

# Latest Neutrino Oscillation Results from the NOvA Experiment

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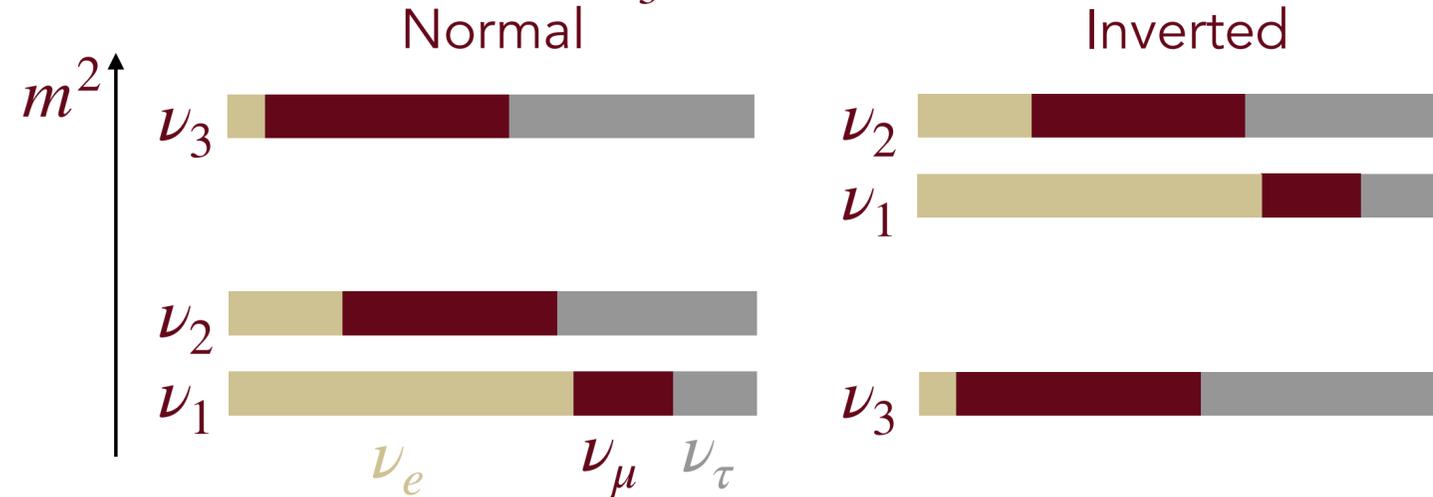


# Neutrino oscillations in NOvA

## 3-flavor Oscillations

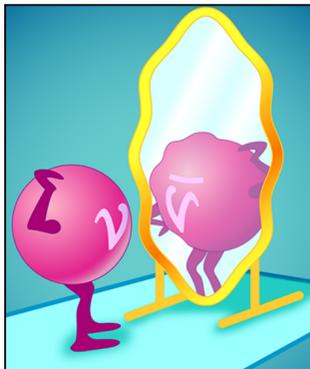
### Mass Ordering

Is  $\nu_3$  the heaviest state?



### CP Violation

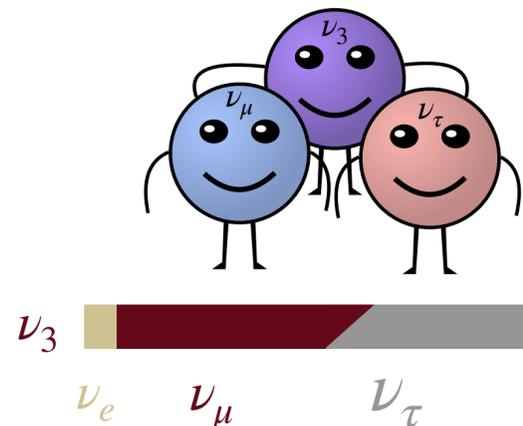
Is there a neutrino-antineutrino asymmetry?



Credit: APS/Carin Cain

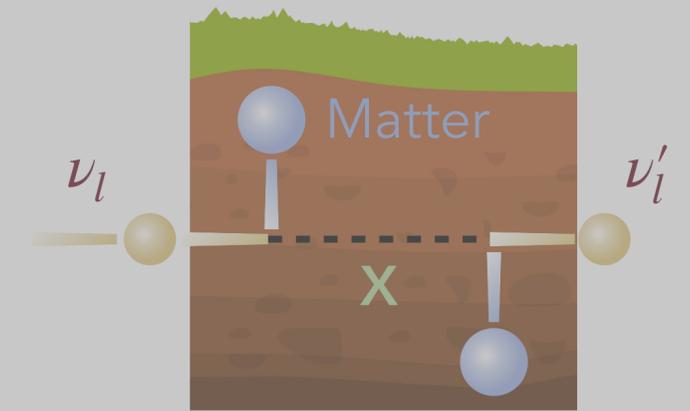
### Maximal Mixing ( $\theta_{23}$ )

Do the  $\nu_\mu$  and  $\nu_\tau$  states couple equally to  $\nu_3$ ?



## Beyond Standard Oscillations

Nonstandard interactions

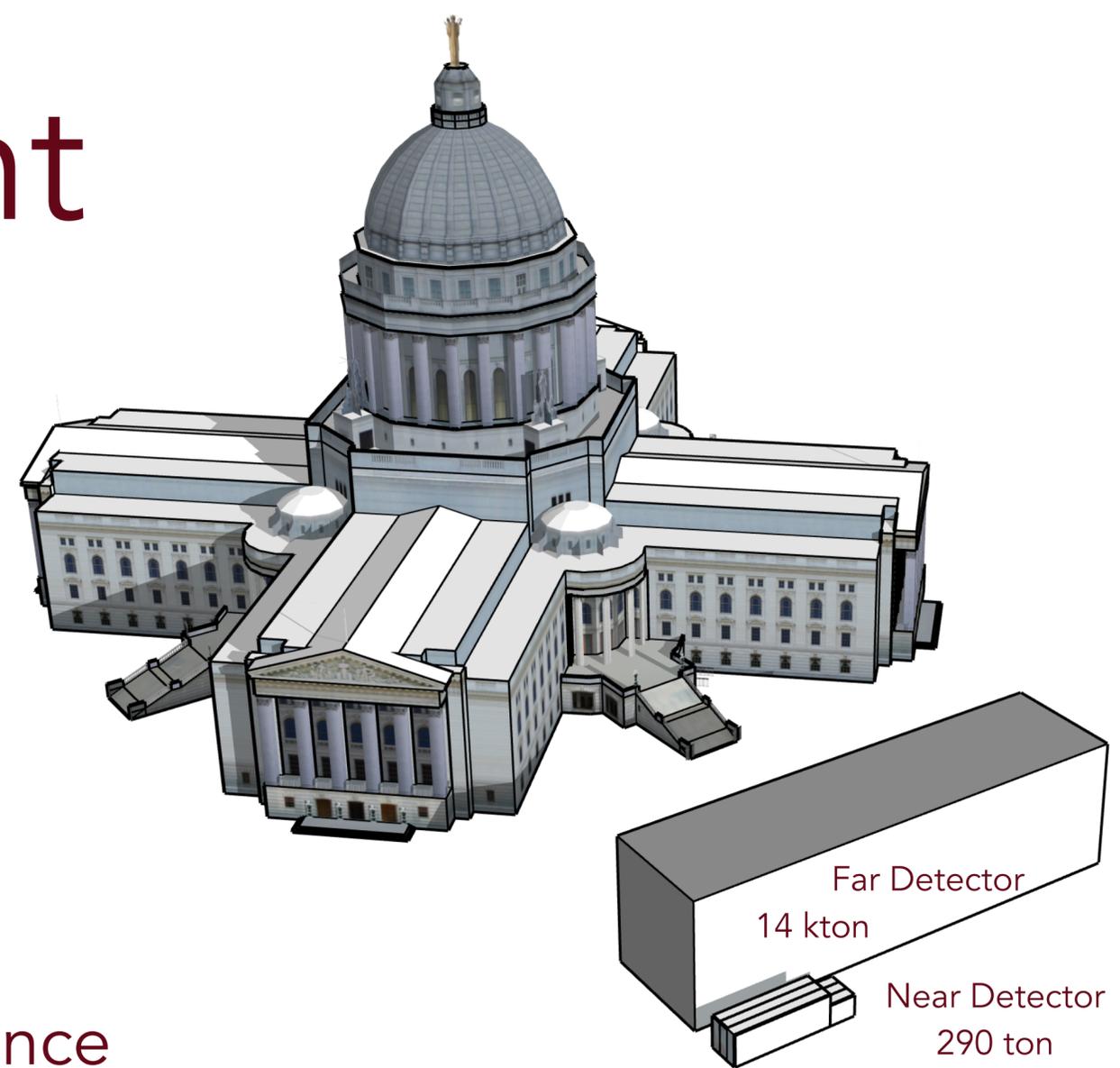
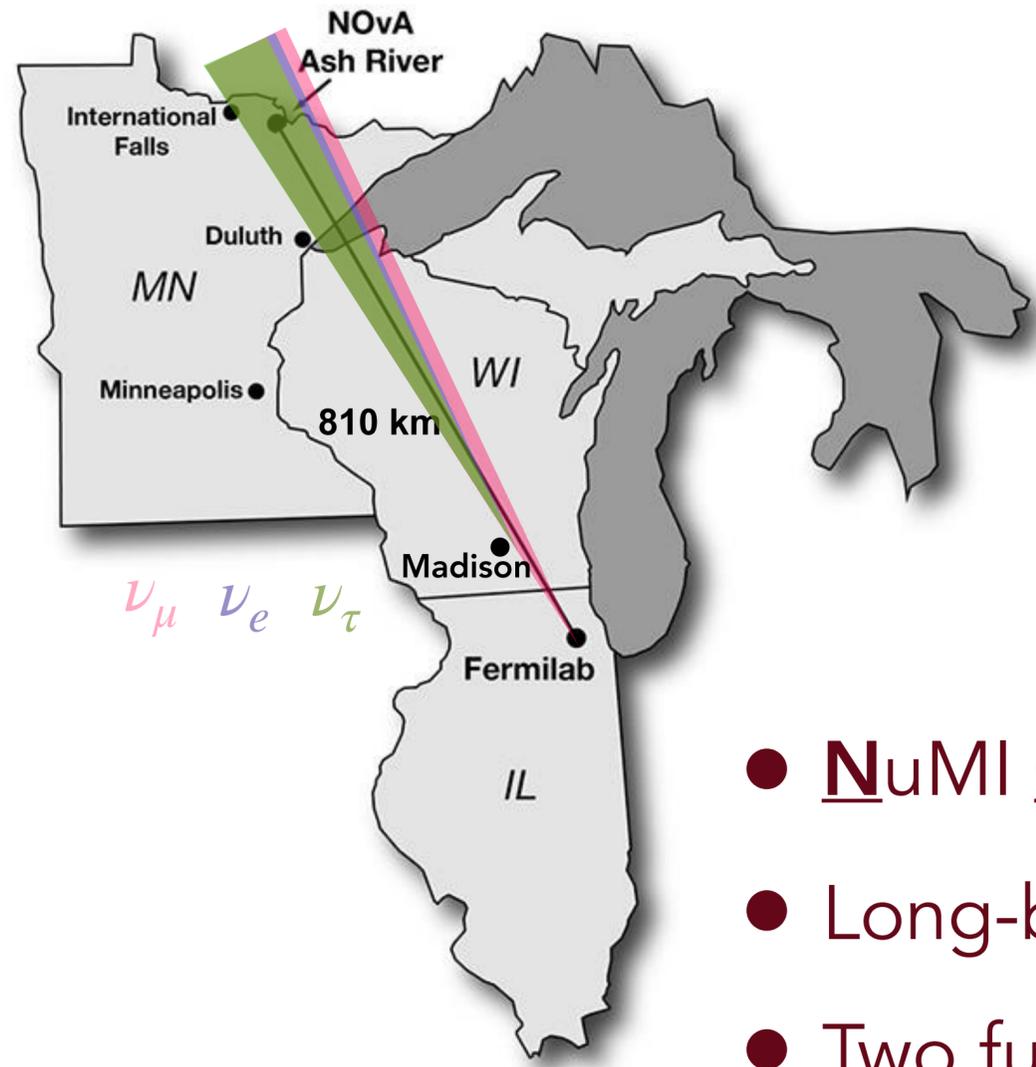


Sterile neutrinos



Credit: Symmetry magazine, Illustration by Sandbox Studio, Chicago with Ana Kova

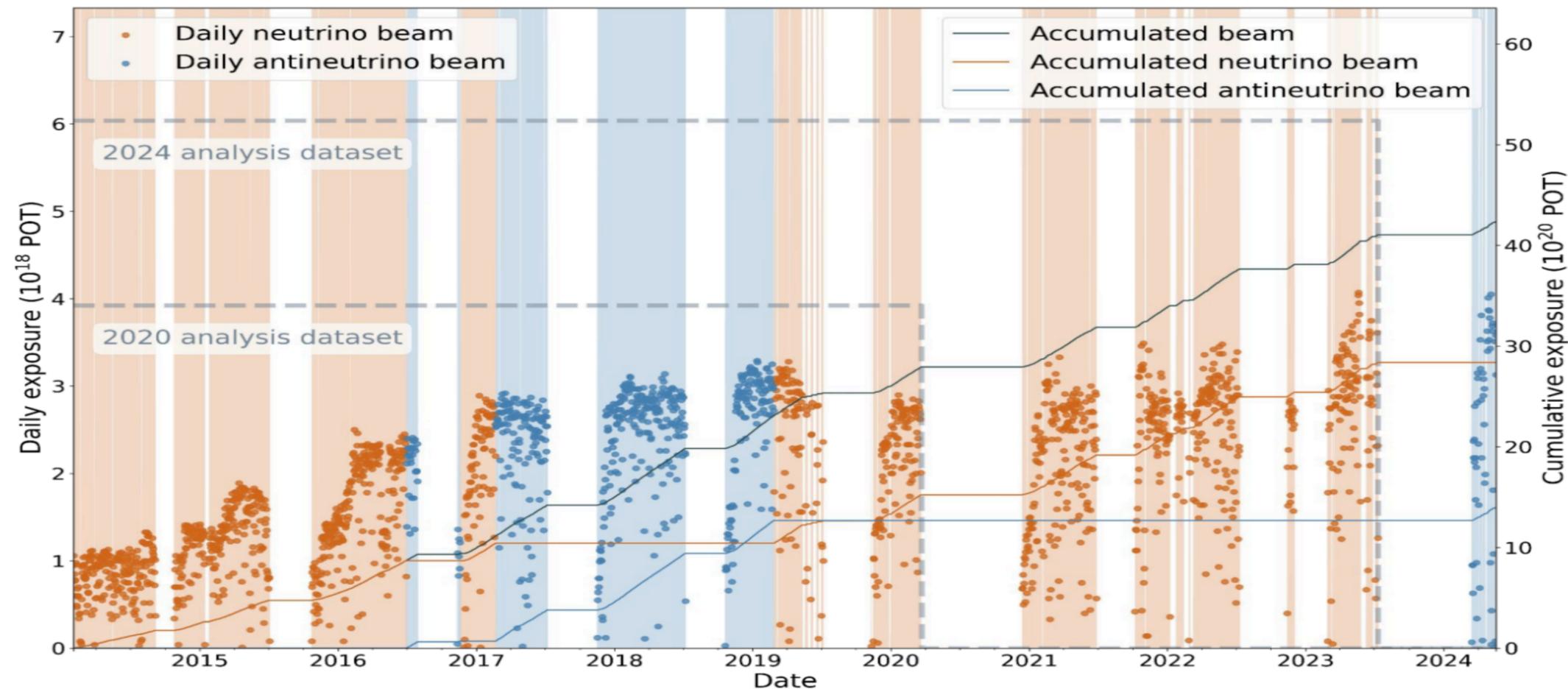
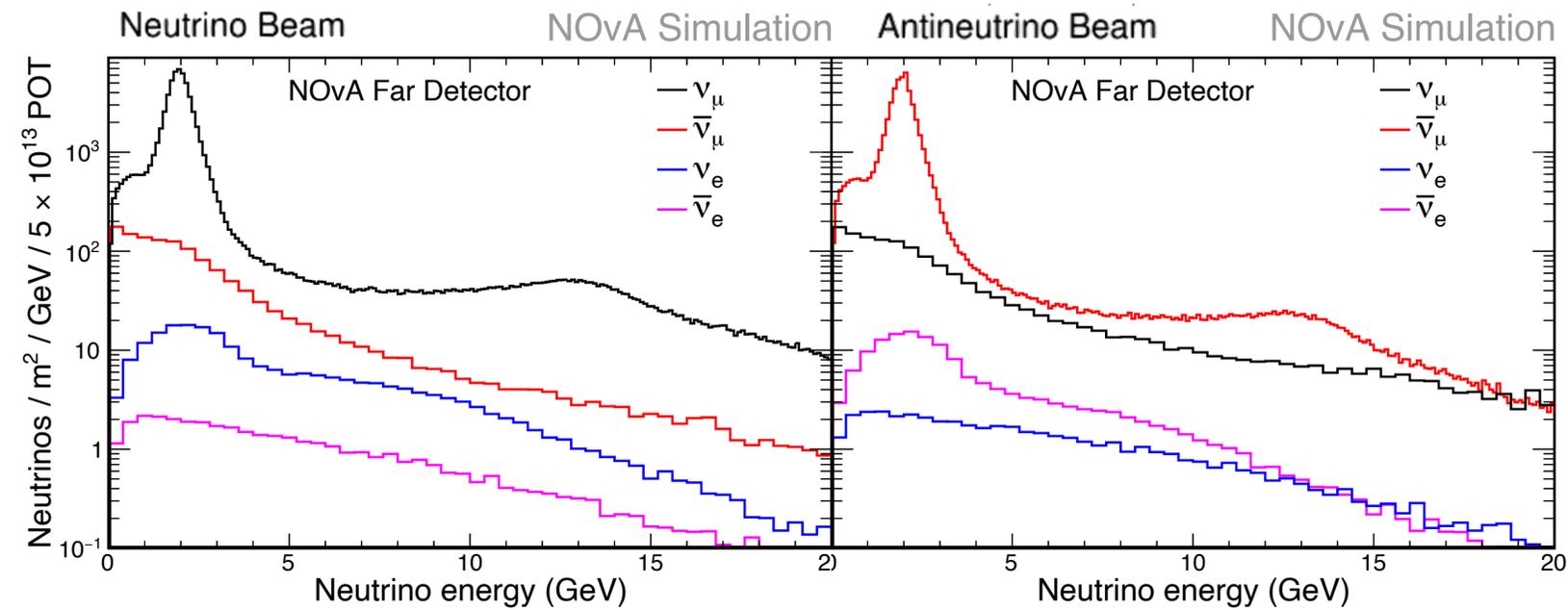
# The NOvA experiment



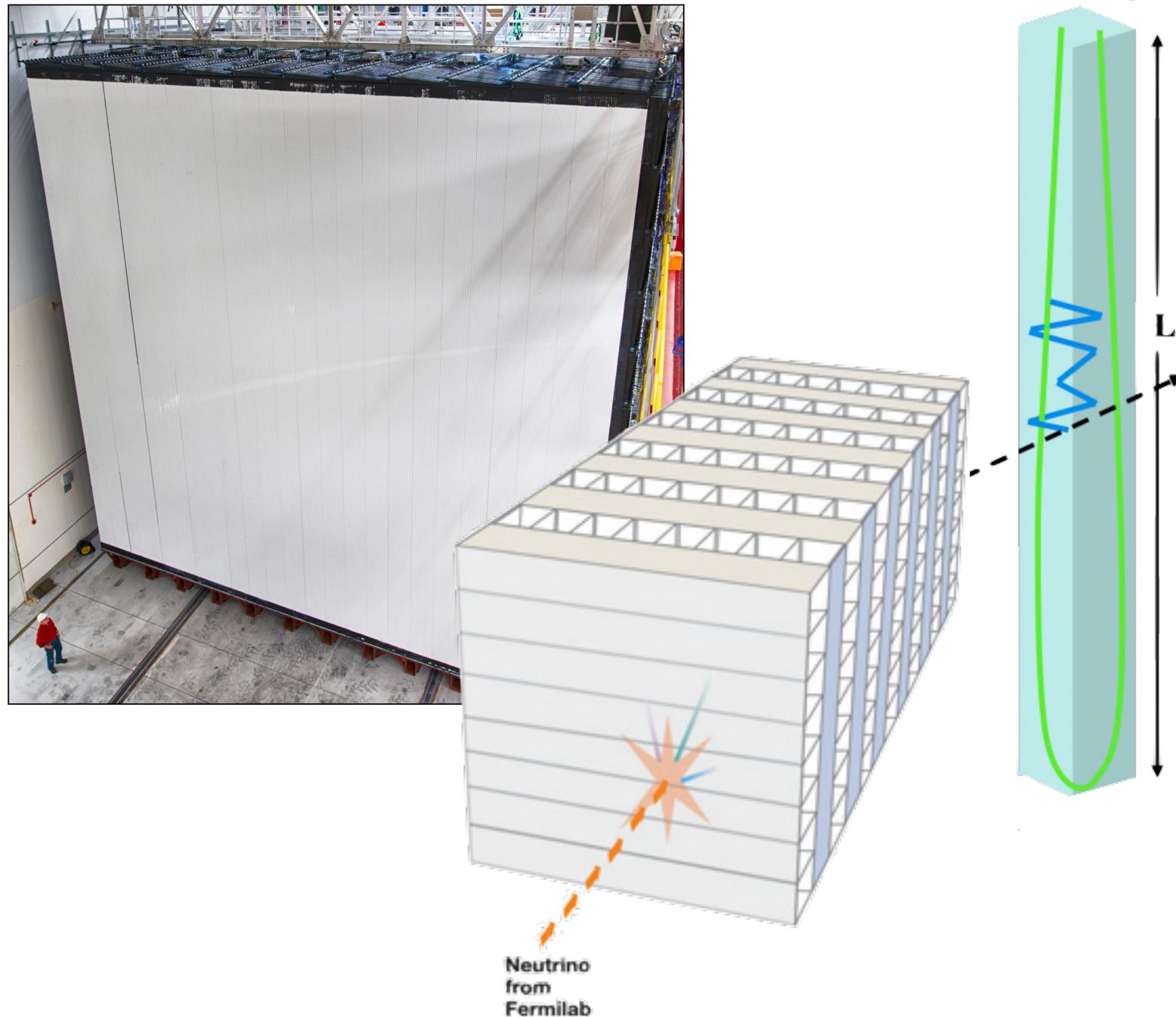
- **NuMI Off-Axis  $\nu_e$  Appearance**
- Long-baseline accelerator neutrino experiment based at Fermilab
- Two functionally equivalent tracking calorimeters in a narrow-band 2 GeV  $\nu_\mu$  or  $\bar{\nu}_\mu$  beam

# The NuMI beam

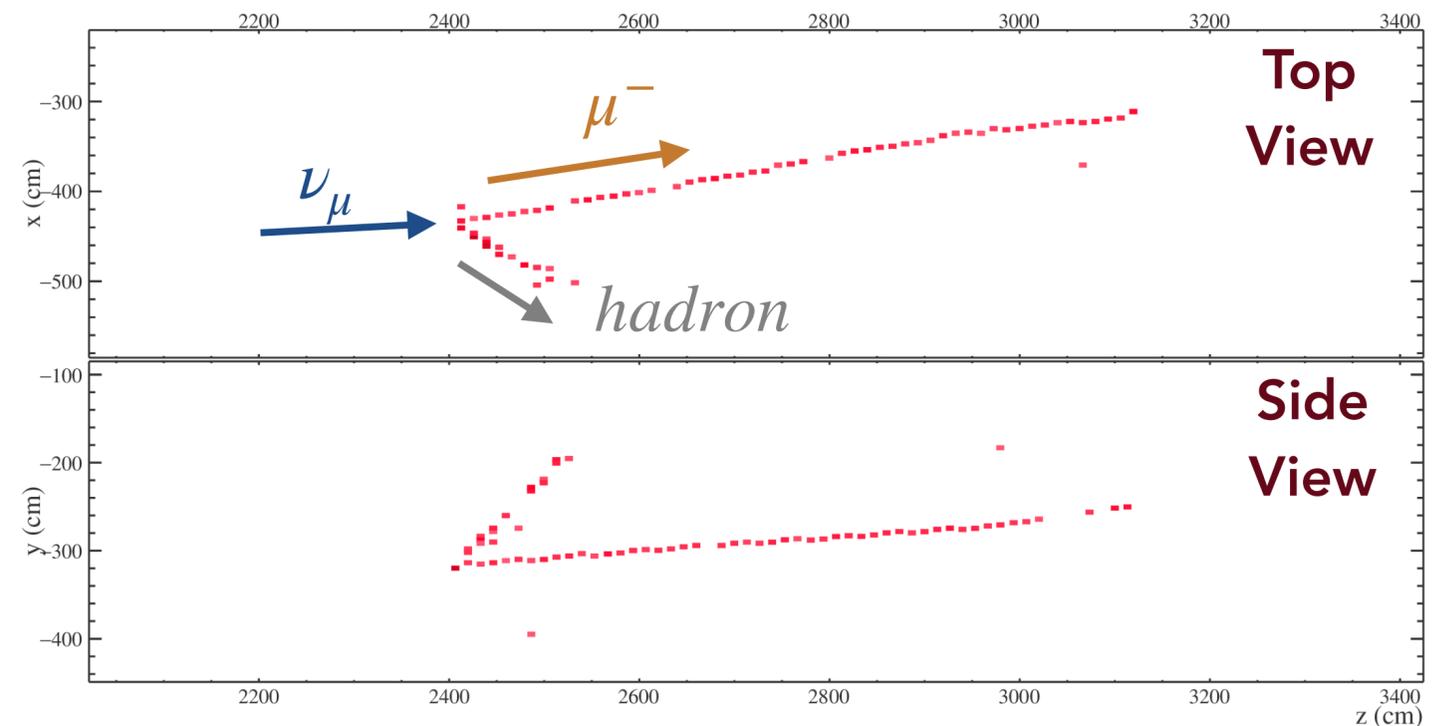
- $39.1 \times 10^{20}$  analyzed POT so far
  - 10 years of running
- Typical beam power of  $\sim 900$  kW
  - Record powers above 1 MW
- Selectable  $\nu_\mu$  or  $\bar{\nu}_\mu$  beam
  - 93-95% purity



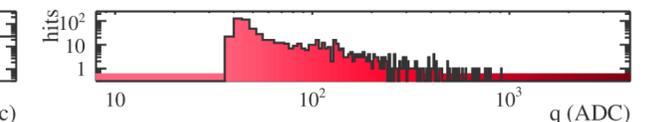
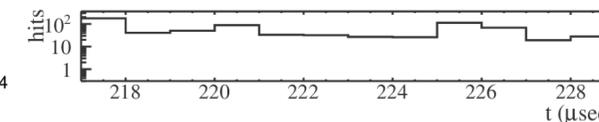
# The NOvA detectors



- Extruded PVC cells filled with scintillating oil
- Looped wavelength shifting fiber viewed by an avalanche photodiode
- Alternating horizontal and vertical planes for 3D reconstruction

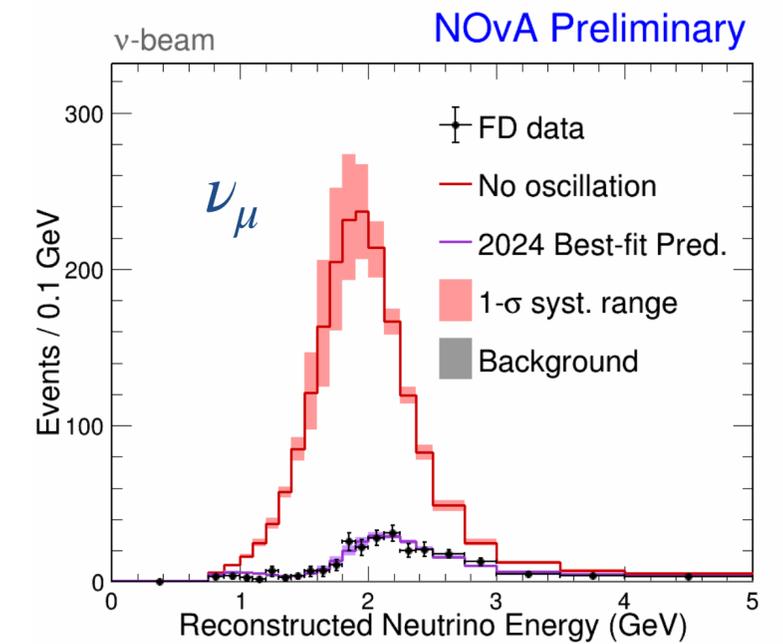
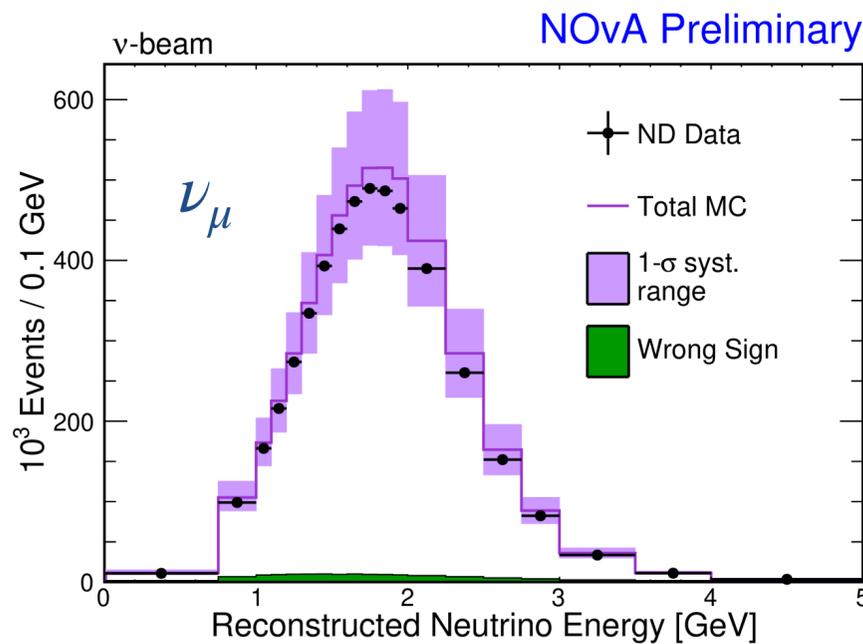


