

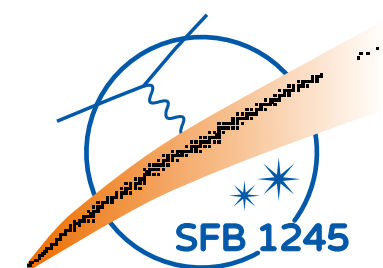
Ab initio overlap integrals for $\mu \rightarrow e$ conversion in nuclei

MH, Hoferichter, Miyagi, Noël, Schwenk, arXiv:2412.04545

Matthias Heinz, ORNL

CIPANP 2025
Madison, WI, June 9, 2025

Work supported by:

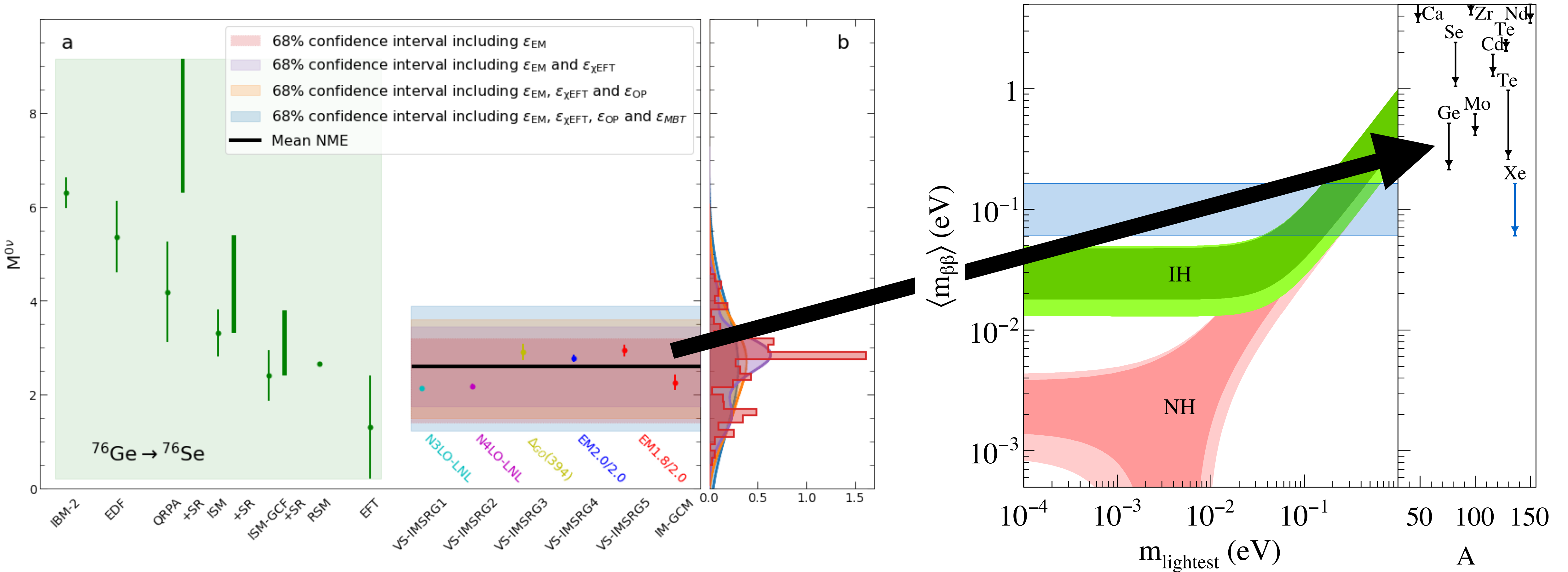


U.S. DEPARTMENT OF
ENERGY

NUCLEI
Nuclear Computational Low-Energy Initiative

This research used resources of the Oak Ridge Leadership Computing Facility located at Oak Ridge National Laboratory, which is supported by the Office of Science of the Department of Energy under contract No. DE-AC05-00OR22725.

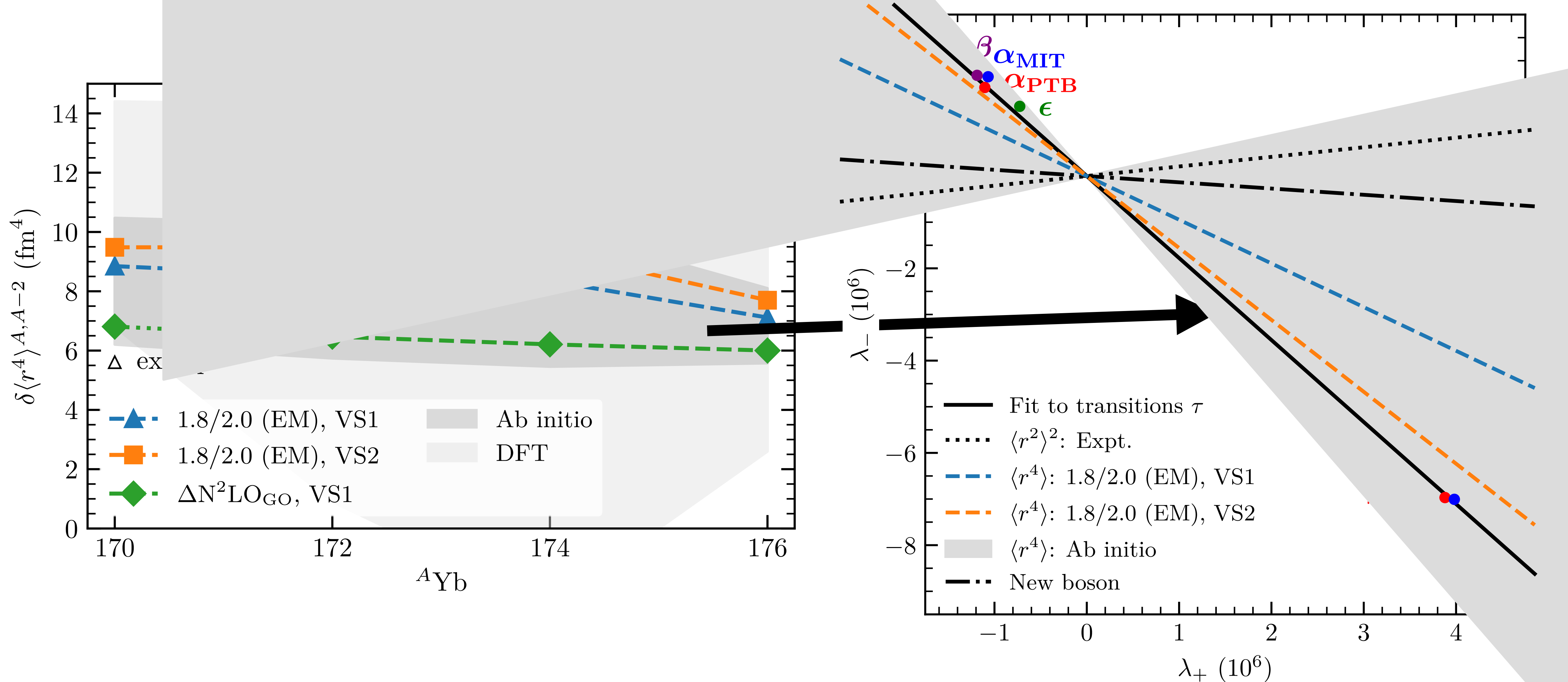
New physics in nuclei requires nuclear theory



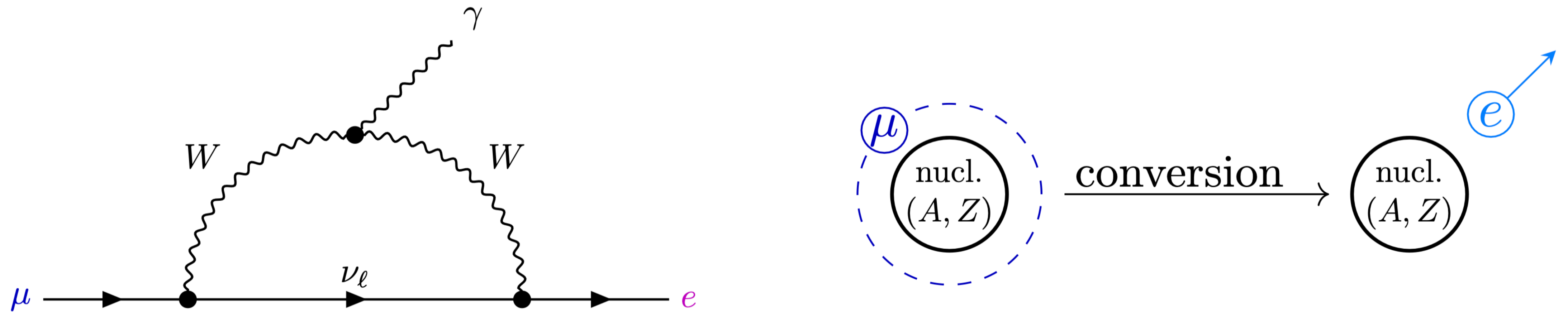
Belley et al., PRL **132** (2024)

Engel, Menéndez, RPP **80** (2016)

Neutron radii requires nuclear theory



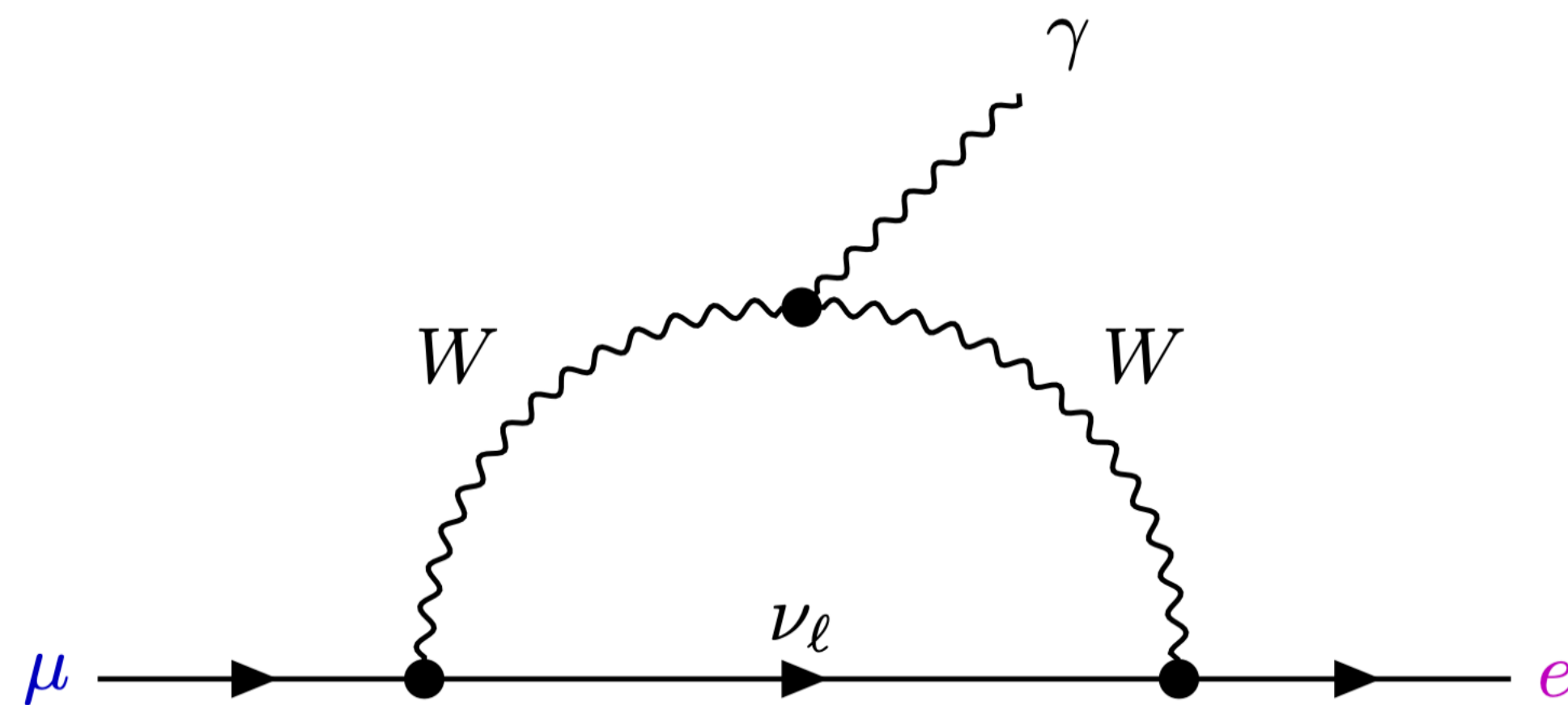
Muon to electron conversion



Figures: F. Noël

- Lepton flavor violation in the standard model suppressed by $(\Delta m_\nu / m_W)^4 \sim 10^{-50}$
- Searches for $\mu \rightarrow e$ conversion constrain new lepton flavor violating interactions
- Complementary channels: $\mu \rightarrow e\gamma$; $\mu \rightarrow 3e$; $\mu \rightarrow e$ [nucl.]

Muon to electron conversion



Figures: F. Noël

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Our focus: ^{27}Al , ^{48}Ti

What has been done so far?

- Microscopic theories mapped to nonrelativistic effective theory operators
See talks by, e.g., Kaori Fuyuto, Evan Rule, William McNulty
- Treatment of leptonic part varies;
Key challenge: **Coulomb distortion**
 - Coulomb Dirac solution
→ Overlap integrals
 - Free Dirac with effective electron momentum q_{eff} → plane-wave formalism
- Neutron density contribution important, but **not well understood** with **unquantified uncertainties**

Method 1. First, we take the proton density from electron scattering experiments given in Appendix A and assume that the neutron density is the same as the proton density. For

*Kitano, Koike, Okada, PRD **66** (2002)*

explore uncertainties associated with alternative schemes for fitting the effective interaction. The uncertainties associated with the SM itself—the basis truncation and the choice of single particle basis—are more difficult to assess. Shell model

*Haxton, Rule, McElvain, Ramsey-Musolf, PRC **107** (2022)*

Revisiting the overlap integrals

$$\mu \rightarrow e \text{ decay rate} \sim \sum_i \left| \bar{C}^{I_i} \times I_i \right|^2$$

- Knowing overlap integrals constrains Wilson coeffs of underlying theories
- Overlap integrals combine nuclear densities and lepton wave functions

$$S^{(n)} \sim \int_0^\infty dr r^2 \rho_n(r) s(r)$$

- Full set for spin-independent formalism: D , $S^{(p)}$, $V^{(p)}$, $S^{(n)}$, $V^{(n)}$
- Charge, point-proton, and point-neutron densities required

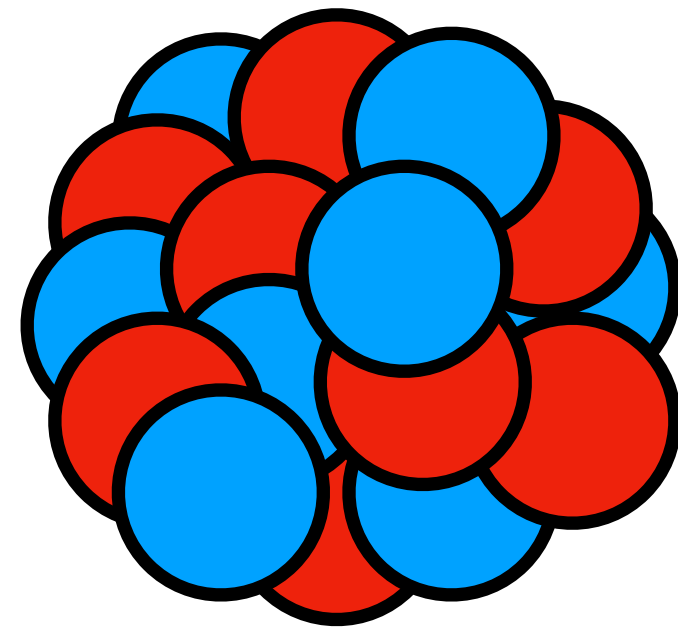
Ab initio nuclear structure theory

Ab initio nuclear structure

● N neutrons

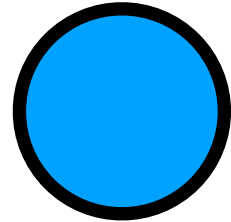
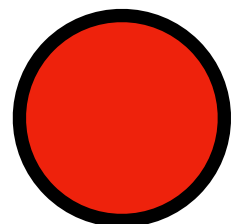
● Z protons

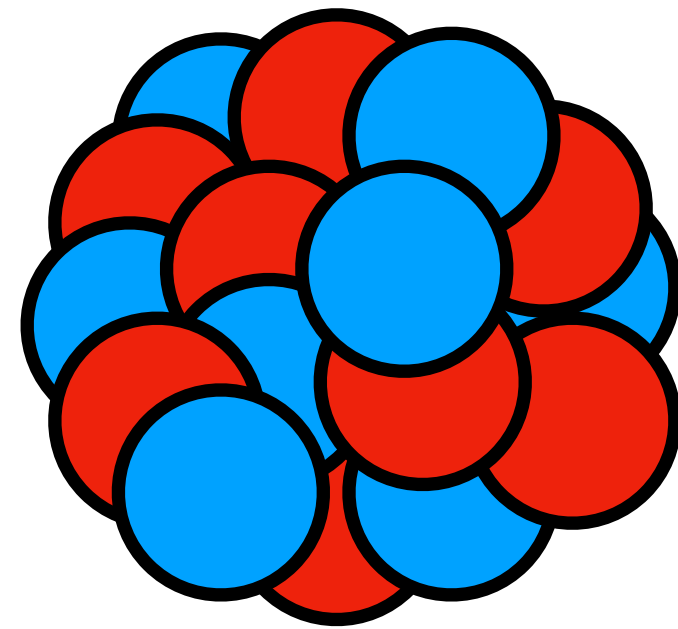
A nucleons



$$H|\Psi\rangle = E|\Psi\rangle$$

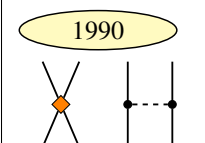
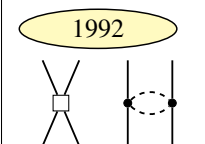
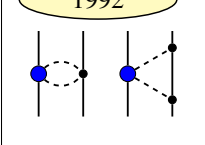
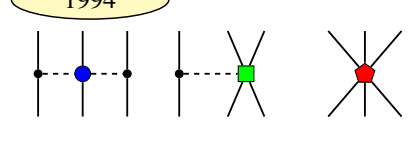
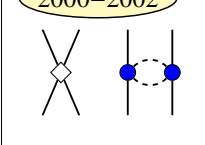
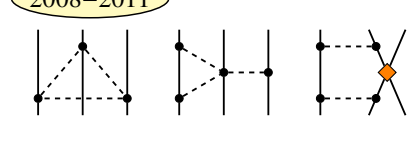
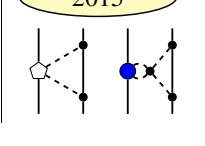
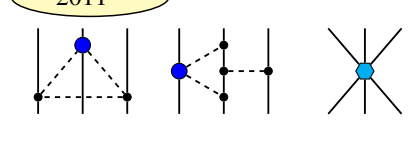
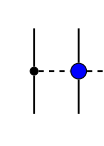
Ab initio nuclear structure

