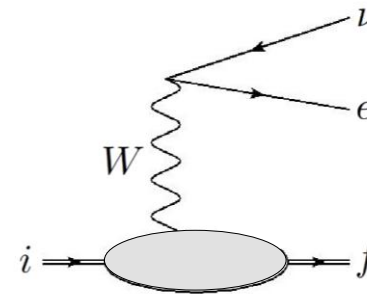
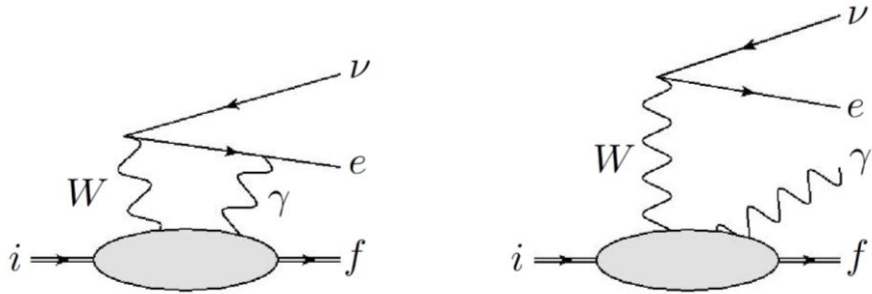
$$|2 - 0| \leq J \leq 2 + 0$$

$$\Delta\pi = -$$

Unique 1st-forbidden

$$d\Gamma \propto |\langle 0^+ || \hat{H}_W || 2^- \rangle|^2 \propto \left\{ 1 + \mathbf{b} \frac{m_e}{E_e} + \mathbf{a} \vec{\beta} \cdot \hat{\nu} + \mathbf{a}_2 [\beta^2 - (\vec{\beta} \cdot \hat{\nu})^2] \right\} |\langle 0^+ || \hat{\mathcal{O}}_J || 2^- \rangle|^2$$

Radiative Corrections:



$$\propto q^{J-1} \xrightarrow{q \rightarrow 0} 0$$

Forbidden!

$$\langle 0^+ || \hat{\mathcal{O}}_J^W \hat{Q}_{J'}^{EM} || 2^- \rangle$$

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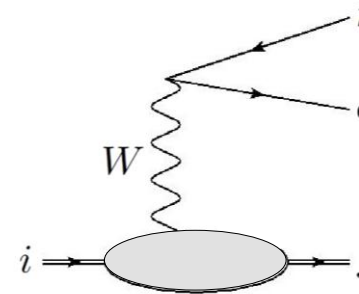
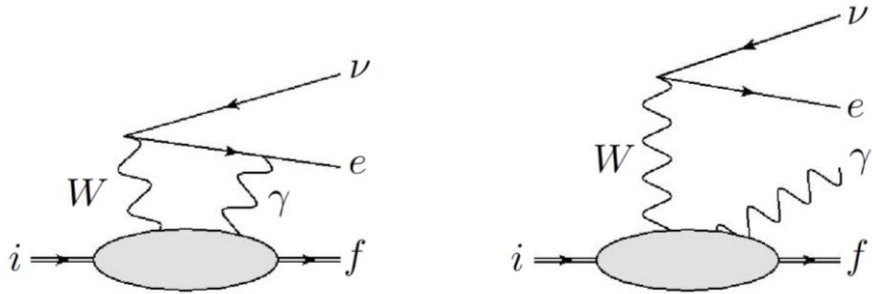
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Forbidden!

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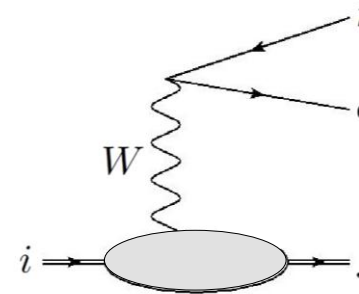
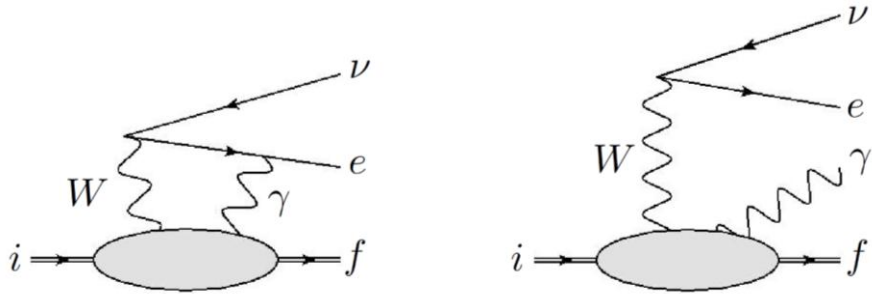
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No longer limited to $J = 2$ operators to satisfy the selection rule!