NOvA - FNAL E929  
 Run: 14828 / 38  
 Event: 192569 / --  
 UTC Tue Apr 22, 2014  
 21:41:51.422846016



# How we measure 3-flavor oscillations

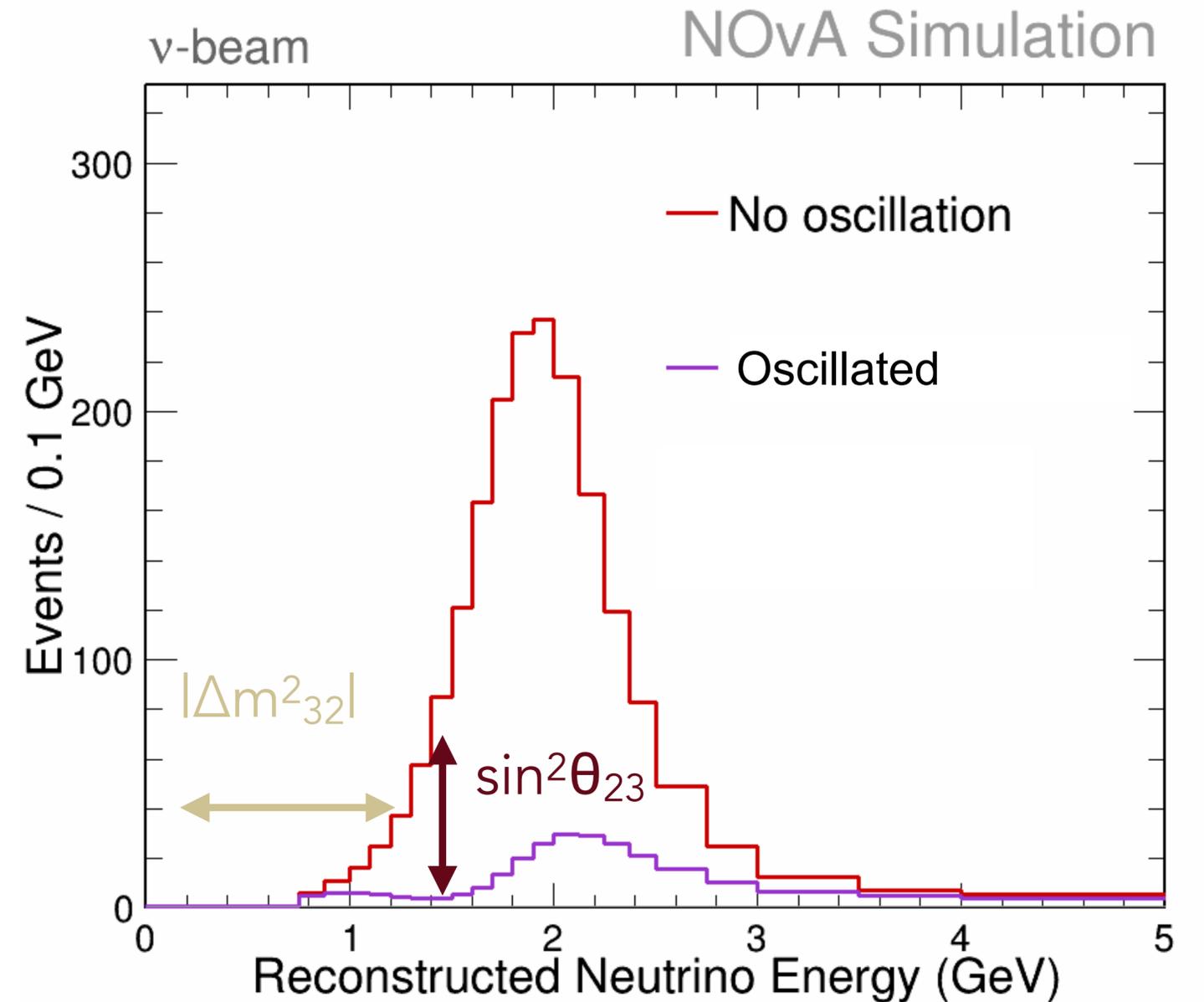
- ND and FD are functionally identical
  - They share many systematic uncertainties (neutrino cross sections and beam flux)
- Use ND data to correct the ND simulation
  - Extrapolate corrected ND simulation to generate data-corrected FD predictions
  - Repeat for different combinations of oscillation parameters and fit



# Decoding our measurements

$\bar{\nu}_\mu$  survival

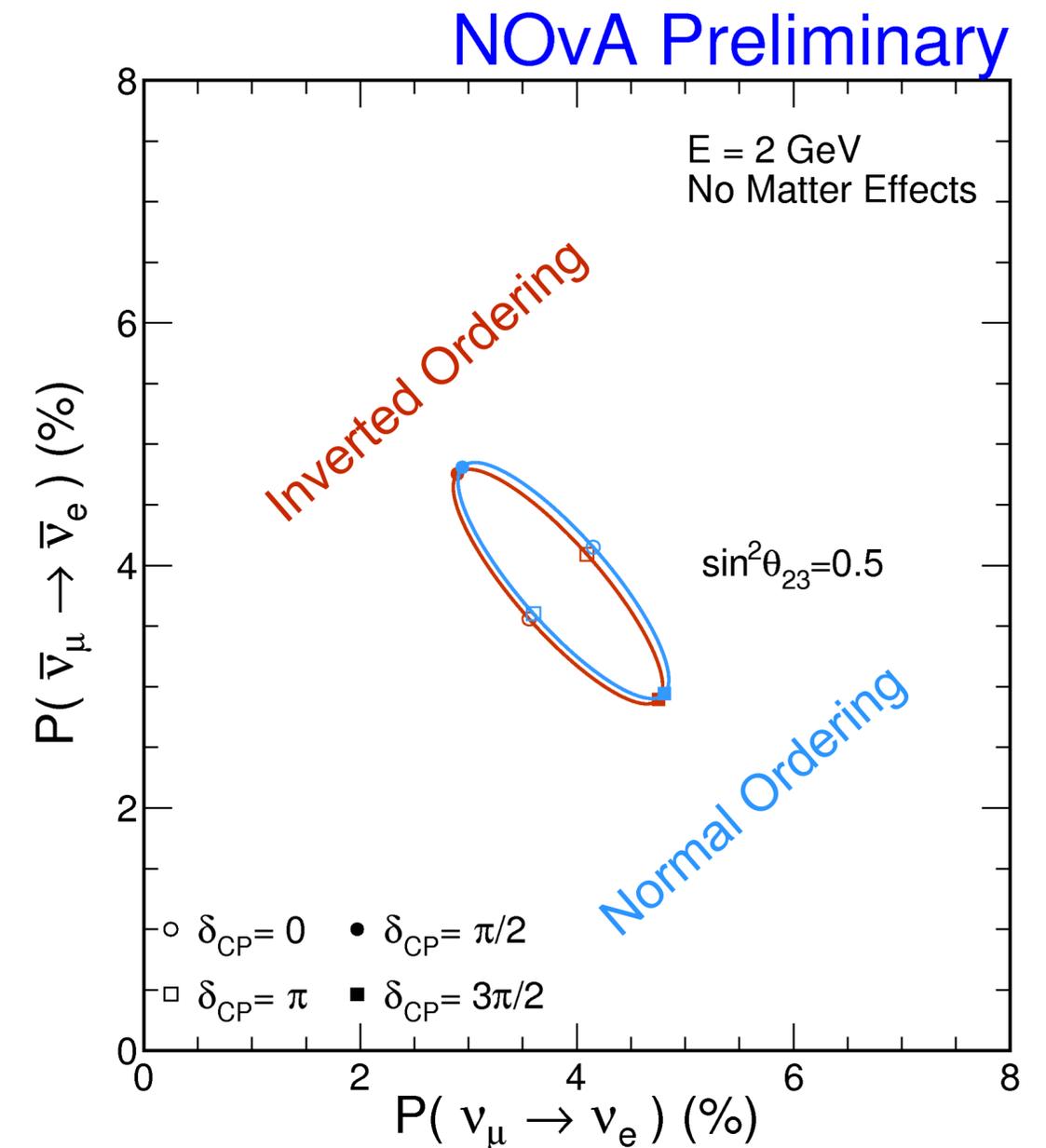
- Dip location depends on  $\Delta m_{32}^2$
- Dip amplitude depends on  $\theta_{23}$



# Decoding our measurements

$(-)$   
 $\bar{\nu}_e$  appearance

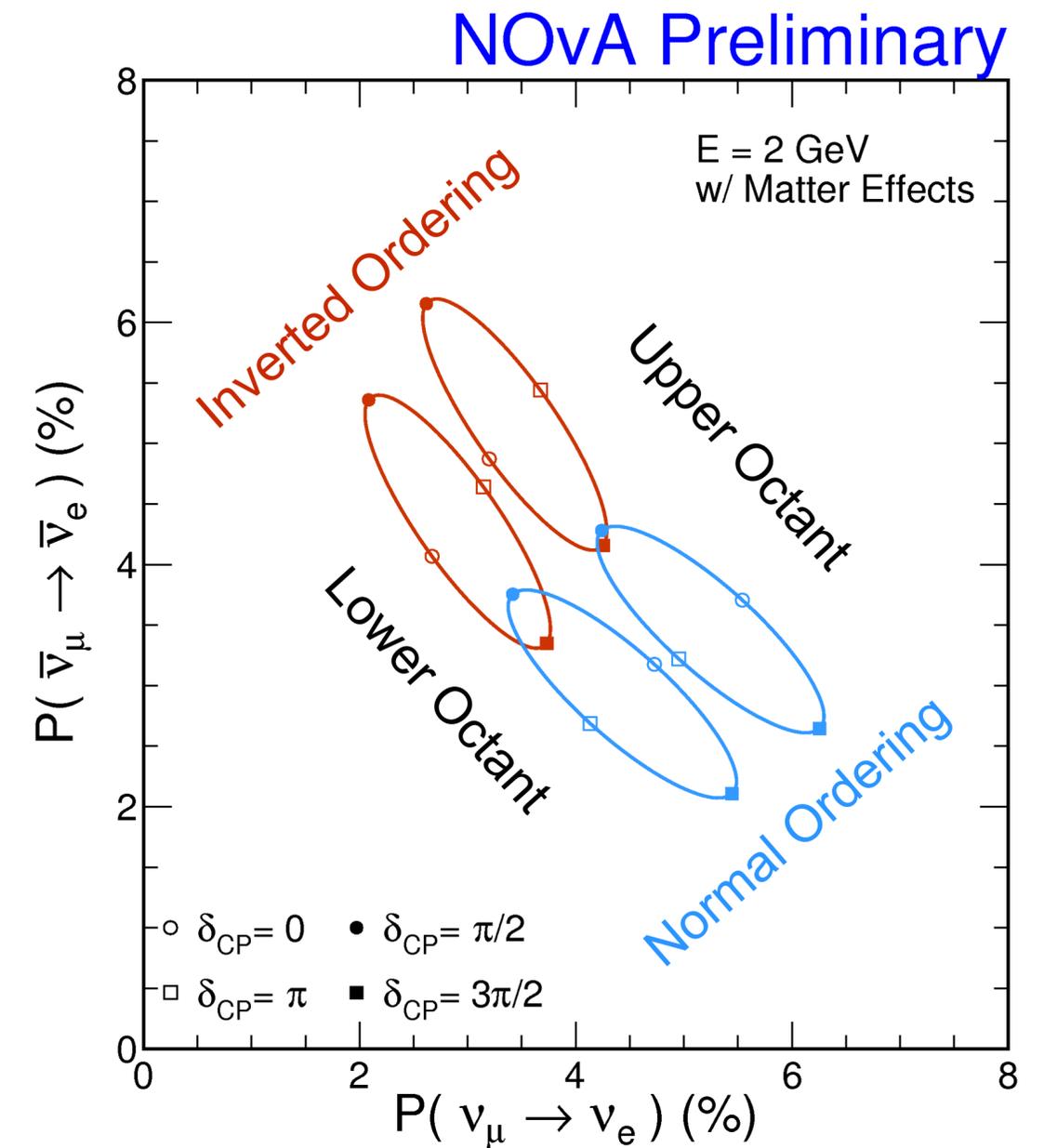
- Compare neutrino and antineutrino oscillations
- Produces ellipses based on  $\delta_{CP}$  and mass ordering
- With maximal mixing and vacuum oscillations the mass ordering and  $\delta_{CP}$  choices create degeneracies



# Decoding our measurements

$\bar{\nu}_e$  appearance

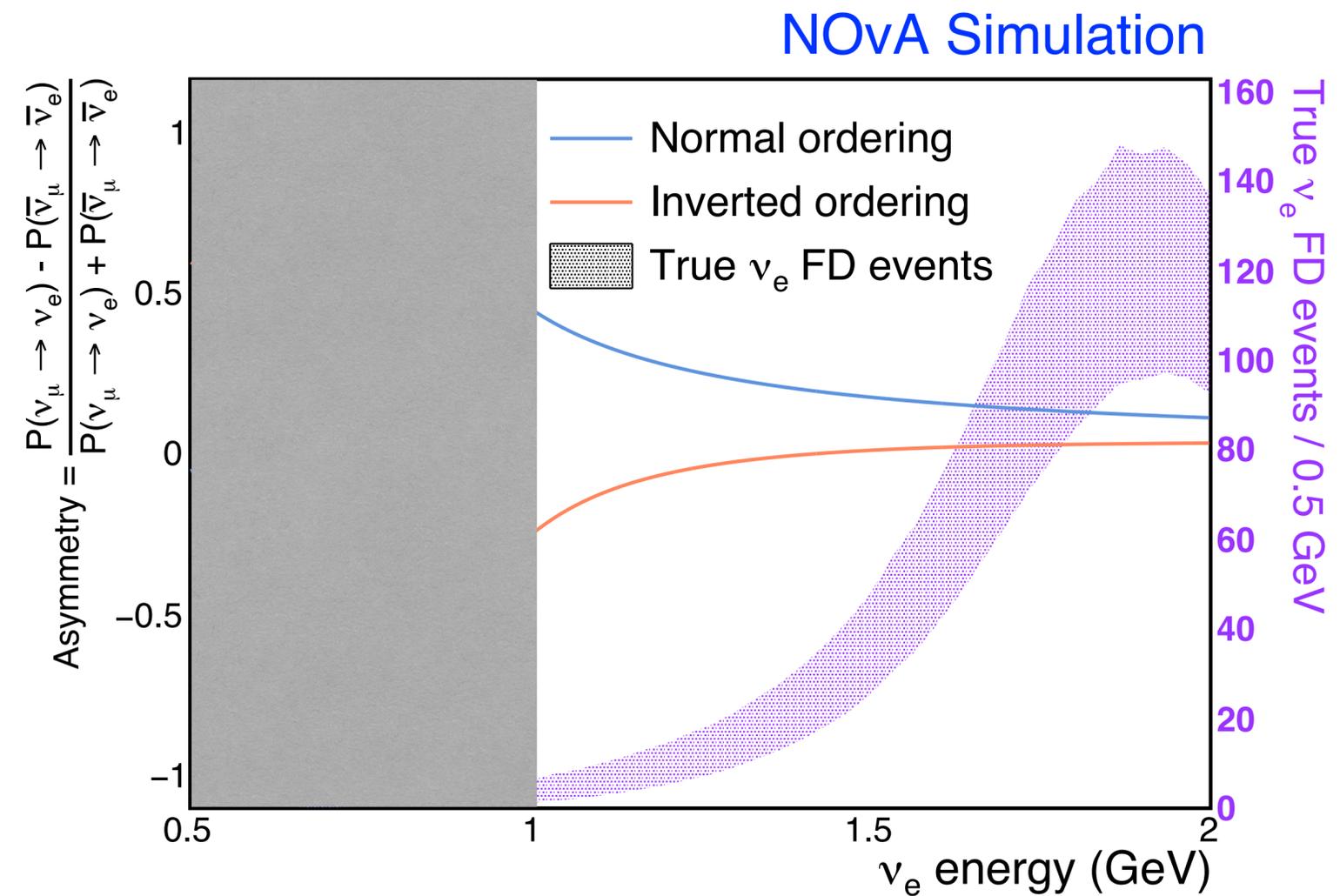
- Matter effects and  $\theta_{23}$  octant pulls the ellipses apart further splits the ellipses
- There are still degenerate regions near the middle



# Latest analysis additions

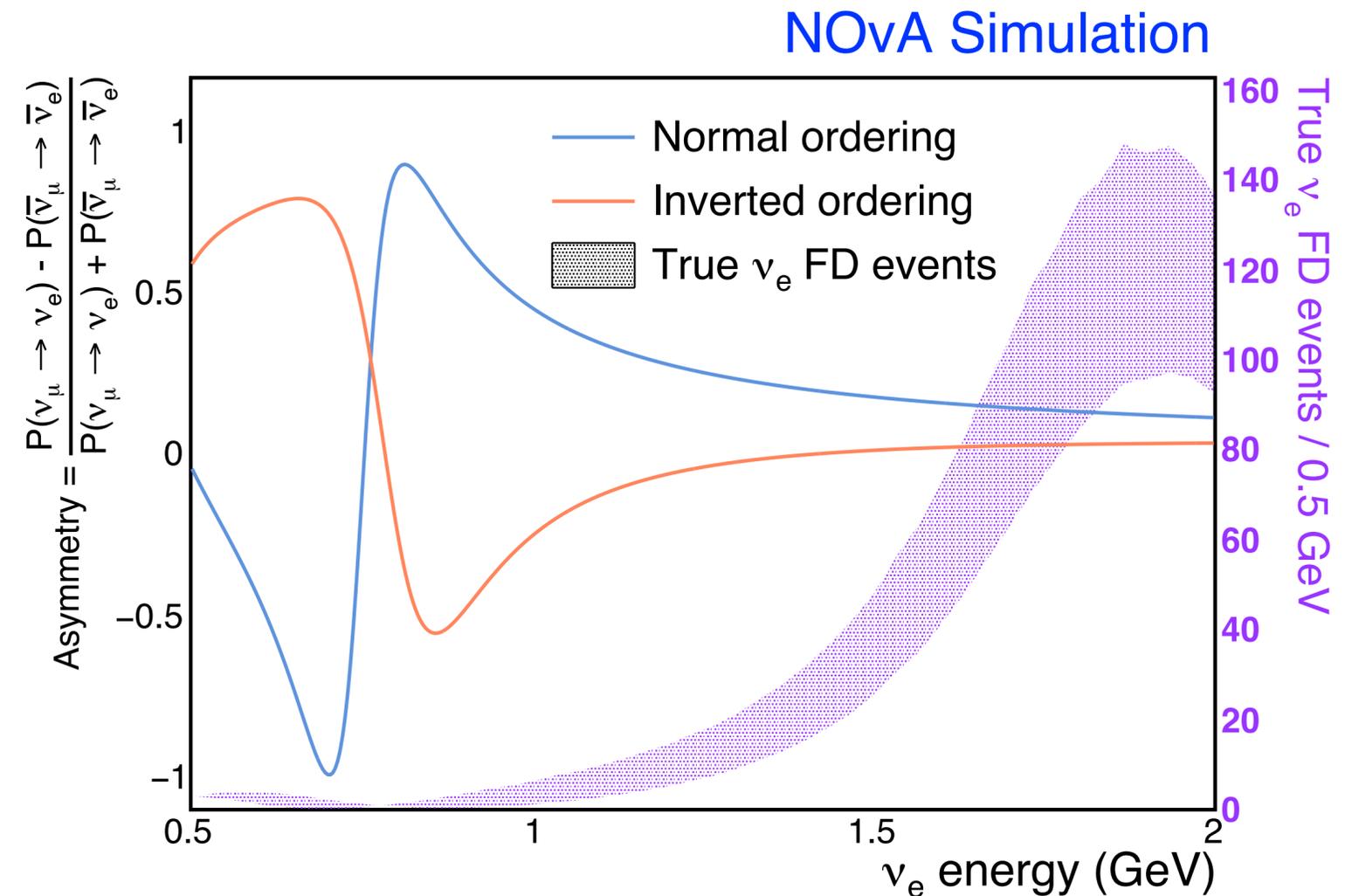
# Low energy $\nu_e$

- Previous analyses were restricted to  $E > 1$  GeV
- $\nu_e$  samples act more like a counting experiment



# Low energy $\nu_e$

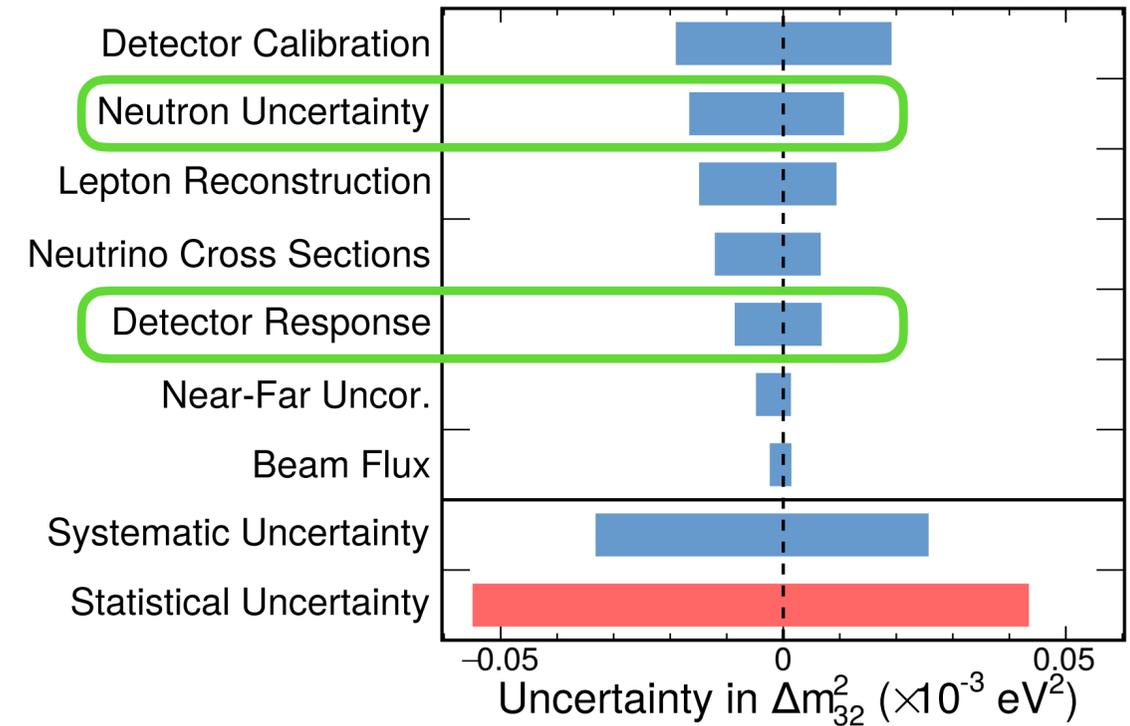
- Previous analyses were restricted to  $E > 1$  GeV
- $\nu_e$  samples act more like a counting experiment
- More shape info available at lower energies
  - Improves mass ordering sensitivity by a few percent
  - But low statistics
    - No low energy events for antineutrino beam



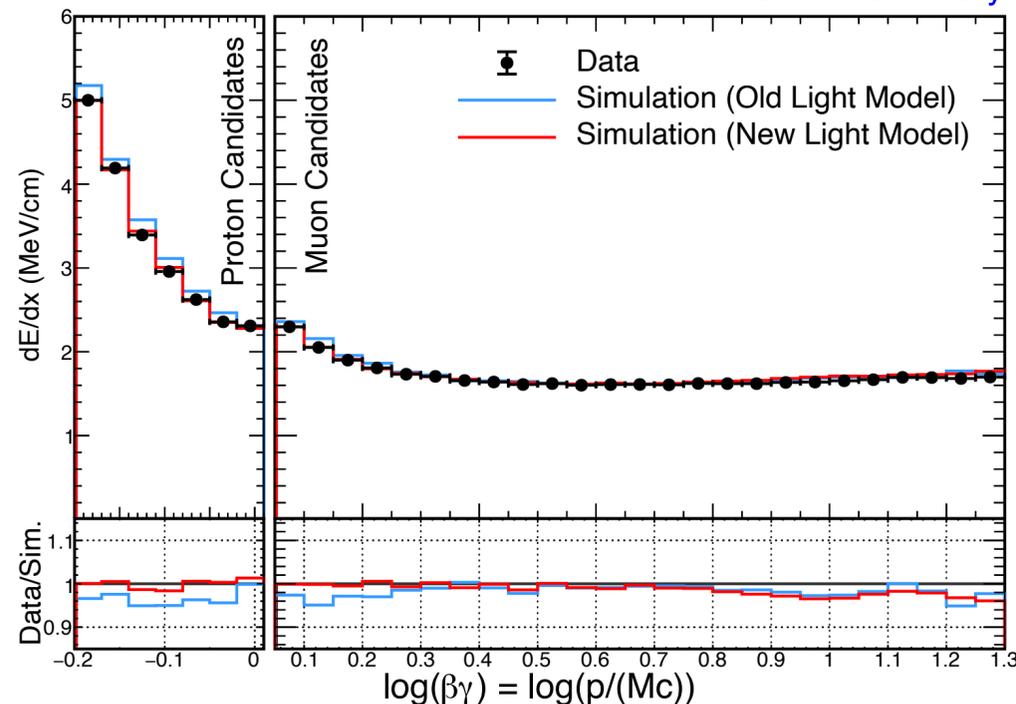
# Detector modeling

- Imperfections in the models for neutron interactions and light production have been leading uncertainties
- New light tuning agrees much better with data
  - Combined fit to multiple ND samples
- Observed an over-simulation of neutron candidates

2020 Analysis Uncertainties NOvA Simulation

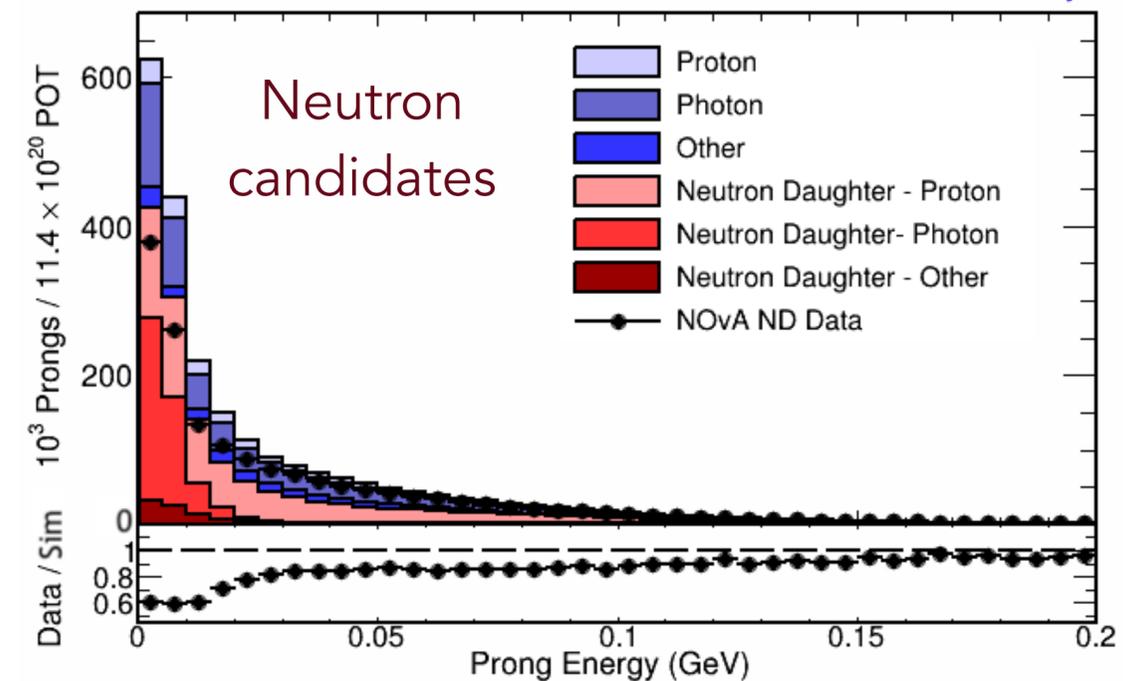


NOvA Preliminary



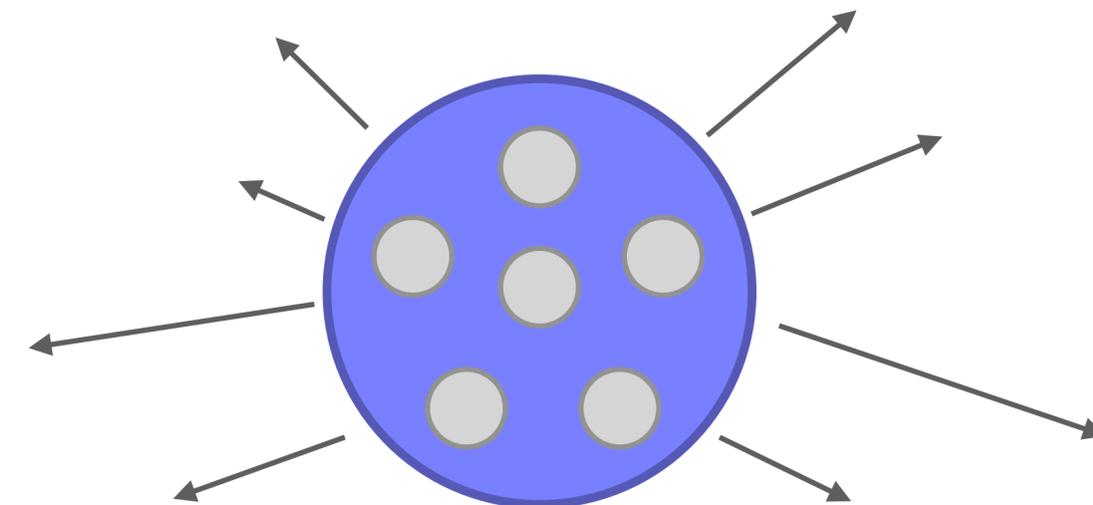
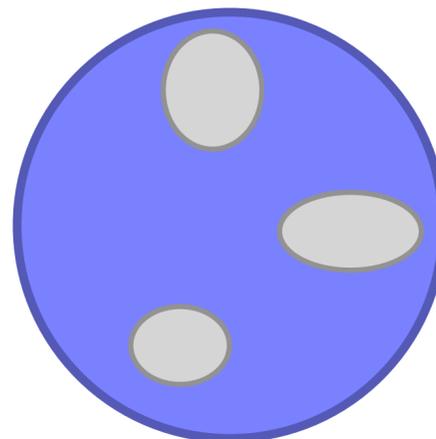
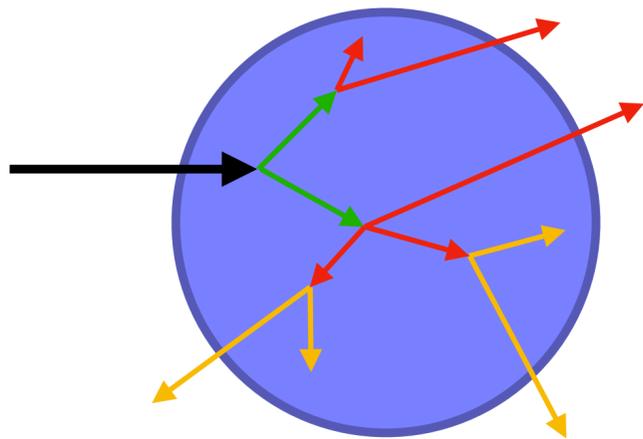
$\bar{\nu}$ -beam