 N neutrons
 Z protons
 A nucleons



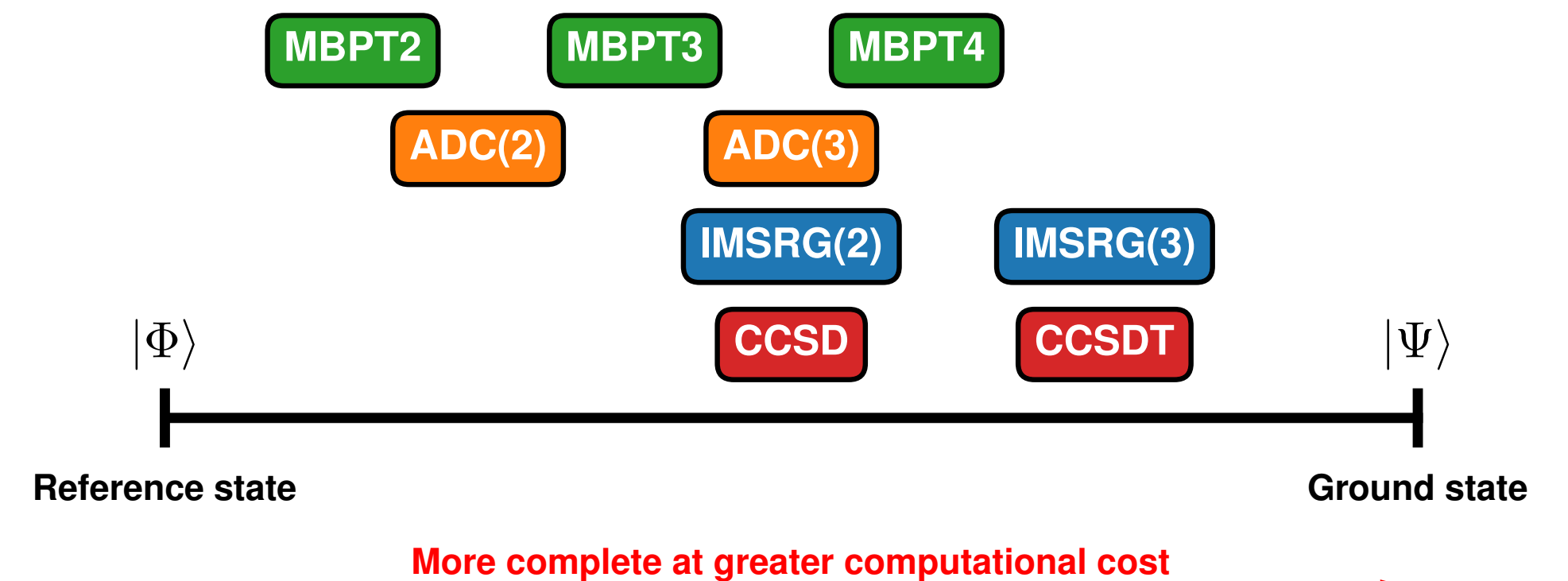
$$H|\Psi\rangle = E|\Psi\rangle$$

Nuclear forces

	NN	3N	4N
LO $\mathcal{O}(Q^0/\Lambda^0)$	 1990 2	—	—
NLO $\mathcal{O}(Q^2/\Lambda^2)$	 1992 7	 1992, 1994	—
N ² LO $\mathcal{O}(Q^3/\Lambda^3)$	 1992 0	 1994 2	—
N ³ LO $\mathcal{O}(Q^4/\Lambda^4)$	 2000–2002 12	 2008–2011 0	 2006 0
N ⁴ LO $\mathcal{O}(Q^5/\Lambda^5)$	 2015 0	 2011– ?	 ?

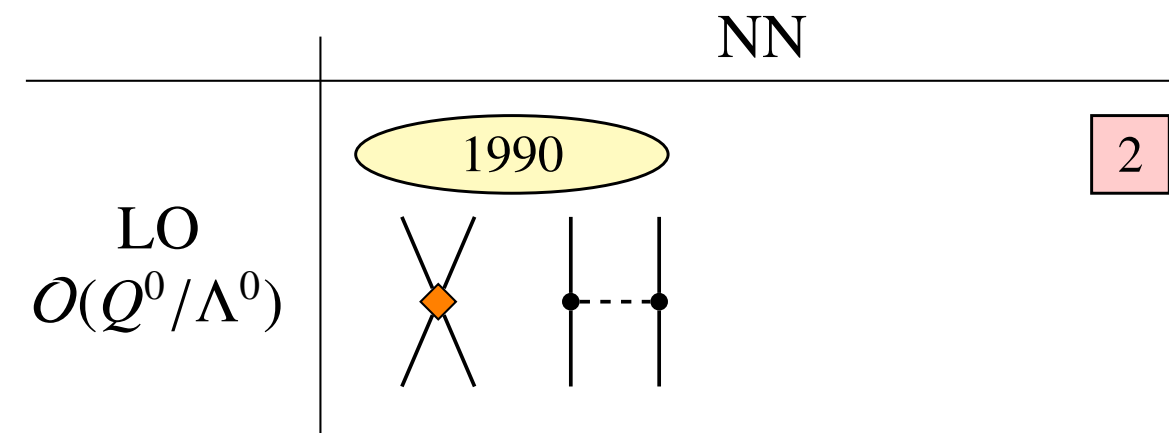
Hebeler, Phys. Rep. **890** (2021)

Many-body methods



Nuclear forces from chiral EFT

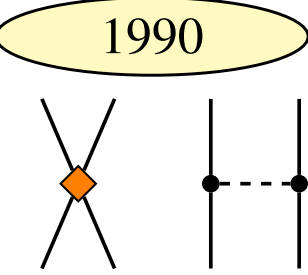
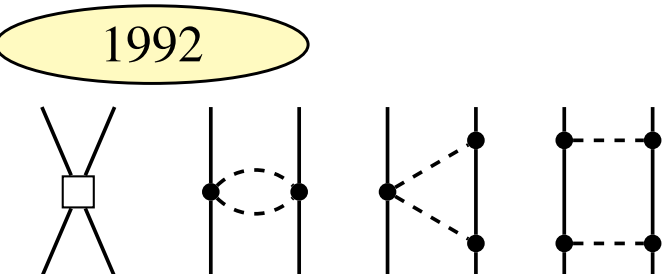
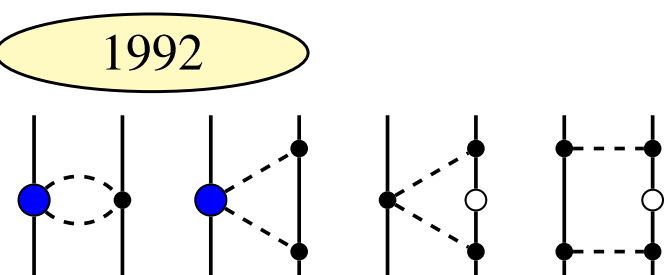
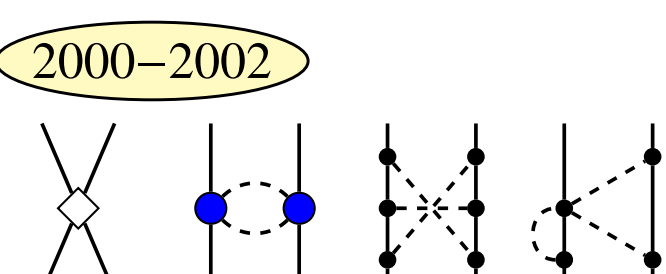
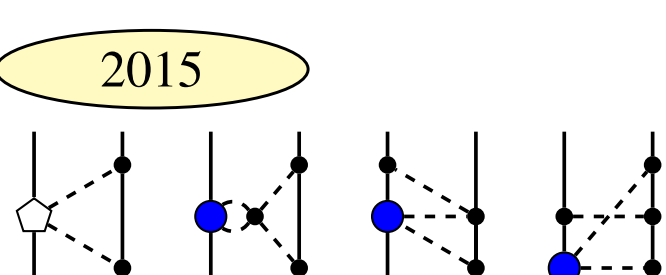
effective field theory



- **Nuclear forces are uncertain**
- Chiral EFT:
Low-energy expansion of QCD
- Long distances: pion exchanges
- Short distances: contact expansion

Nuclear forces from chiral EFT

effective field theory

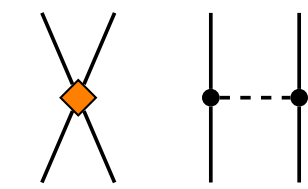
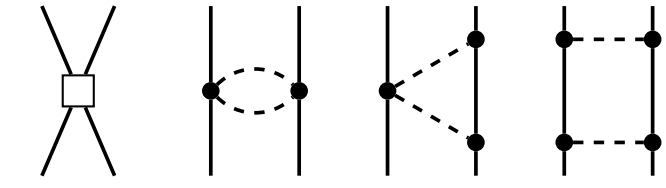
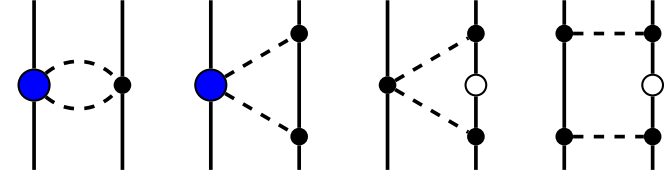
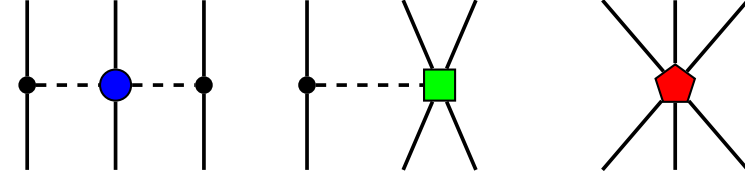
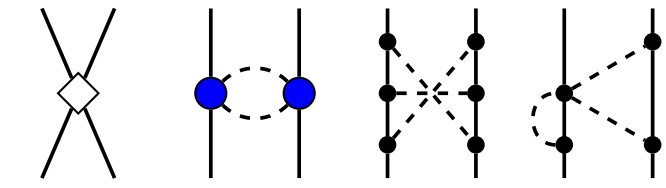
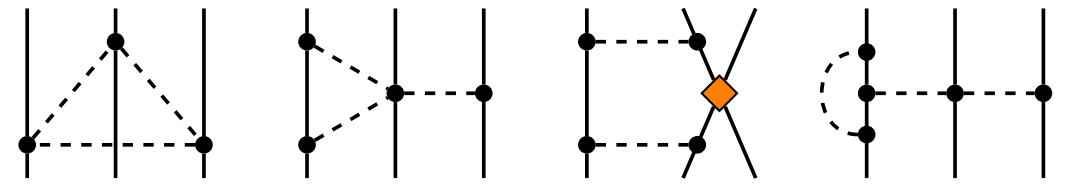
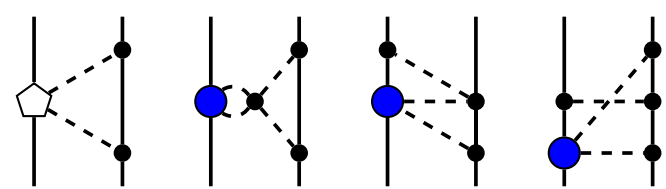
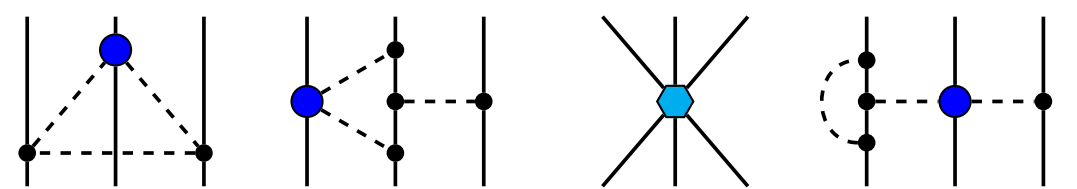
	NN	
LO $O(Q^0/\Lambda^0)$		2
NLO $O(Q^2/\Lambda^2)$		7
N ² LO $O(Q^3/\Lambda^3)$		0
N ³ LO $O(Q^4/\Lambda^4)$		12
N ⁴ LO $O(Q^5/\Lambda^5)$		0

Hebeler, Phys. Rep. **890** (2021)

- **Nuclear forces are uncertain**
- Chiral EFT:
Low-energy expansion of QCD
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Nuclear forces from chiral EFT

effective field theory

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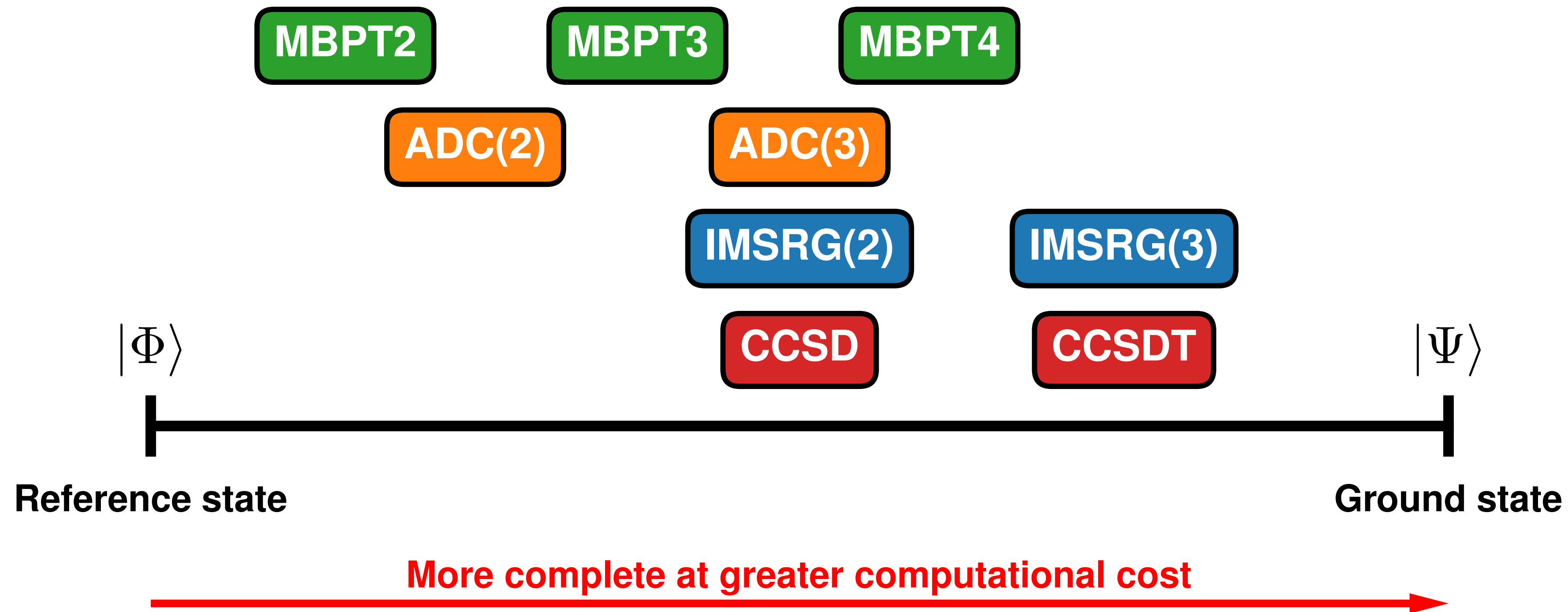
- **Nuclear forces are uncertain**
- Chiral EFT:
 - **Low-energy expansion** of QCD
 - Long distances: pion exchanges
 - Short distances: contact expansion
- Natural inclusion of 3B forces
- Truncation & **systematic improvement**
- **Uncertainty quantification**
- **Consistent prediction of currents to describe external probes**

Hebeler, Phys. Rep. 890 (2021)

Many-body expansion methods



Many-body expansion methods

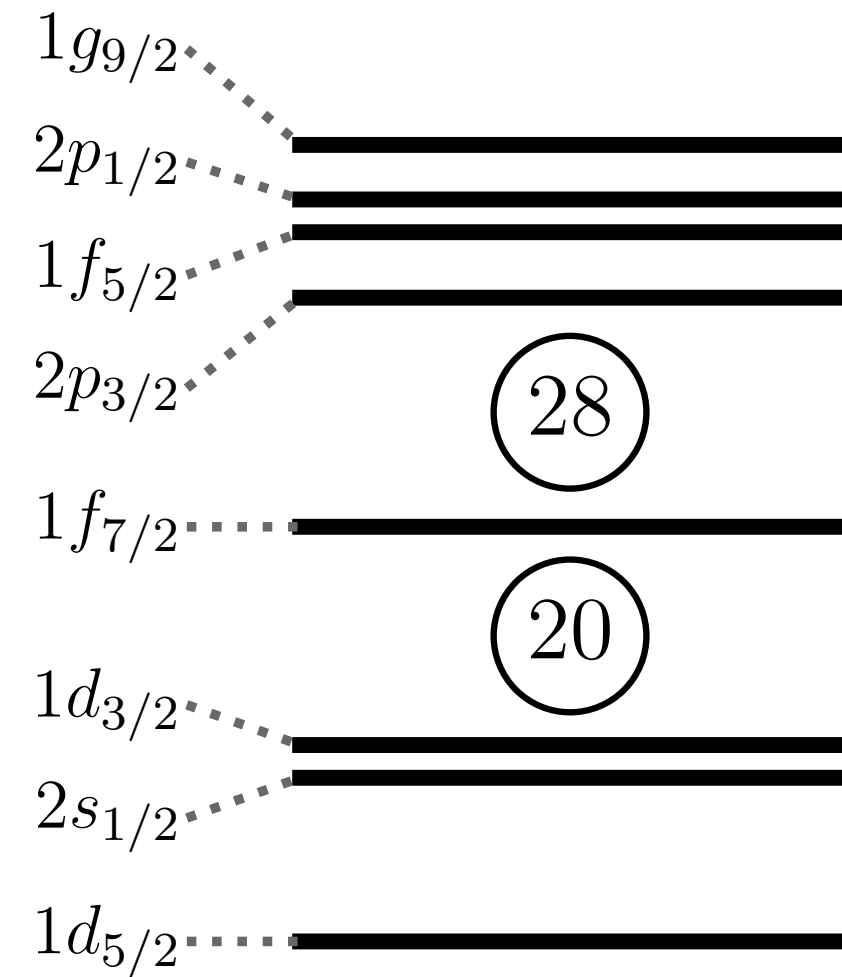


- **Systematically improvable expansion** around reference state $|\Phi\rangle$
- **Tractable computational cost** in larger nuclei
- Approximate many-body solution with **quantifiable uncertainty**