$$\begin{aligned} |J_X - 0| &\leq J \leq J_X + 0 \\ |2 - J_X| &\leq J' \leq 2 + J_X \end{aligned}$$

angular momentum

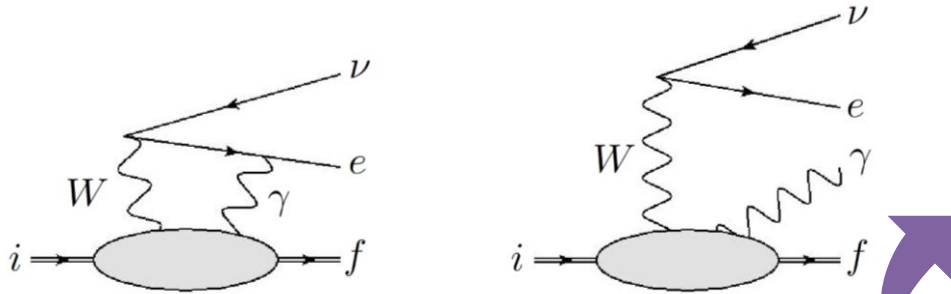
parity

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$$d\Gamma \propto |\langle 0^+ || \hat{H}_W || 2^- \rangle|^2 \propto \left\{ 1 + \mathbf{b} \frac{m_e}{E_e} + \mathbf{a} \vec{\beta} \right\}.$$

Radiative Corrections:



$$\langle 0^+ || \hat{O}_J^W \hat{Q}_{J'}^{EM} || 2^- \rangle = \sum_X \langle 0^+ || \hat{O}_J^W || \psi_X \rangle \langle \psi_X || \hat{Q}_{J'}^{EM} || 2^- \rangle$$

➤ **Choose wisely the intermediate states and operators:**

No longer limited to $J = 2$ operators to satisfy the selection rule!

$$\begin{aligned} |J_X - 0| &\leq J \leq J_X + 0 \\ |2 - J_X| &\leq J' \leq 2 + J_X \end{aligned}$$

	$\Delta\pi$	transition	multipoles
$J = 0$	+	Fermi	$\hat{C}_0^V \sim 1$
	-	1 st forbidden	$\hat{L}_0^A \sim qR$ $\hat{C}_0^A \sim p_f/m_N$
$J > 0$	$(-)^J$	J^{th} forbidden	$\hat{C}_J^V \sim (qR)^J$ $\hat{M}_J^A \sim (qR)^J$
	$(-)^{J-1}$	Gamow Teller ($J = 1$) unique $(J - 1)^{\text{th}}$ forbidden	$\hat{L}_J^A \sim (qR)^{J-1}$ $\hat{C}_J^A \sim p_f/m_N (qR)^J$ $\hat{M}_J^V \sim p_f/m_N (qR)^J$

[AGM & Gazit, J.Phys.G 2022](#)