NOvA Preliminary



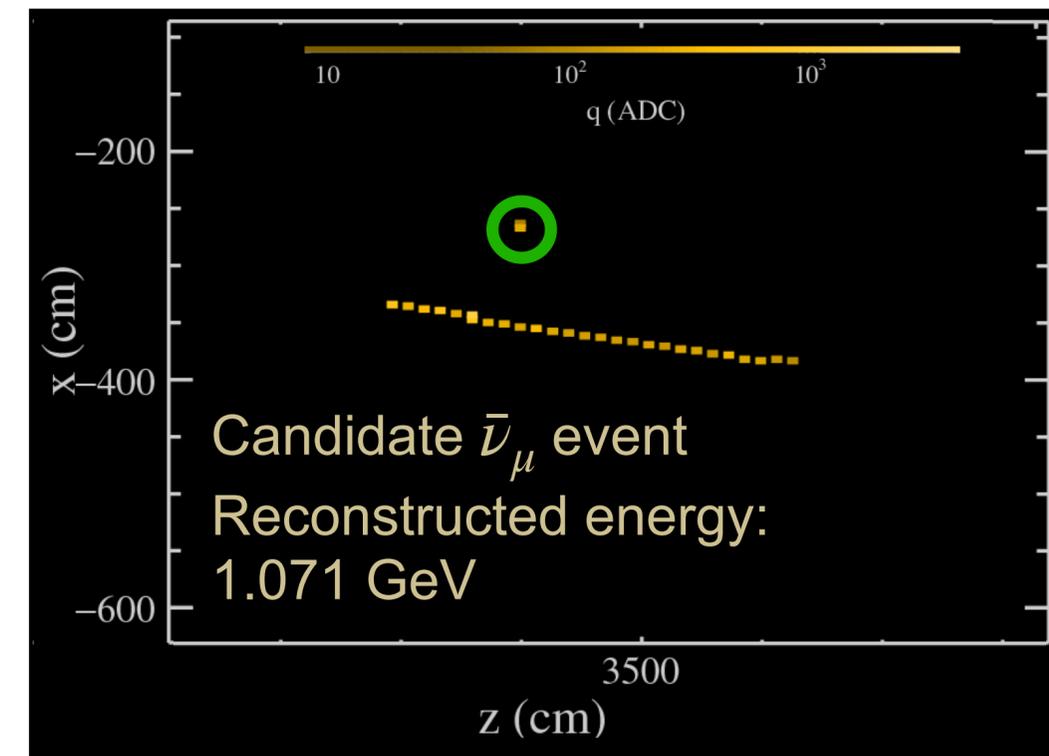
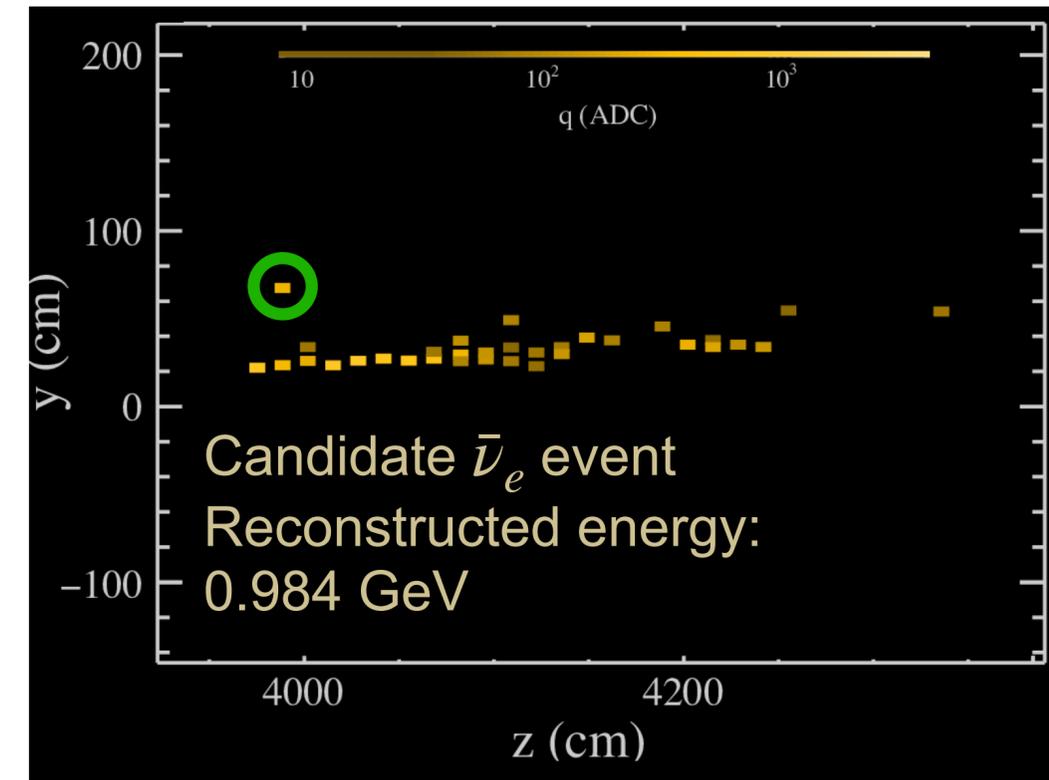
# The problem with neutrons

- To study neutrinos we need to know their incoming energy
- Neutrinos are neutral  $\rightarrow$  invisible until they interact  $\rightarrow$  sum energy of outgoing particles
- Neutrons are neutral  $\rightarrow$  can carry energy away unseen
- Our nuclear colleagues have done a great job with them, but generally at lower energies (from thermal to  $\sim 20$  MeV)



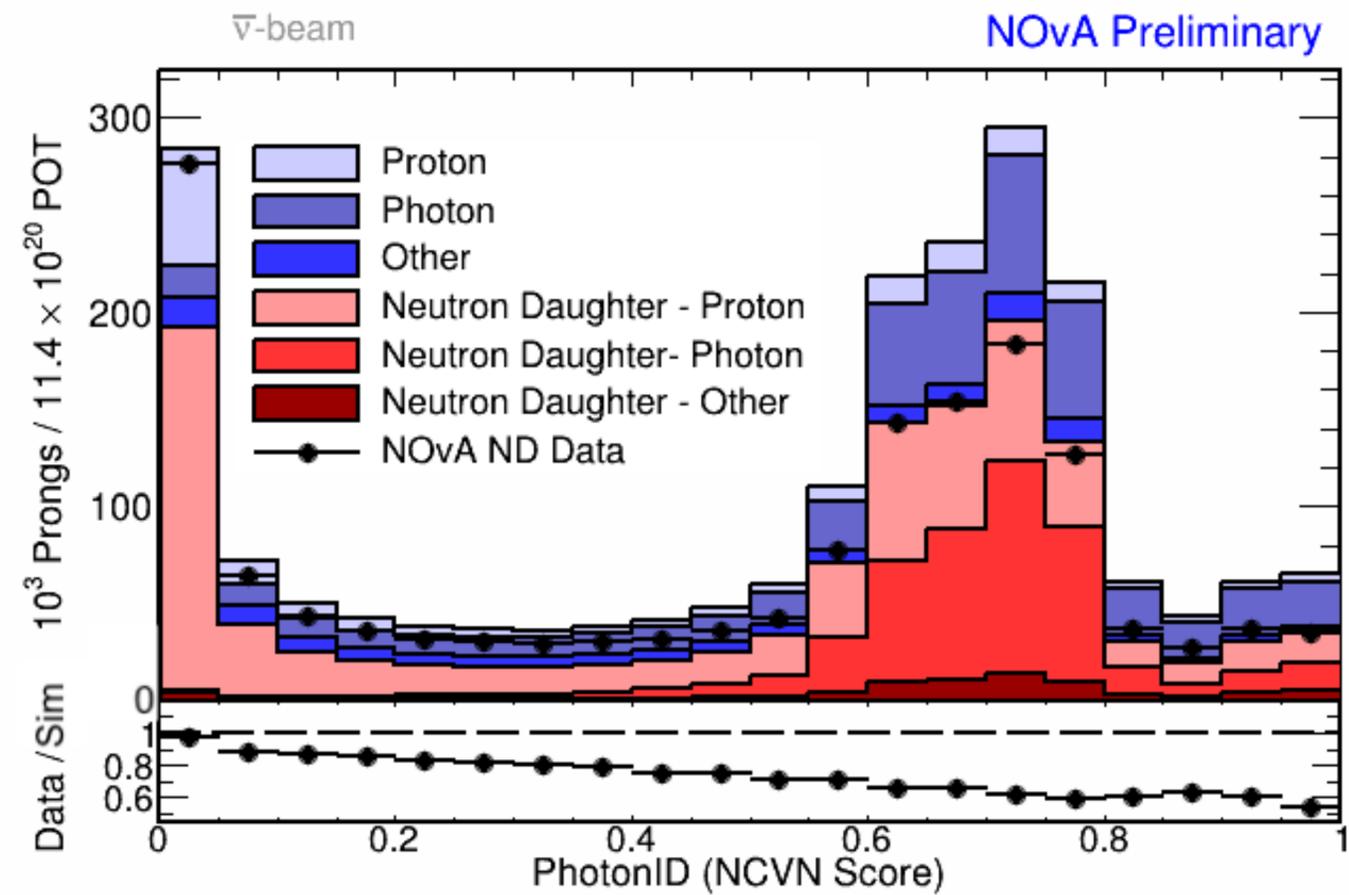
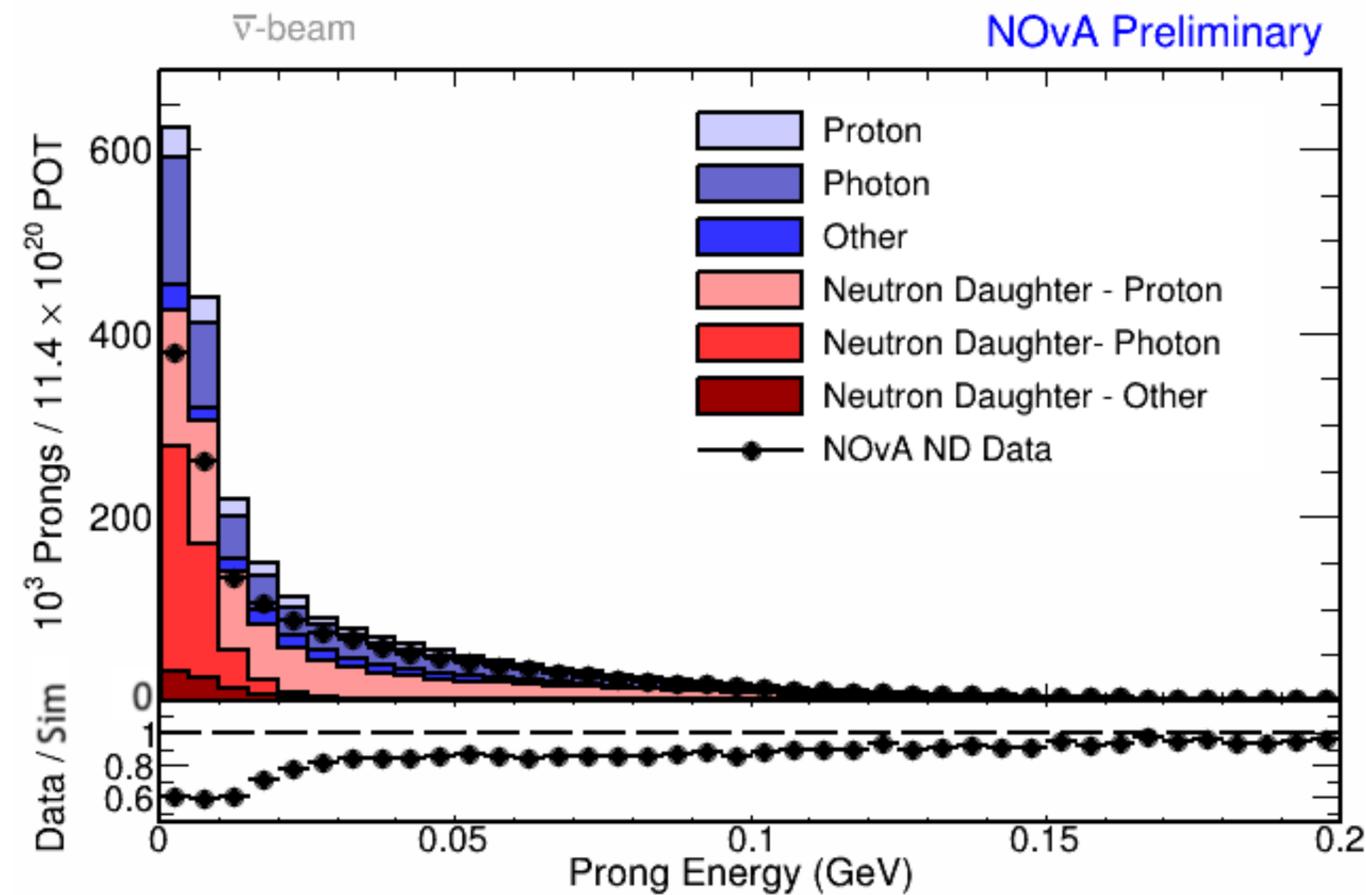
# Neutron selection

- Three main interaction modes
  - Elastic scattering: nuclear recoil energy is typically low
  - Inelastic scattering: produces photons and charged particles
  - Neutron capture: if they thermalize, on chlorine or hydrogen
- Developed a simple criteria to select neutron hit candidates
  - Many elastic scatters before producing visible particles
    - Hits should be  $> 20$  cm from the neutrino interaction
    - A candidate should contain fewer than 6 hits
  - Selects visible neutrons with 71% efficiency and 61% purity



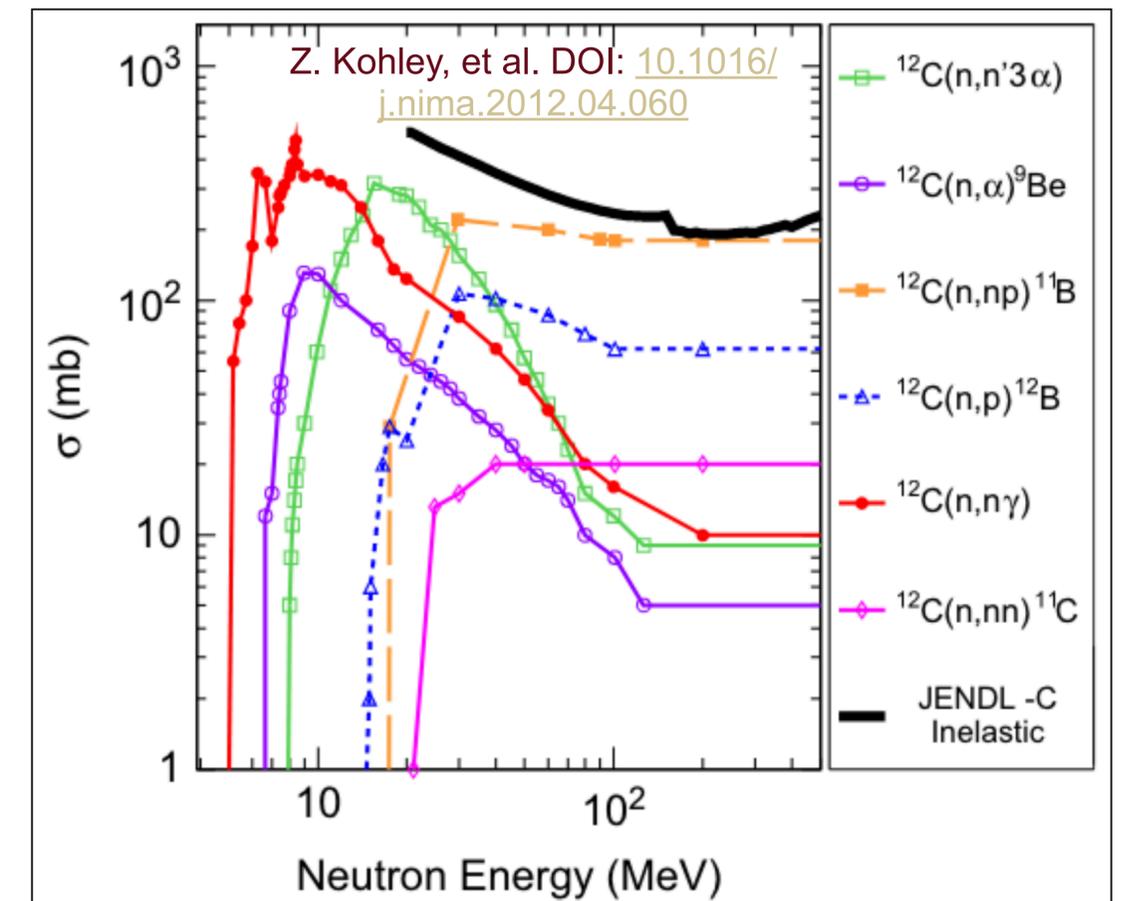
# Neutron over-simulation

- Base simulation produces up to 40% more neutrons as compared to data
  - De-excitation photons produced by neutron inelastic scattering (medium red) are concentrated at low energy where the simulation excess is the greatest



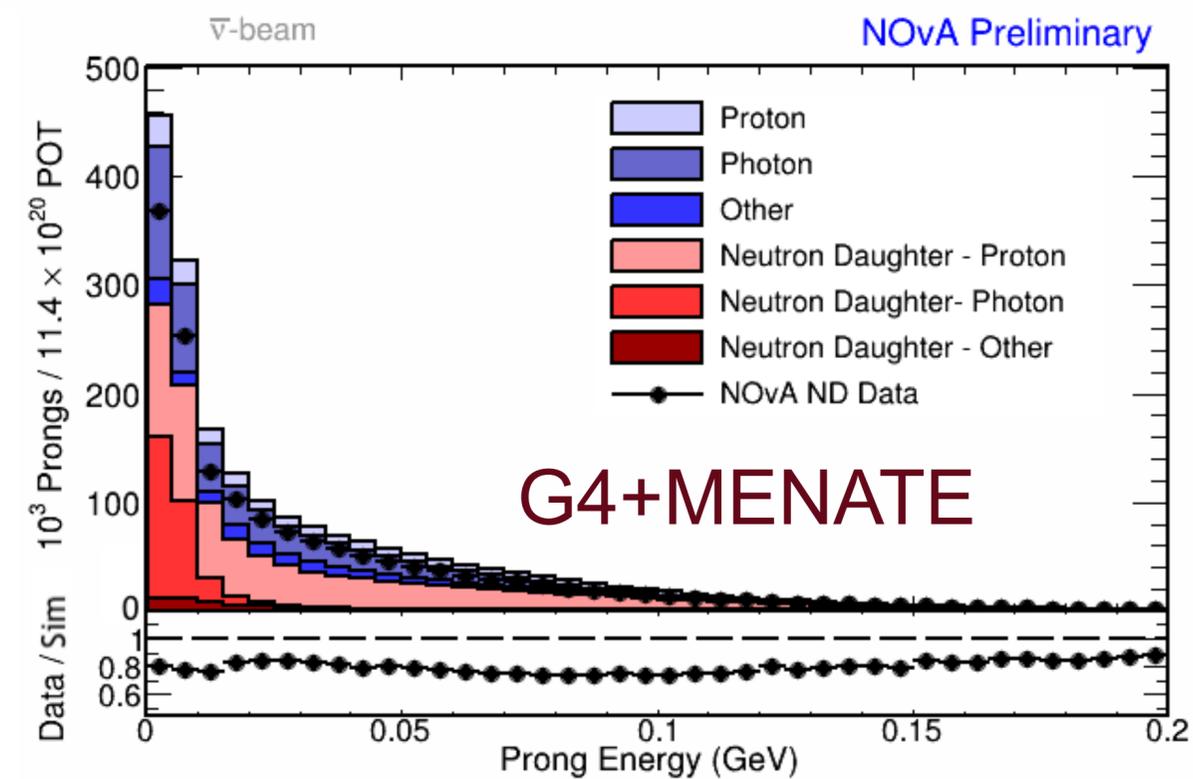
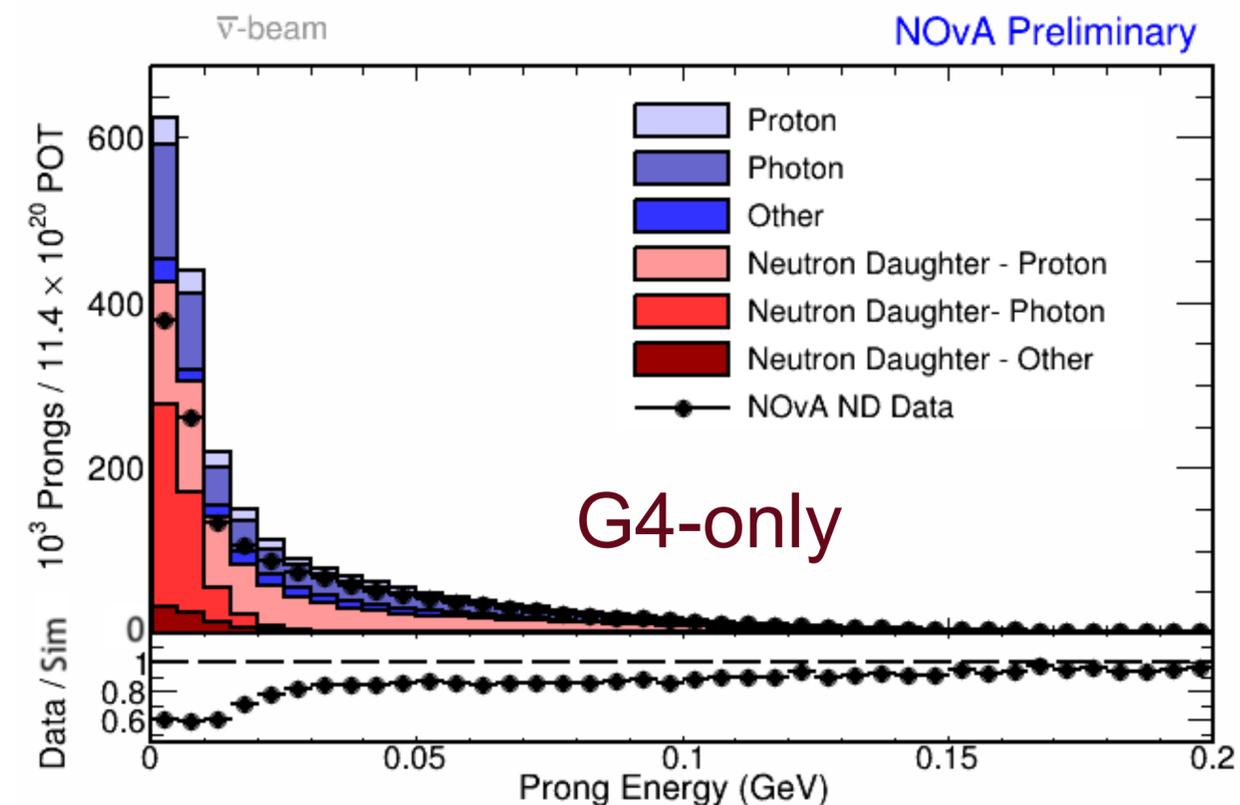
# Alternative neutron inelastic model

- De-excitation photons produced by Geant4:
  - The HP model: determines outgoing particles based on measurements (limited data)
  - Intranuclear cascade: theory-driven and statistically determines outgoing particles while using the overall cross-section
- MENATE extends the ethos of the HP model to higher energies
  - But only for neutron-on-carbon interactions (that's what we have data for)
  - Limits  $n+^{12}\text{C}$  to only 6 final inelastic states
  - Our implementation integrates directly with Geant4



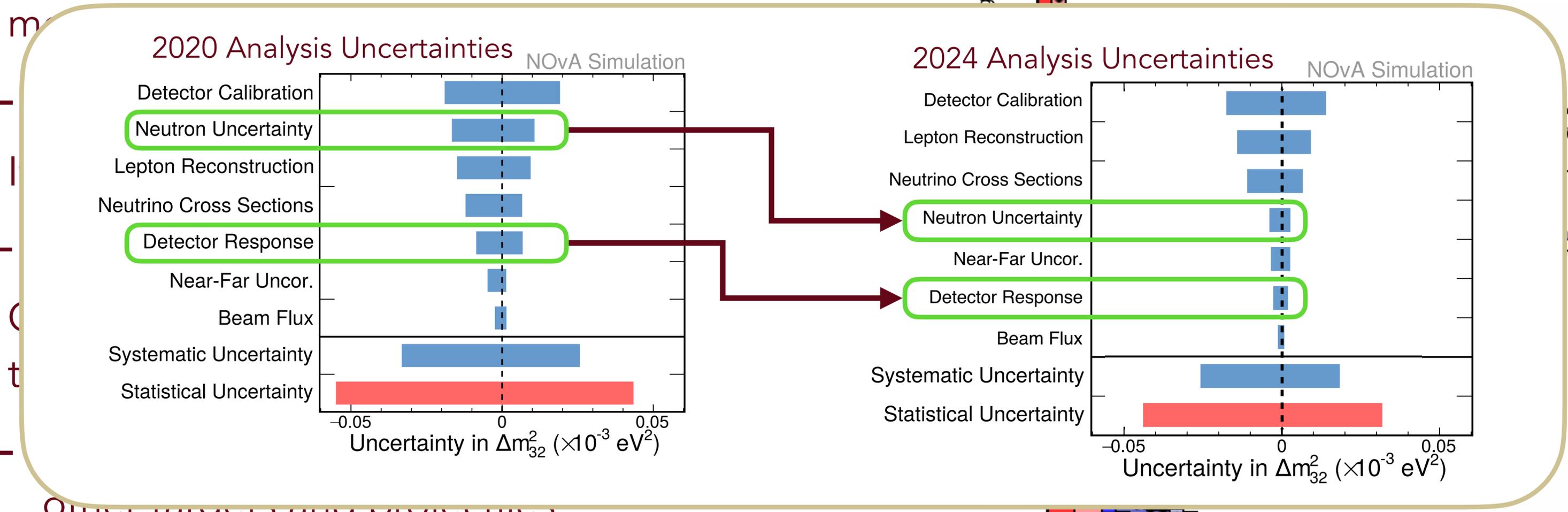
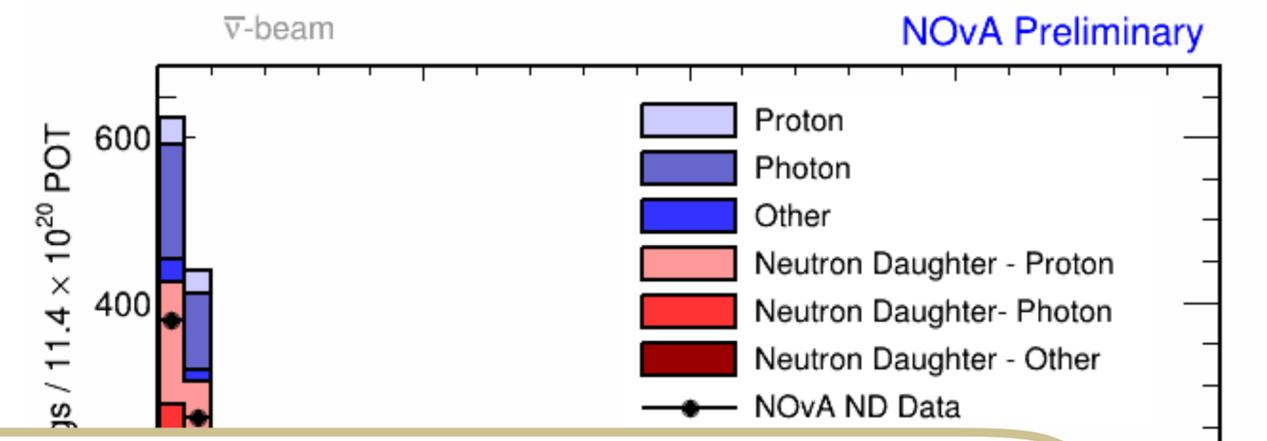
# Comparing to data

- MENATE reduces the over-simulation in the most populous bins
  - Photon production is significantly decreased
- It increases the over-simulation at higher energies
  - Proton production is slightly increased
- Overall, the data-simulation ratio is flatter, but there is still a significant over-simulation
  - Only addressing neutron-on-carbon; there are other targets and projectiles