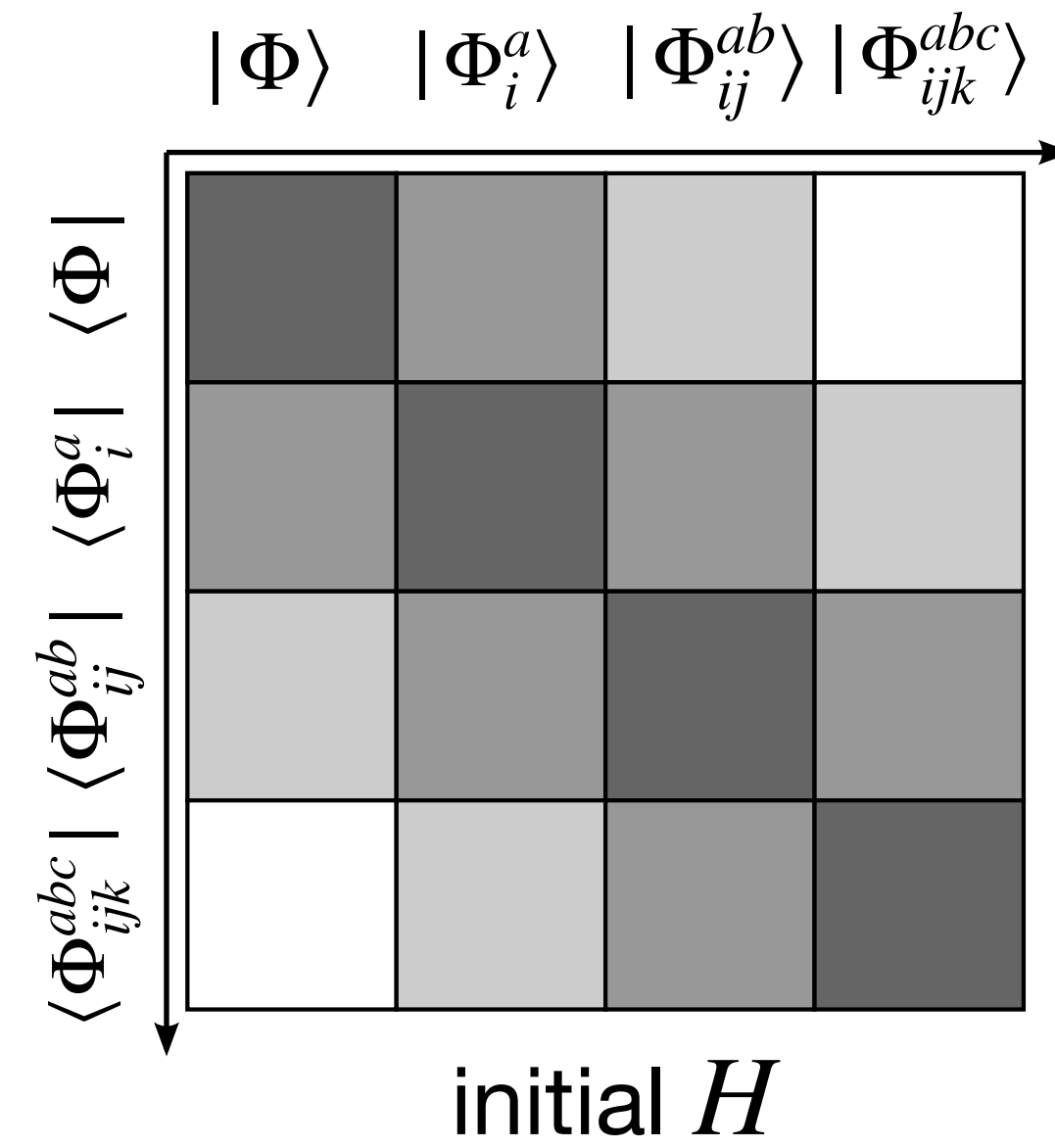
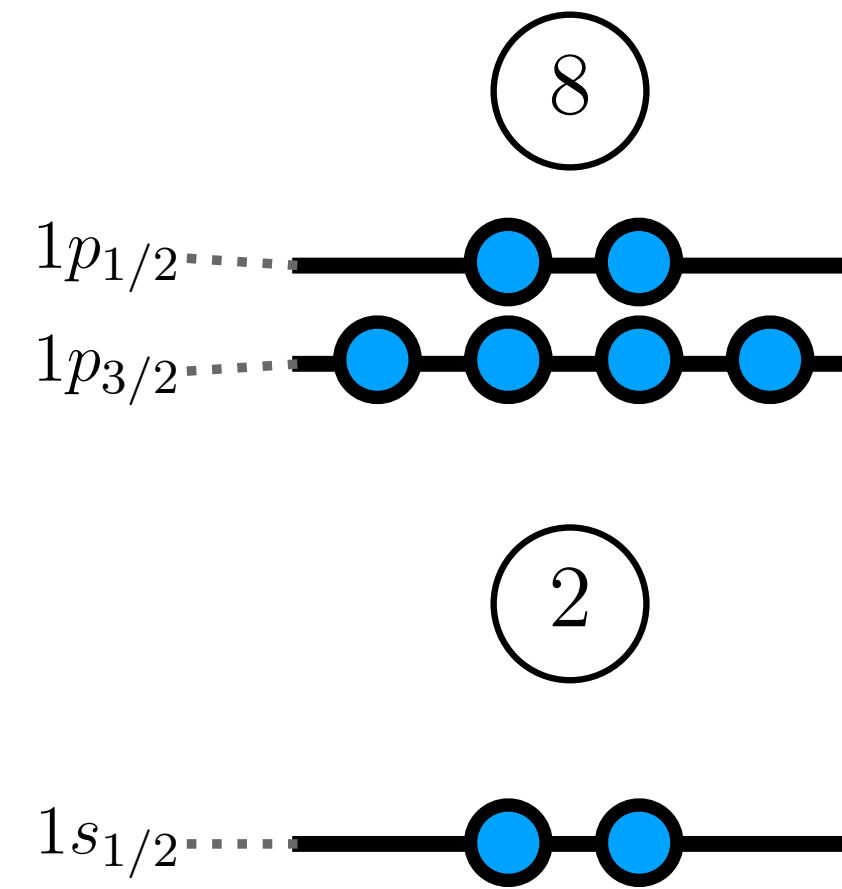
The IMSRG

in-medium similarity renormalization group

excitations



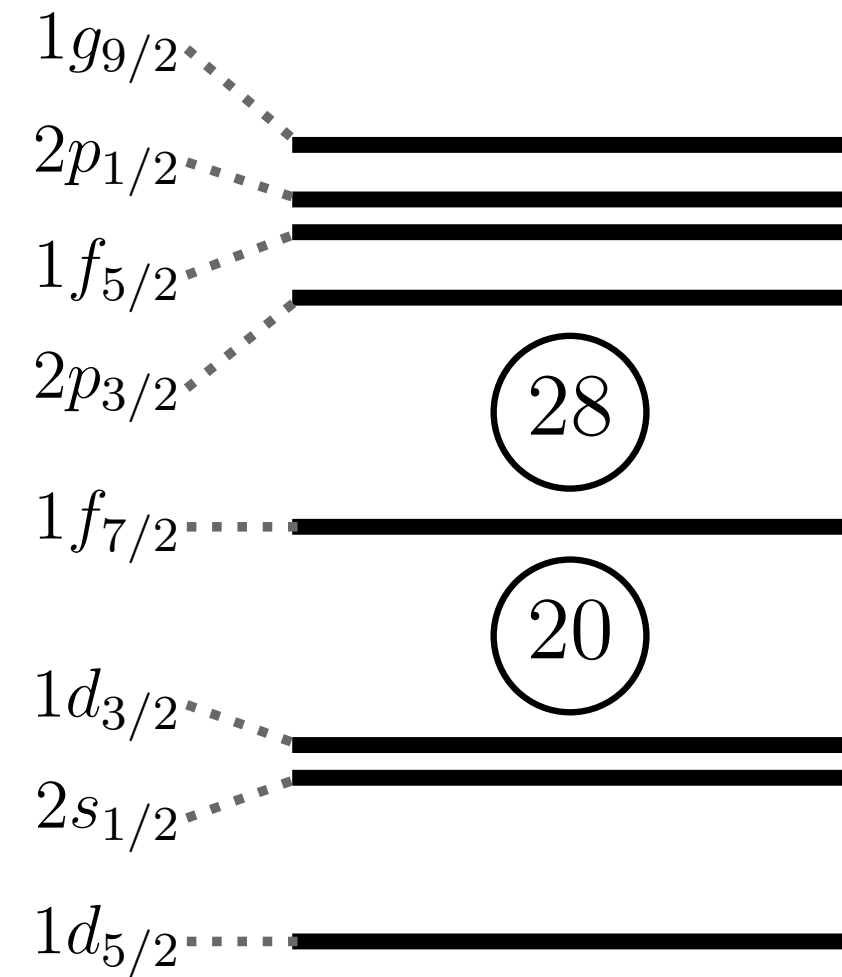
reference state



The IMSRG

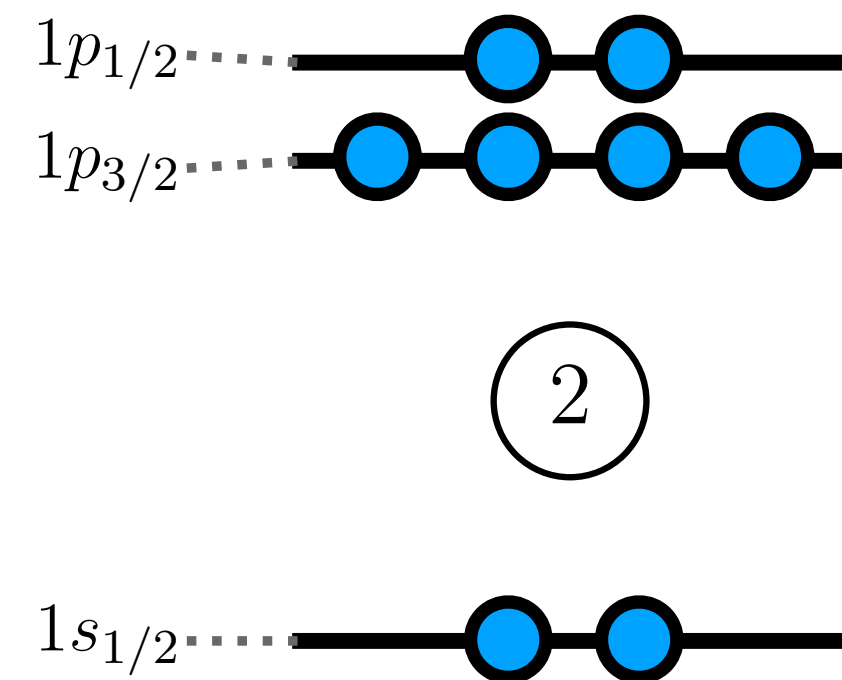
in-medium similarity renormalization group