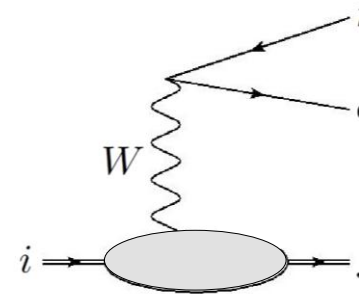
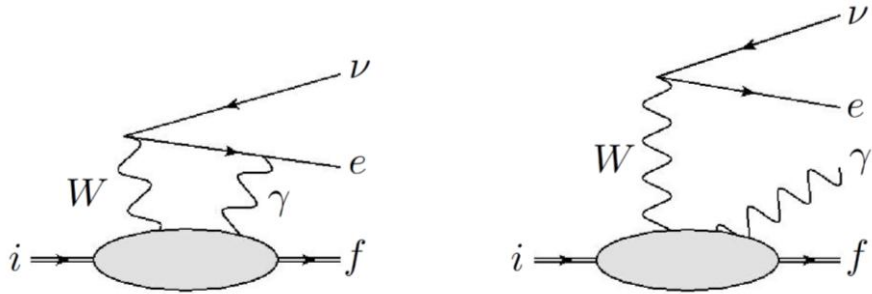
Decay process	Intermediate states	Multipole operators
$i(2^-) \rightarrow f(0^+)$	$i(2_g^-), f(0_g^+)$	$C_0^{\text{em}}, L_2^A, E_2^A$
	$i(1_X^+), f(1_X^-)$	$C_1^{\text{em}}, L_1^{\text{em}}, E_1^{\text{em}}, L_1^A, E_1^A$
$i(0^+) \rightarrow f(2^-)$	$i(0_g^+), f(2_g^-)$	$C_0^{\text{em}}, L_2^A, E_2^A$
	$i(1_X^-), f(1_X^+)$	$C_1^{\text{em}}, L_1^{\text{em}}, E_1^{\text{em}}, L_1^A, E_1^A$

Unique 1st-forbidden

$$d\Gamma \propto |\langle 0^+ \| \hat{H}_W \| 2^- \rangle|^2 \propto \left\{ 1 + \mathbf{b} \frac{m_e}{E_e} + \mathbf{a} \vec{\beta} \cdot \hat{\mathbf{v}} + \mathbf{a}_2 [\beta^2 - (\vec{\beta} \cdot \hat{\mathbf{v}})^2] \right\} |\langle 0^+ \| \hat{\mathcal{O}}_J \| 2^- \rangle|^2$$

$$\propto q^{J-1} \xrightarrow{q \rightarrow 0} 0$$

Radiative Corrections:



Forbidden!

$$\langle 0^+ \| \hat{\mathcal{O}}_J^W \hat{Q}_{J'}^{EM} \| 2^- \rangle = \sum_X \langle 0^+ \| \hat{\mathcal{O}}_J^W \| \psi_X \rangle \langle \psi_X \| \hat{Q}_{J'}^{EM} \| 2^- \rangle$$

No longer limited to $J = 2$ operators to satisfy the selection rule!