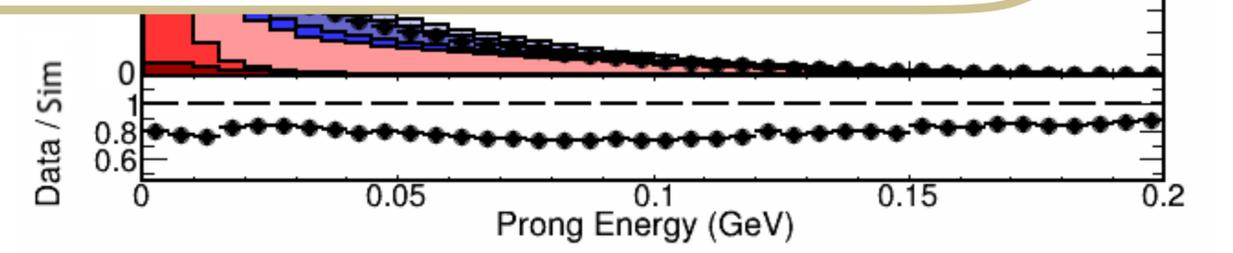


# Comparing to data

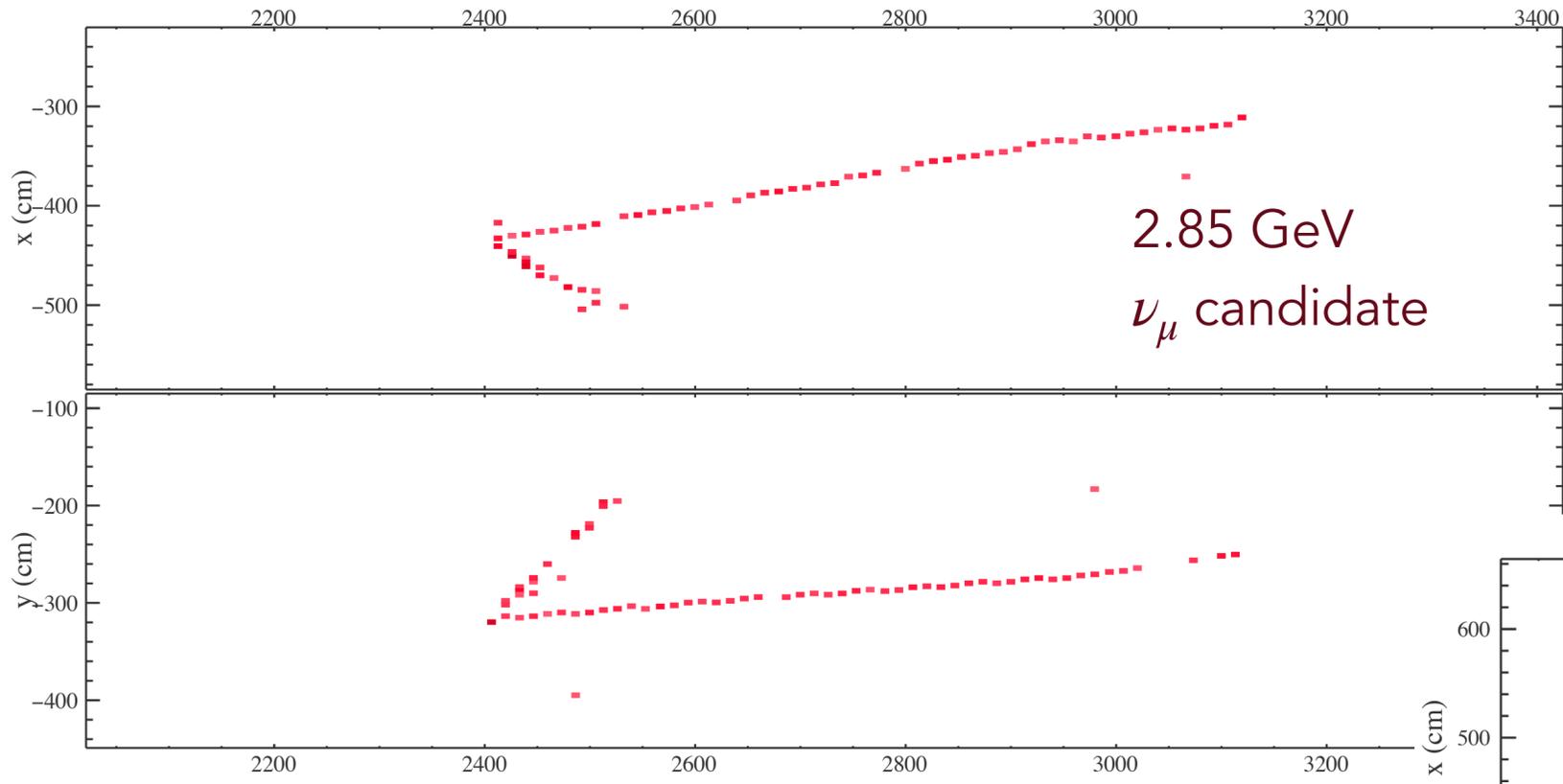
- MENATE reduces the over-simulation in the



Other targets and projects



# Results



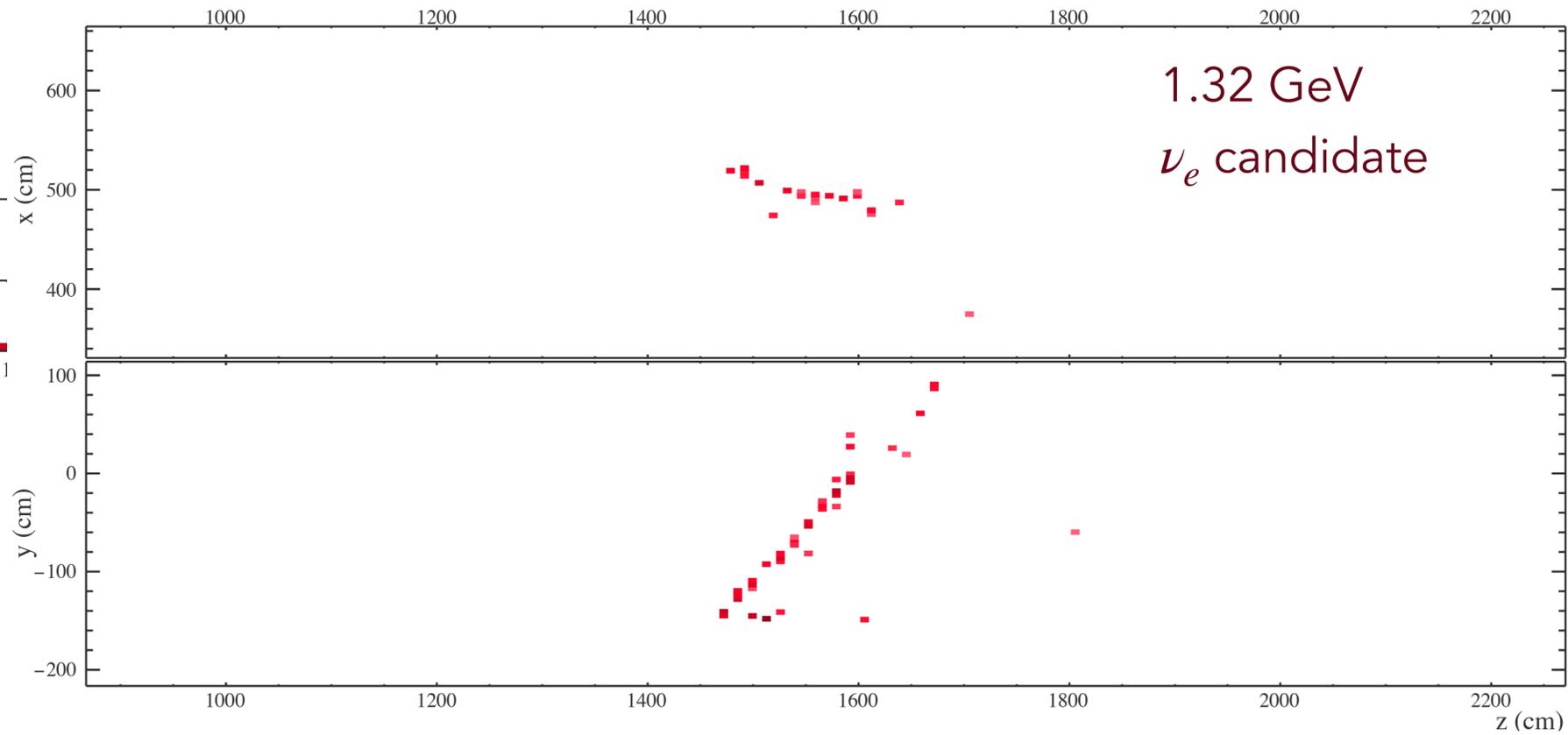
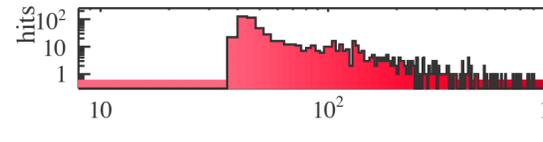
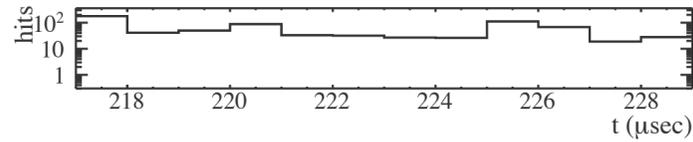
NOvA - FNAL E929

Run: 14828 / 38

Event: 192569 / NuMI

UTC Tue Apr 22, 2014

21:41:51.422846016



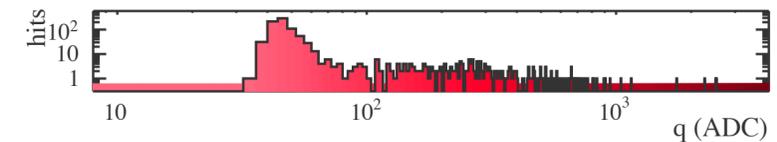
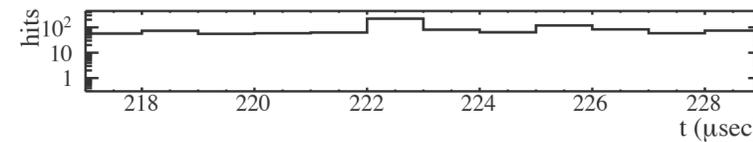
NOvA - FNAL E929

Run: 40072 / 11

Event: 6764 / NuMI

UTC Fri Jun 4, 2021

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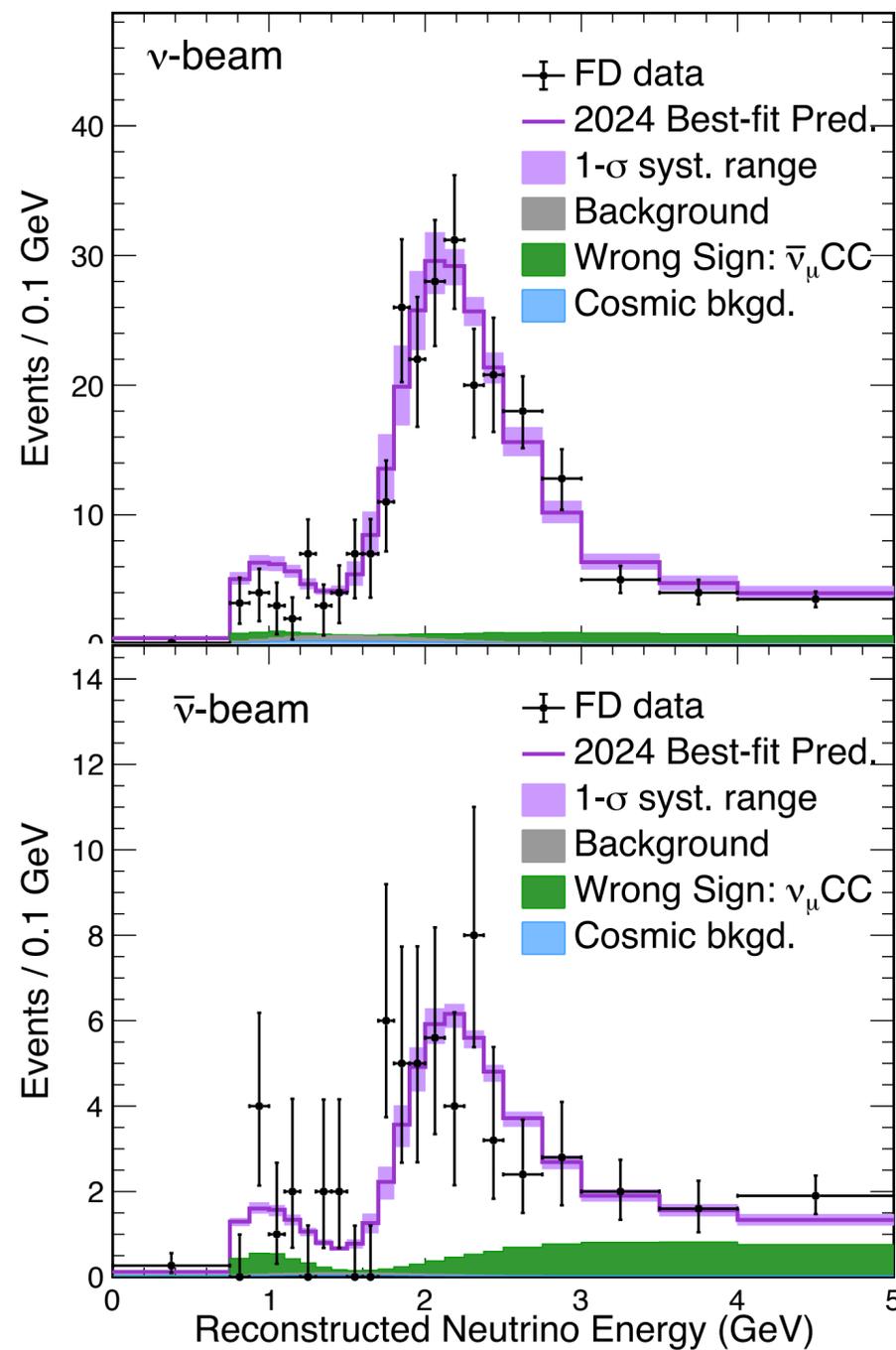


# FD data

Sample	Num. Events	Background
$\nu_\mu$	384	11.3
$\bar{\nu}_\mu$	106	1.7
$\nu_e$	181	61.7
$\bar{\nu}_e$	32	12.2

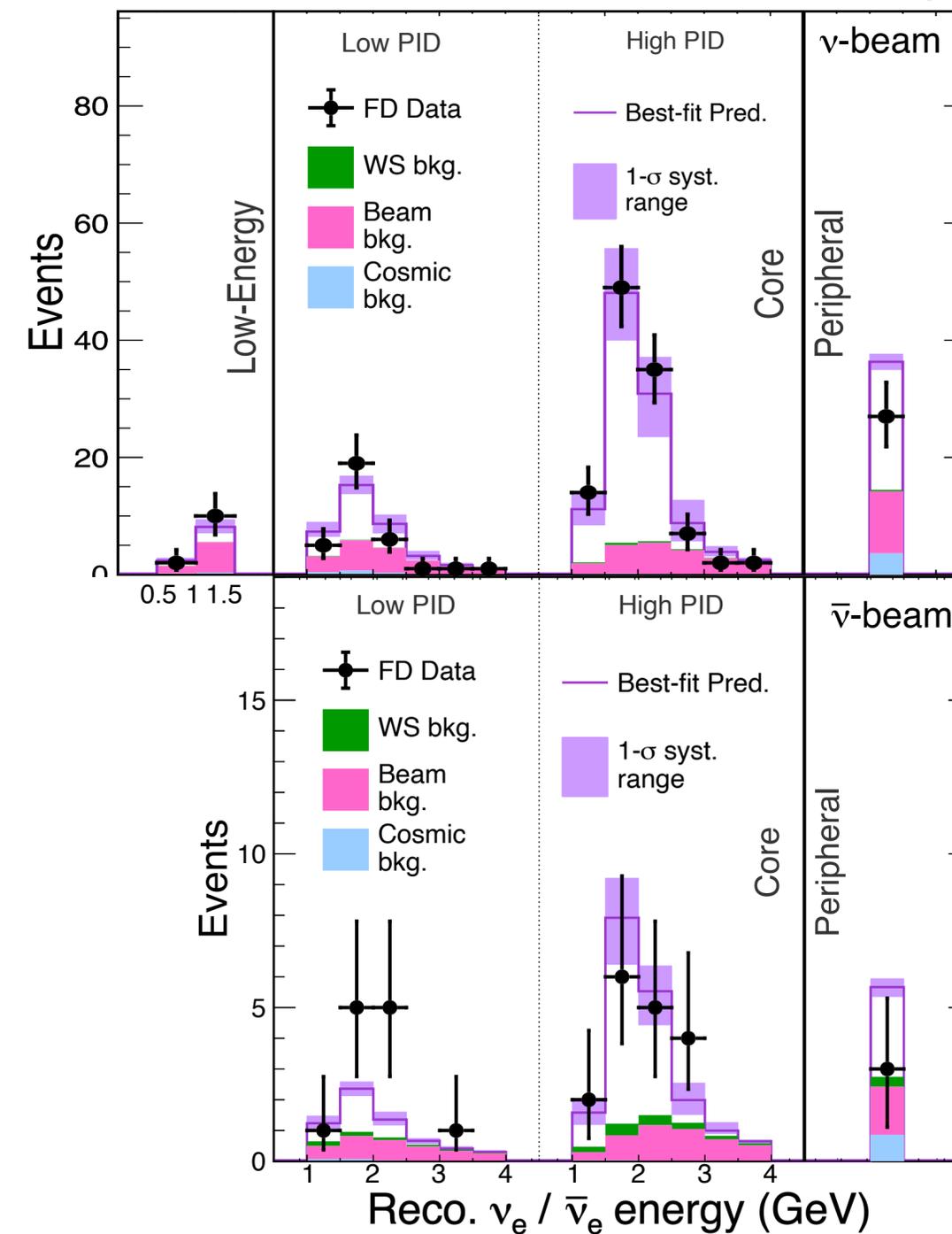
$\nu_\mu$

NOvA Preliminary



$\nu_e$

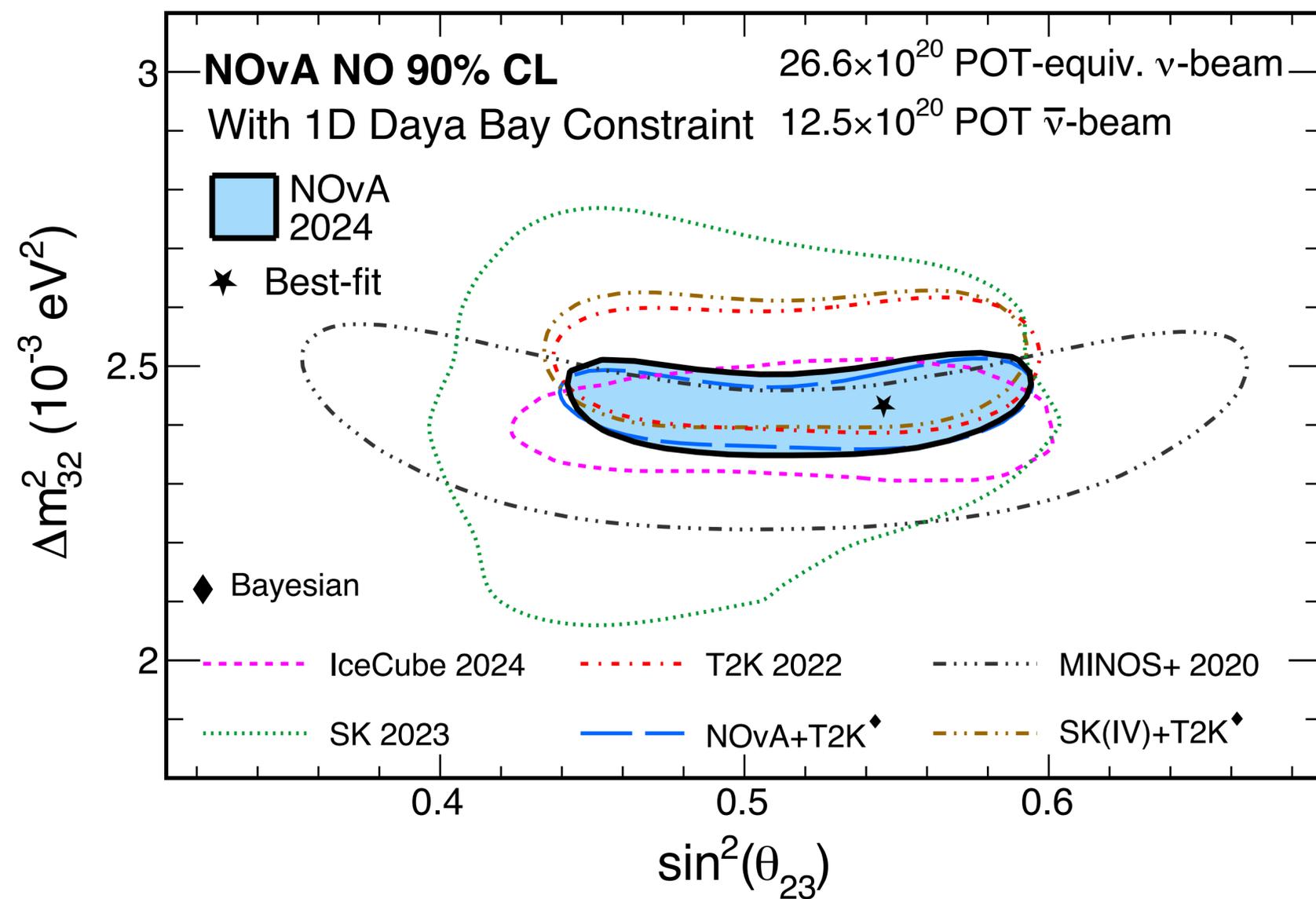
NOvA Preliminary



# Frequentist results

NOvA Preliminary

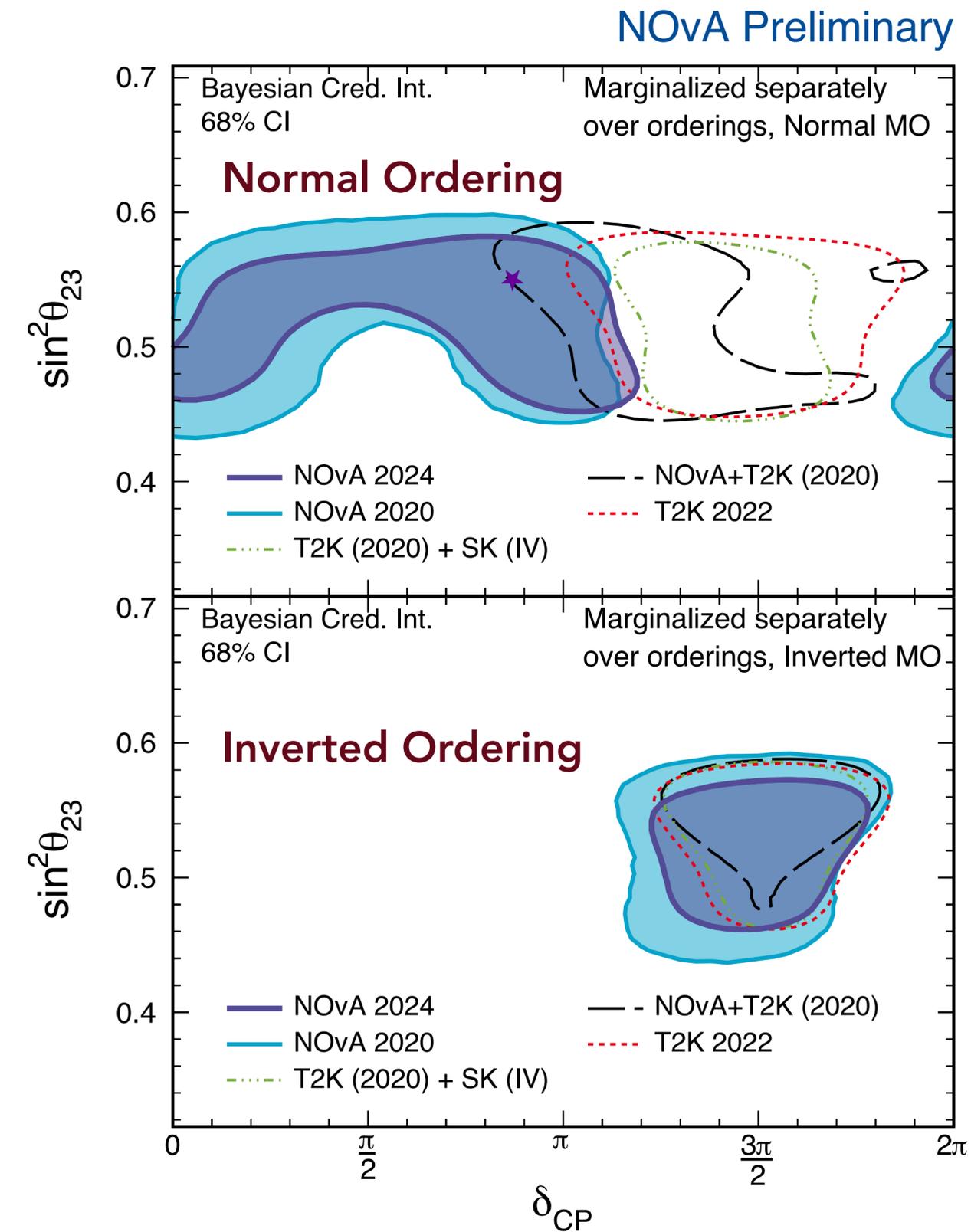
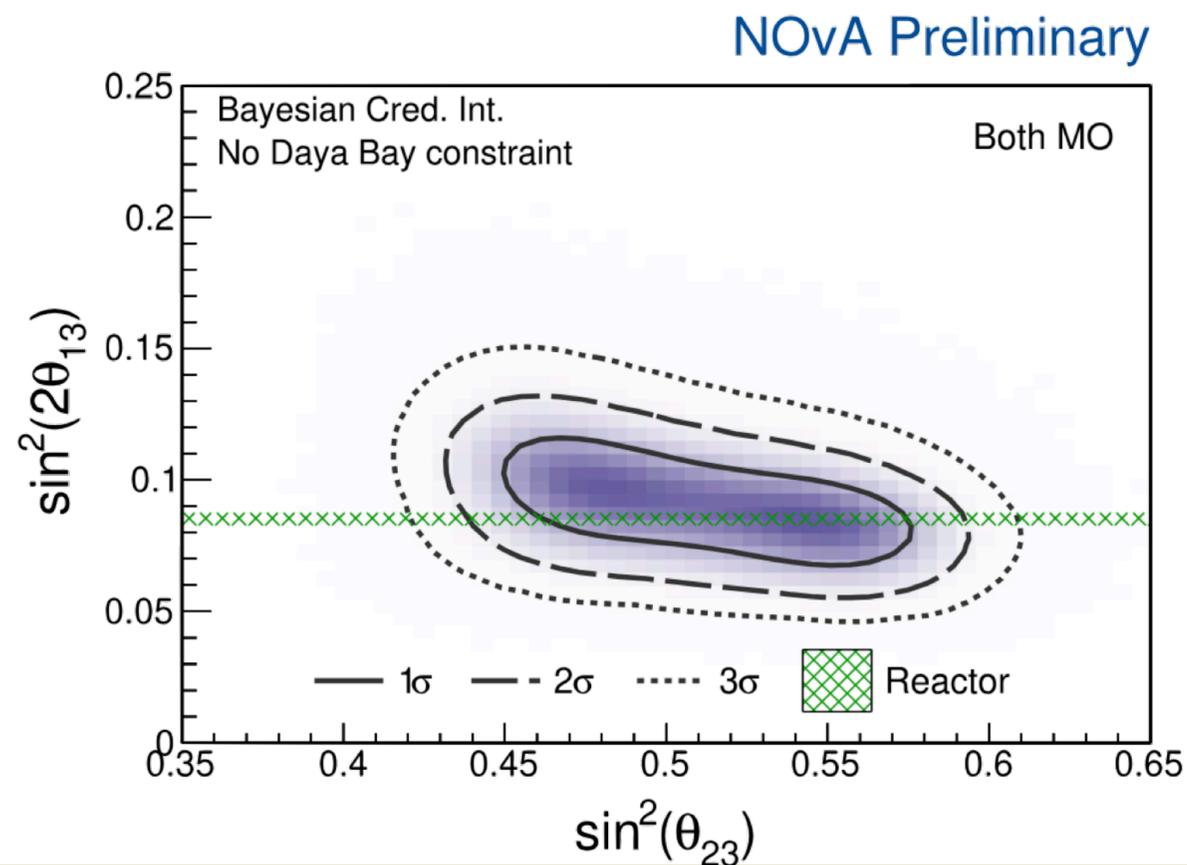
Parameter	Best-fit	Normal Ordering Preference ( $\sigma$ )	
$\sin^2(\theta_{23})$	$0.546^{+0.032}_{-0.075}$	W/ 1D Daya Bay constraint	p-value 0.1731
$\Delta m_{32}^2$ ( $10^{-3} \text{ eV}^2$ )	$2.433^{+0.035}_{-0.036}$		$1.36 \sigma$
$\delta_{CP}$ ( $\pi$ )	0.875	W/ 2D Daya Bay constraint	p-value 0.1158
			$1.57 \sigma$



- New result is consistent with previous measurements and other experiments
- Prefer the Upper  $\theta_{23}$  octant and Normal Mass Ordering
- Most precise single-experiment measurement of  $\Delta m_{32}^2$