excitations

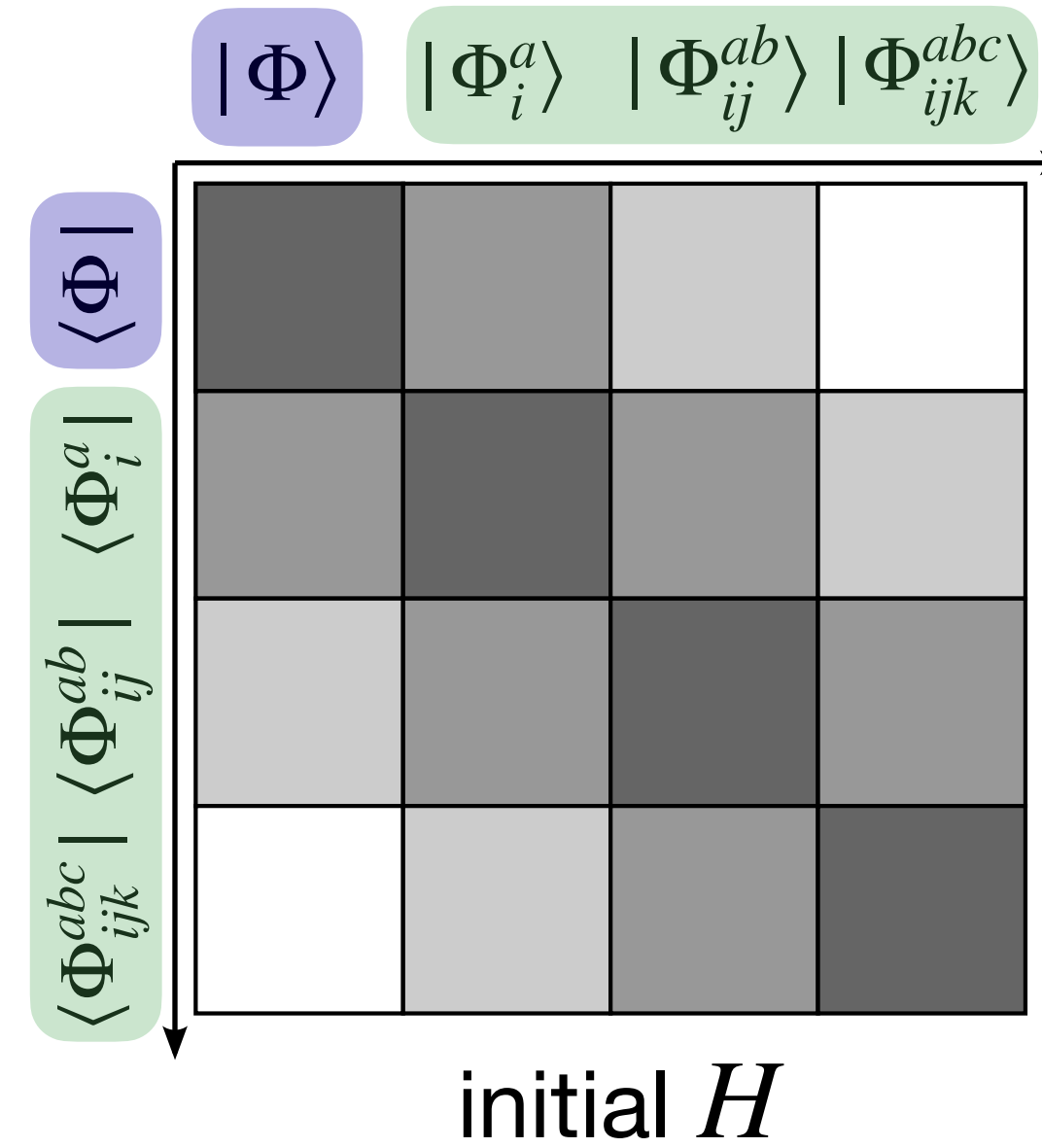


8

reference state



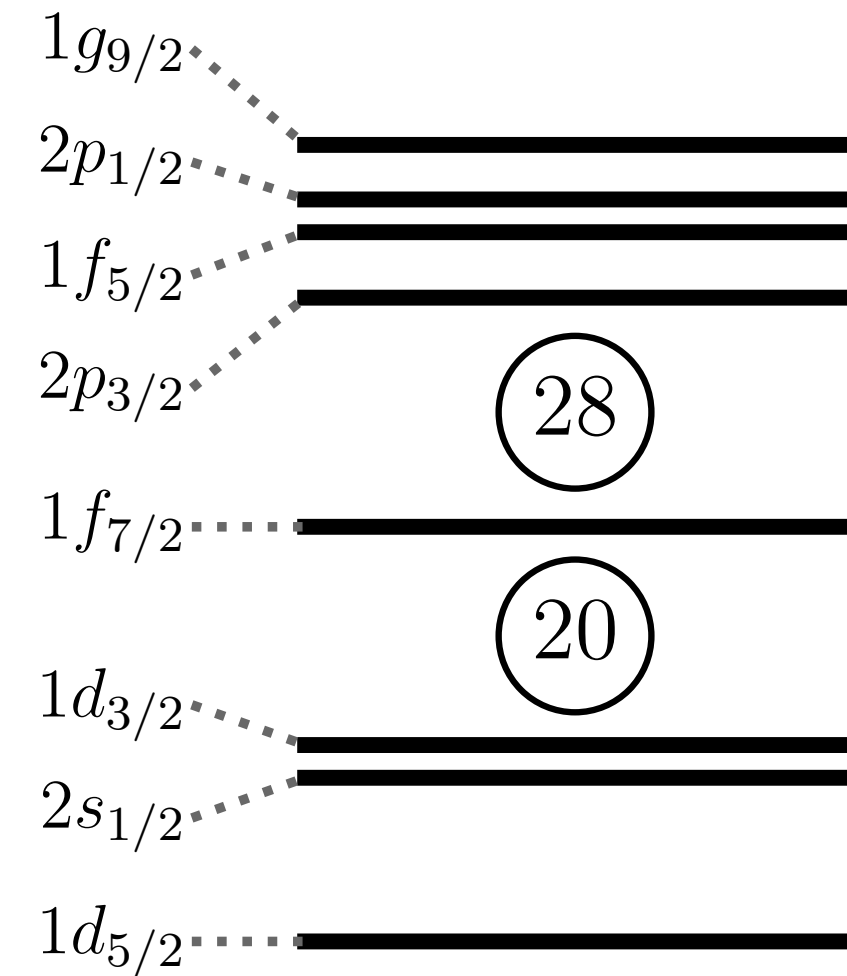
2



The IMSRG

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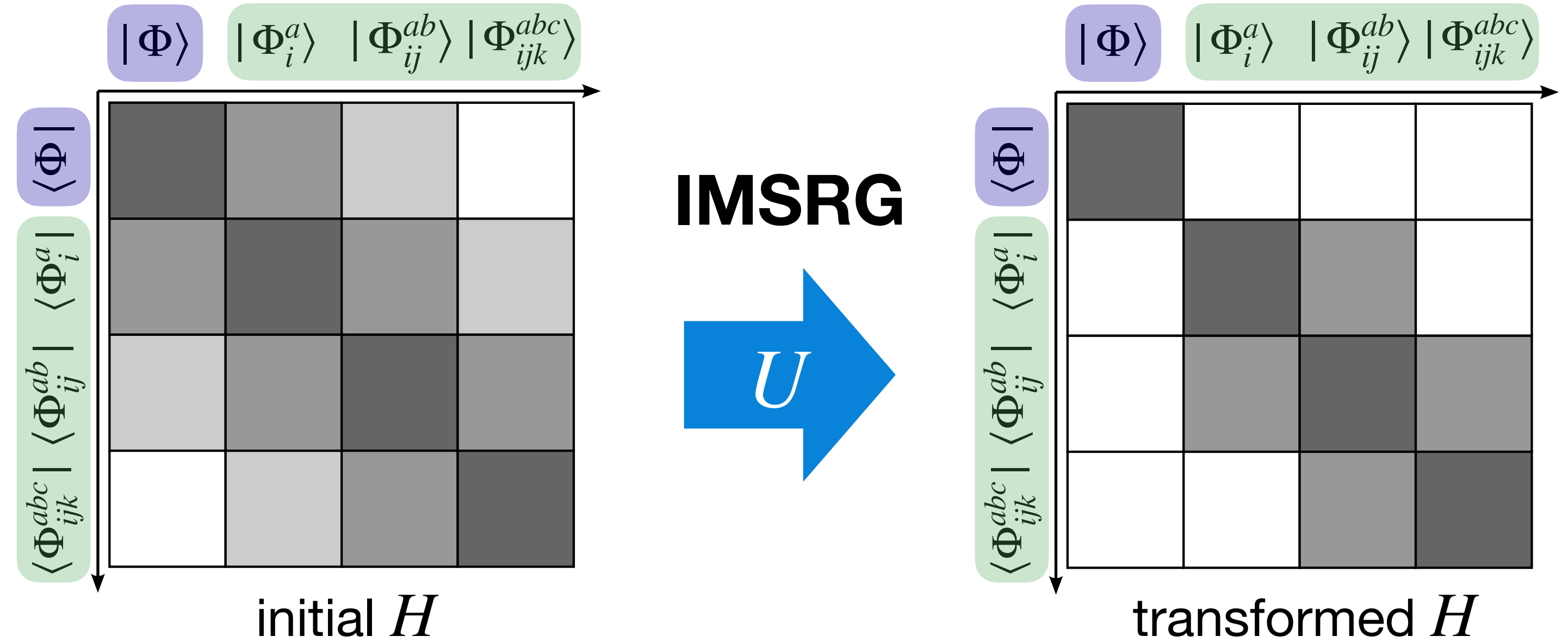
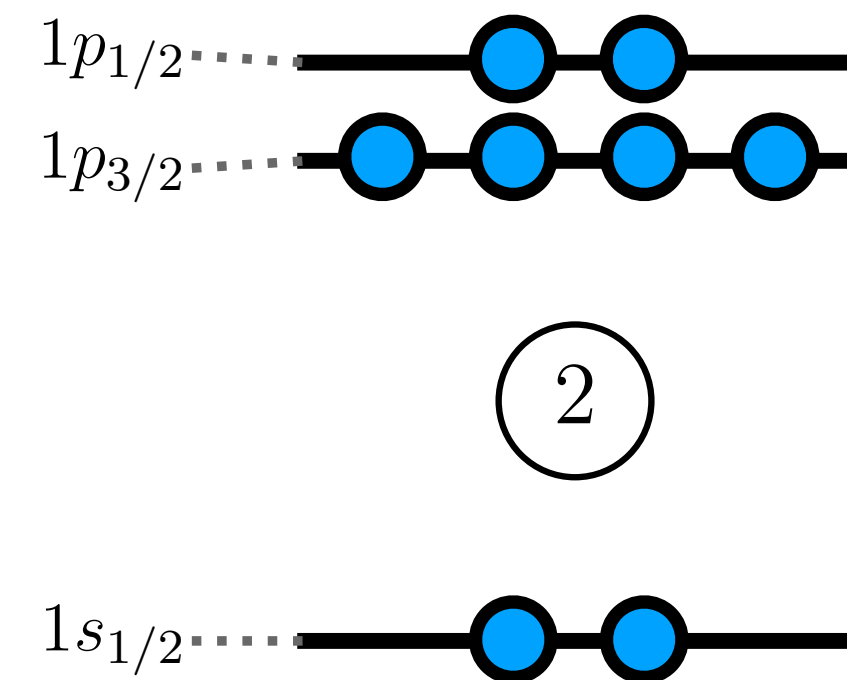
excitations



decouple

8

reference state



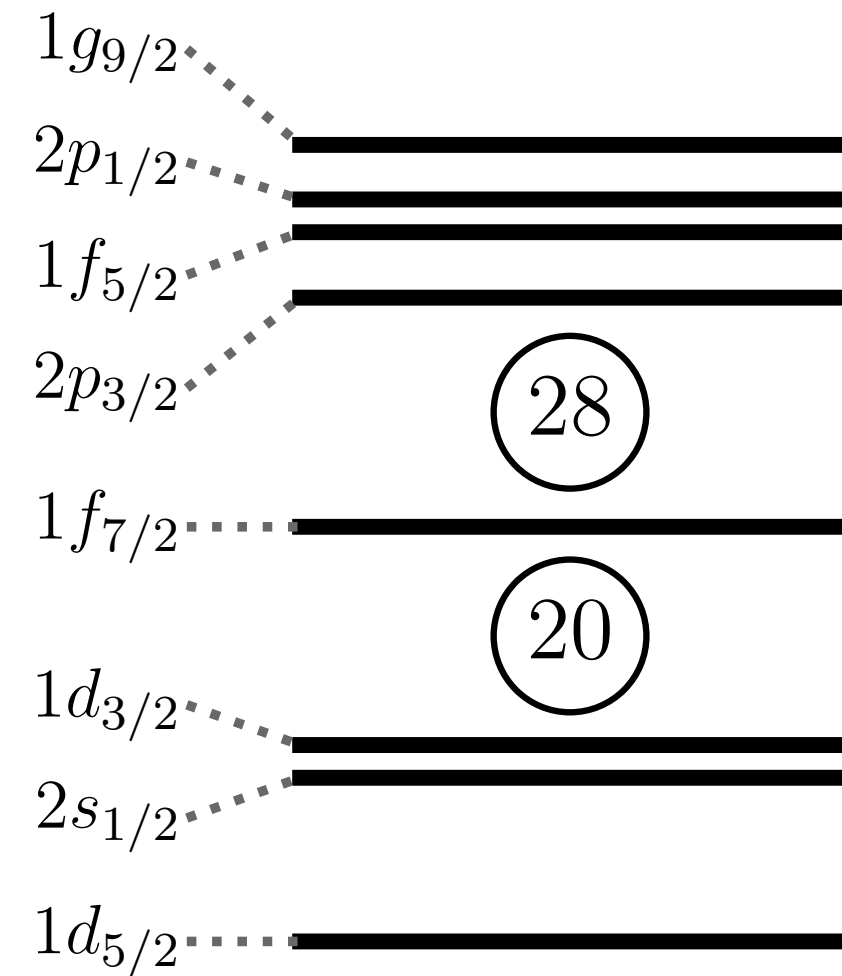
Hergert et al., Phys. Rep. **621** (2016)

- IMSRG:** Unitary transformation $U = e^{\Omega}$ to decouple reference state from excitations

The IMSRG

in-medium similarity renormalization group

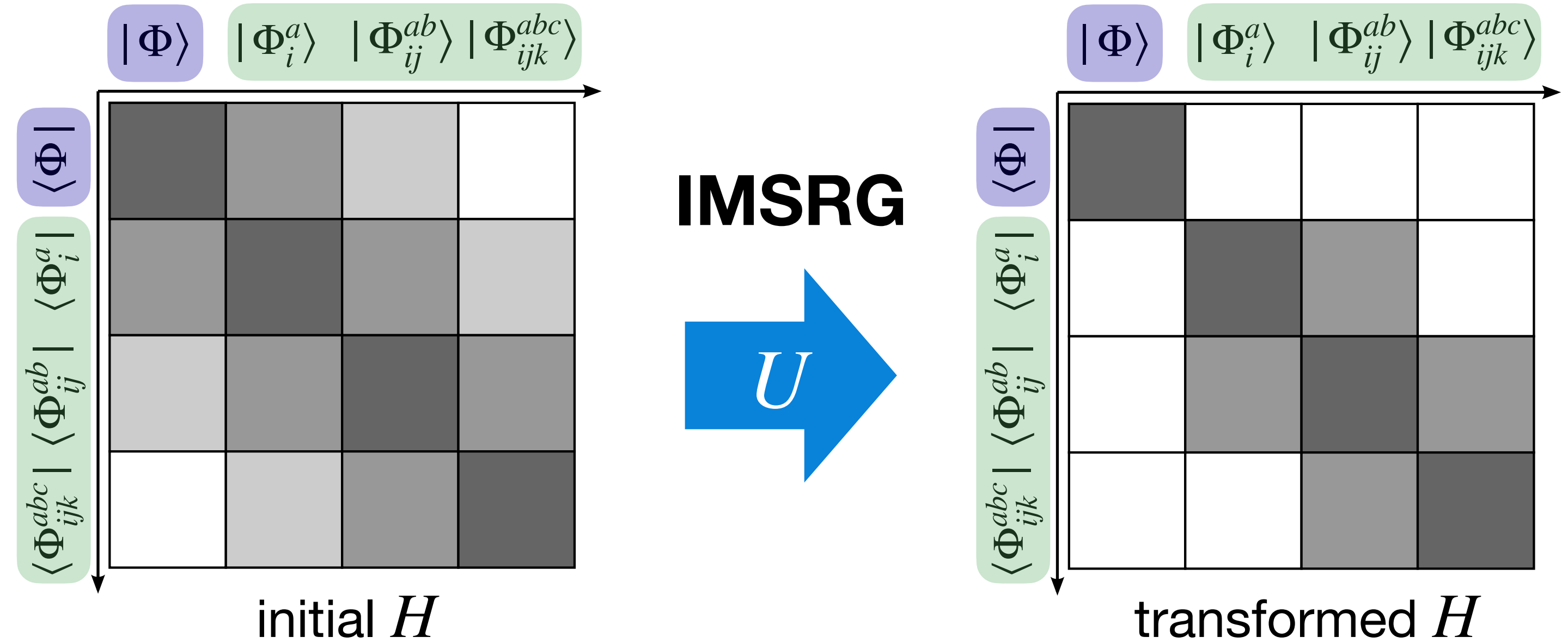
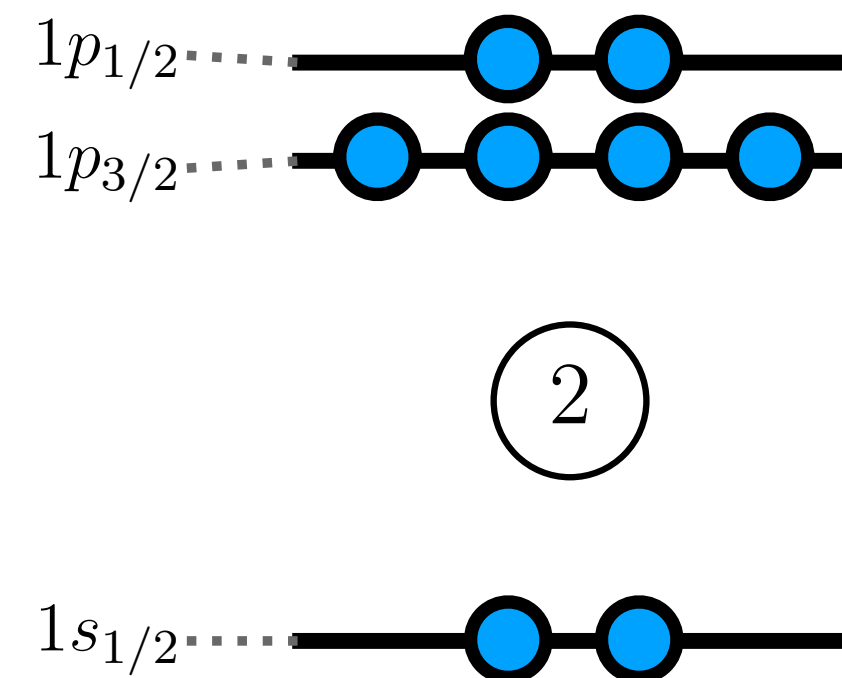
excitations



decouple

8

reference state



Hergert et al., Phys. Rep. **621** (2016)

- **IMSRG**: Unitary transformation $U = e^{\Omega}$ to decouple reference state from excitations
- Expansion and truncation in **many-body operators**

$$U = e^{\Omega} = e^{\Omega_1 + \Omega_2 + \Omega_3 + \dots}$$

MH et al., PRC **103** (2021)
PRC **111** (2025)
Stroberg, He (2024)

- **IMSRG(3)** for precision and uncertainty quantification

Computed nuclear responses

Nuclear response notation: $X_{J,p/n}(q)$

Spin-independent:

- $M_{0,p}, M_{0,n}$: coherent, related to point-proton, point-neutron densities

Semi-coherent:

- $\Phi''_{0,p}, \Phi''_{0,n}$: spin-orbit effects

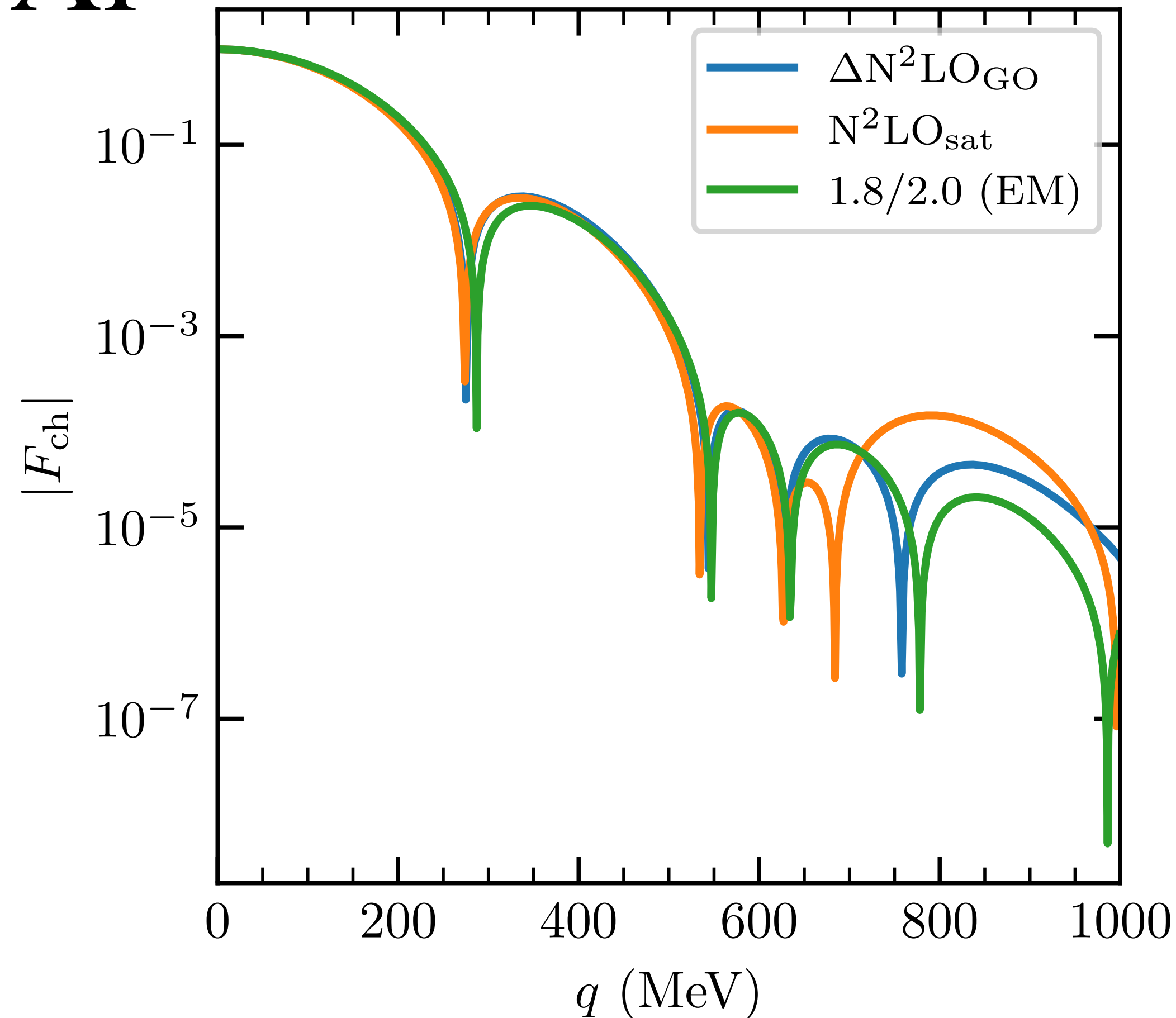
Incoherent:

- $M_{J,p/n}, \Phi''_{J,p/n}, J = 2, \dots; \overset{\text{orb. ang. mom.}}{\Delta_{J,p/n}}, \overset{\text{spin}}{\Sigma'_{J,p/n}}, \Sigma''_{J,p/n}, J = 1, \dots$
- No obvious hierarchy for incoherent responses $\rightarrow$ need up to $J = 5$ for ^{27}Al