$$d\Gamma \propto \underset{\text{interference}}{f_0} q^0 + f_1 q^1 + f_2 q^2$$

angular momentum

parity

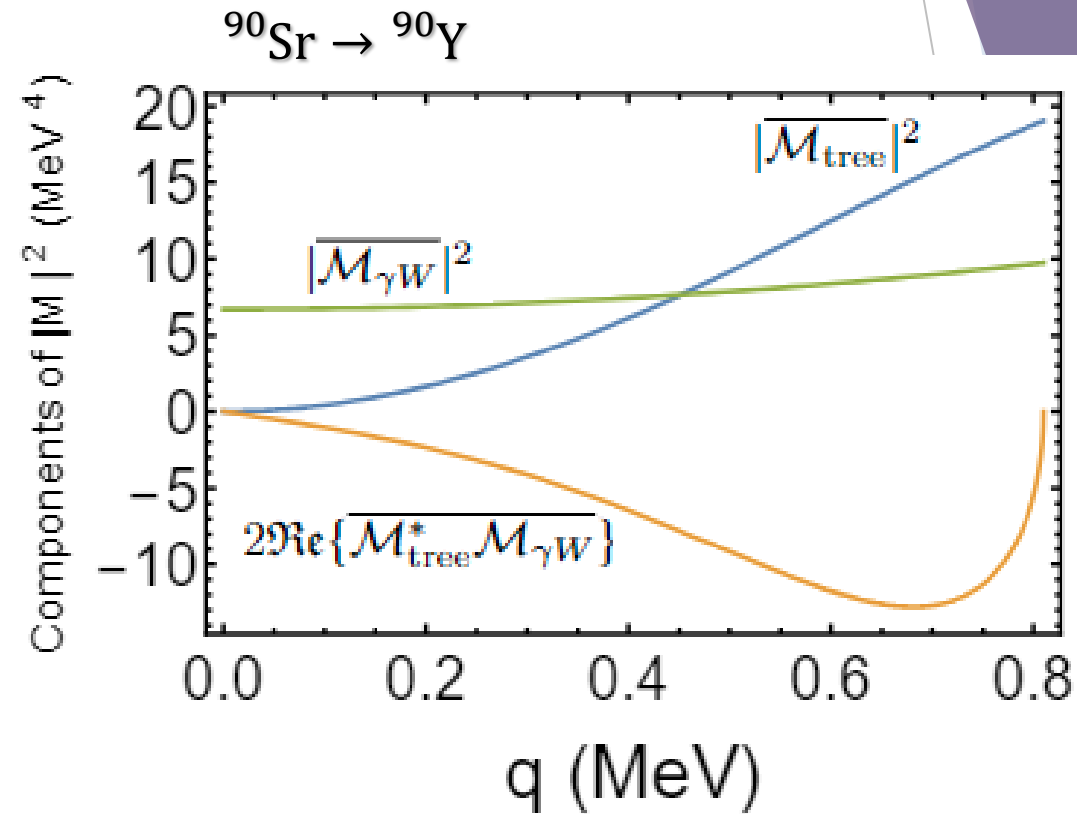
$$J_i^{\pi_i} \rightarrow J_f^{\pi_f}$$

$$|J_i - J_f| \leq J \leq J_i + J_f$$

$$\Delta\pi = \pi_i \cdot \pi_f$$

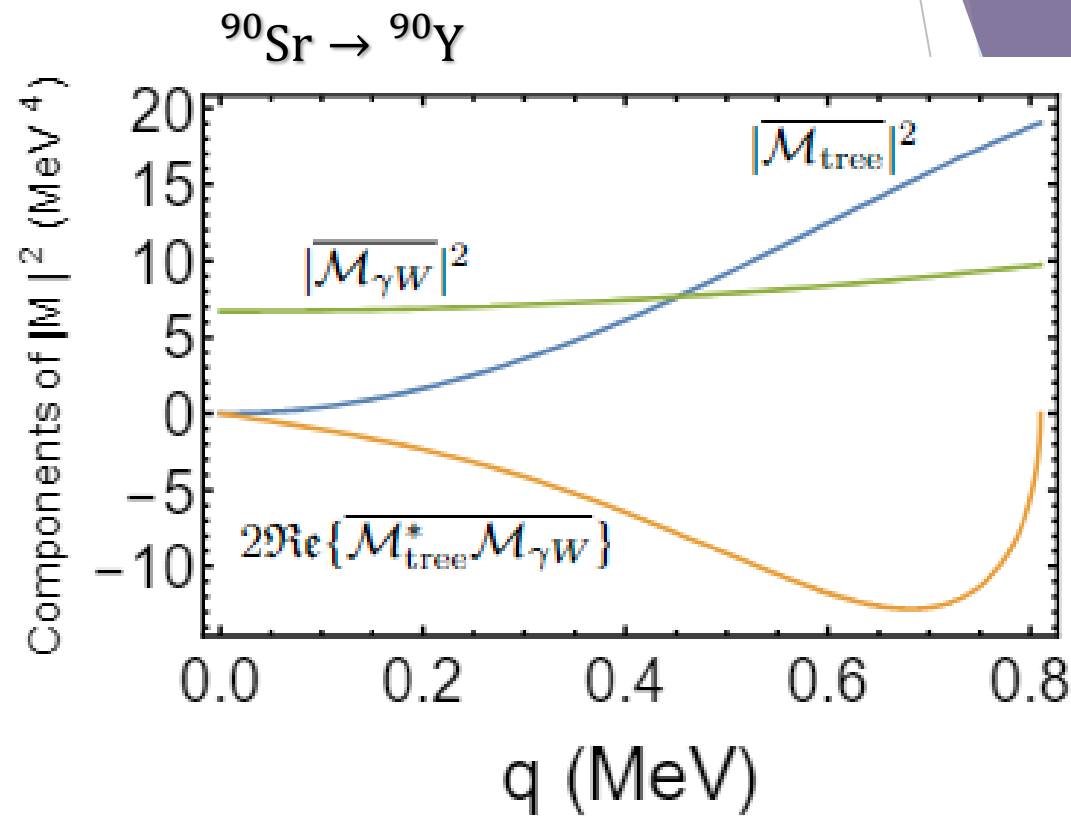