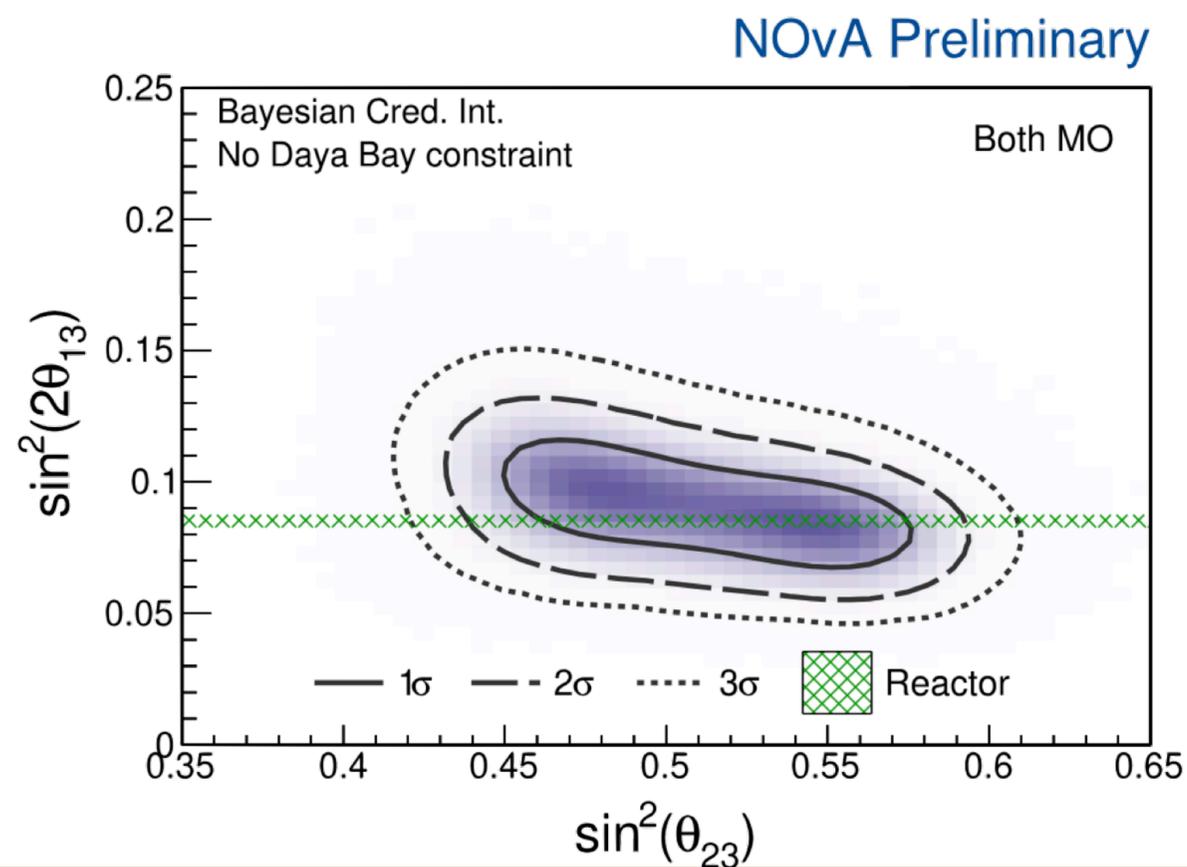
# Bayesian results

- Prefer CP conserving values of  $\delta_{CP}$  in the normal mass ordering and CP violating values in the inverted
- Preference for Normal Mass Ordering and Upper  $\theta_{23}$  Octant is enhanced by reactor constraints

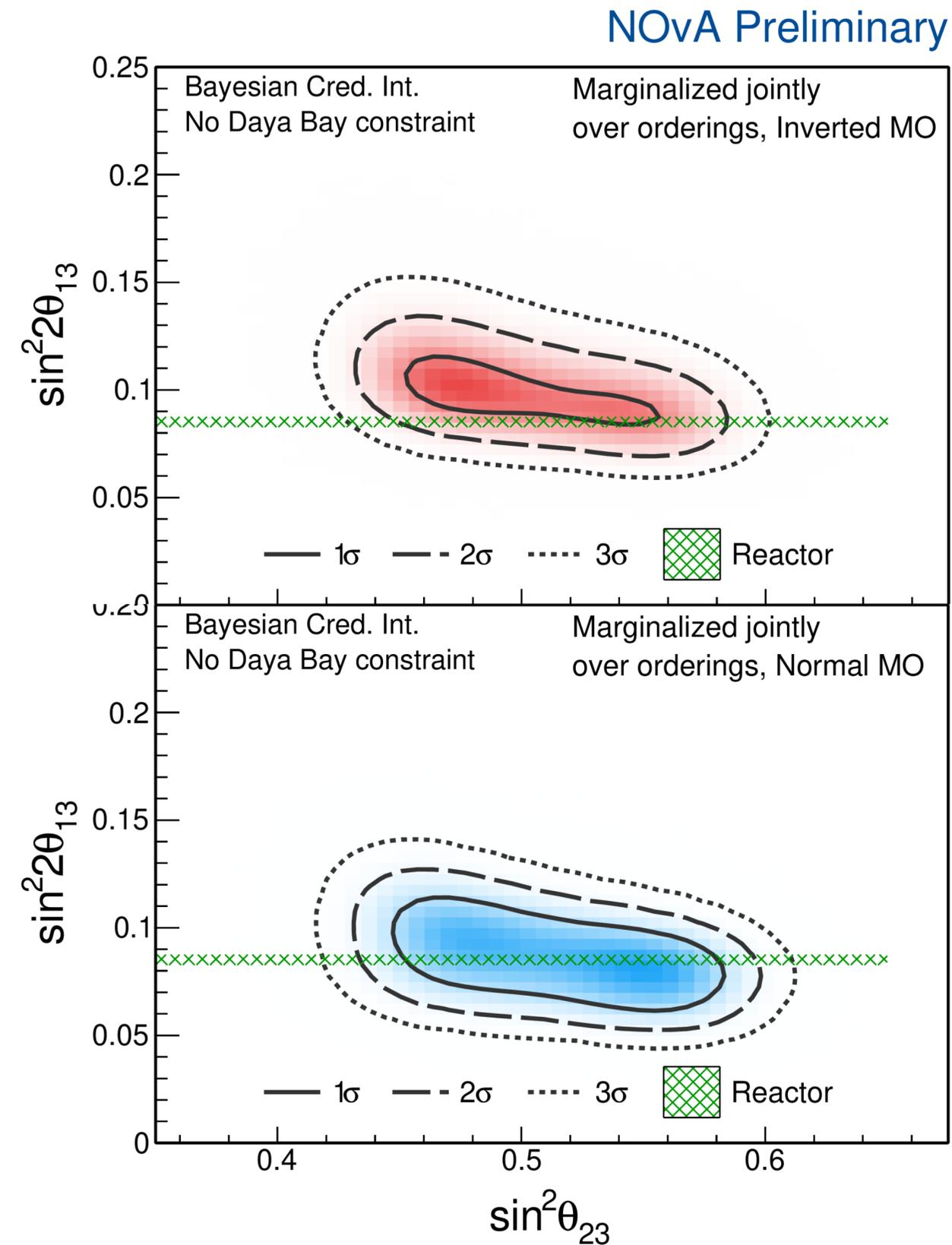


# Bayesian results

- Prefer CP conserving values of  $\delta_{CP}$  in the normal mass ordering and CP violating values in the inverted
- Preference for Normal Mass Ordering and Upper  $\theta_{23}$  Octant is enhanced by reactor constraints

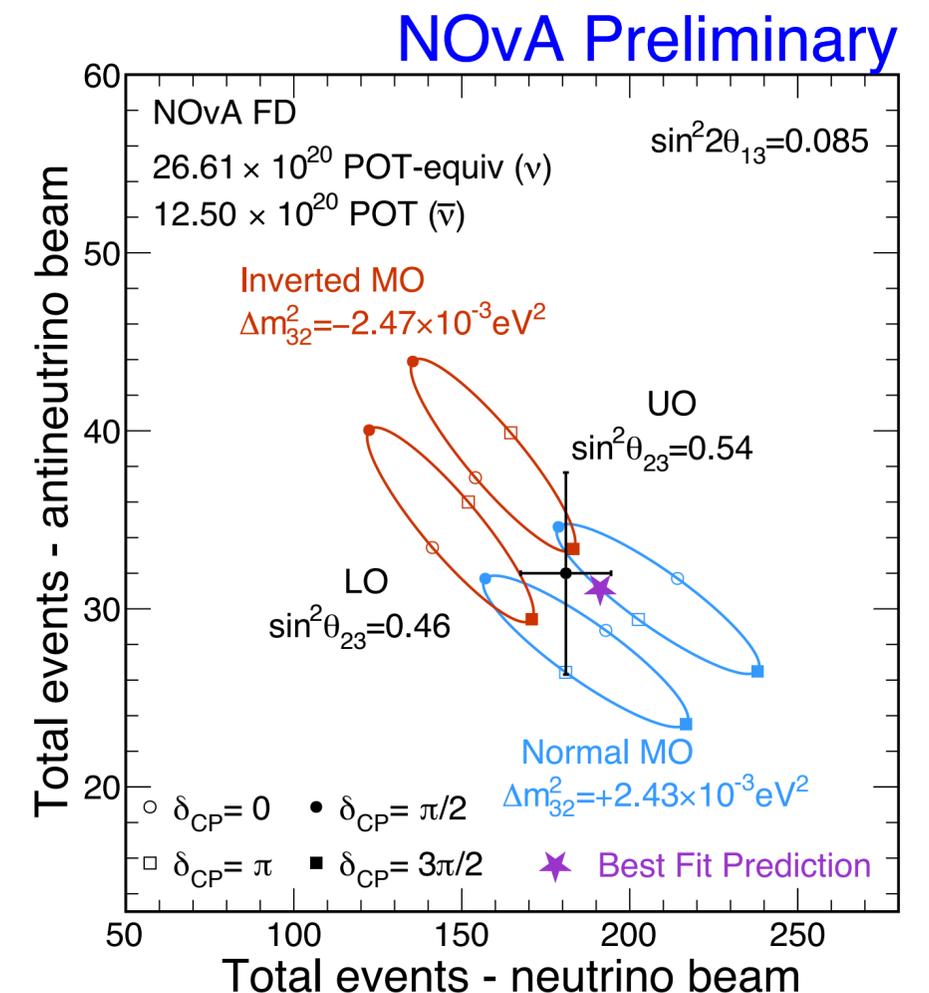


- NO preference at 87% posterior probability



# Summary

- Newest NOvA results contain 2X more neutrino-mode data from 10 years of running
- Most most precise single-experiment measurement of  $\Delta m_{32}^2$ 
  - Now the most precisely measured oscillation parameter
- Slight preference for normal mass ordering and upper  $\theta_{23}$  octant, but in a highly degenerate region
- We are analyzing results from a Test Beam run to help constrain our largest systematic uncertainty



# Summary

# Thanks!

This document was prepared by NOvA using the resources of the Fermi National Accelerator Laboratory (Fermilab), a U.S. Department of Energy, Office of Science, Office of High Energy Physics HEP User Facility. Fermilab is managed by FermiForward Discovery Group, LLC, acting under Contract No. 89243024CSC000002.

- New
- Most
- N
- Slight
- octa
- We
- cons



50 100 150 200 250  
Total events - neutrino beam

# Backups

# Neutrino Oscillations

- Neutrino are observed in "flavor states" ( $\nu_e, \nu_\mu, \nu_\tau$ )
- Which are superpositions of "mass states" ( $\nu_1, \nu_2, \nu_3$ )
  - The mixing is described by the unitary matrix:  $U_{PMNS}$
- As a neutrino propagates, the mass states interfere causing the flavor to change from one type to another
- Oscillations depend on:
  - Elements of  $U_{PMNS}$
  - Distance of travel / neutrino energy
  - Mass (squared) splittings

$$|\nu_\alpha(\mathbf{x}, t)\rangle = \sum_{i=1}^3 U_{\alpha,i}^* |\nu_i(\mathbf{x}, t)\rangle$$

$$P(\nu_\alpha \rightarrow \nu_\beta) = f(U_{PMNS}, L/E, \Delta m_{ij}^2)$$

$$\Delta m_{ij}^2 = m_i^2 - m_j^2$$

# The PMNS Matrix

- With three flavor states and three mass states,  $U_{PMNS}$  is a 3 X 3 matrix

$$U_{PMNS} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \times \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta_{CP}} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta_{CP}} & 0 & c_{13} \end{pmatrix} \times \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

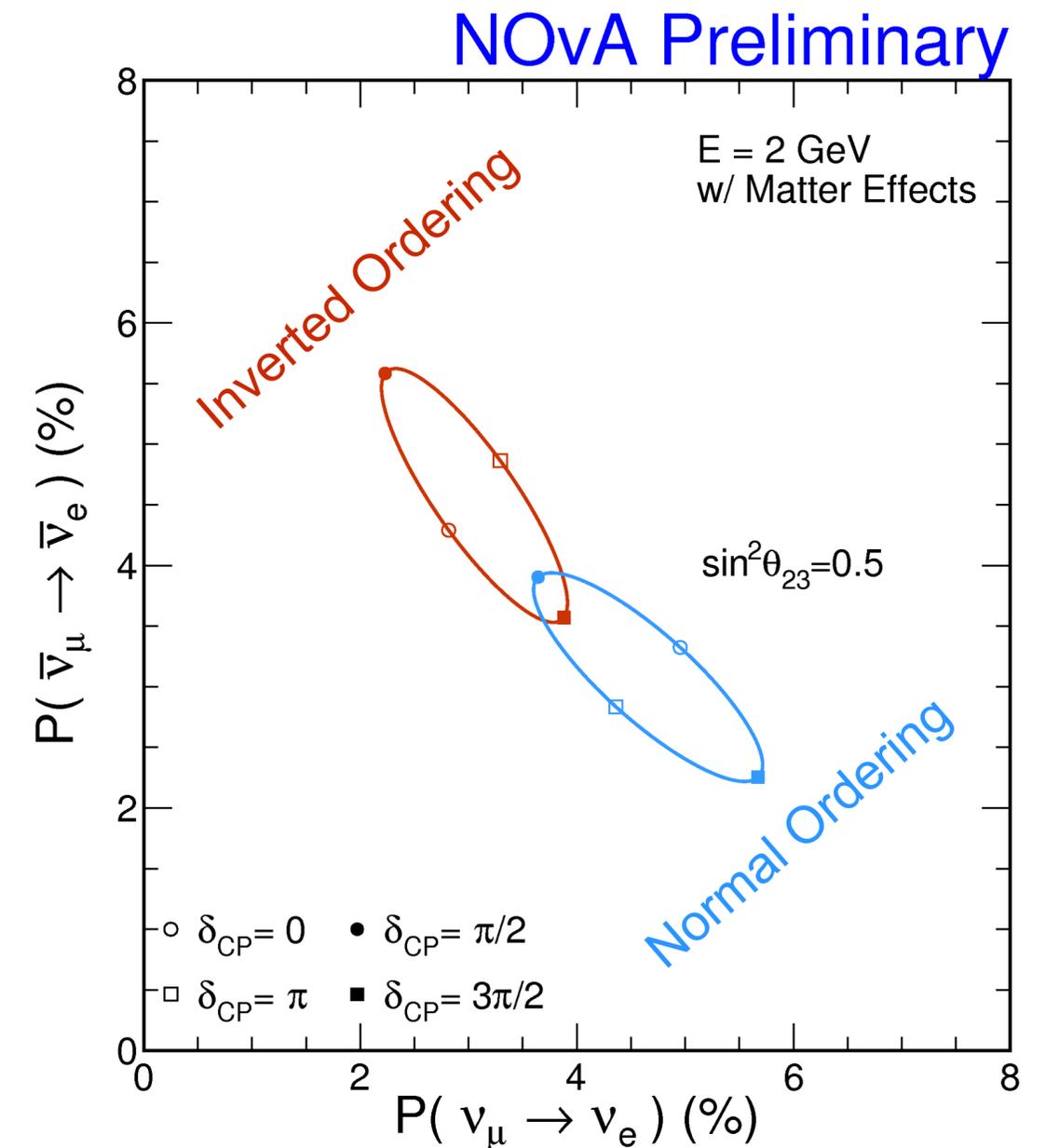
$$c_{ij} = \cos \theta_{ij}, \quad s_{ij} = \sin \theta_{ij}$$

- Three mixing angles:  $\theta_{12}$ ,  $\theta_{13}$ ,  $\theta_{23}$ ; and a (potentially) CP-violating phase:  $\delta_{CP}$ 
  - CP = charge inversion + parity inversion
    - Charge transformation: change particle to antiparticle (and vice versa)
    - Parity transformation: mirror all spatial coordinates
  - CP violation in the neutrino sector might explain why the universe is matter-dominated

# Decoding our measurements

$\bar{\nu}_e$  appearance

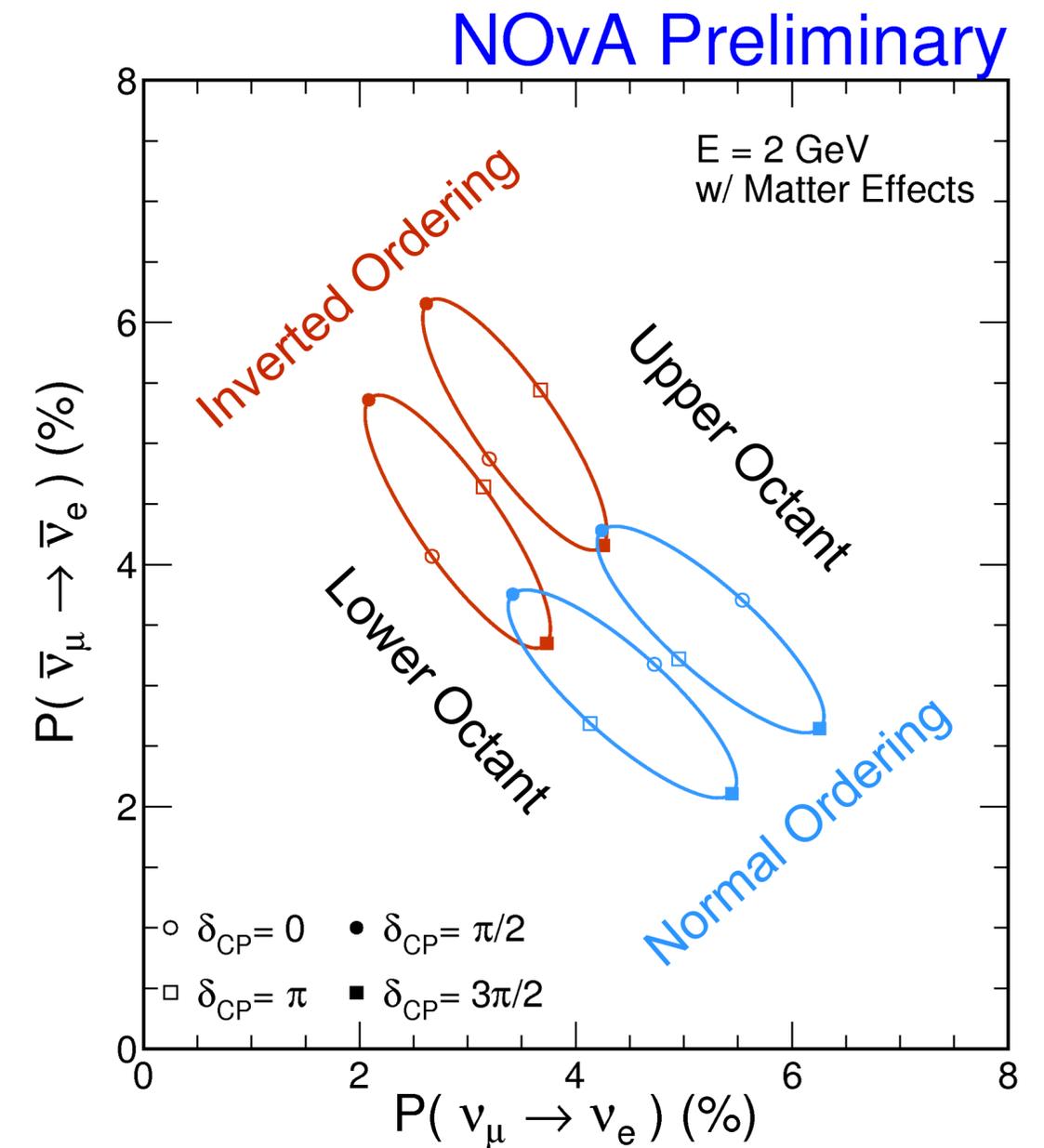
- Oscillations in matter pull the ellipses apart
  - Due to coherent forward scattering on electrons
  - NO increases the neutrino oscillation probability and suppresses antineutrino oscillations
  - IO is reversed



# Decoding our measurements

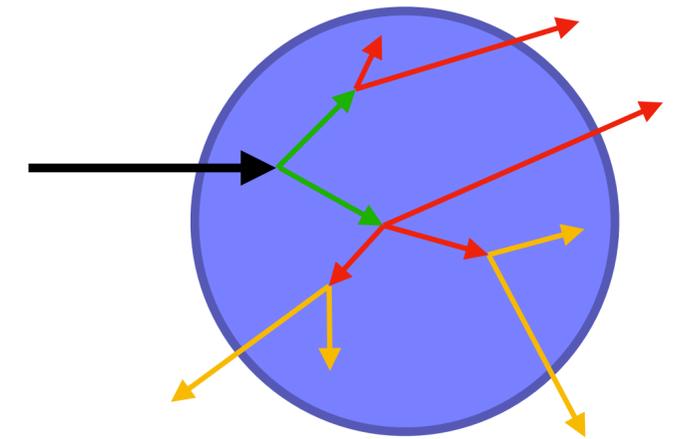
$\bar{\nu}_e$  appearance

- If  $\theta_{23}$  does not produce maximal mixing then the choice of octant ( $<$  or  $>$   $45^\circ$ ) further splits the ellipses
- There are still degenerate regions near the middle



# Neutron inelastic scattering

- Geant4 uses two models based on KE
  - Below 20 MeV: Data-driven "High precision"
  - 20 MeV-10 GeV: Intranuclear cascade
- NOvA uses the Bertini intranuclear cascade, which is split into sub-models
  - Nucleon-nucleon interactions: basically two-body



Beam flux  
simulation

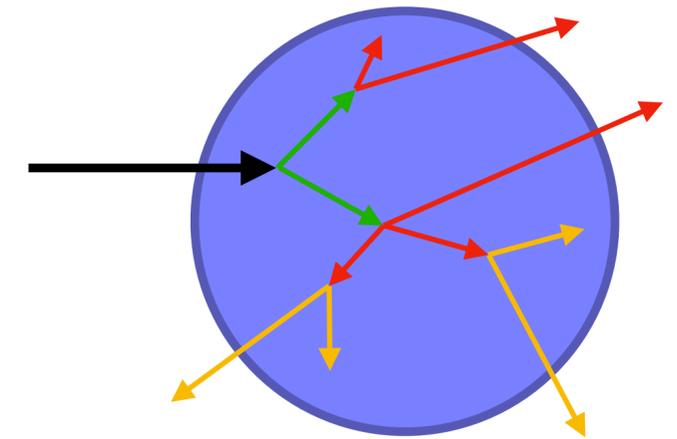
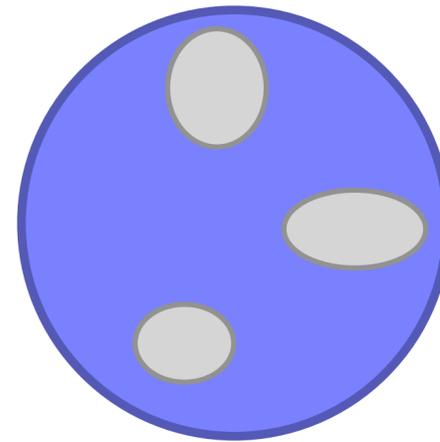
Neutrino interactions:  
GENIE

Particle propagation:  
Geant4

Custom light model  
and electronics

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- NOvA uses the Bertini intranuclear cascade, which is split into sub-models
  - Nucleon-nucleon interactions: basically two-body
  - Pre-equilibrium: balances the residual excitons



Beam flux simulation

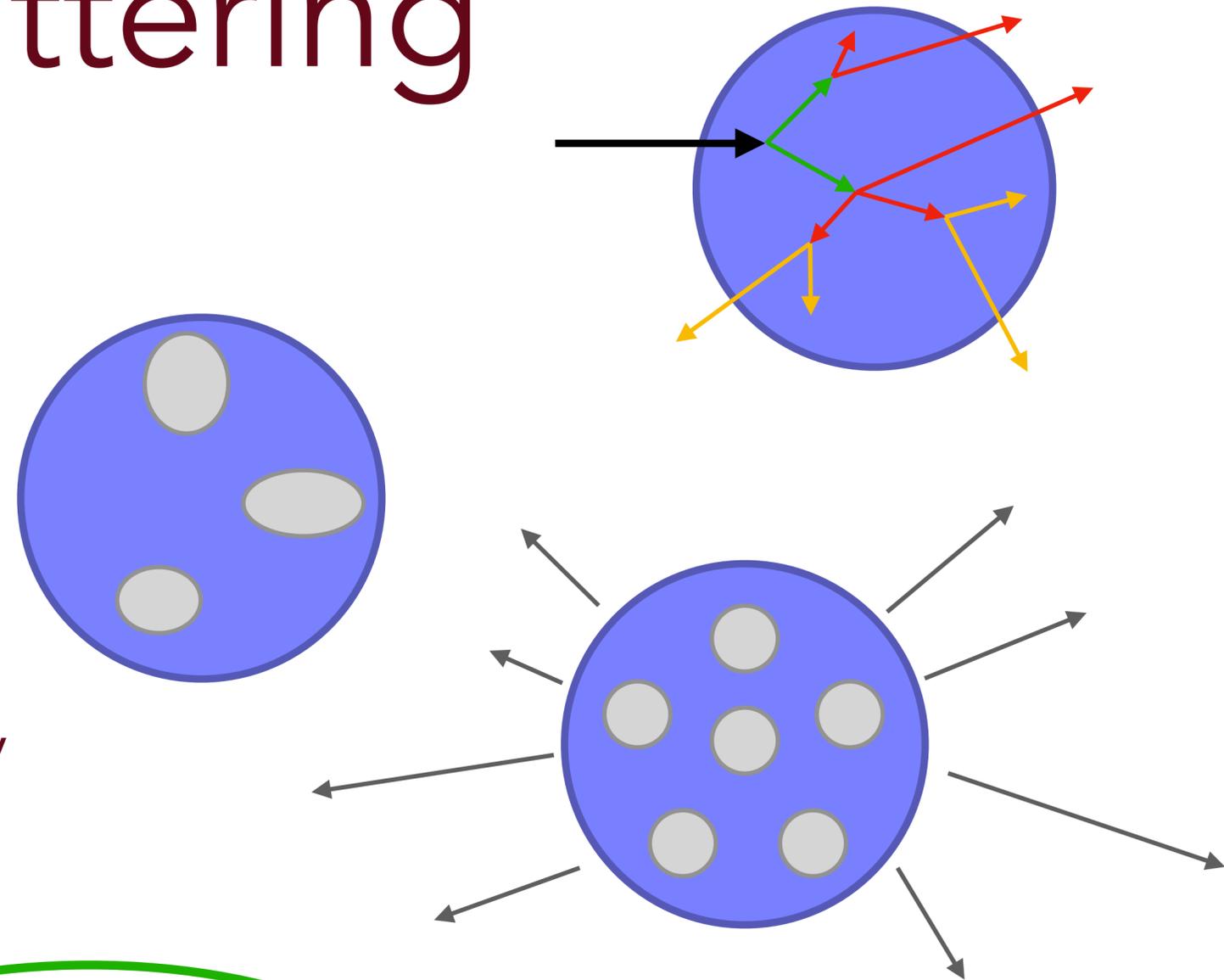
Neutrino interactions:  
GENIE

Particle propagation:  
Geant4

Custom light model  
and electronics

# Neutron inelastic scattering

- Geant4 uses two models based on KE
  - Below 20 MeV: Data-driven "High precision"
  - 20 MeV-10 GeV: Intranuclear cascade
- NOvA uses the Bertini intranuclear cascade, which is split into sub-models
  - Nucleon-nucleon interactions: basically two-body
  - Pre-equilibrium: balances the residual excitons
  - Evaporation: heavy particles first, then photons