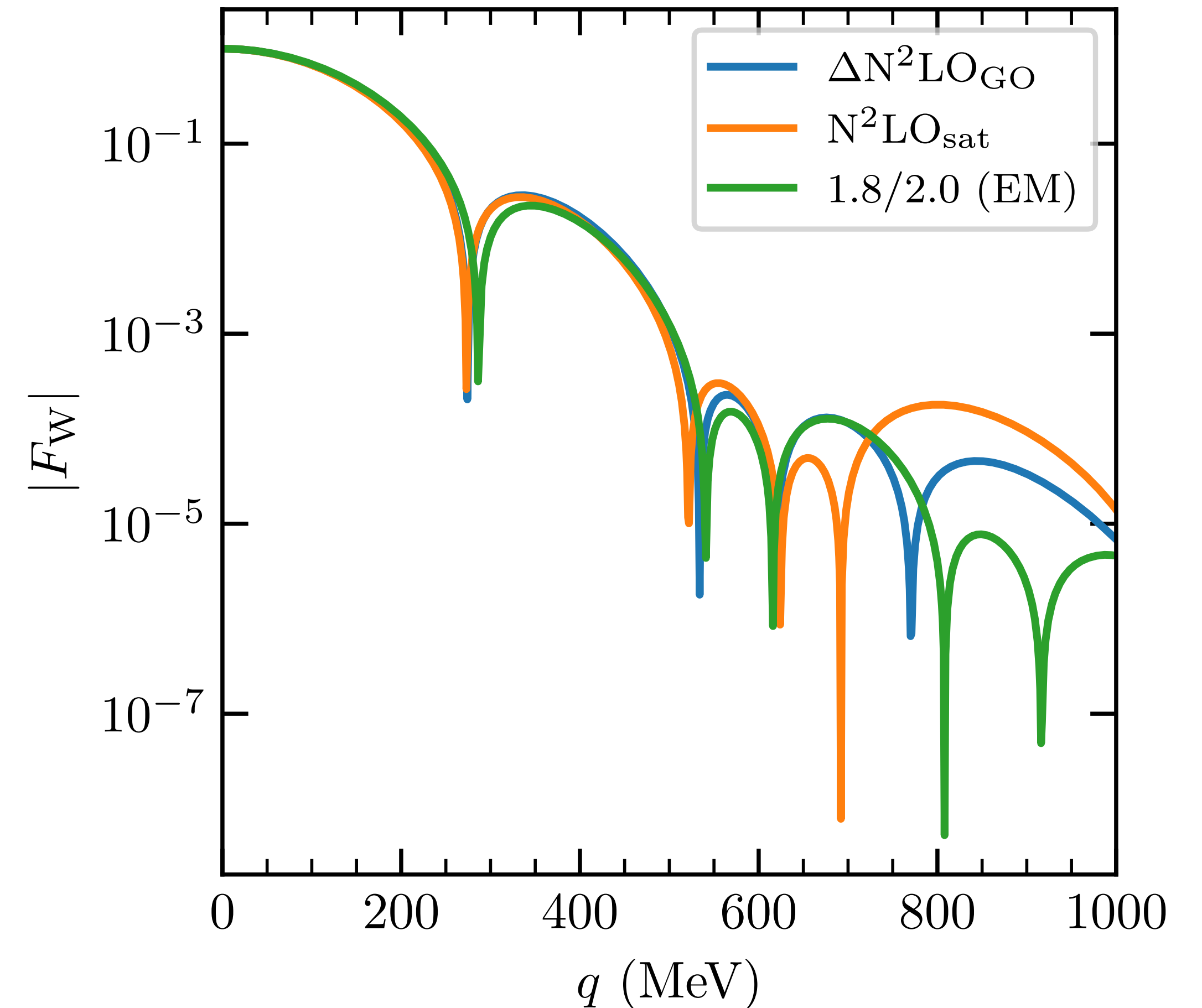
Ab initio $\mu \rightarrow e$ overlap integrals

Charge and weak responses are correlated

^{27}Al charge = easy to measure



weak = hard to measure



- **Approach:** Correlations between $\rho_{\text{ch}}(r)$, $\rho_n(r)$, $\rho_p(r)$ constrain overlap integrals

The correlation

Ensemble of 42 Hamiltonians:

- “Magic” 4 [1.8/2.0 (EM), ...]

Hebeler et al., PRC **83** (2011)

- N^2LO_{sat}

Ekström et al., JPG **42** (2015)

- ΔN^2LO_{GO} ($\Lambda = 394$ MeV)

Jiang et al., PRC **102** (2020)

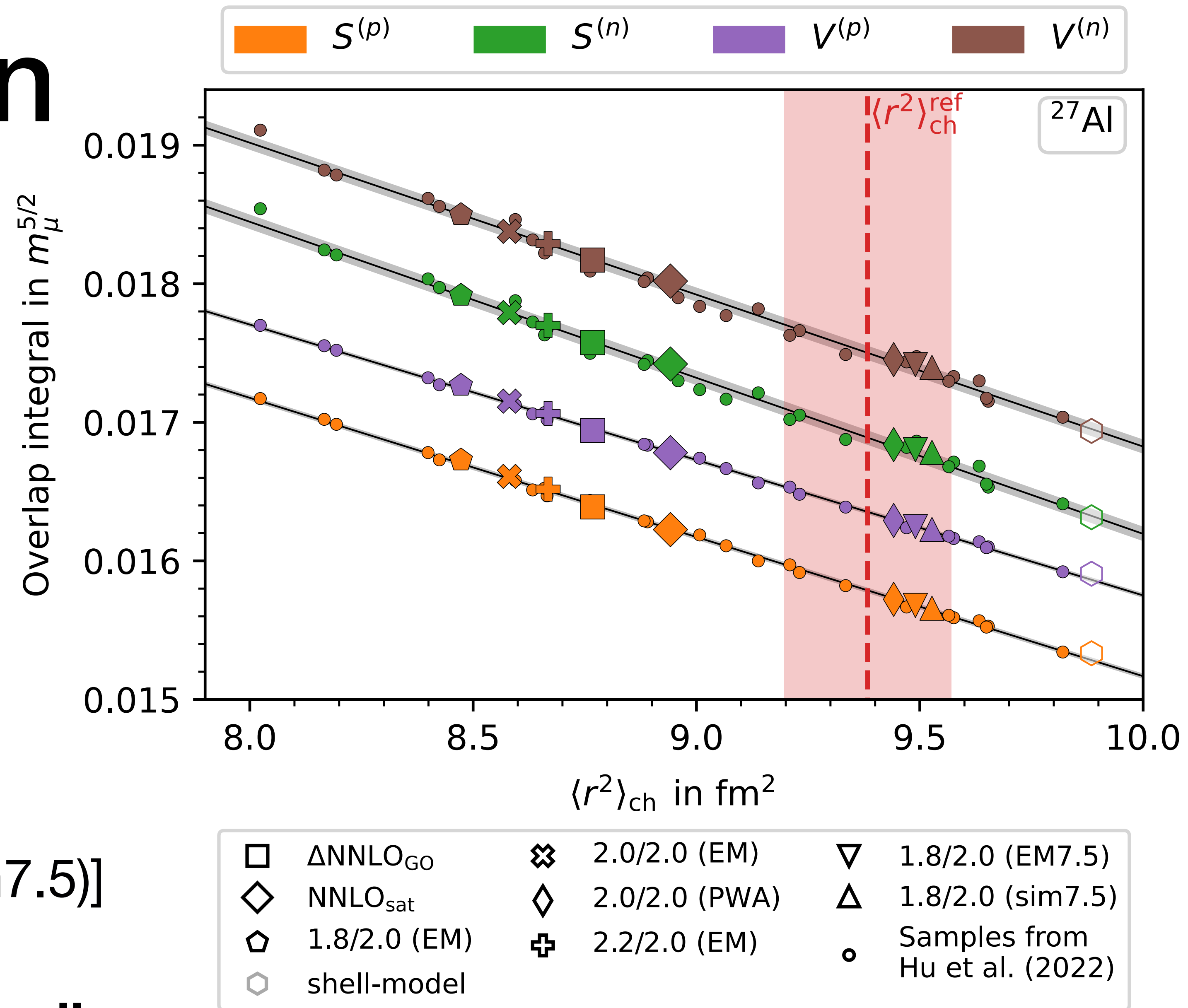
- 34 nonimplausible samples

Hu et al., Nat. Phys. **18** (2022)

- Refit “magic” [1.8/2.0 (EM7.5/sim7.5)]

Arthuis et al., arXiv:2401.06675

Tight correlations with R_{ch}^2 observed!



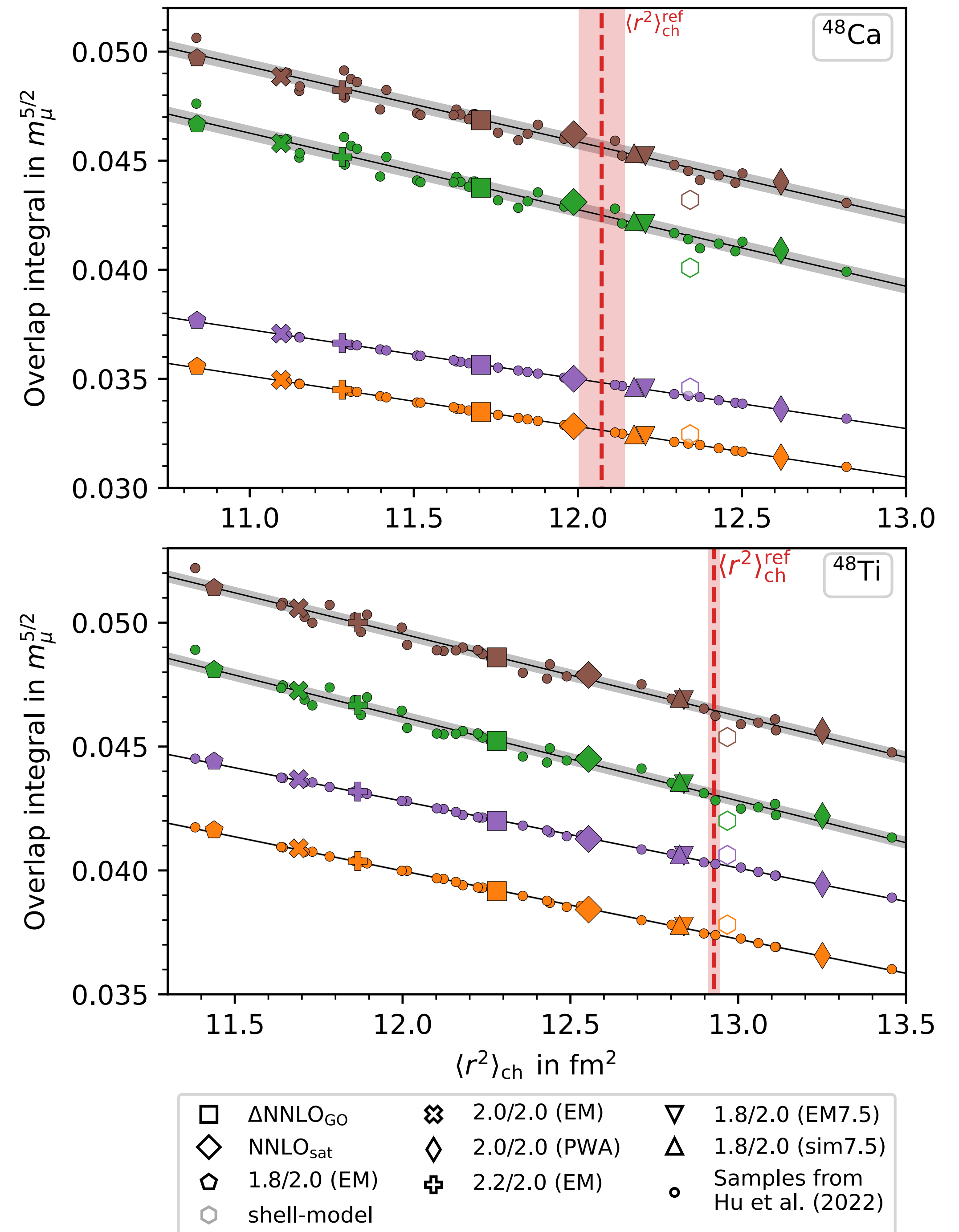
MH, Hoferichter, Miyagi, Noël, Schwenk, arXiv:2412.04545

The correlation

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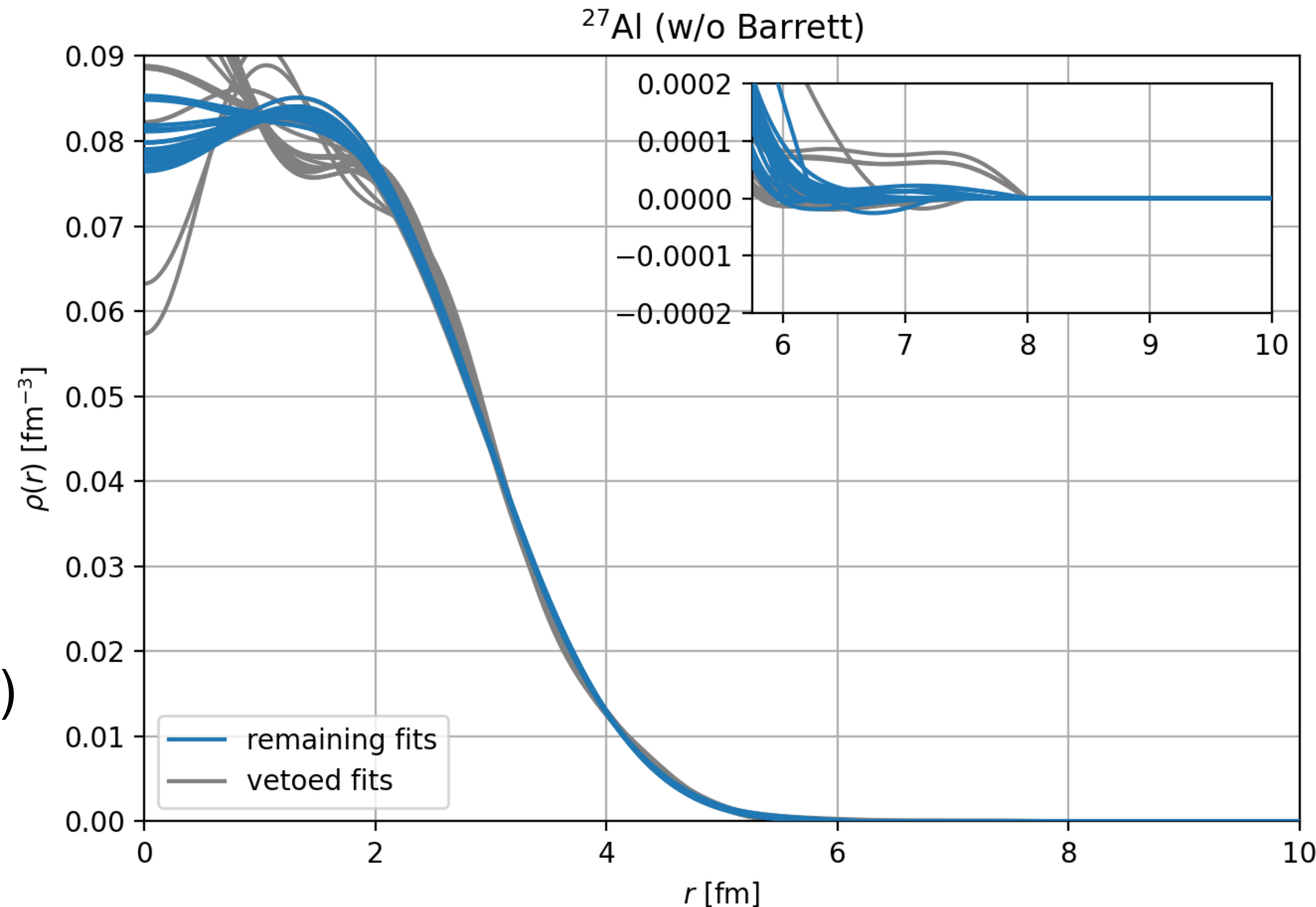
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Charge distribution uncertainties

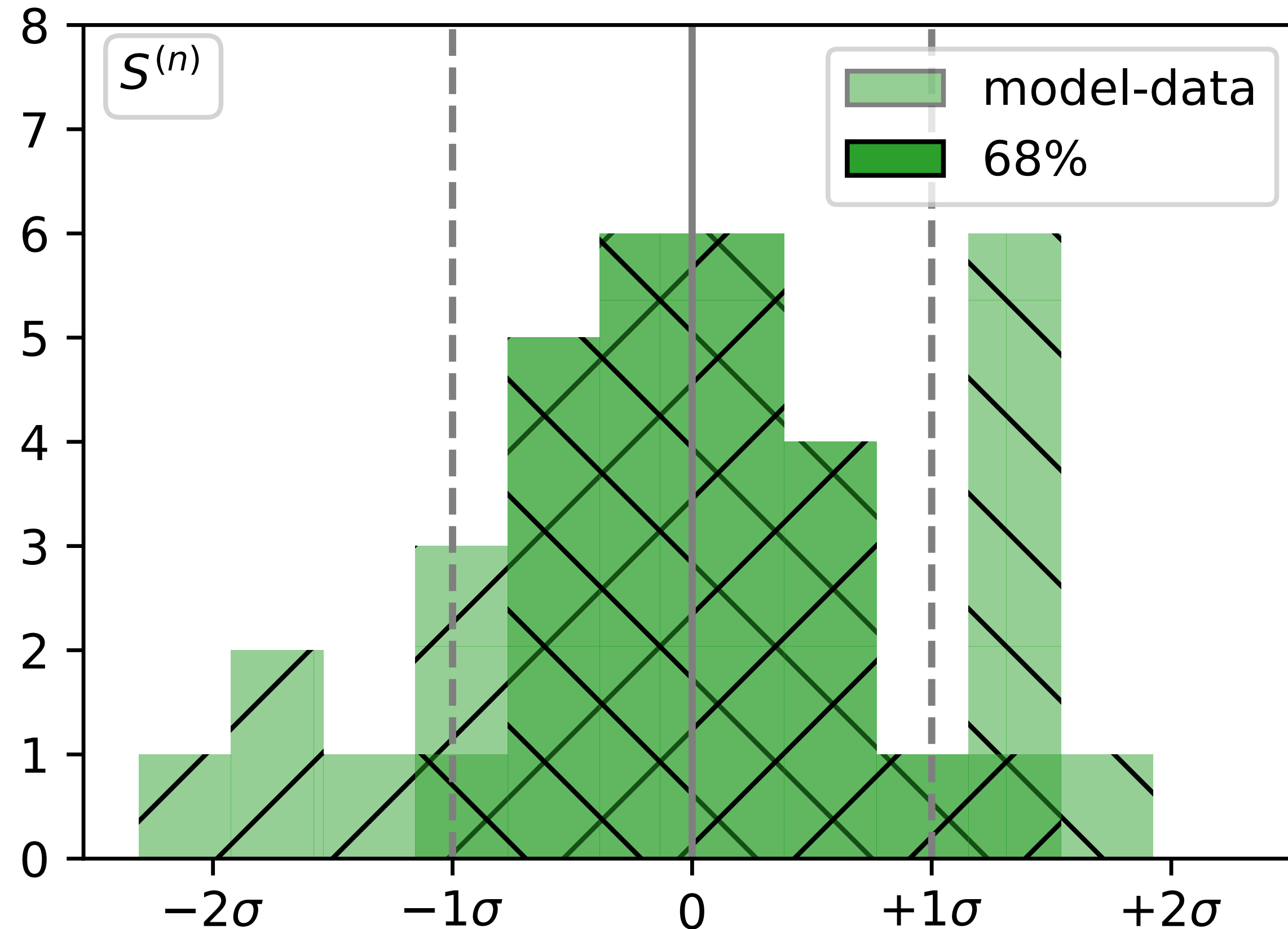
- Re-analysis of electron-scattering data
- Improved evaluation of uncertainties (more conservative for ^{27}Al)
- Goes into uncertainty for R_{ch}^2 and D (overlap integral)