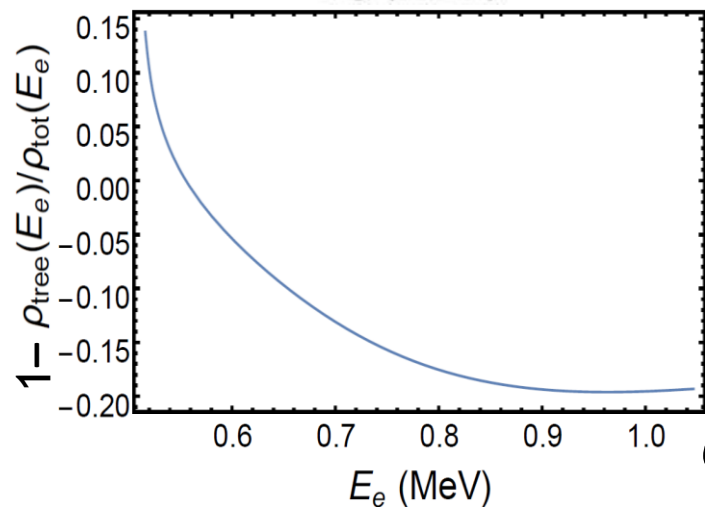
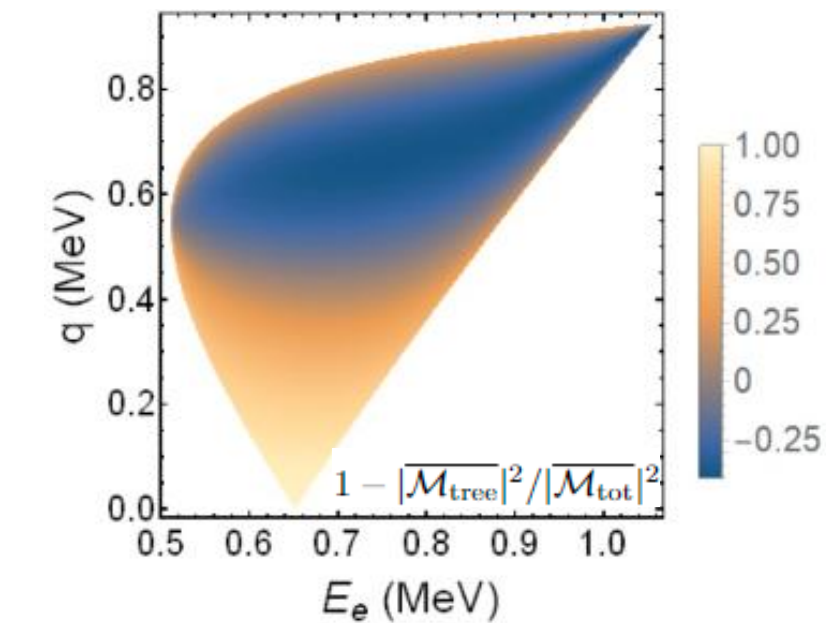
Unique Forbidden decays

No decay @ $q \rightarrow 0$



$$d\Gamma \propto \frac{|\mathcal{M}_{\gamma W}|^2}{f_0} q^0 + \frac{2\Re\{\mathcal{M}_{\text{tree}}^* \mathcal{M}_{\gamma W}\}}{f_1} q^1 + \frac{|\mathcal{M}_{\text{tree}}|^2}{f_2} q^2$$