Beam flux simulation

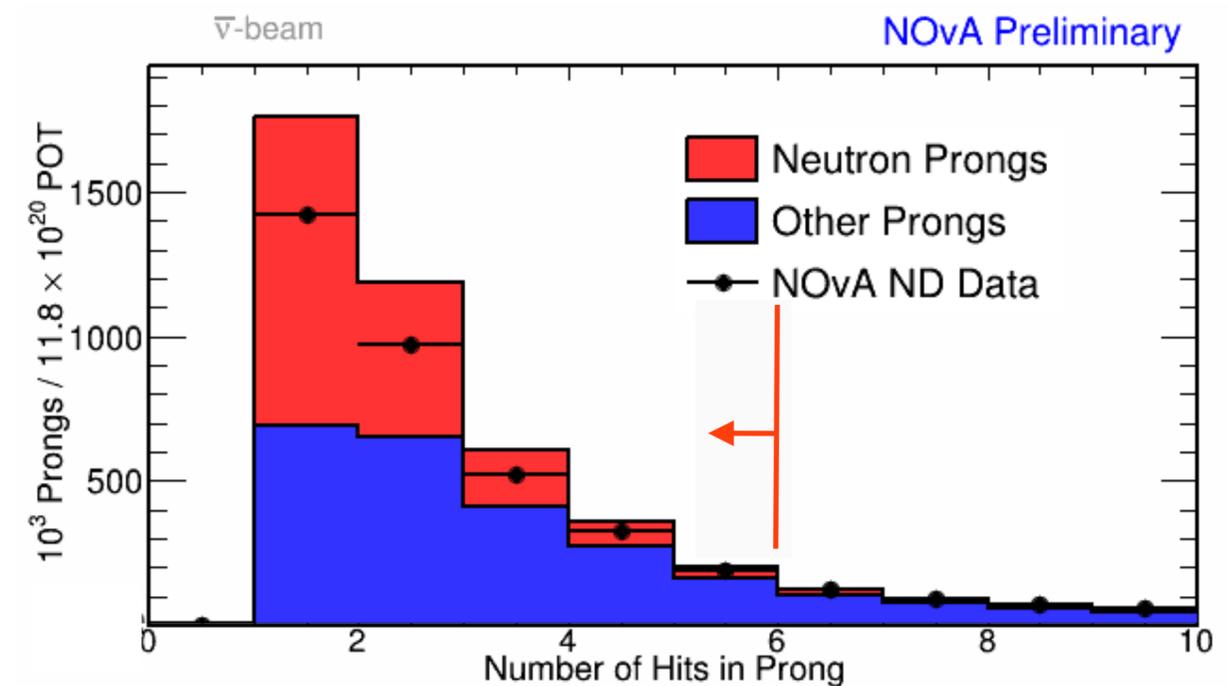
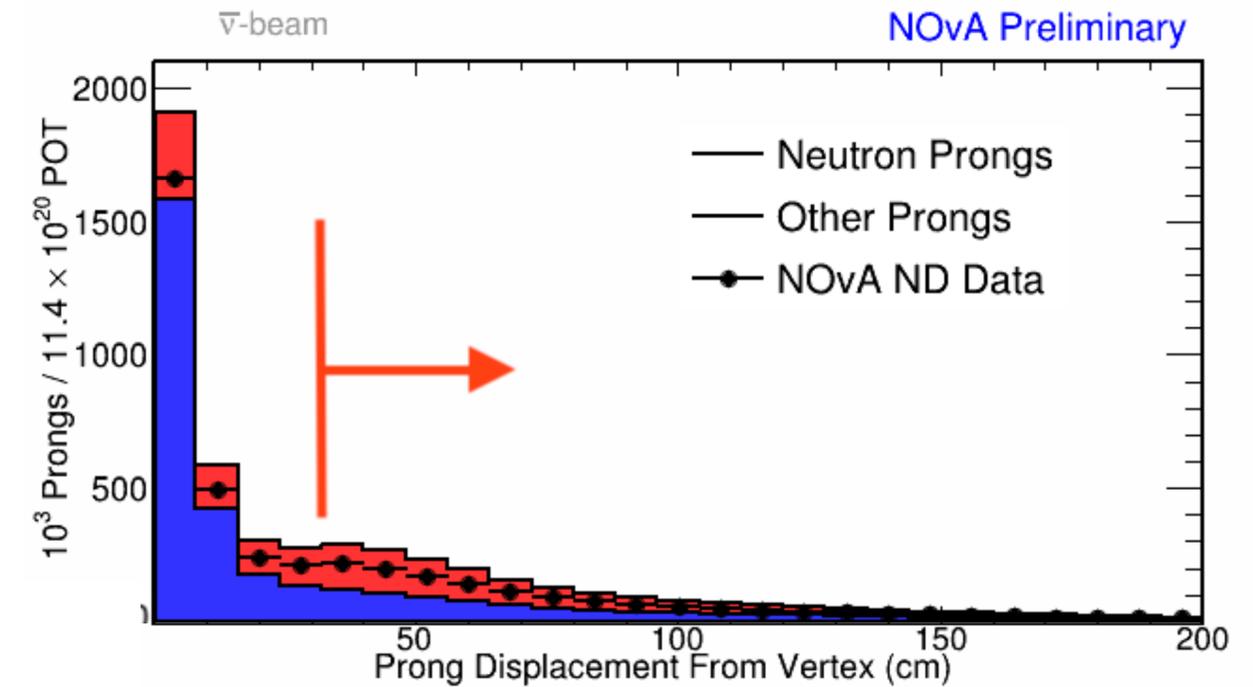
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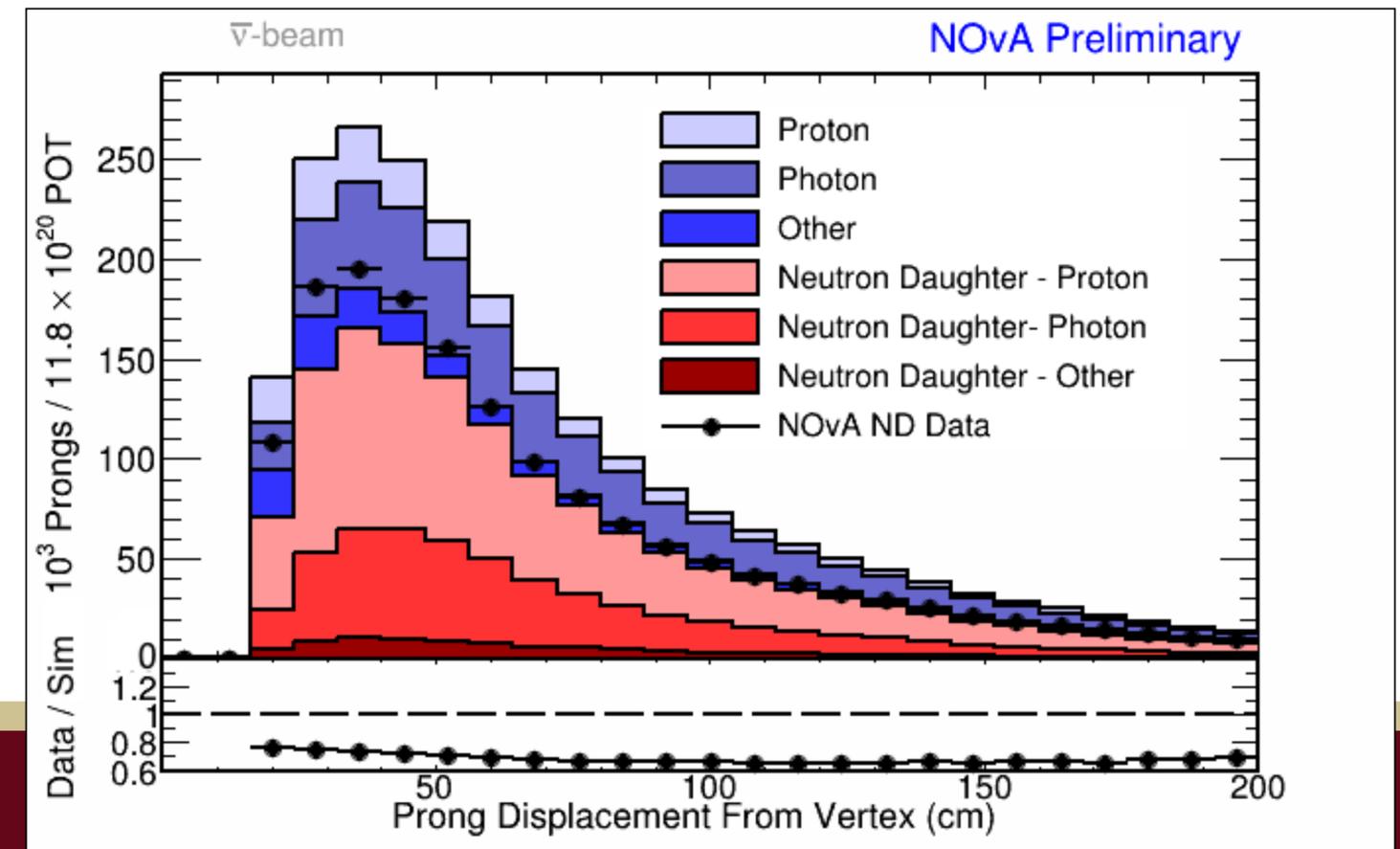
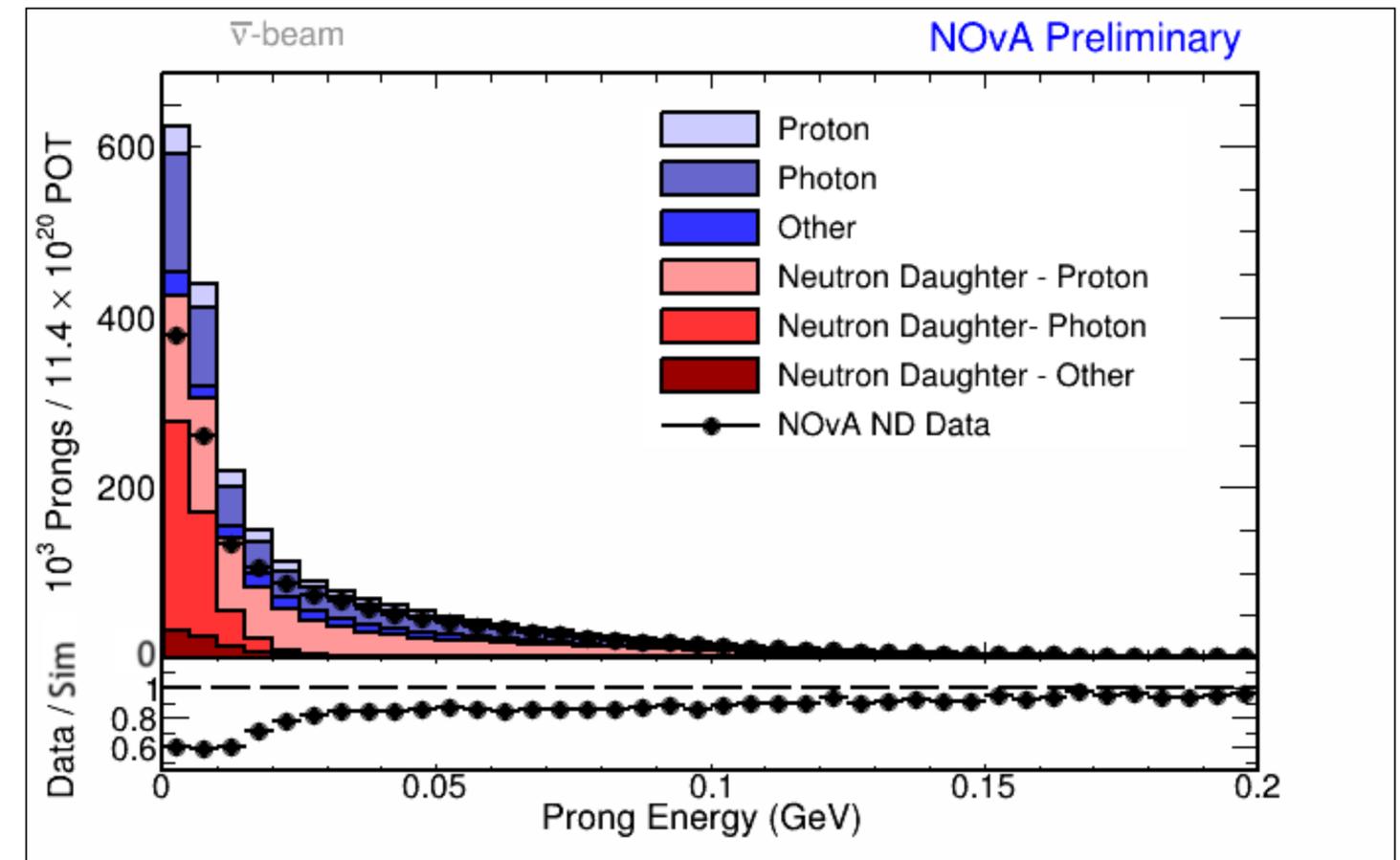
# Selecting neutrons

- Frequent elastic scattering means neutron-related-hits are typically far from the neutrino interaction vertex
  - Select only hit clusters that are  $> 20$  cm away
- Neutron energy depositions are not highly correlated to their kinetic energies
  - Typically leave only a few visible hits
  - Require hit clusters to have fewer than 6 hits



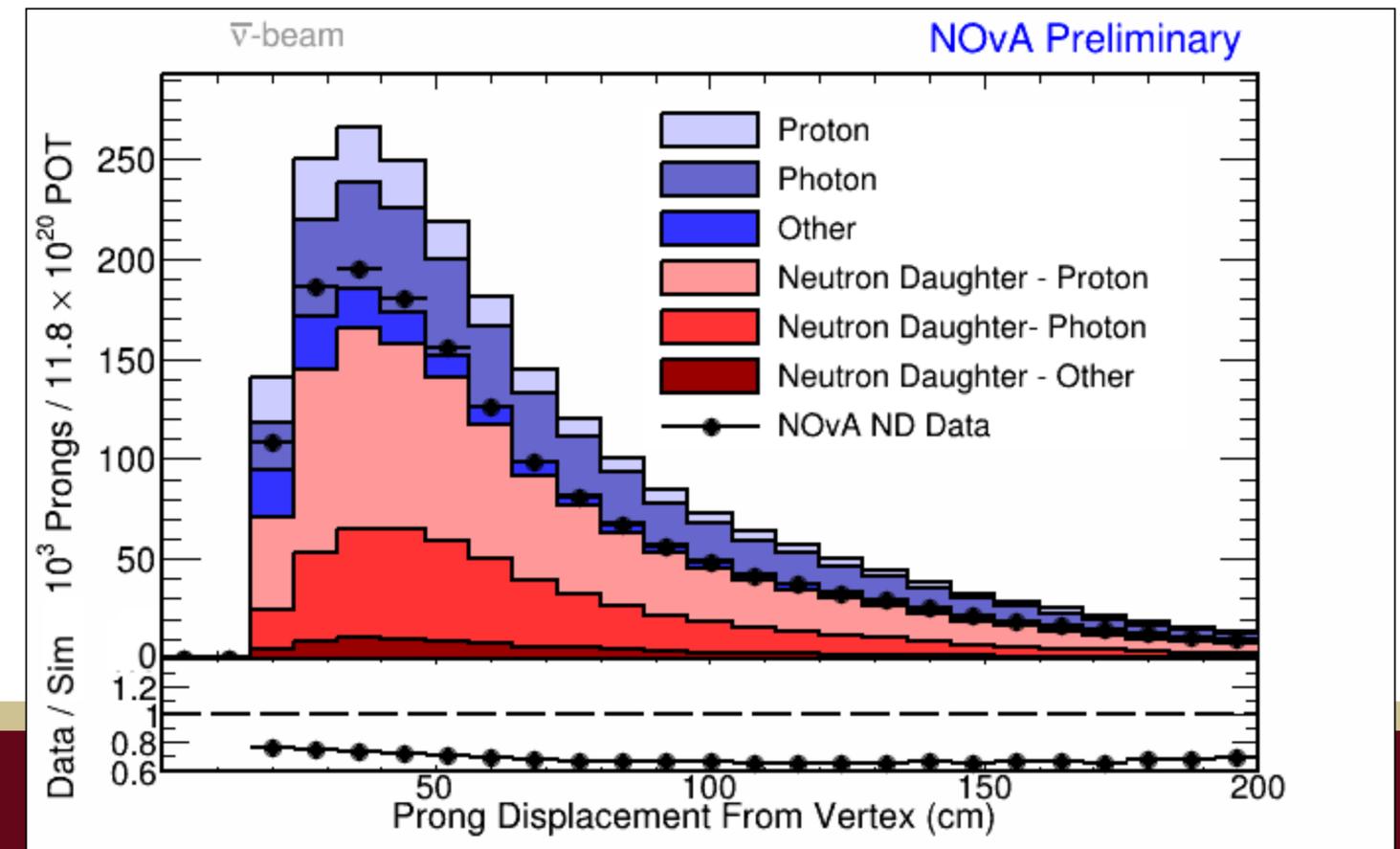
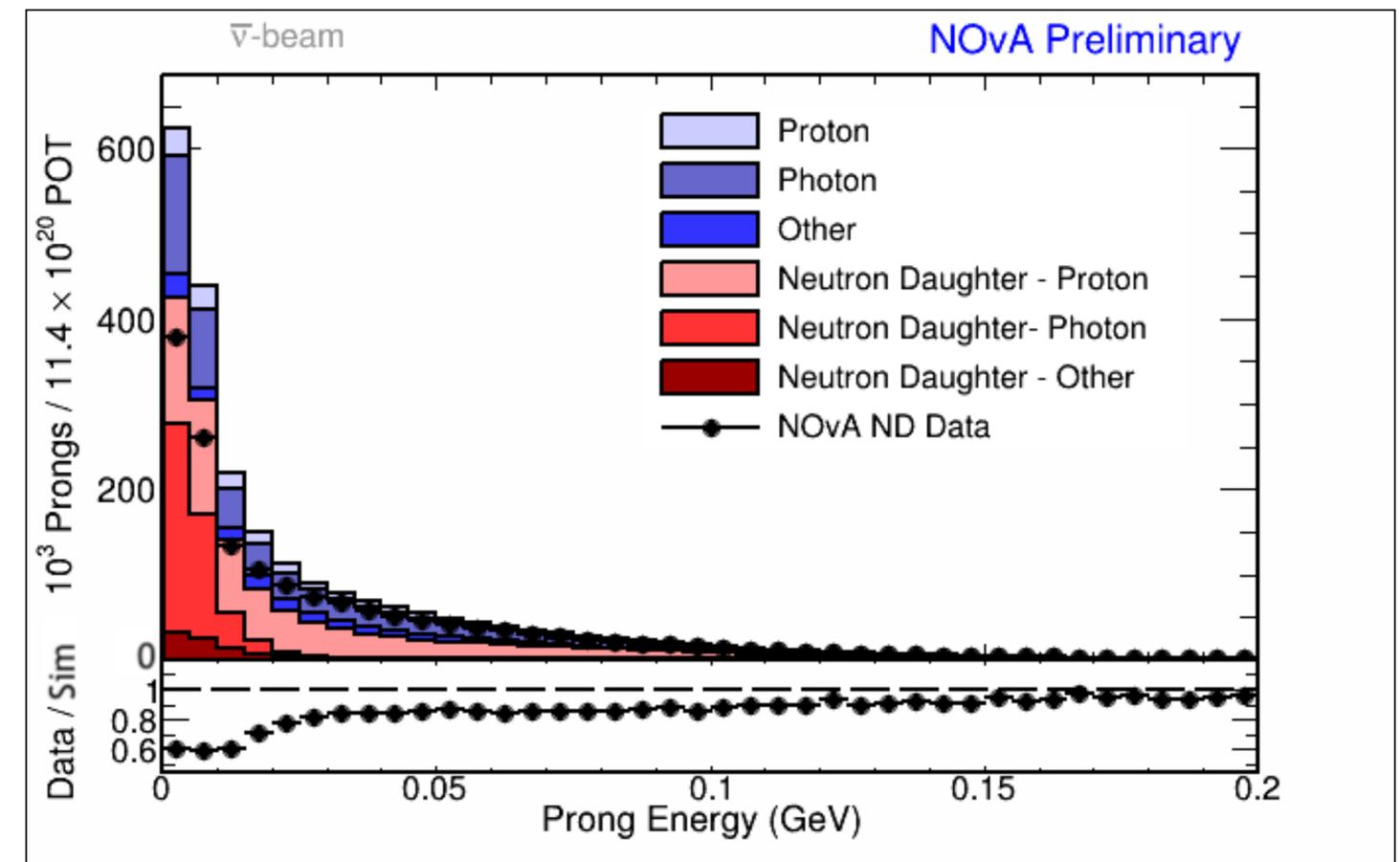
# Selection results

- Apply NOvA's standard  $\bar{\nu}_\mu$  selection at the ND then perform neutron selection
  - 71% efficiency for selecting visible neutrons
  - Selected sample is 61% pure



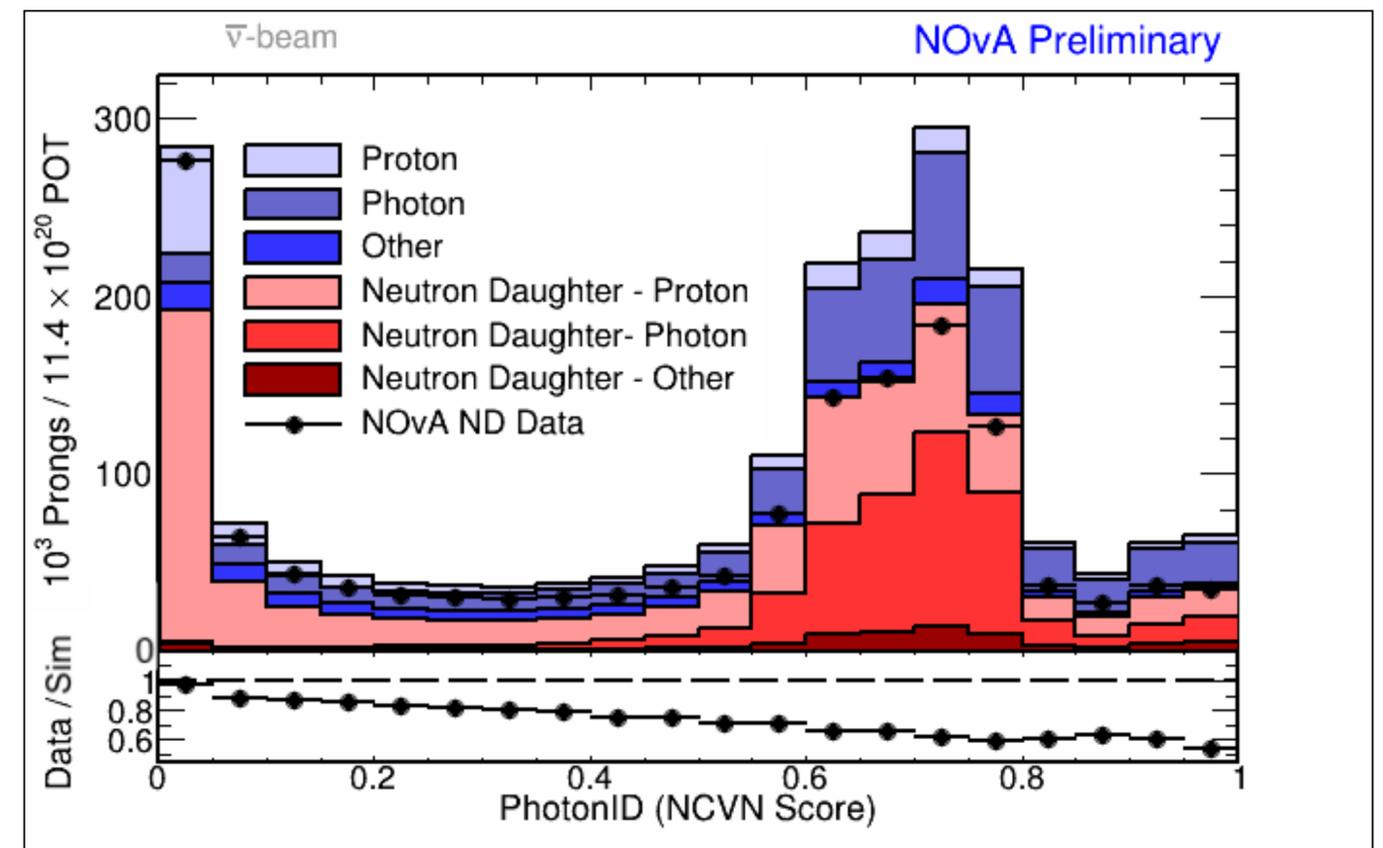
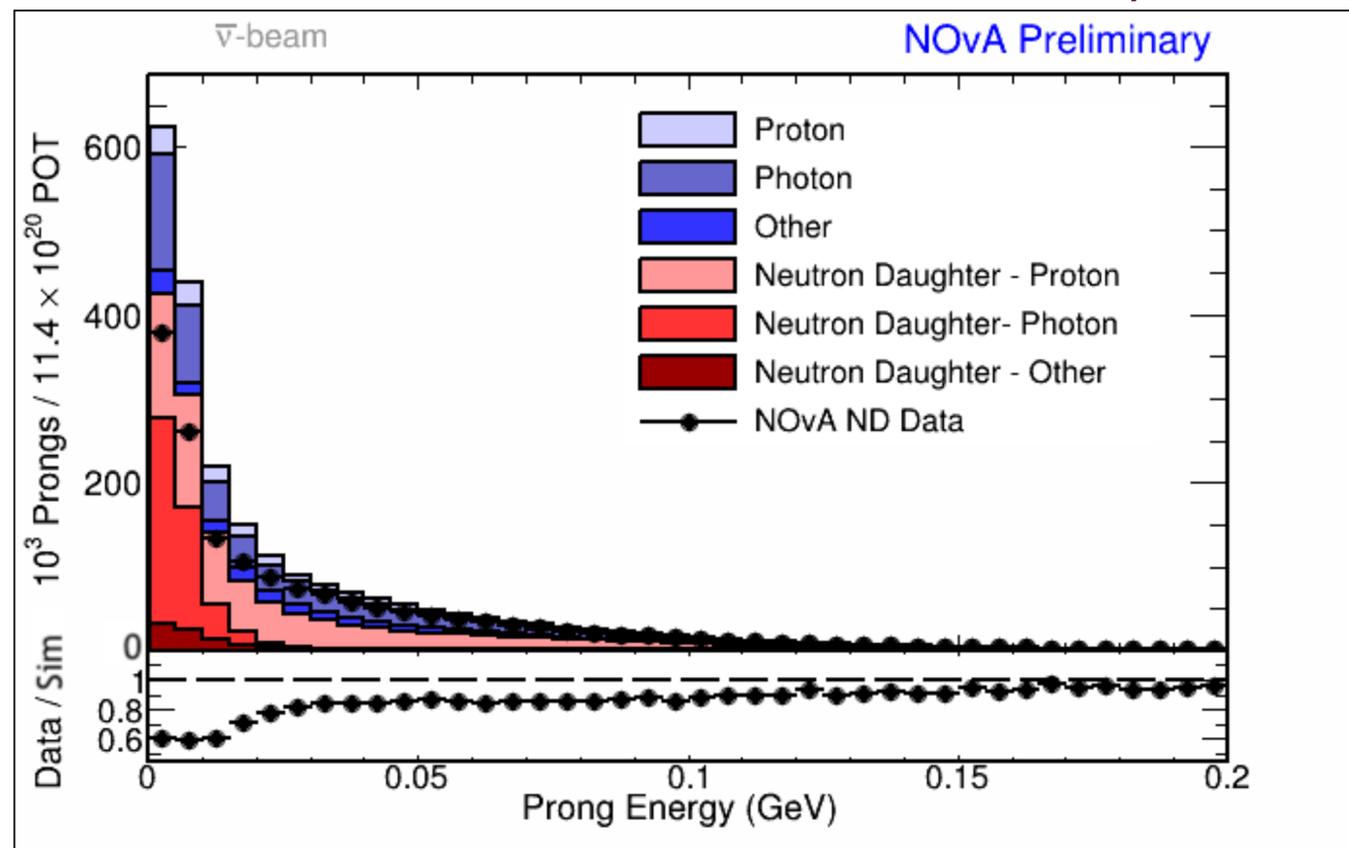
# Selection results

- Apply NOvA's standard  $\bar{\nu}_\mu$  selection at the ND then perform neutron selection
  - 71% efficiency for selecting visible neutrons
  - Selected sample is 61% pure
- Simulation is broken down by particle type
  - Red shades are particles associated with a primary neutron
  - Blue shades are contamination from other primary particles
- Significant data-simulation discrepancy
  - Up to 40% in the lowest energy bins



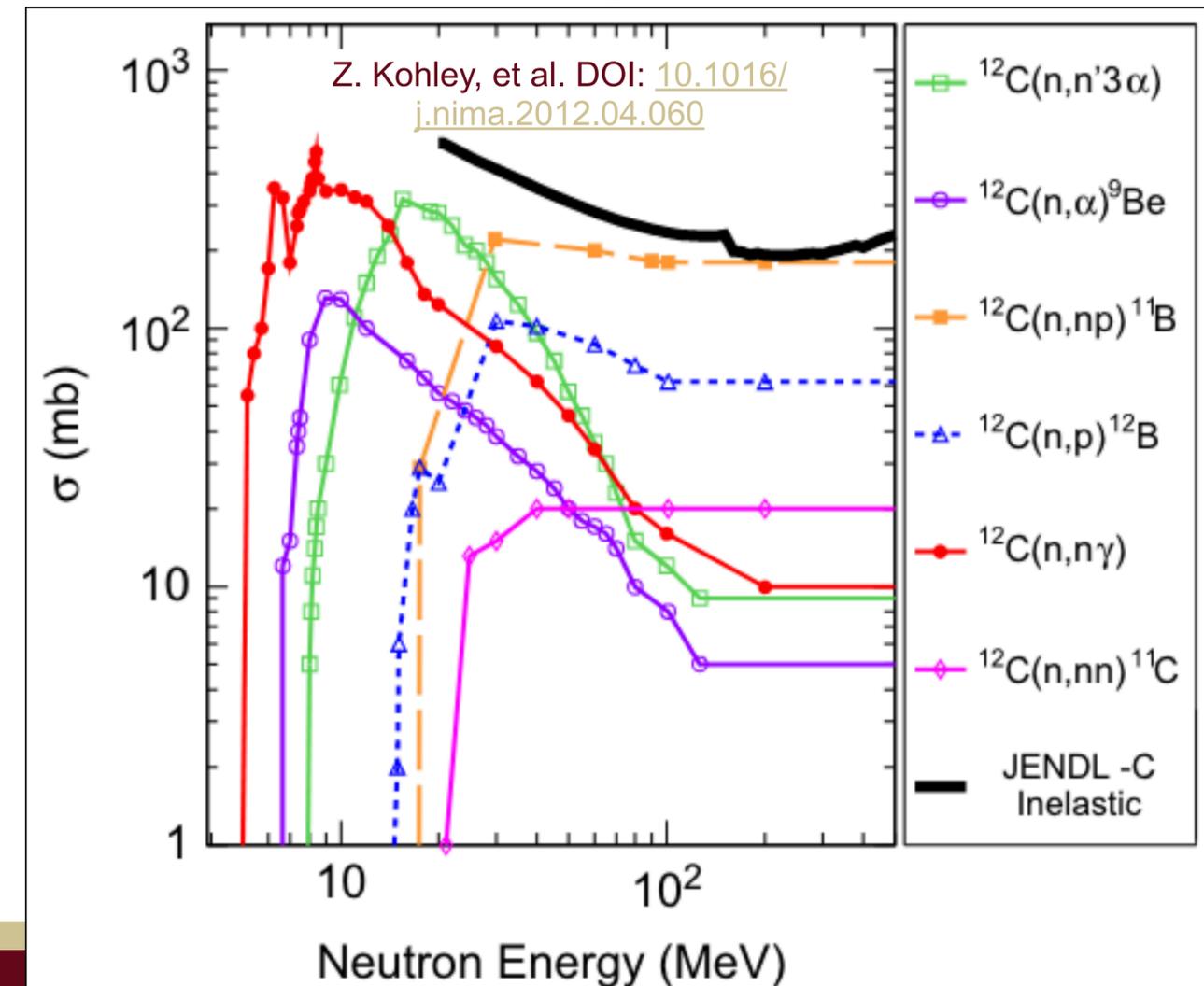
# What's causing the discrepancy?