Nuclear structure uncertainties

^{27}Al

Remaining scatter around correlation

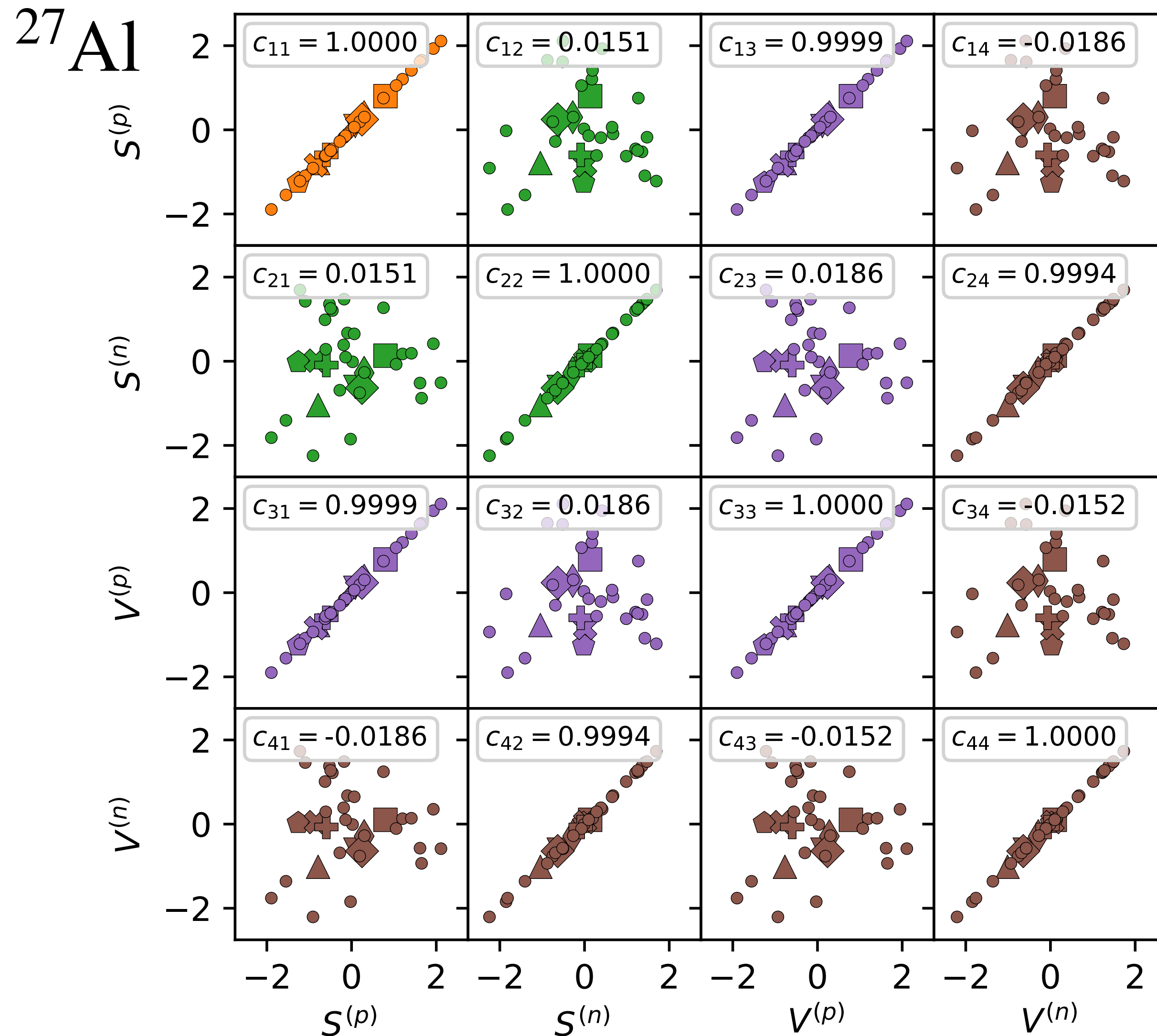


- Correlation $\rightarrow$ simple linear fit

$$S_{i,\text{corr}}^{(n)} = a (R_{\text{ch},i}^2 - R_{\text{ch,ref}}^2) + b$$

- **How to assess correlation uncertainties?**
- Residuals $S_i^{(n)} - S_{i,\text{corr}}^{(n)}$ approximately normally distributed
- Include variance σ^2 in b
- Can also consider **covariance** between overlap integrals

Nuclear structure uncertainties



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Ab initio overlap integral predictions

Predictions + uncertainties

	Our result	Kitano et al.
D	0.0359(2)	0.0362
$S^{(p)}$	0.01579(2)(19)	0.0155
$^{27}\text{Al } S^{(n)}$	0.01689(5)(21)	0.0167
$V^{(p)}$	0.01635(2)(18)	0.0161
$V^{(n)}$	0.01750(5)(21)	0.0173

Covariance matrix

	^{27}Al				
	D	$S^{(p)}$	$S^{(n)}$	$V^{(p)}$	$V^{(n)}$
D	1.0000	0.7205	0.7030	0.7210	0.7028
$S^{(p)}$		1.0000	0.9656	1.0000	0.9645
$S^{(n)}$			1.0000	0.9664	1.0000
$V^{(p)}$				1.0000	0.9654
$V^{(n)}$					1.0000

- ^{27}Al : Uncertainties from e scattering dominate $\rightarrow$ strong correlations among overlap integrals
- ^{48}Ti : Correlation analysis uncertainties dominate for neutrons

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Ab initio overlap integral predictions

Predictions + uncertainties

	Our result	Kitano et al.
D	0.08640(11)	0.0864
$S^{(p)}$	0.03742(05)(5)	0.0368
$^{48}\text{Ti } S^{(n)}$	0.04305(25)(6)	0.0435
$V^{(p)}$	0.04029(04)(5)	0.0396
$V^{(n)}$	0.04646(24)(5)	0.0468

Covariance matrix

	^{48}Ti				
	D	$S^{(p)}$	$S^{(n)}$	$V^{(p)}$	$V^{(n)}$
D	1.0000	0.4657	0.1169	0.5003	0.1163
$S^{(p)}$		1.0000	0.1118	0.9991	0.0916
$S^{(n)}$			1.0000	0.1176	0.9997
$V^{(p)}$				1.0000	0.0978
$V^{(n)}$					1.0000

- ^{27}Al : Uncertainties from e scattering dominate $\rightarrow$ strong correlations among overlap integrals
- ^{48}Ti : Correlation analysis uncertainties dominate for neutrons

Ab initio overlap integral predictions

Predictions + uncertainties

	Our result	Kitano et al.
D	0.08640(11)	0.0864
$S^{(p)}$	0.03742(05)(5)	0.0368
$^{48}\text{Ti } S^{(n)}$	0.04305(25)(6)	0.0435
$V^{(p)}$	0.04029(04)(5)	0.0396
$V^{(n)}$	0.04646(24)(5)	0.0468

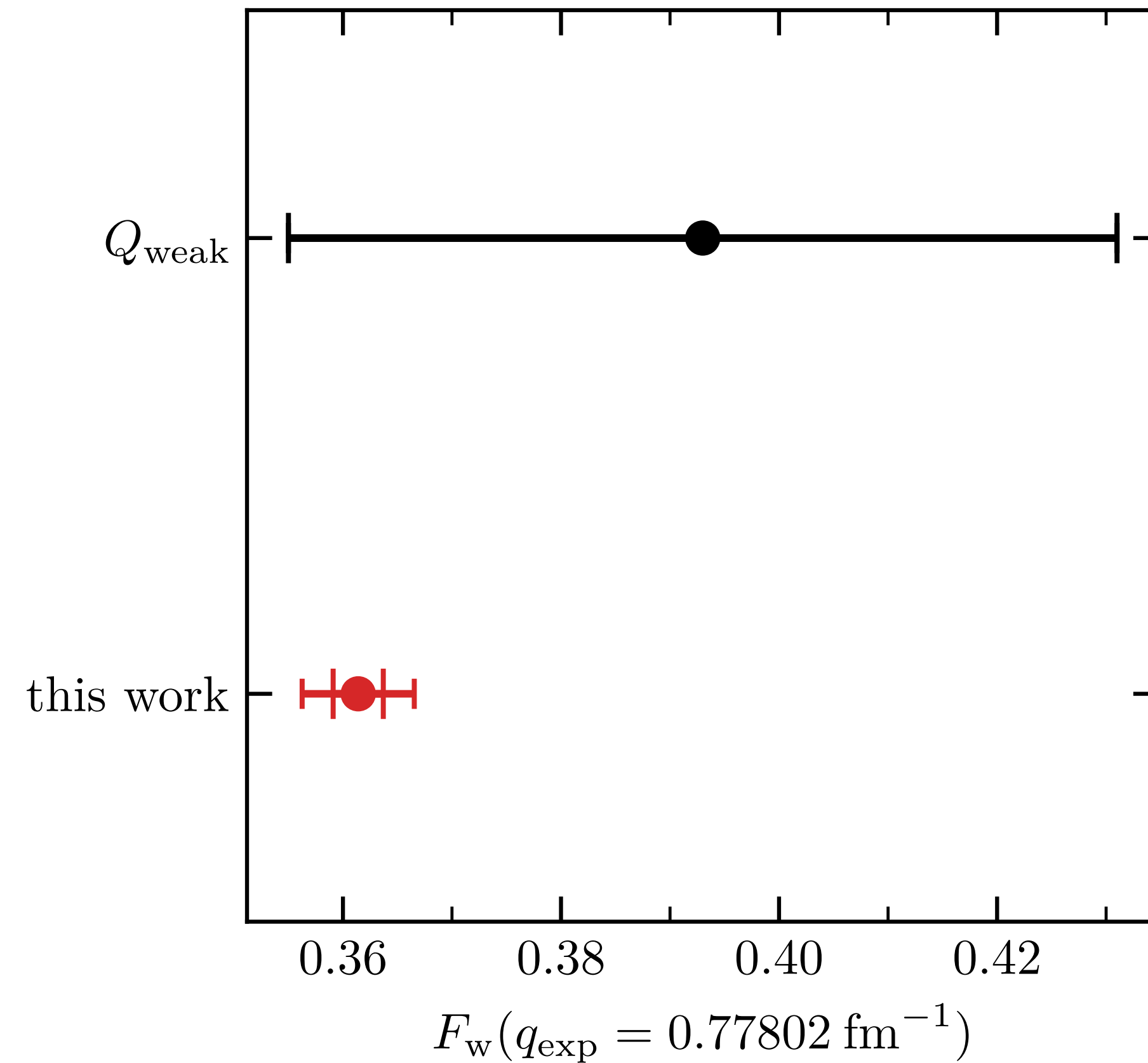
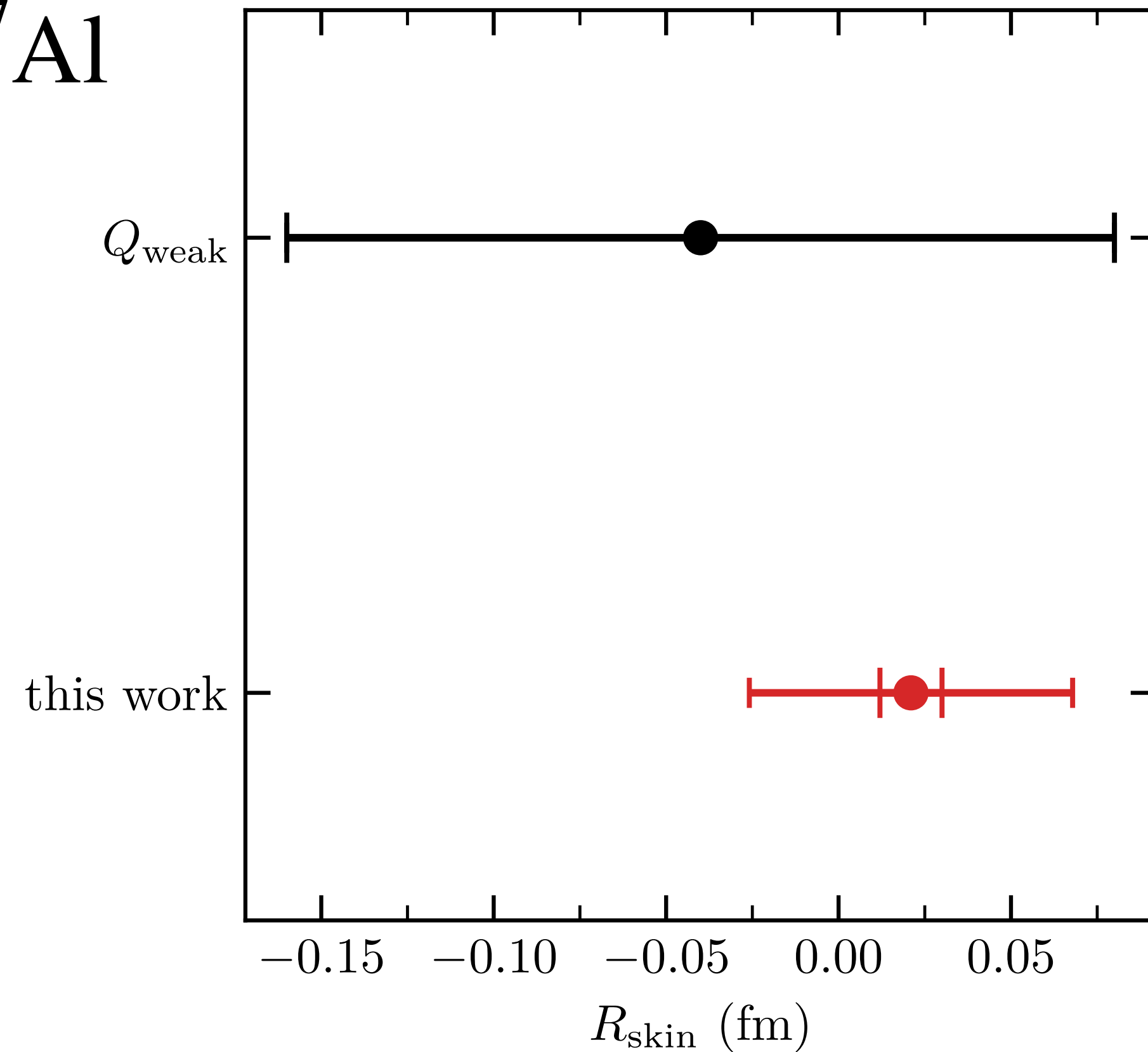
Covariance matrix

	^{48}Ti				
	D	$S^{(p)}$	$S^{(n)}$	$V^{(p)}$	$V^{(n)}$
D	1.0000	0.4657	0.1169	0.5003	0.1163
$S^{(p)}$		1.0000	0.1118	0.9991	0.0916
$S^{(n)}$			1.0000	0.1176	0.9997
$V^{(p)}$				1.0000	0.0978
$V^{(n)}$					1.0000

- ^{27}Al : Uncertainties from e scattering dominate $\rightarrow$ strong correlations among overlap integrals
- ^{48}Ti : Correlation analysis uncertainties dominate for neutrons