Unique Forbidden decays

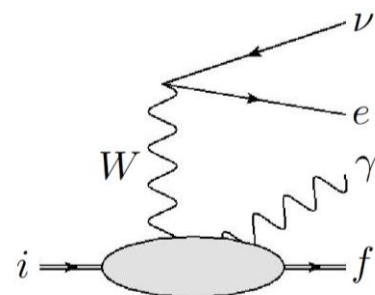
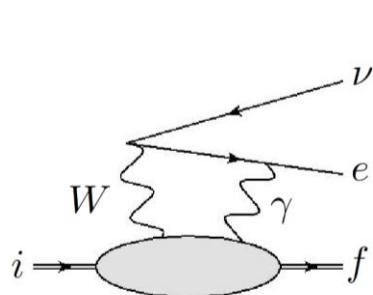


$$d\Gamma \propto \overline{|\mathcal{M}_{\gamma W}|^2} q^0 + \overline{|\mathcal{M}_{\text{tree}}|^2} q^1 + 2\Re\{\overline{\mathcal{M}_{\text{tree}}^*} \mathcal{M}_{\gamma W}\} q^2$$

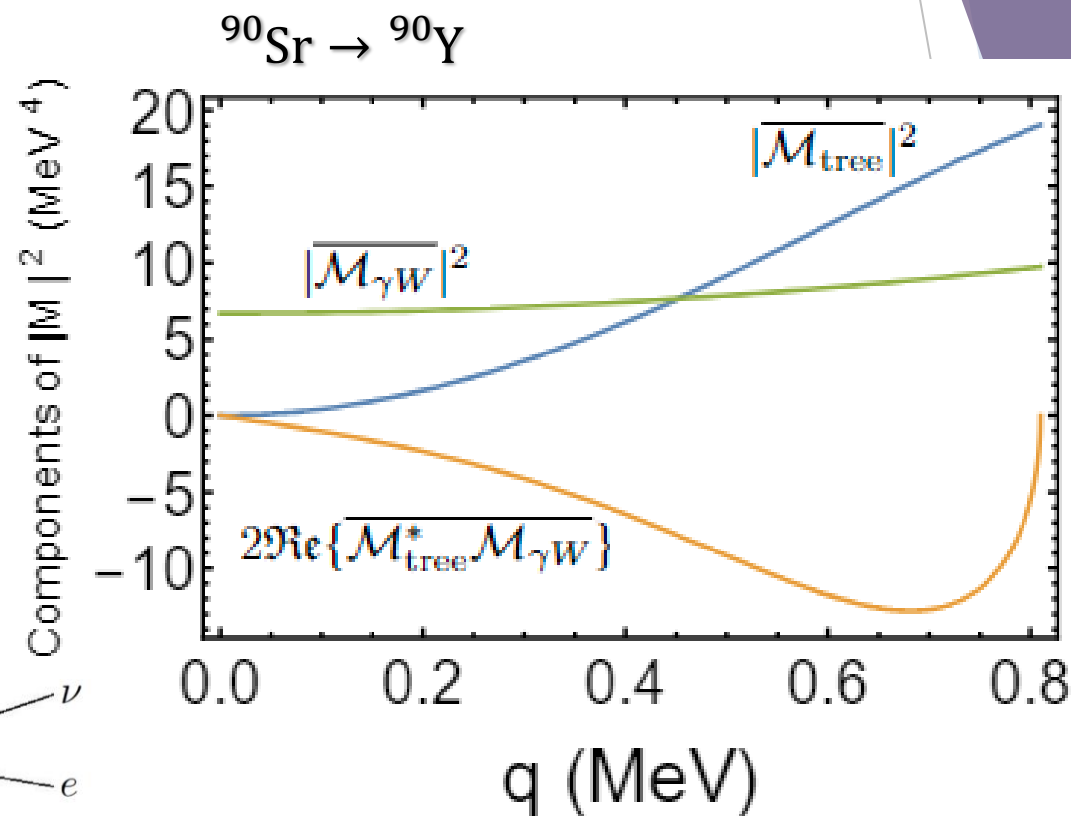