- De-excitation photons produced by neutron inelastic scattering (medium red) are concentrated at low energy where the simulation excess is the greatest
- We trained a neural network to identify neutron-daughter particle types
  - Over-simulation "follows" the photons



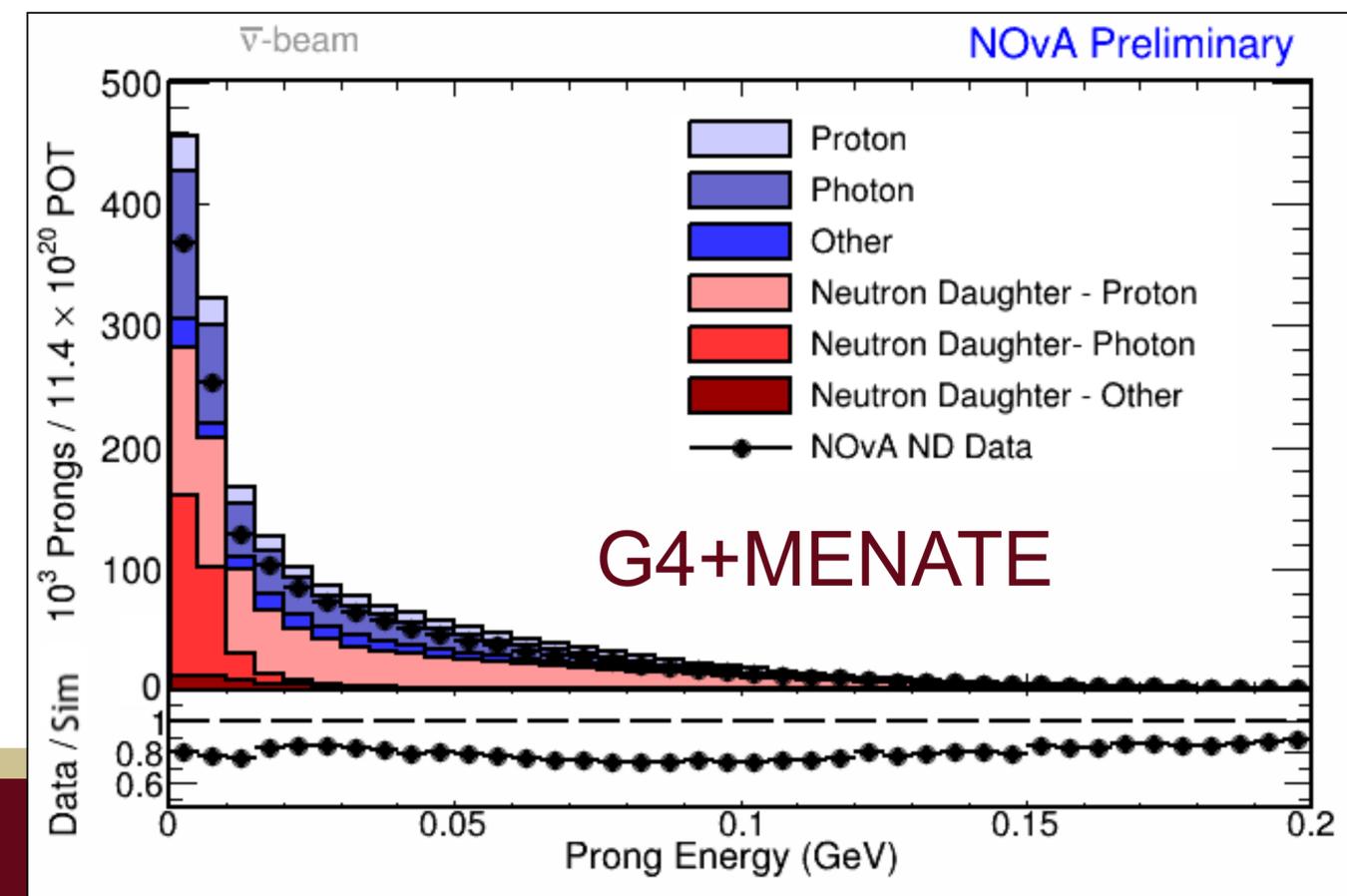
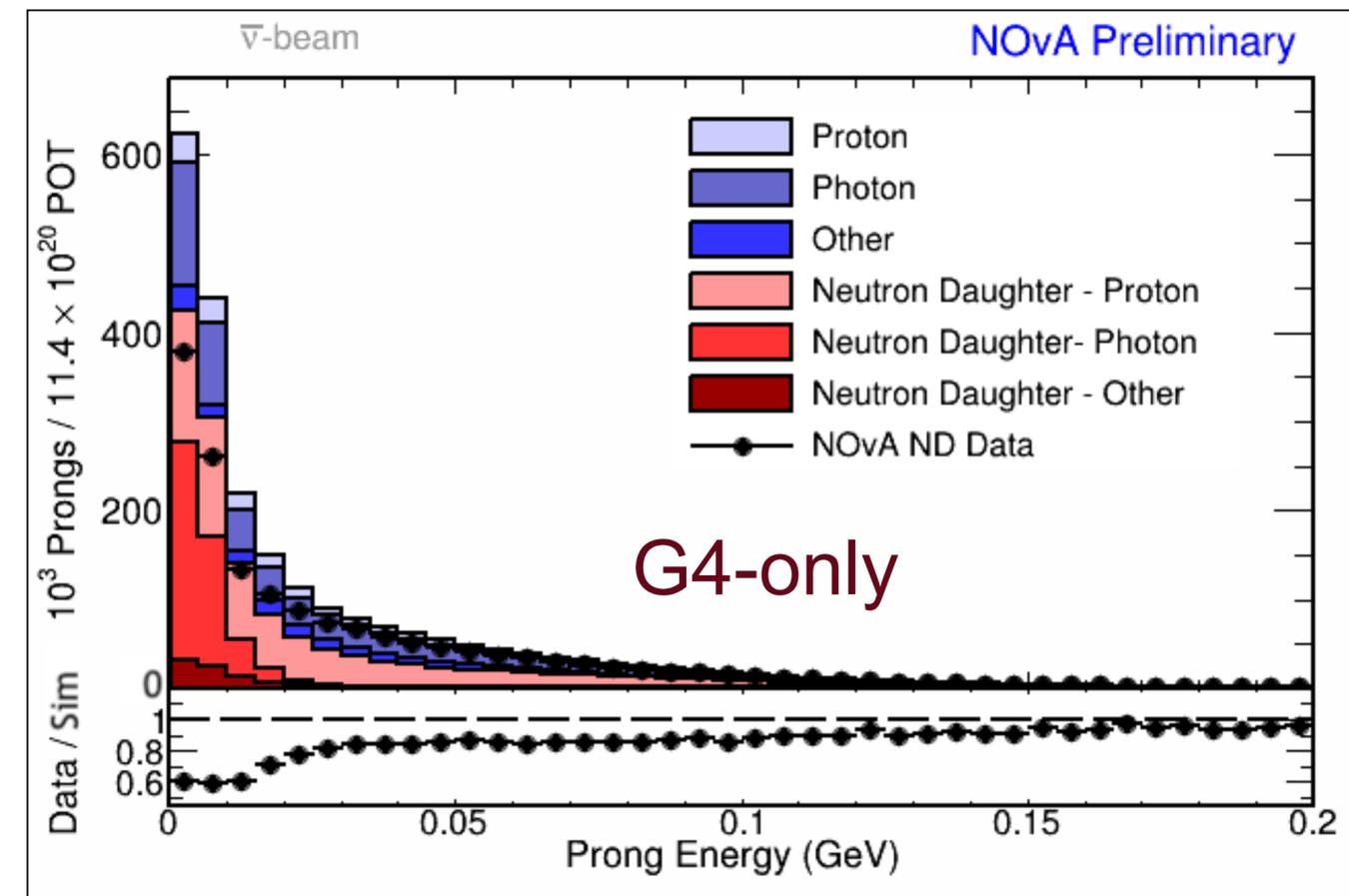
# Alternative neutron inelastic model

- De-excitation photons produced by Geant4:
  - The HP model: determines outgoing particles based on measurements (limited data)
  - Intranuclear cascade: theory-driven and statistically determines outgoing particles while using the overall cross-section
- “MENATE” extends the ethos of the HP model to higher energies
  - But only for neutron-on-carbon interactions (that’s what we have data for)
  - Limits  $n+^{12}\text{C}$  to only 6 final states
  - Our implementation integrates directly with Geant4



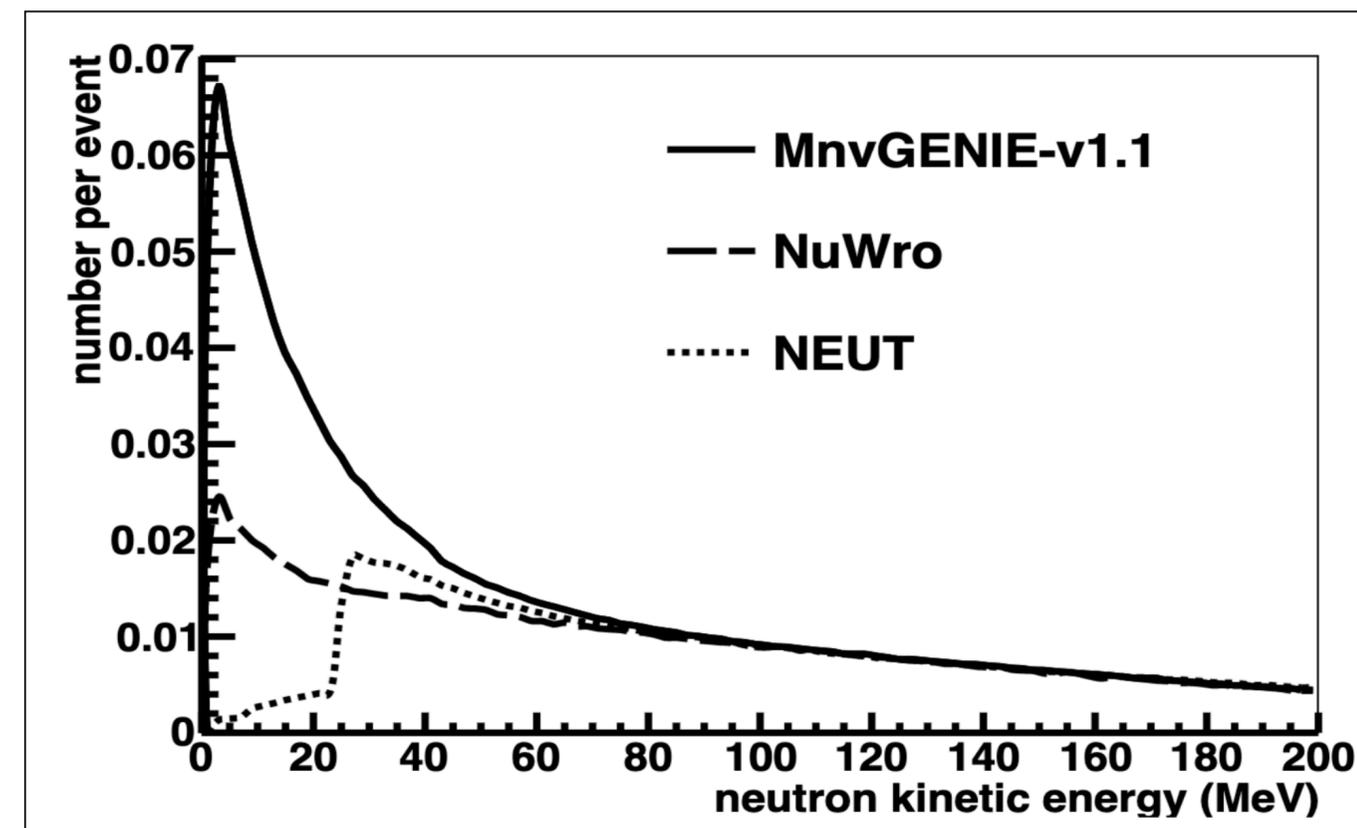
# Comparing to data

- MENATE reduces the over-simulation in the most populous bins
  - Photon production is significantly decreased
- It increases the over-simulation at higher energies
  - Proton production is slightly increased
- Overall, the data-simulation ratio is flatter, but there is still a significant over-simulation
  - Only addressing neutron-on-carbon; there are other targets and projectiles



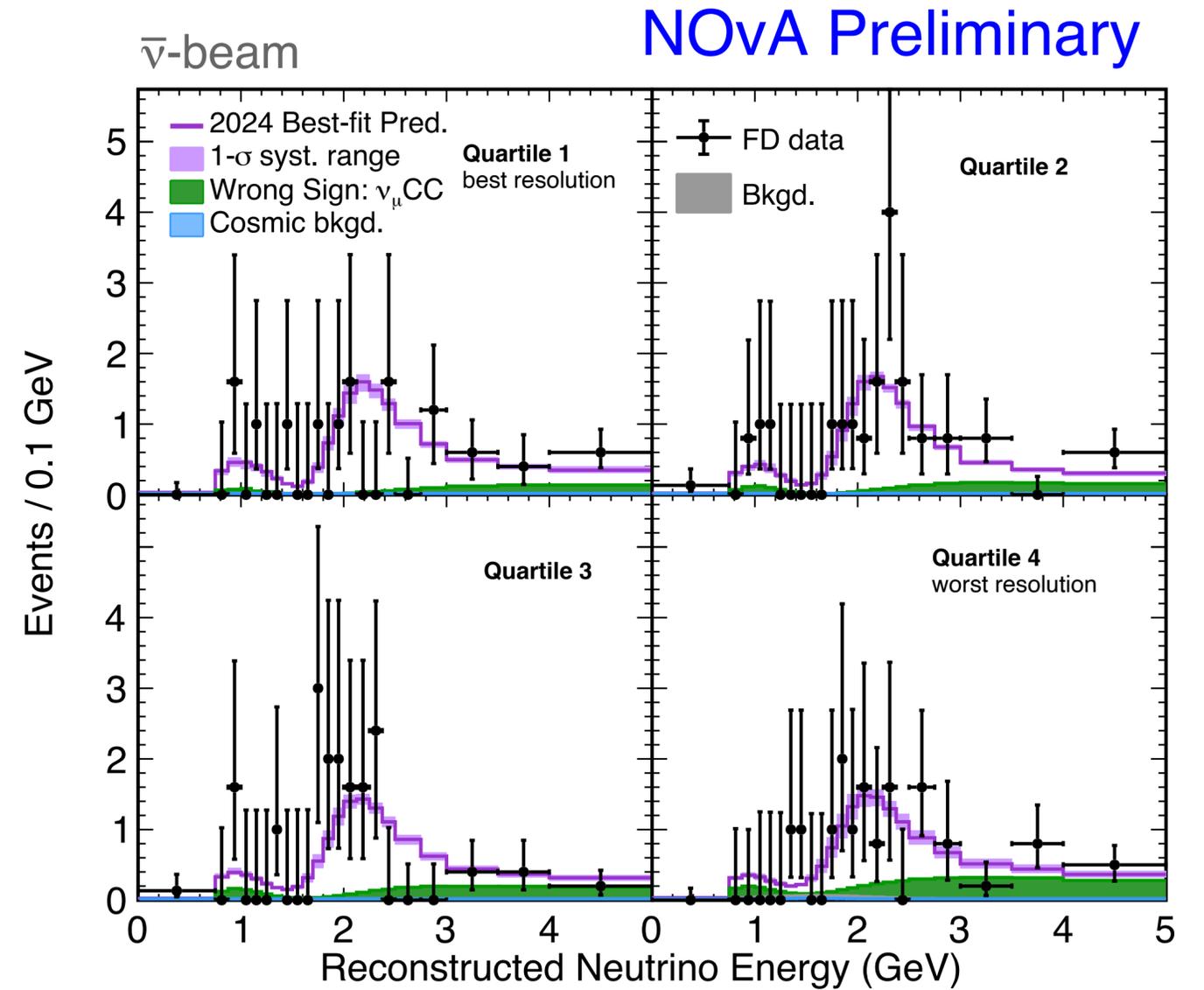
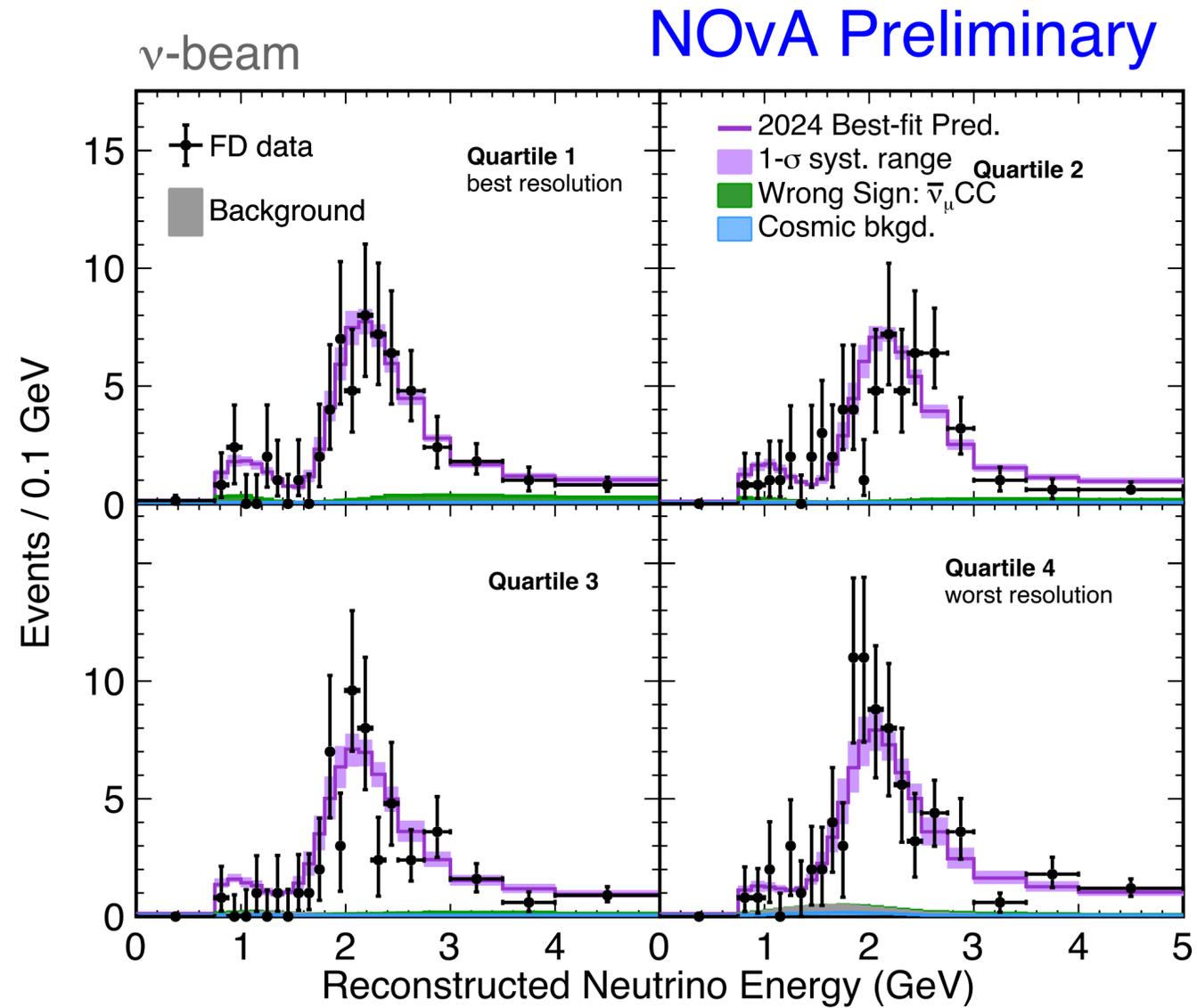
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- Overall, the data-simulation ratio is flatter, but there is still a significant over-simulation
  - Only addressing neutron-on-carbon; there are other targets and projectiles
  - Large uncertainties on neutron production
    - This would produce a “normalization” discrepancy



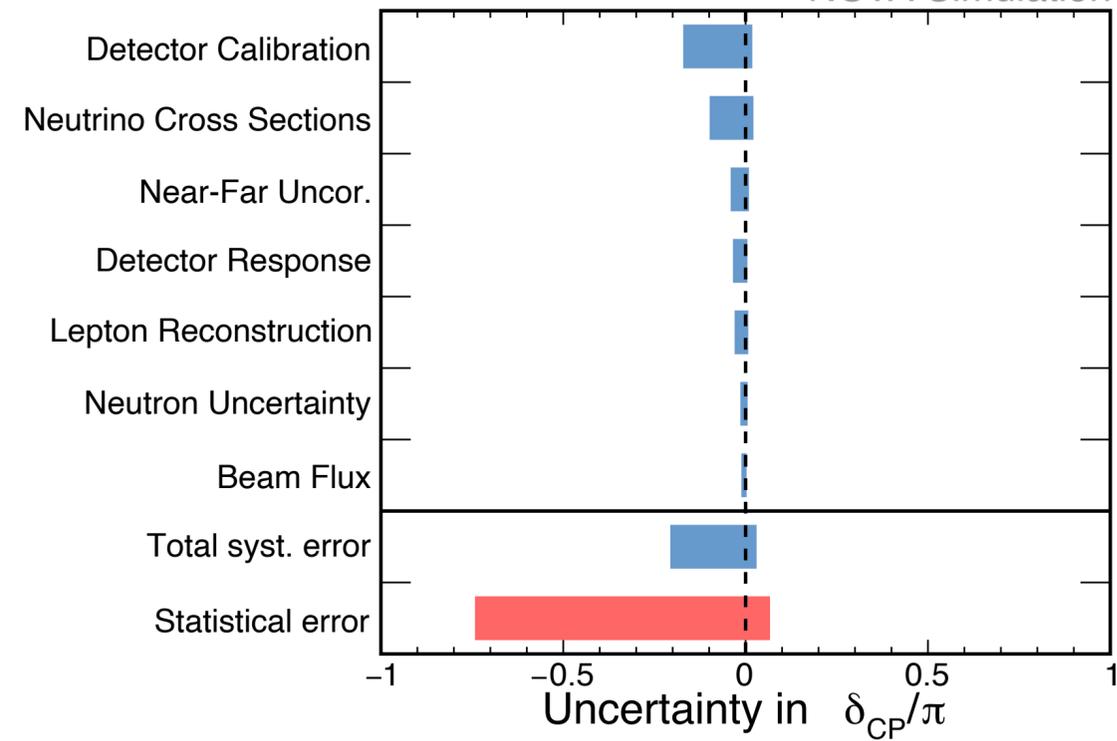
M. Elkins, et al. DOI: [10.1103/PhysRevD.100.052002](https://doi.org/10.1103/PhysRevD.100.052002)

# Numu Ehad Quants

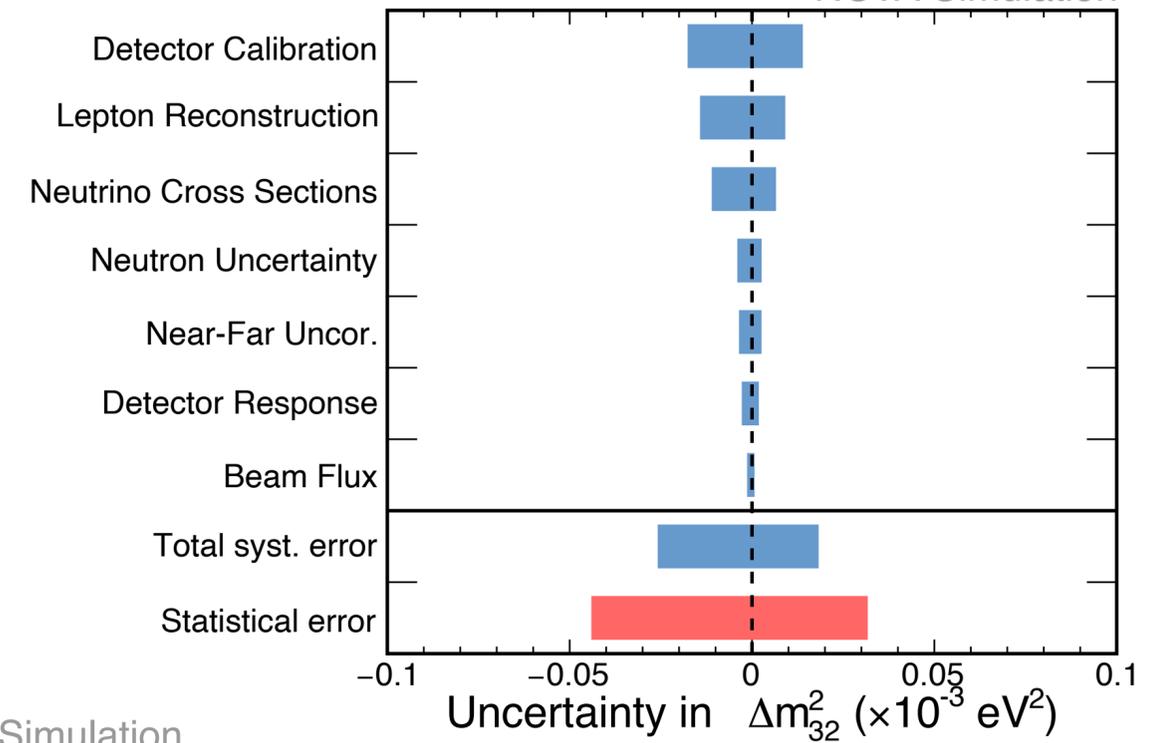


# Freq. Uncertainties

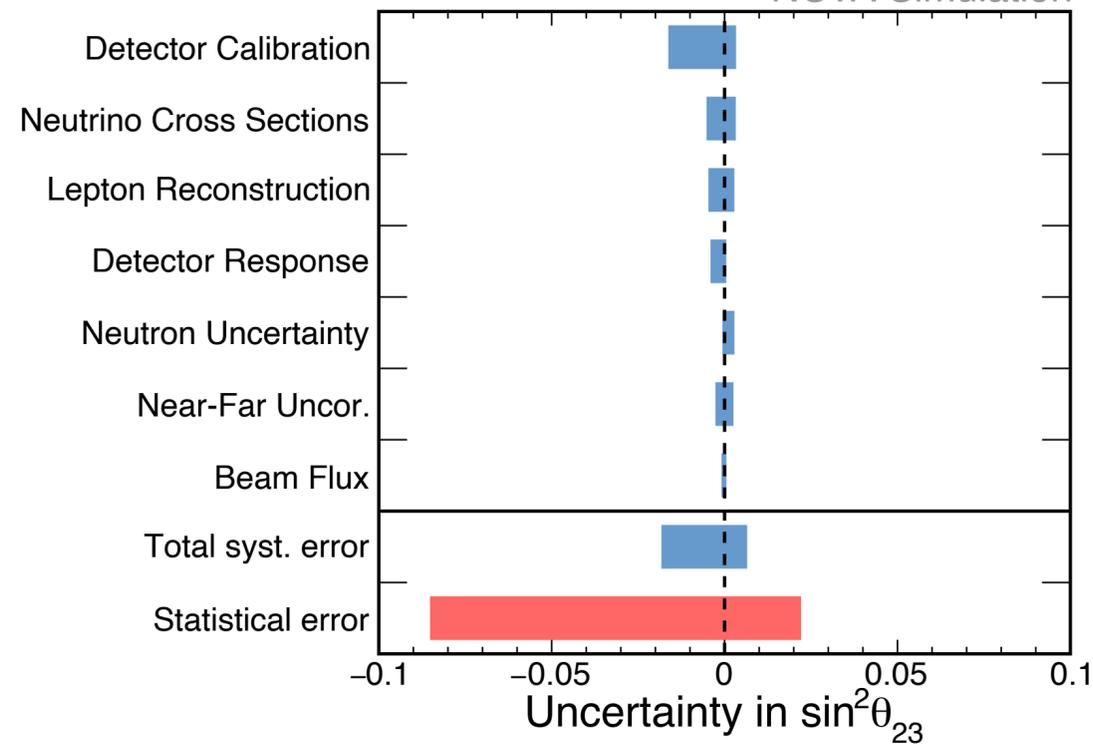
NOvA Simulation



NOvA Simulation



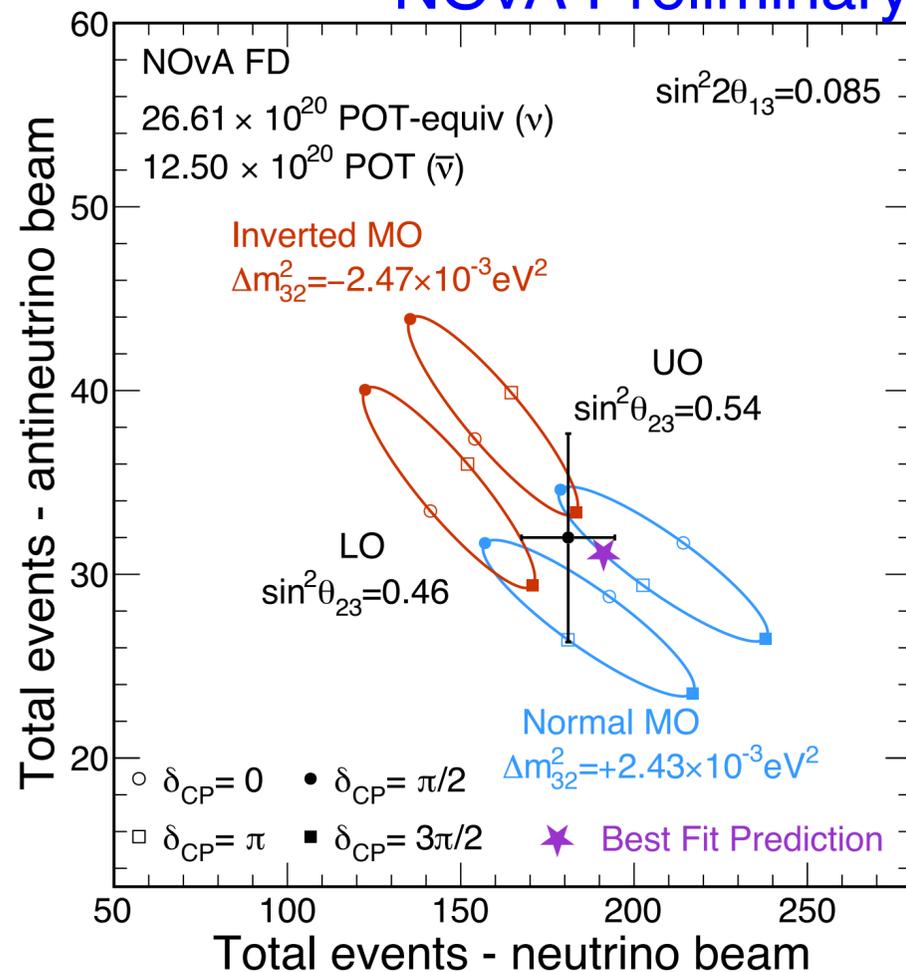
NOvA Simulation



# Bayesian results

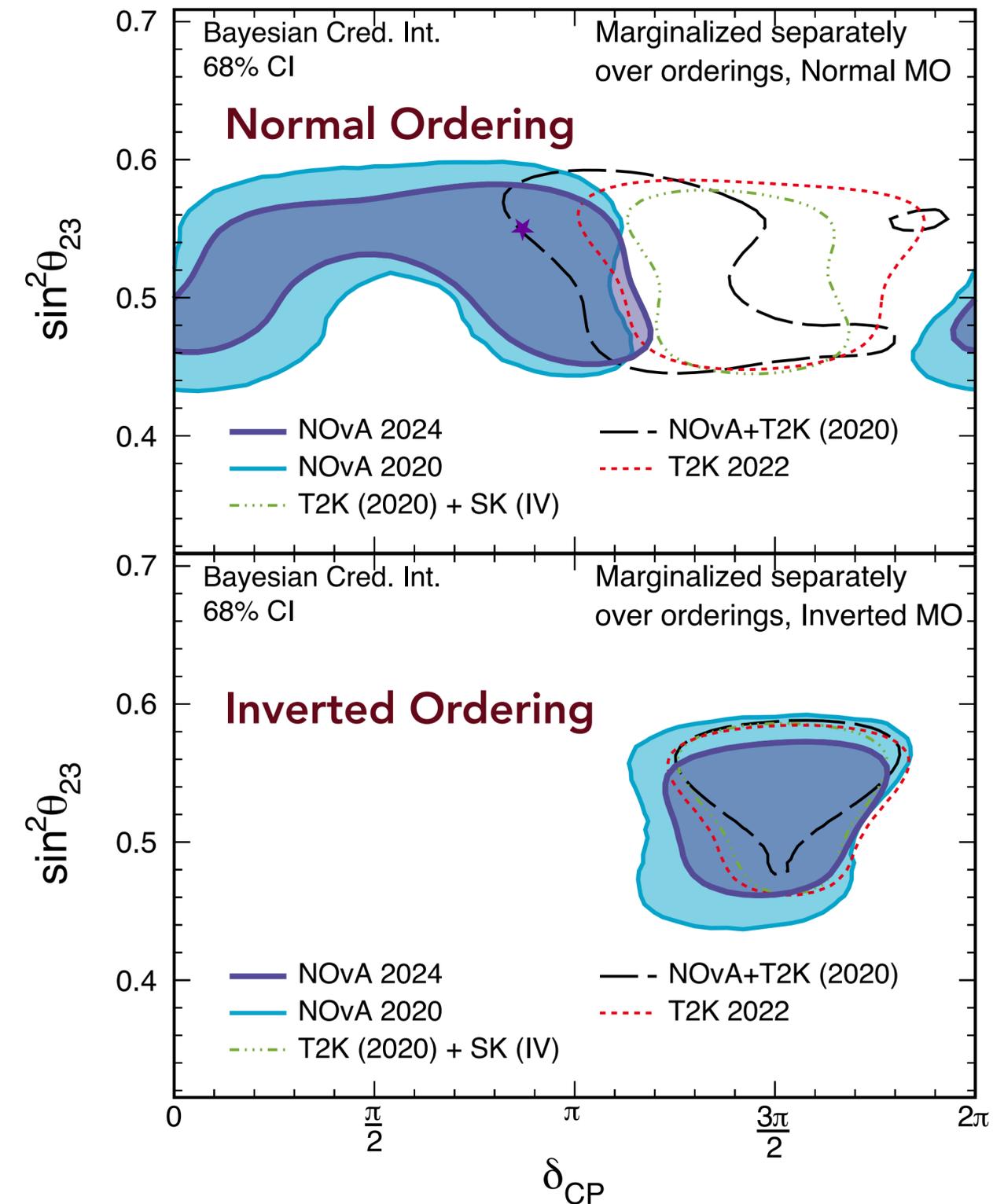
- Prefer CP conserving values of  $\delta_{CP}$  in the normal mass ordering and CP violating values in the inverted