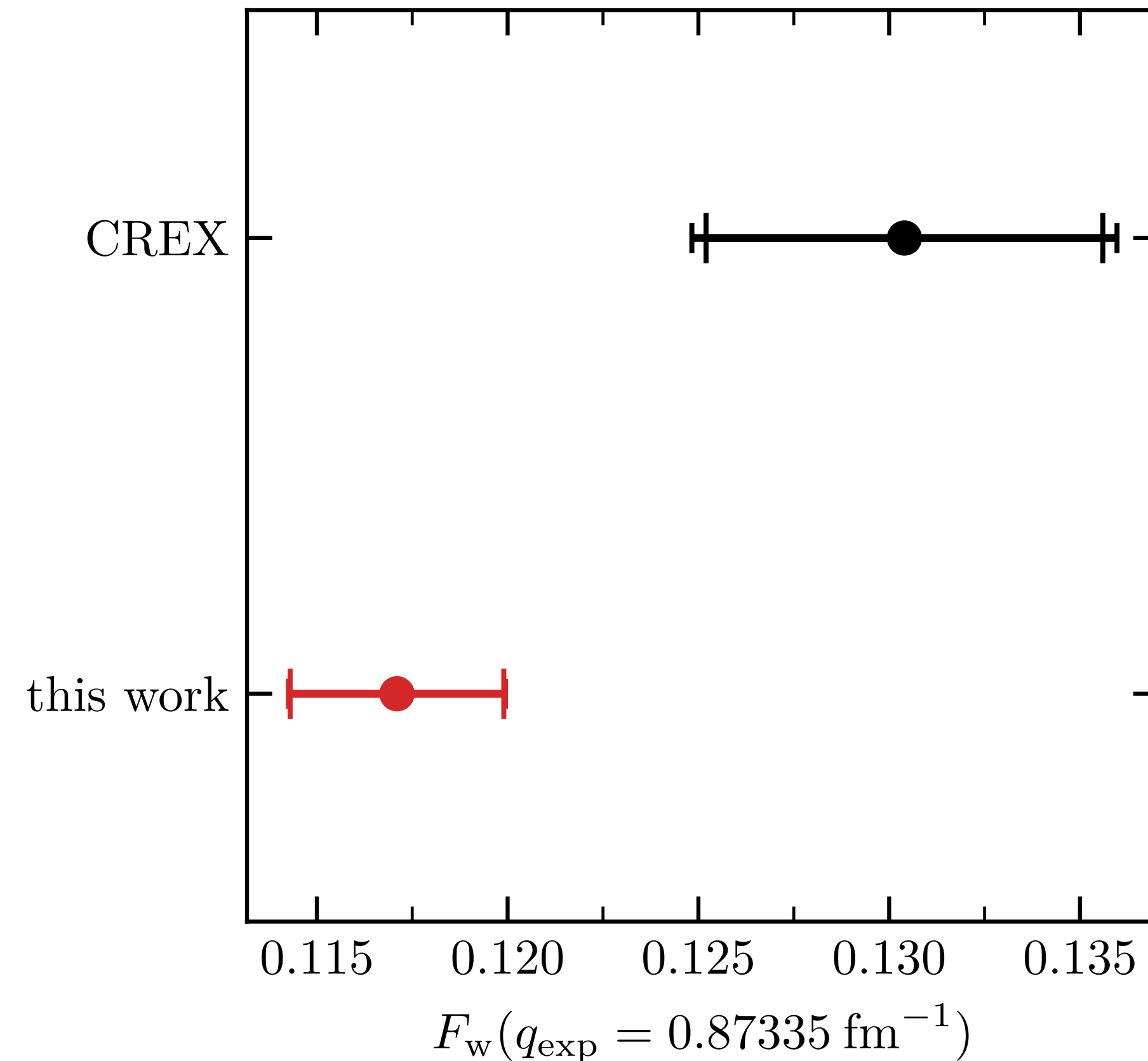
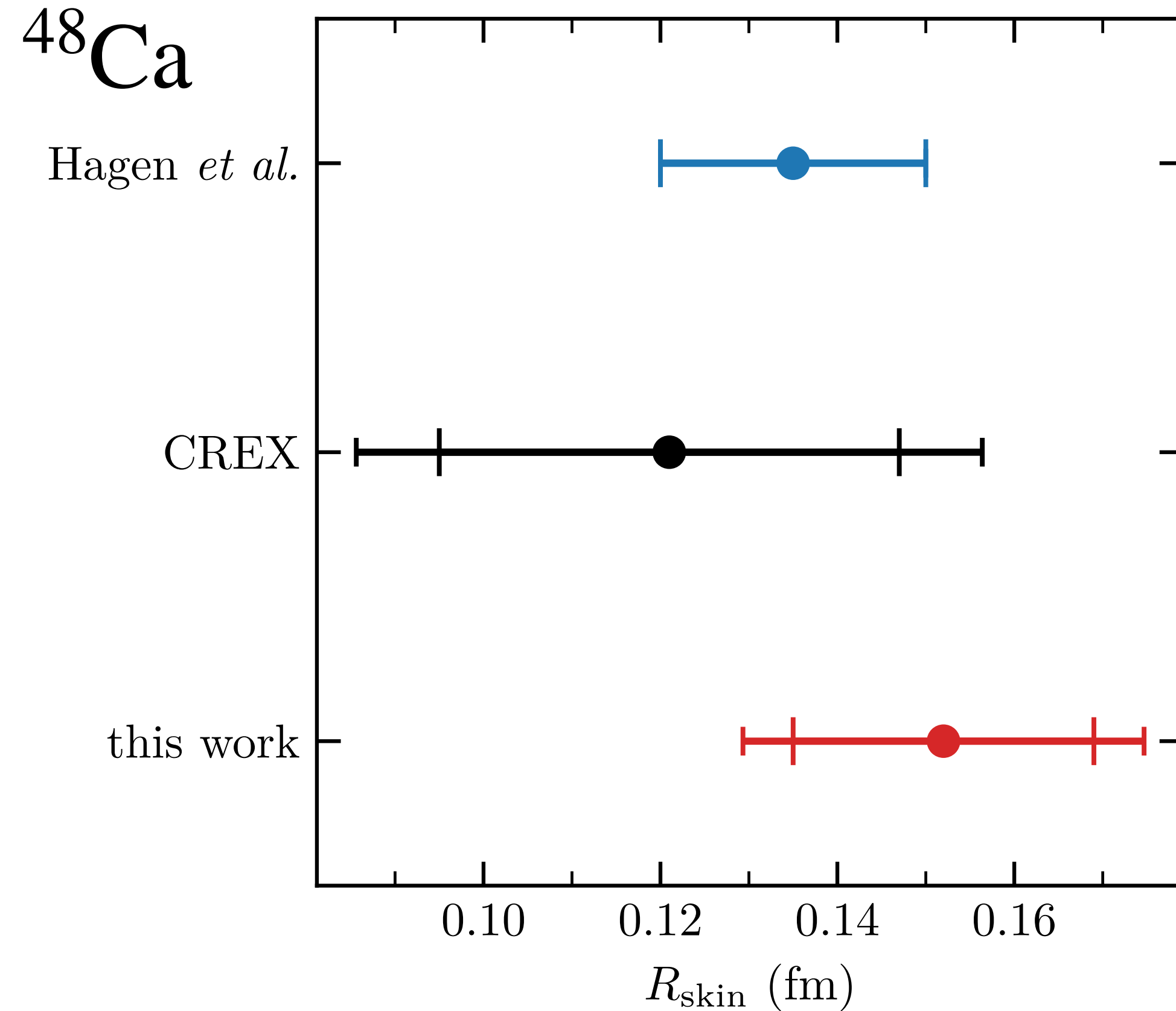
Bonus: Predictions for weak scattering

^{27}Al



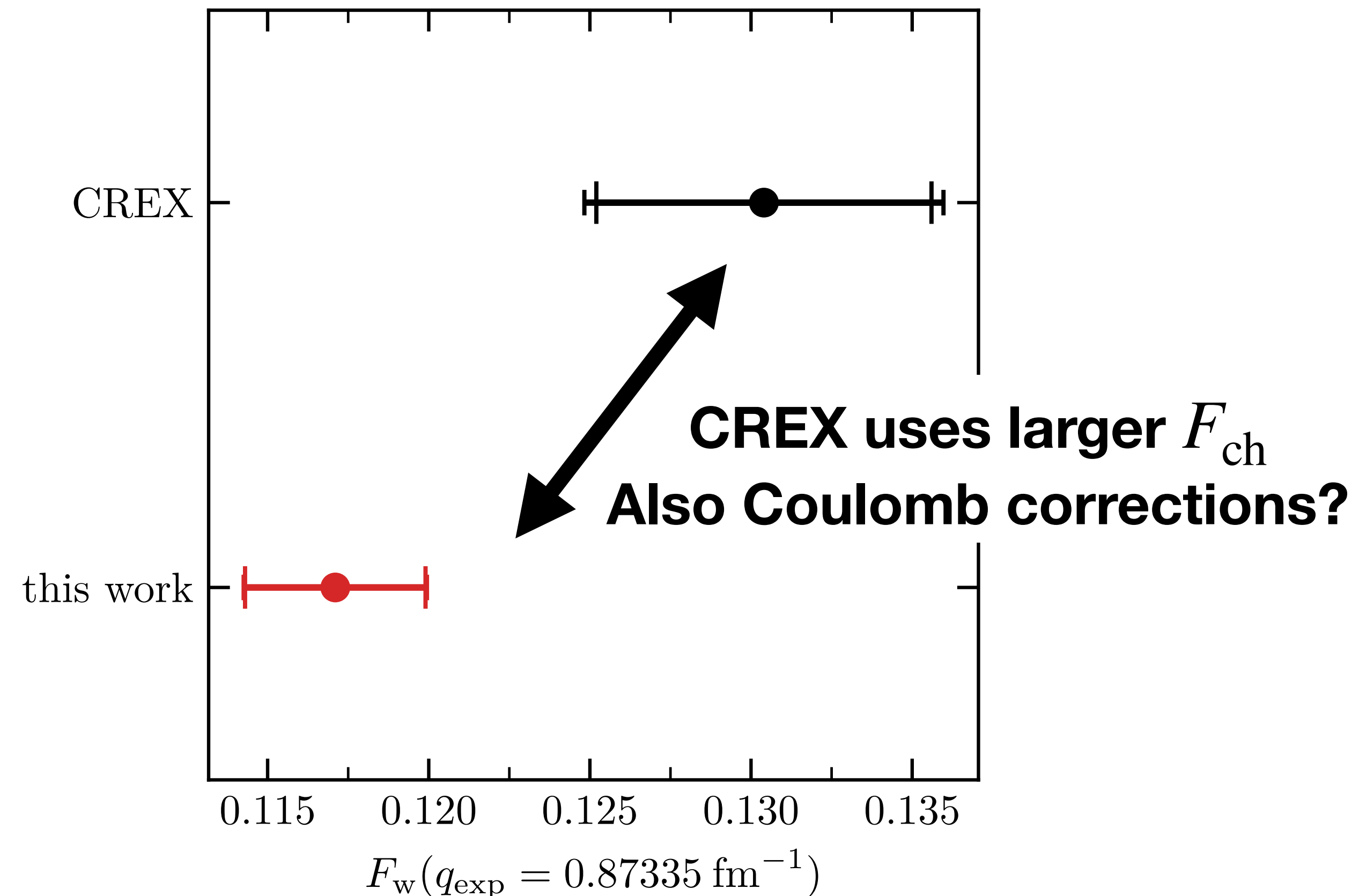
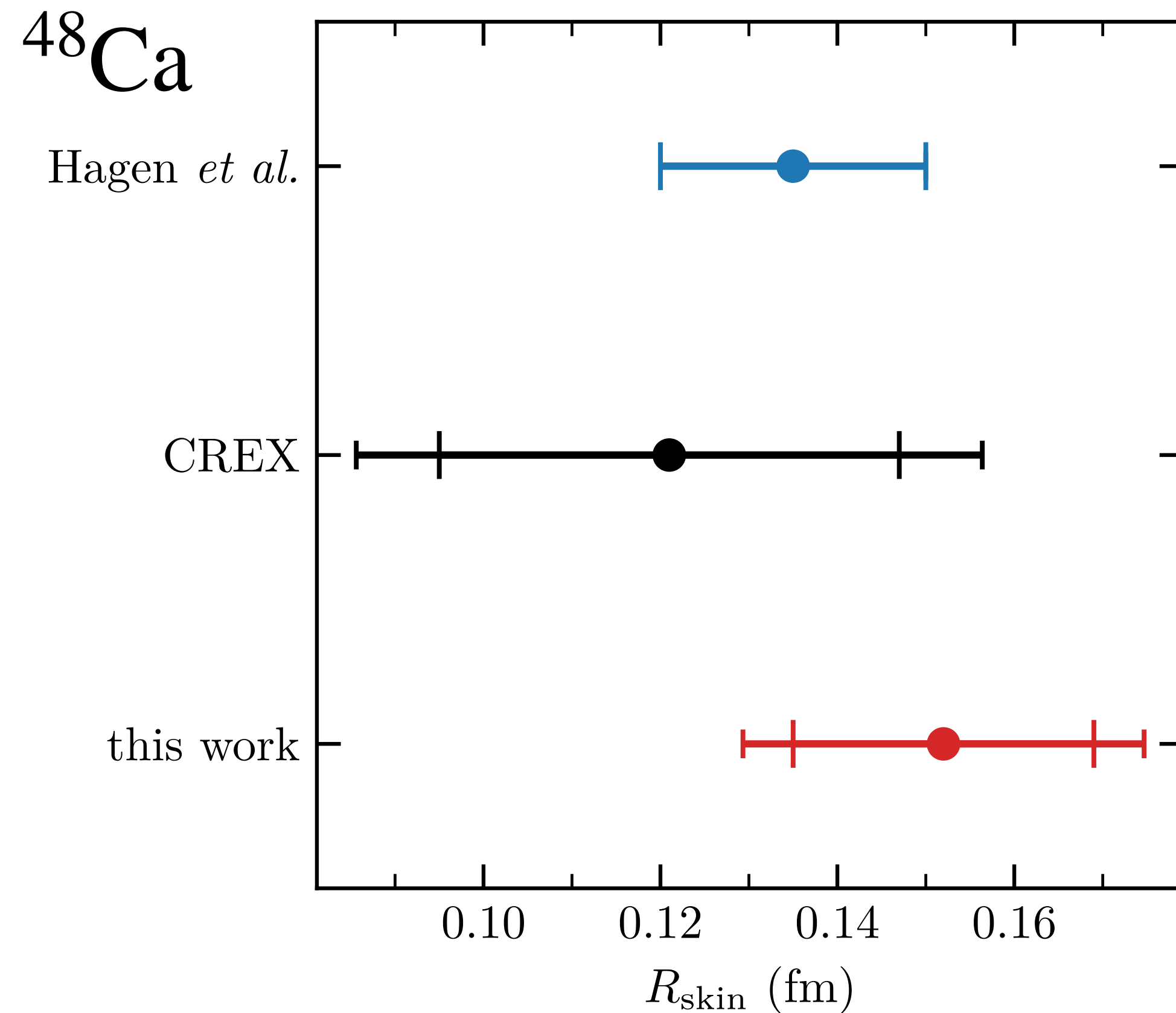
Q_{weak} , PRL **128** (2022), CREX, PRL **129** (2022), Hagen et al., Nat. Phys. **12** (2016)
MH, Hoferichter, Miyagi, Noël, Schwenk, arXiv:2412.04545

Bonus: Predictions for weak scattering

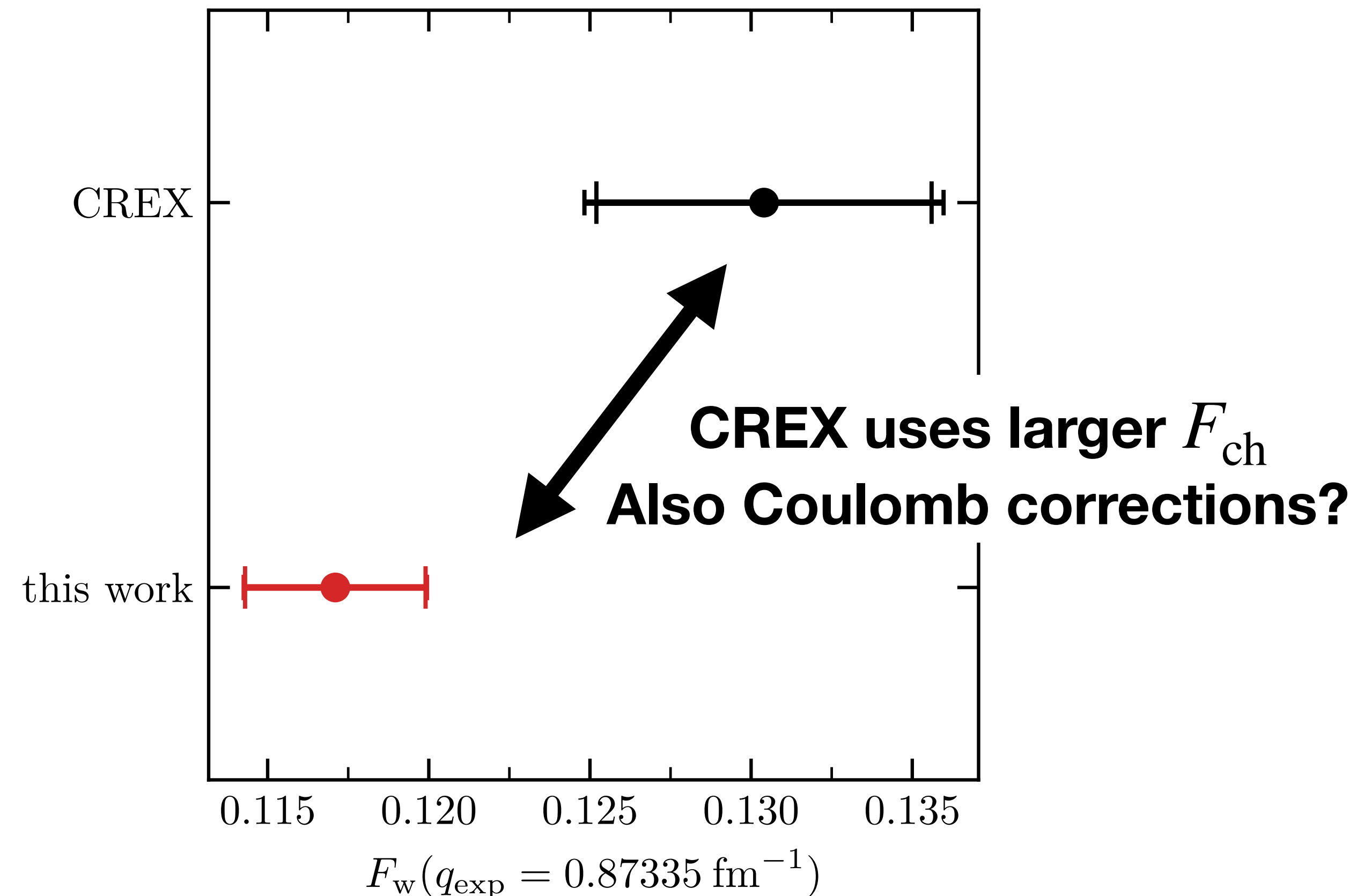
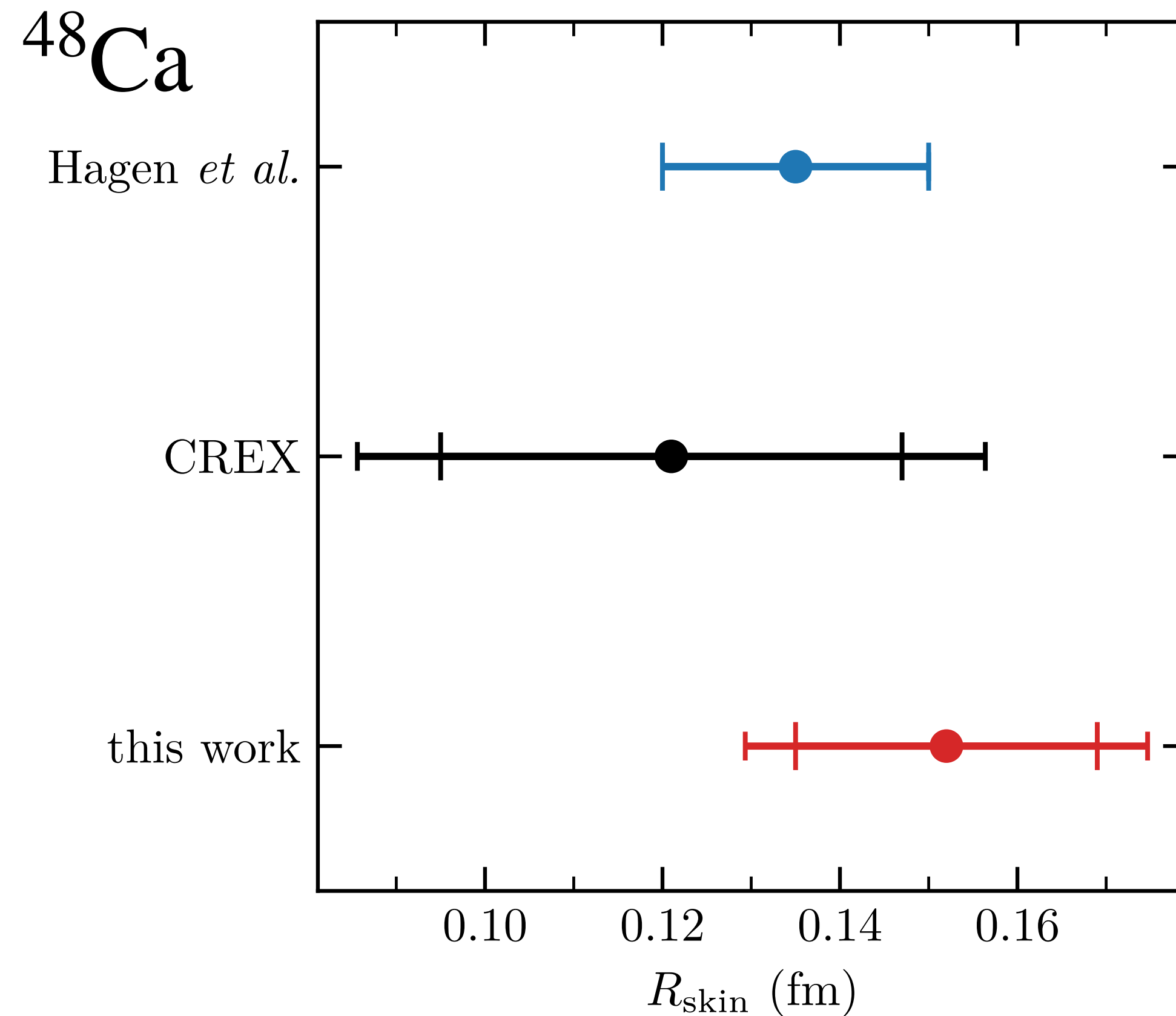