Unique Forbidden decays

No decay @ $q \rightarrow 0$

Light New Physics



$$d\Gamma \propto f_0 q^0 + f_1 q^1 + f_2 q^2$$



Unique Forbidden decays: Summary

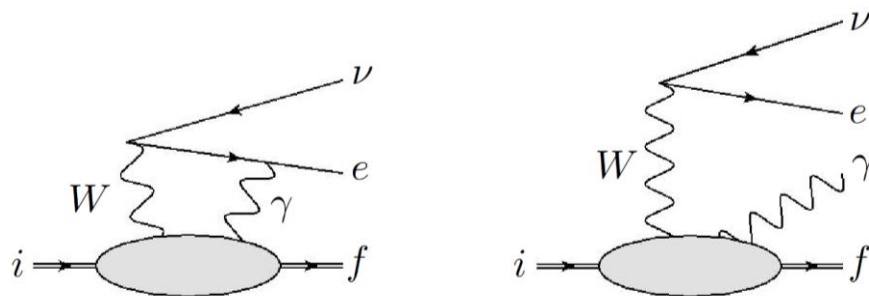
No decay @ $q \rightarrow 0$

Experiments

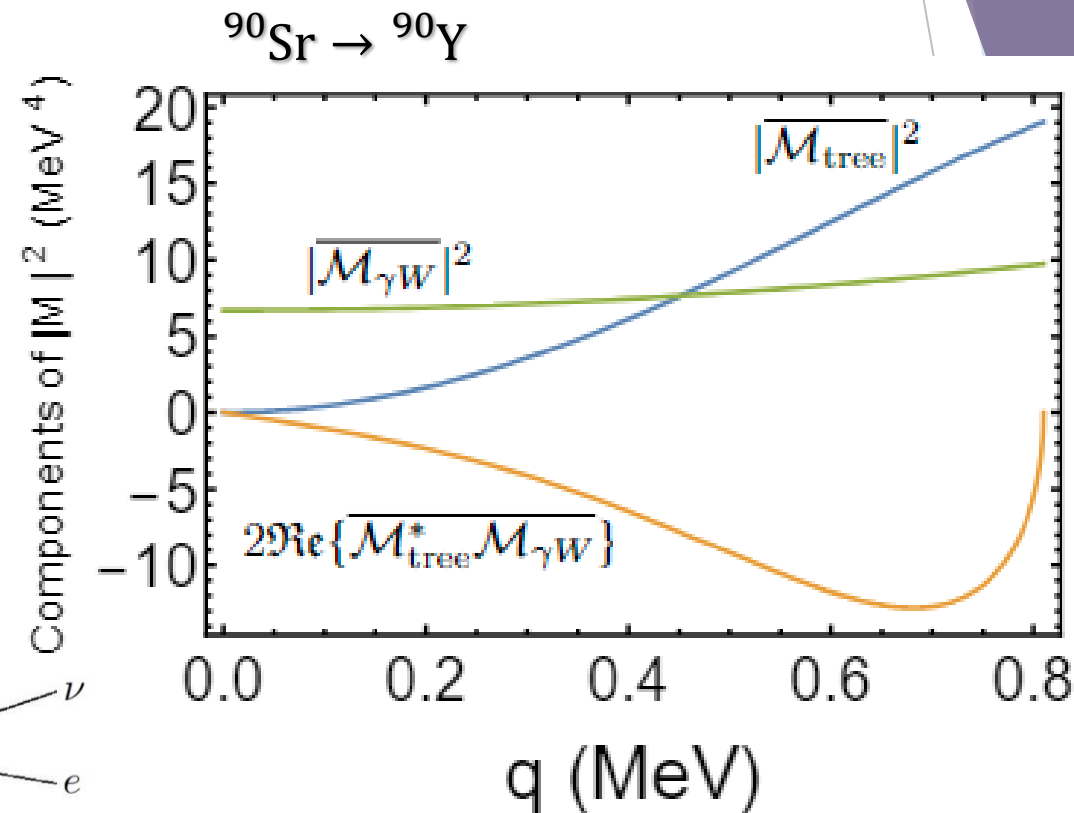
@ ORNL, HUJI, SOREQ (FRIB?)

Spectrum is sensitive to:

- ↙ Exotic Weak Interactions
- ↙ Light New Physics



$$d\Gamma \propto f_0 q^0 + f_1 q^1 + f_2 q^2$$



Unique Forbidden decays: Summary & Outlook

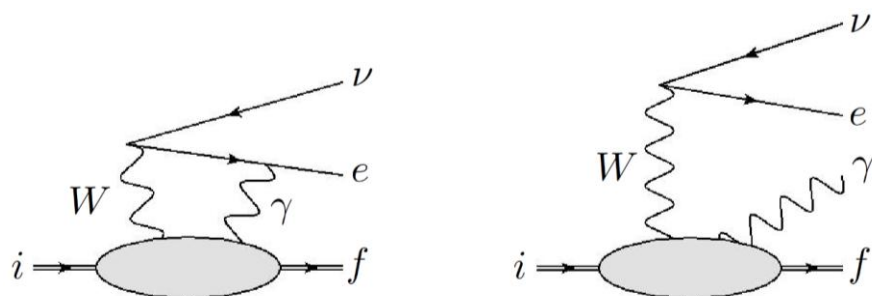
No decay @ $q \rightarrow 0$

Experiments

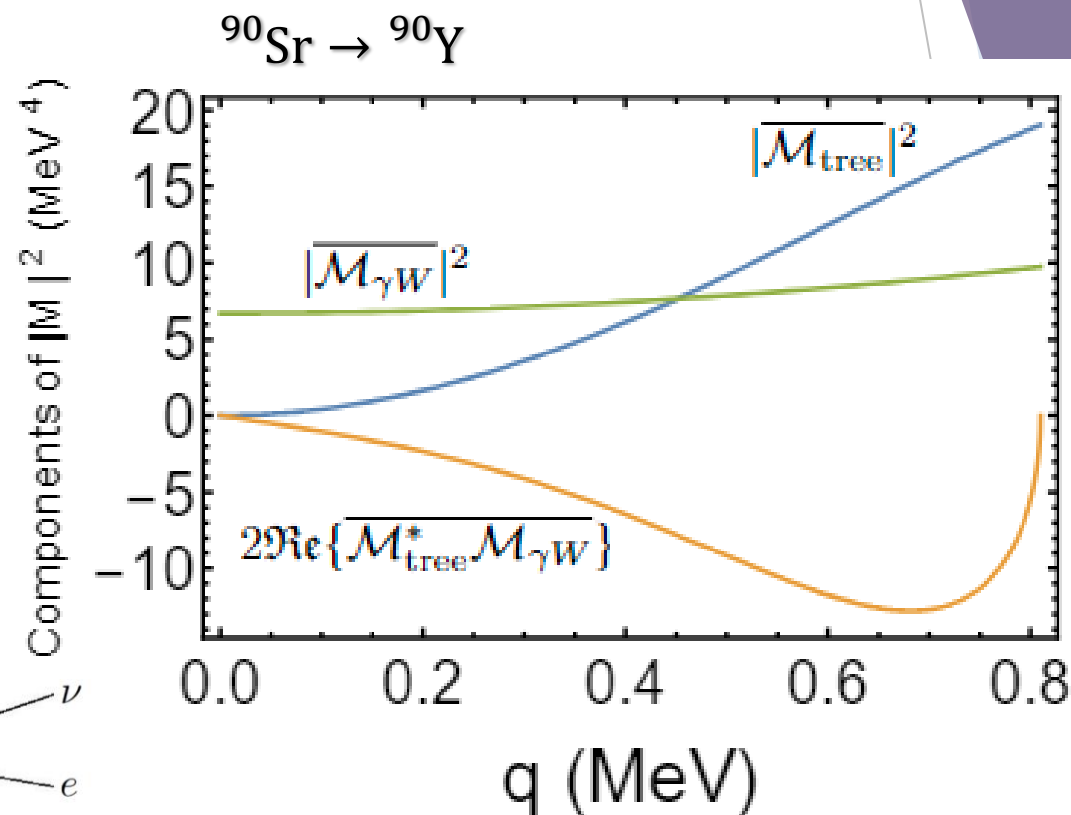
@ ORNL, HUJI, SOREQ (FRIB?)

Spectrum is sensitive to:

- ↙ Exotic Weak Interactions
- ↙ Light New Physics



$$d\Gamma \propto f_0 q^0 + f_1 q^1 + f_2 q^2$$



The new radiative regime:

- Beyond $q \rightarrow 0$ power counting
- With nuclear matrix elements
- Light new physics possibilities

Thanks!

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LPC Caen

Charlie Rasco
ORNL

Guy Ron
Hebrew U.

DOE Topical Collaboration “Nuclear Theory for New Physics”

FRIB Theory Alliance

U.S. DOE Office of Science, Office of Nuclear Physics

Nuclear Theory for New Physics



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