NOvA Preliminary

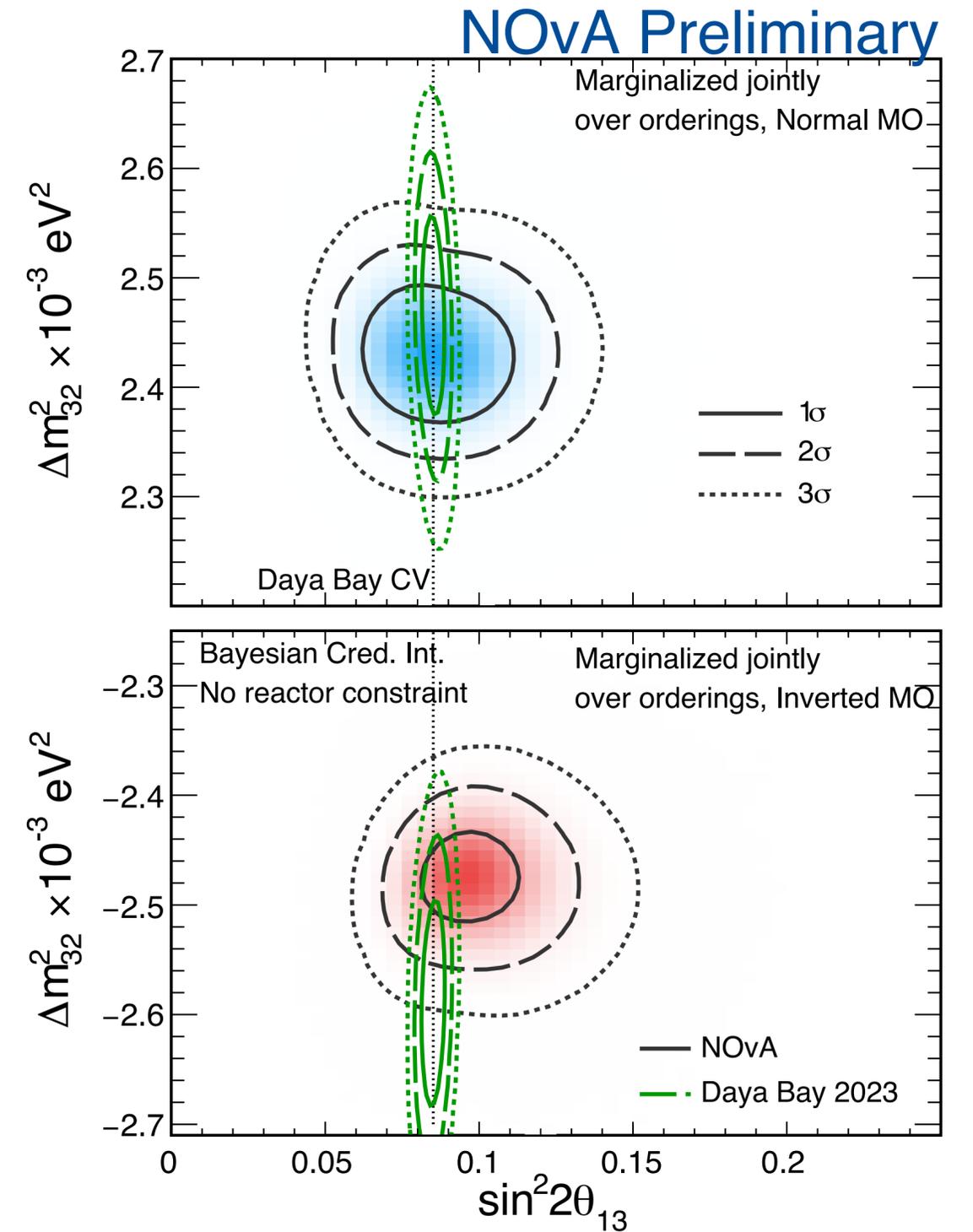
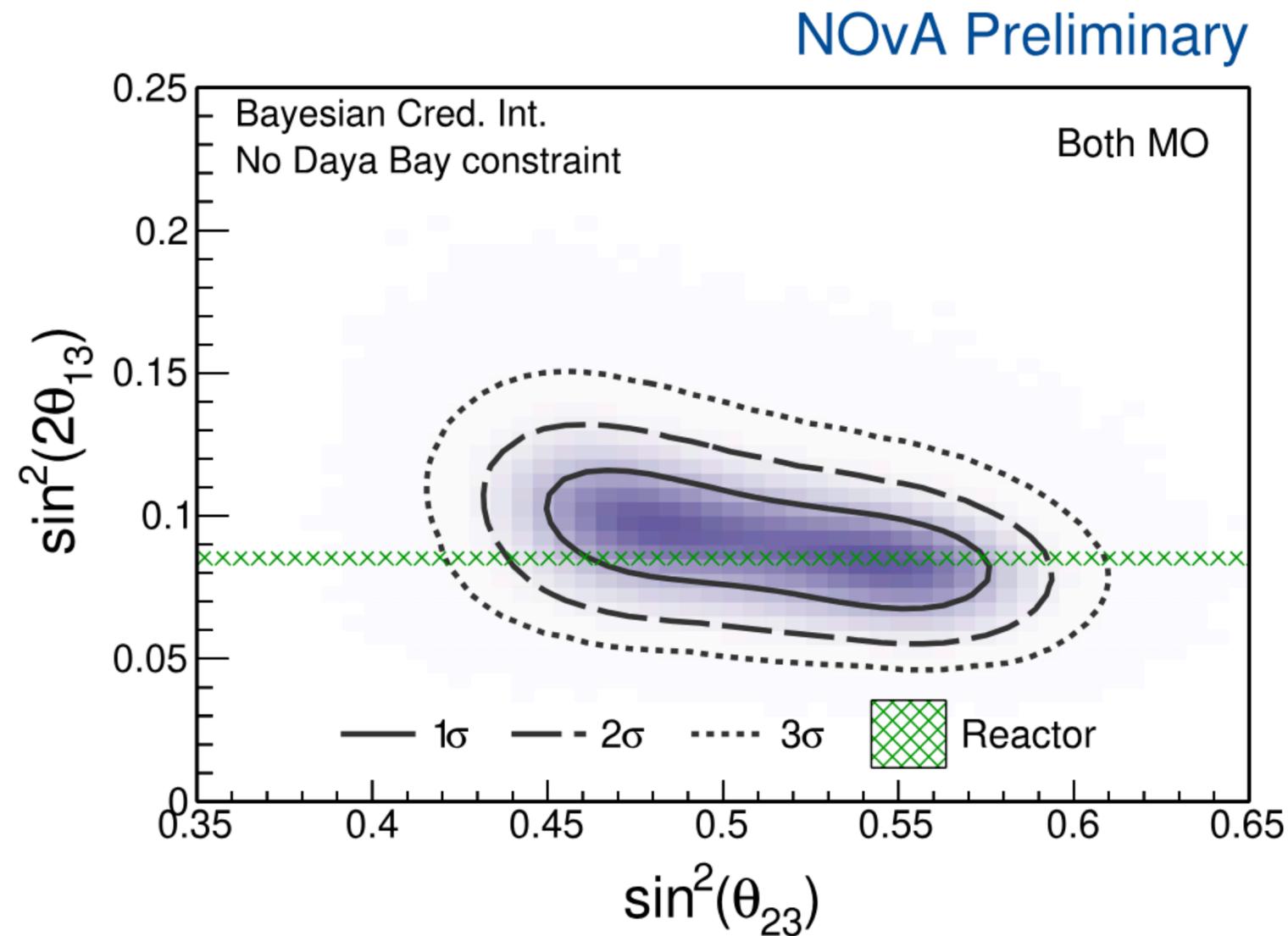


- Our data lie in the highly degenerate region
- Preference for normal mass ordering is enhanced by reactor constraints

NOvA Preliminary

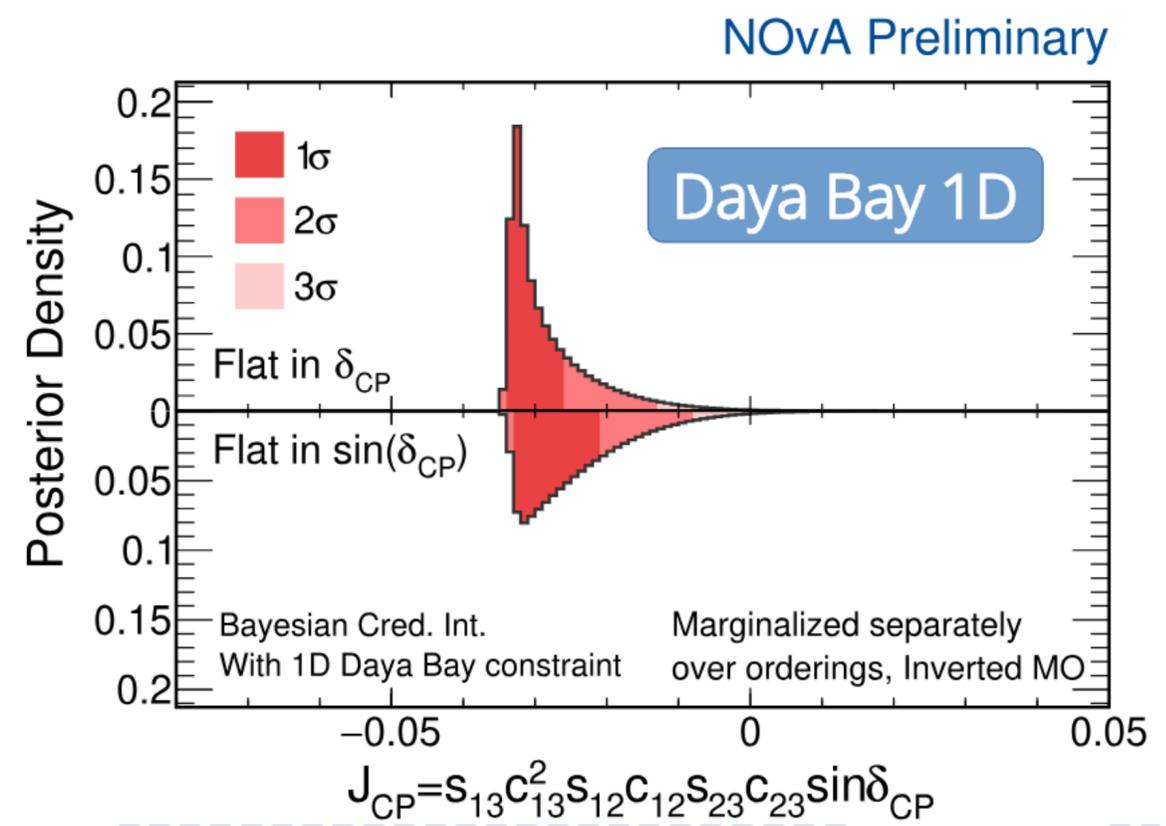
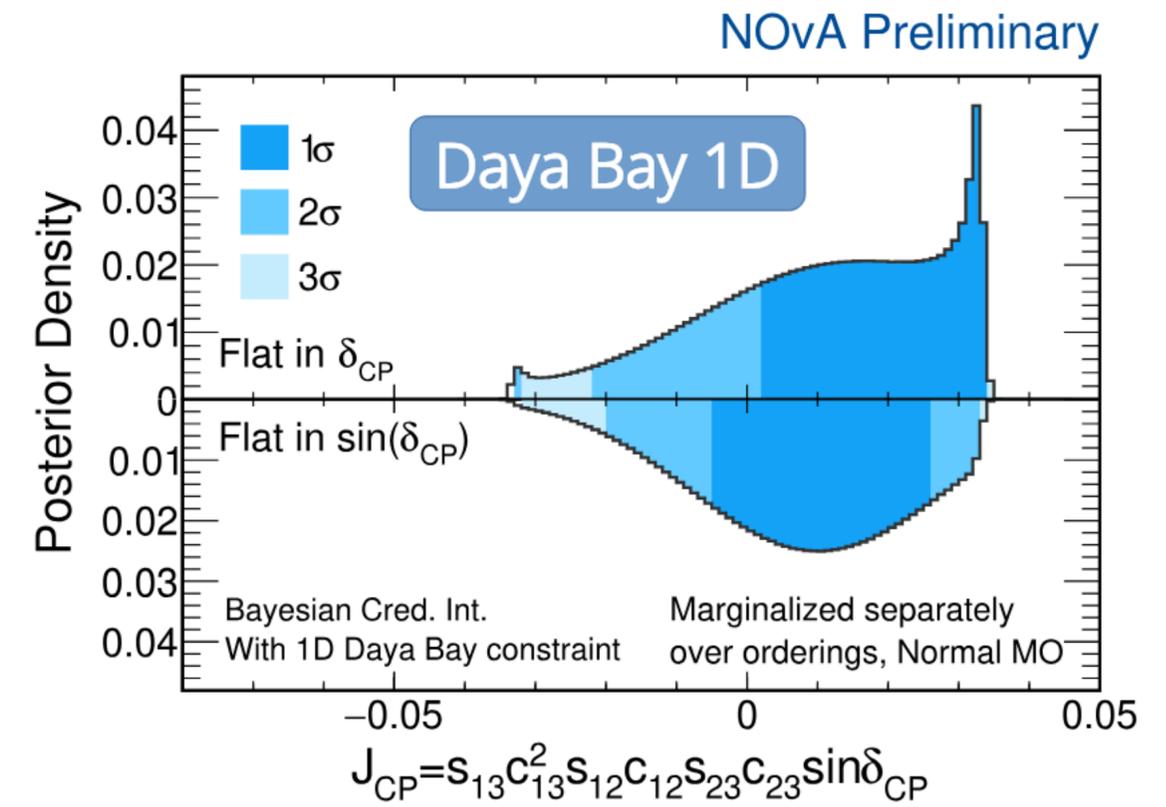


# Daya Bay's constraint



# CP Violation

- The Jarlskog invariant is a parameterization-independent measure of CP violation
  - $J = 0 \rightarrow$  CP-conservation
  - $J \neq 0 \rightarrow$  CP-violation
- Choice of prior
  - Flat in  $\sin\delta_{CP}$  provides data-only preference
  - Flat in  $\delta_{CP}$  (theoretically motivated) is biased away from minimal-CP-violation
  - Doesn't impact interpretation of result

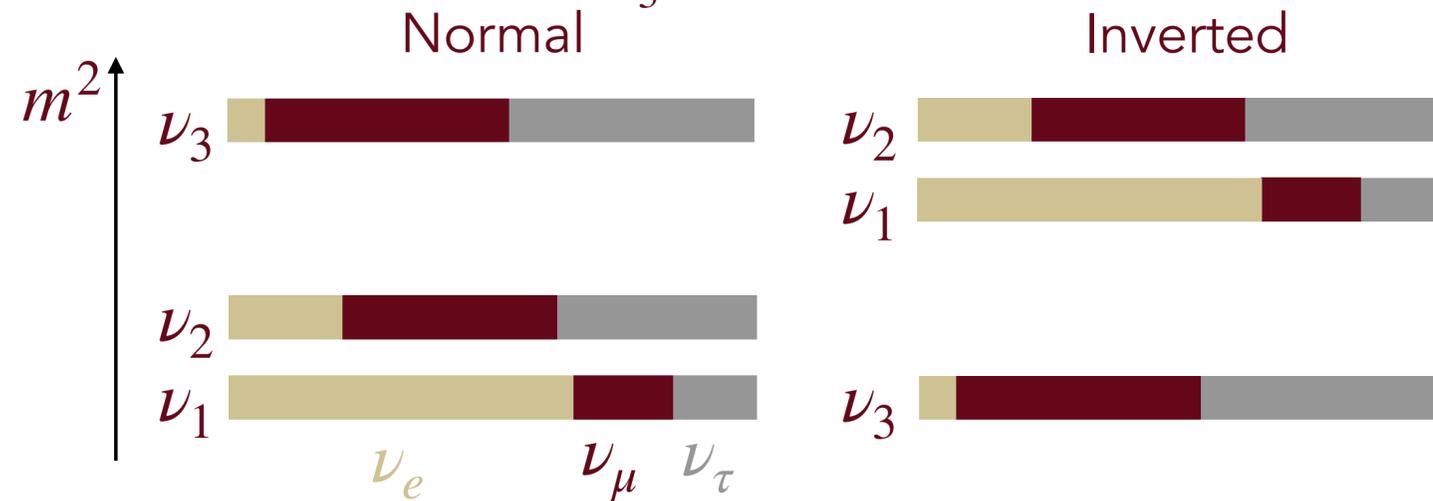


# Neutrino oscillations in NOvA

## 3-flavor Oscillations

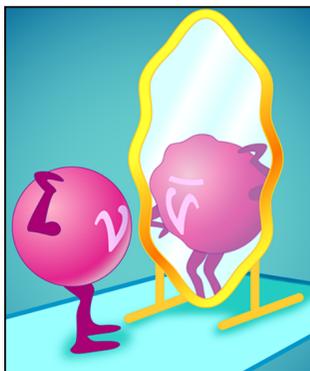
### Mass Ordering

Is  $\nu_3$  the heaviest state?



### CP Violation

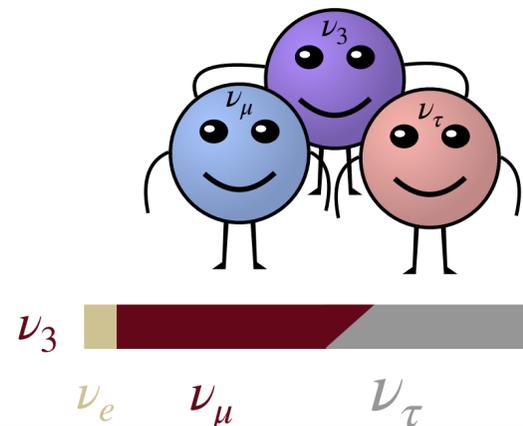
Is there a neutrino-antineutrino asymmetry?



Credit: APS/Carin Cain

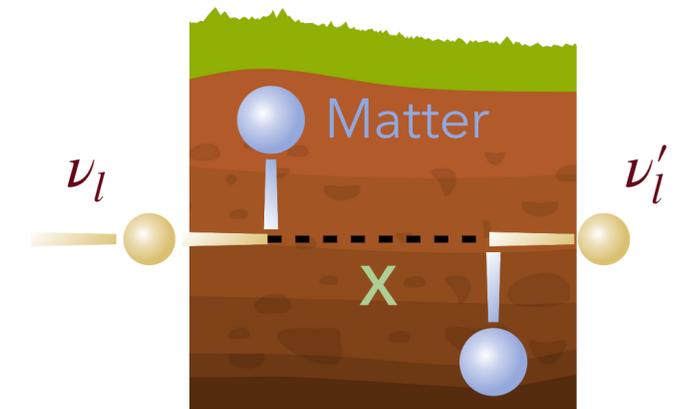
### Maximal Mixing ( $\theta_{23}$ )

Do the  $\nu_\mu$  and  $\nu_\tau$  states couple equally to  $\nu_3$ ?

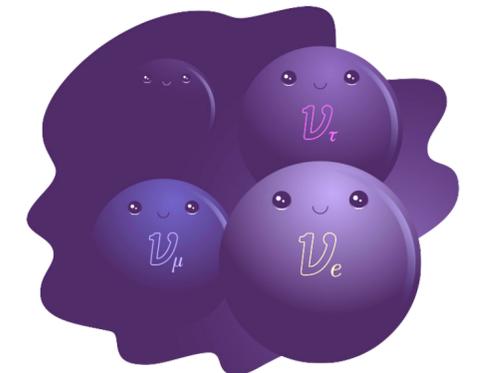


## Beyond Standard Oscillations

Nonstandard interactions



Sterile neutrinos

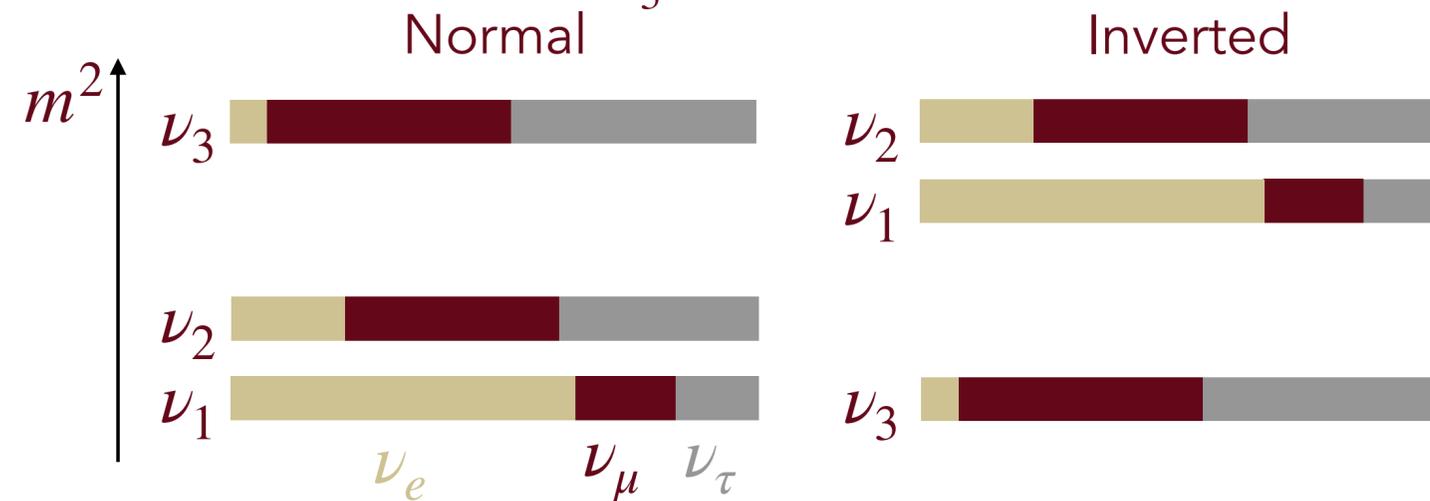


Credit: Symmetry magazine, Illustration by Sandbox Studio, Chicago with Ana Kova

# 3-Flavor Oscillations

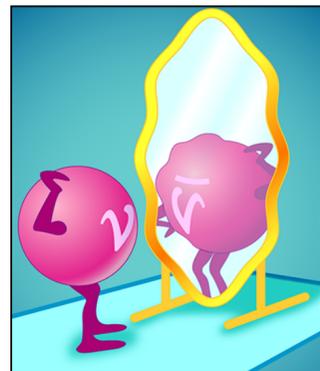
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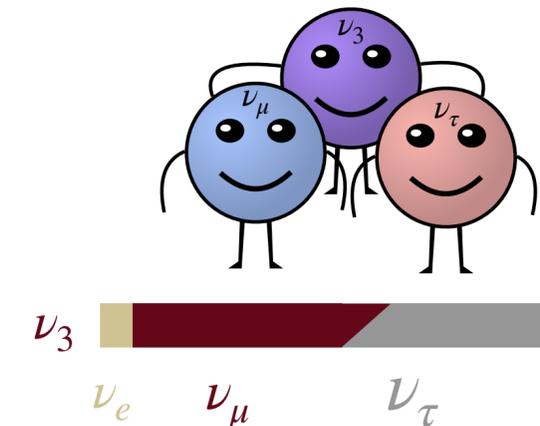
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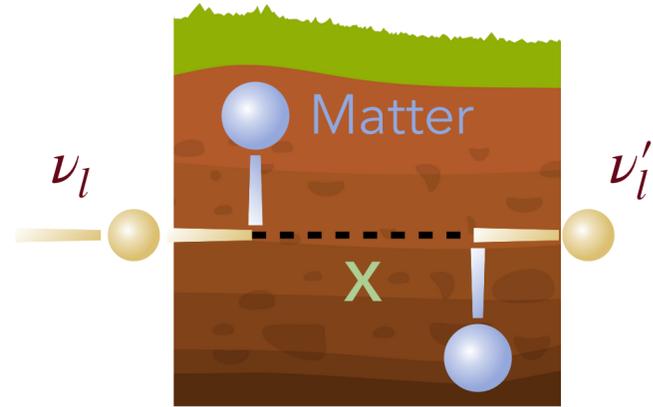
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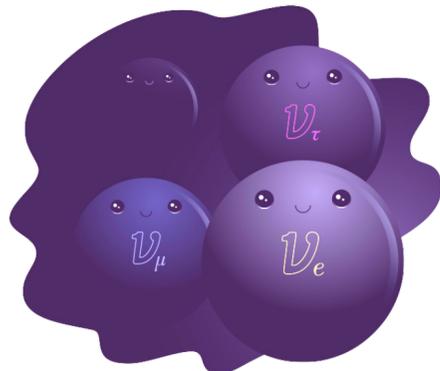
# Beyond Standard Oscillations

Nonstandard interactions



$$\mathcal{H} = U\mathcal{H}_0U^\dagger + \sqrt{2}G_F n_e \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} 1 + \varepsilon_{ee} & \varepsilon_{e\mu} & \varepsilon_{e\tau} \\ \varepsilon_{e\mu}^* & \varepsilon_{\mu\mu} & \varepsilon_{\mu\tau} \\ \varepsilon_{e\tau}^* & \varepsilon_{\mu\tau}^* & \varepsilon_{\tau\tau} \end{pmatrix}$$

Sterile neutrinos



Credit: [Symmetry magazine](#), Illustration by [Sandbox Studio](#), Chicago with Ana Kova

$$\begin{bmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \\ \nu_s \end{bmatrix} = \begin{bmatrix} U_{e1} & U_{e2} & U_{e3} & U_{e4} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} & U_{\mu4} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} & U_{\tau4} \\ U_{s1} & U_{s2} & U_{s3} & U_{s4} \end{bmatrix} \begin{bmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \\ \nu_4 \end{bmatrix}$$

3+1 model