Q_{weak} , PRL **128** (2022), CREX, PRL **129** (2022), Hagen *et al.*, Nat. Phys. **12** (2016)
MH, Hoferichter, Miyagi, Noël, Schwenk, arXiv:2412.04545

Bonus: Predictions for weak scattering



Bonus: Predictions for weak scattering



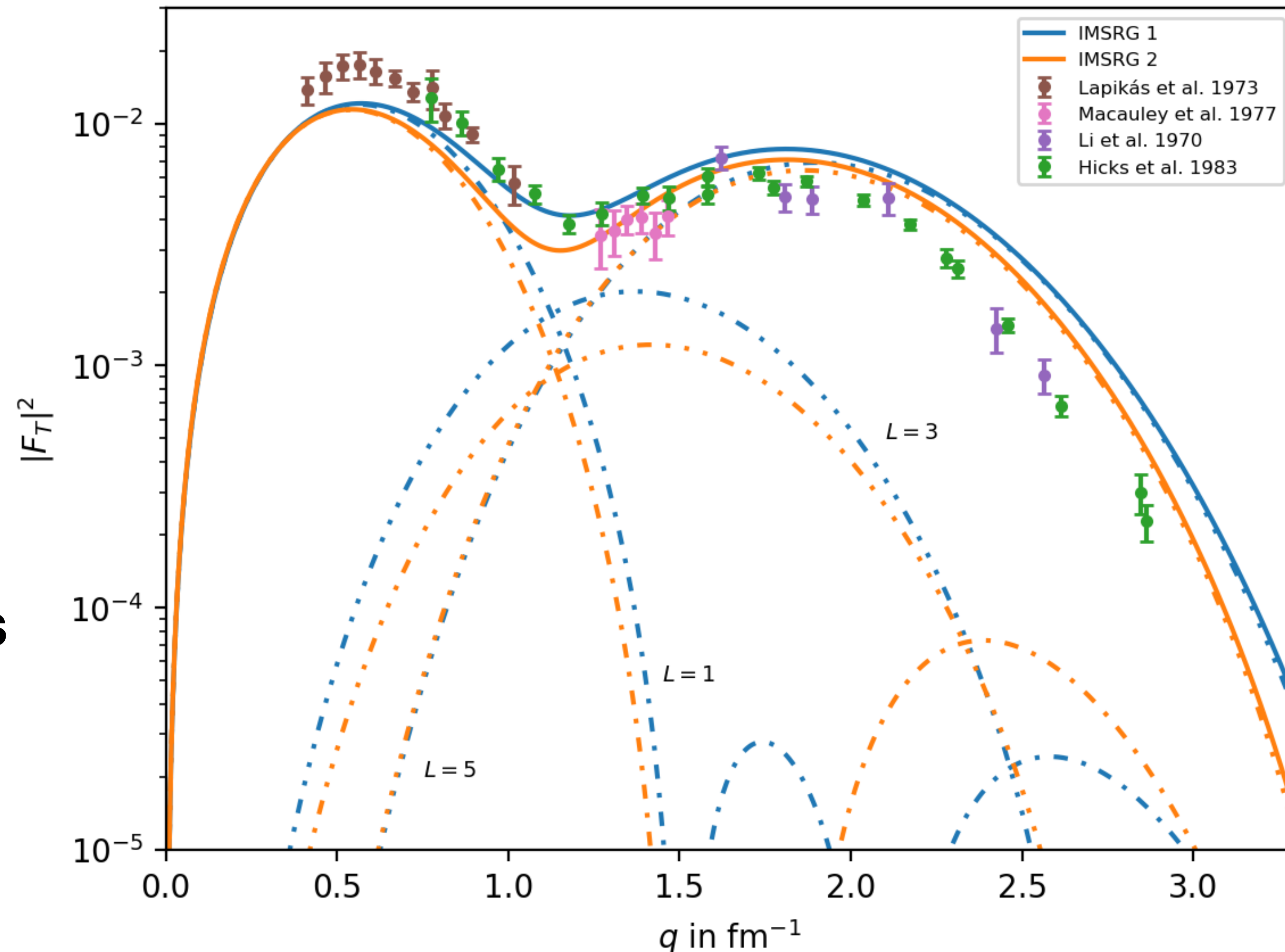
- Apparent tension with CREX partially resolved by accounting for Coulomb corrections and comparing with measured A_{PV}

Q_{weak} , PRL **128** (2022), CREX, PRL **129** (2022), Hagen *et al.*, Nat. Phys. **12** (2016)

MH, Hoferichter, Miyagi, Noël, Schwenk, arXiv:2412.04545

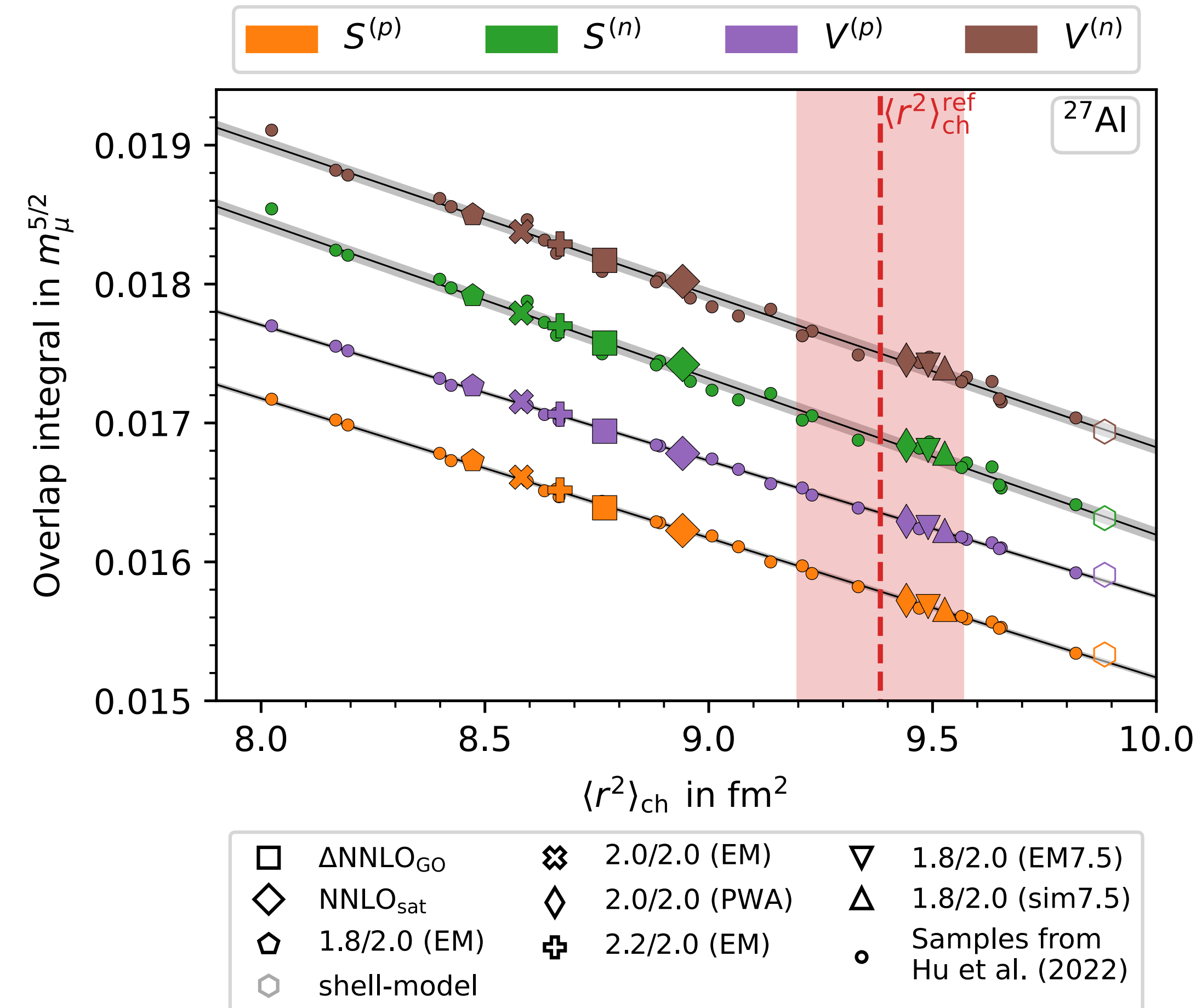
What next?

- **Incoherent spin-dependent** responses also contribute
 - Naively subleading contribution, but may be important for some microscopic theories
- **Understand impact of uncertainties** on analyses constraining Wilson coefficients of SMEFT operators



Conclusion

- Ab initio predictions of overlap integrals for spin-independent $\mu \rightarrow e$ conversion
- **Comprehensive treatment** of Hamiltonian (and many-body) uncertainties
- **Correlation analysis** accounts for correlated uncertainties
- Consistent with past work on weak scattering
- **Key input for inferences** of implications for BSM physics from $\mu \rightarrow e$ decay rate



MH, Hoferichter, Miyagi, Noël, Schwenk, arXiv:2412.04545

Acknowledgments

Coauthors:

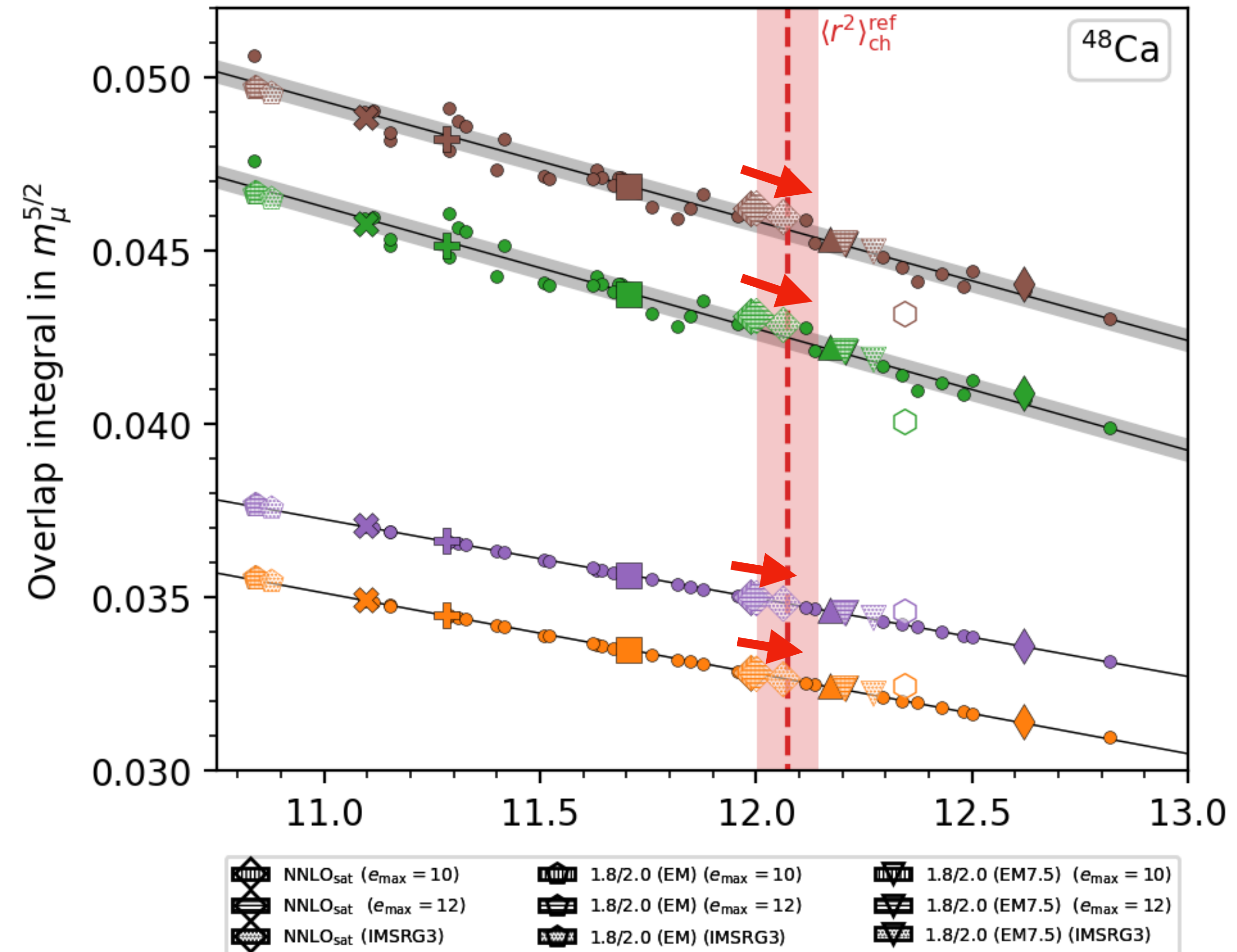
- **Uni Bern:** Frederic Noël,
Martin Hoferichter
- **TU Darmstadt:** Achim Schwenk
- **Univ. of Tsukuba:** Takayuki Miyagi



Thank you for your attention!

What about many-body uncertainties?

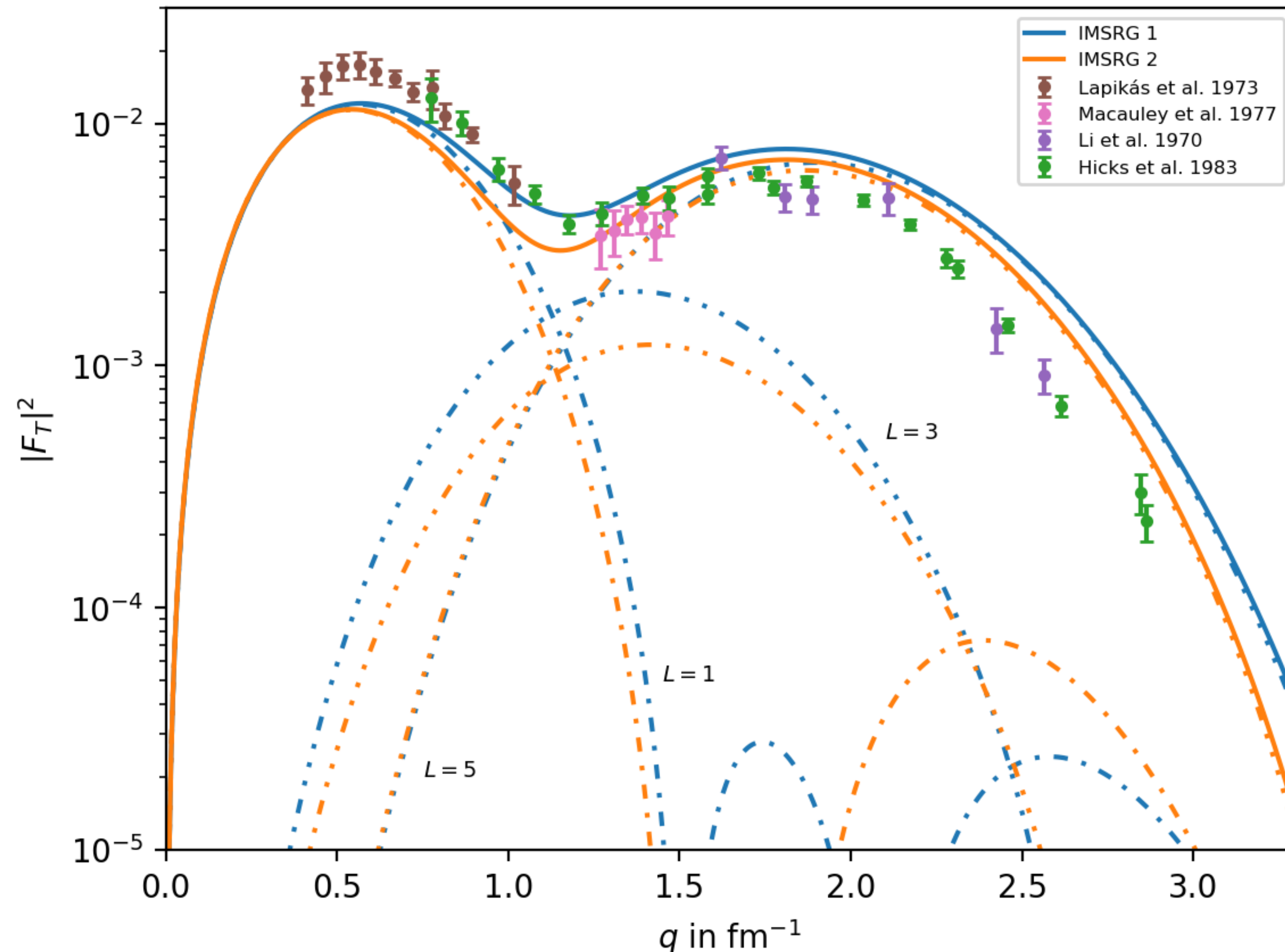
- Adjusted both model-space and many-body truncations
- Resulting points still lie on the correlation
- Correlation encompasses both diverse Hamiltonian uncertainties and many-body uncertainties



IMSRG(3) = IMSRG(3)- N^7 with restricted 3B operators

Can we resolve subleading responses?

- Incoherent spin-dependent responses also contribute
- Challenges:
 - Many more responses to evaluate: $M, \Phi'', \Delta, \Sigma', \Sigma''$ for $J \leq 5$
 - Overlap integral formalism not yet established
 - **Which observable to correlate with?**
 - **Can we actually resolve this given current uncertainties?**
 - **2BCs might be relevant?**



Many-body calculation uncertainties

