



Photo Credit: Bryce Richter



# Exclusive Reactions at COMPASS

**CIPANP 2025**

Madison, Wisconsin

June 10, 2025

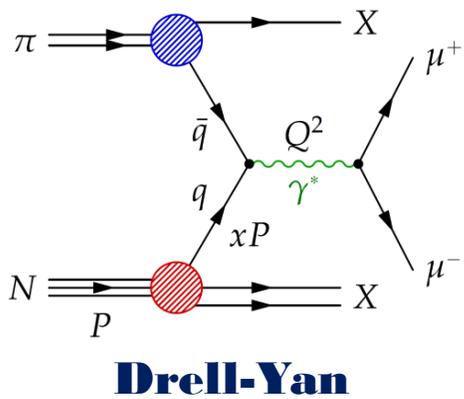
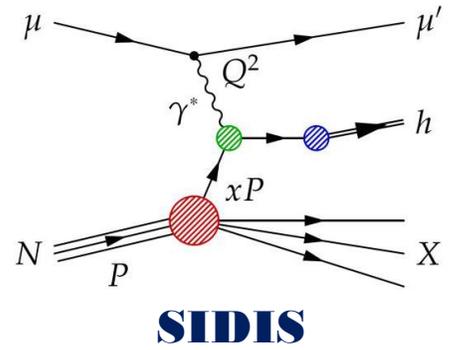
Po-Ju Lin

Department of Physics, National Central University

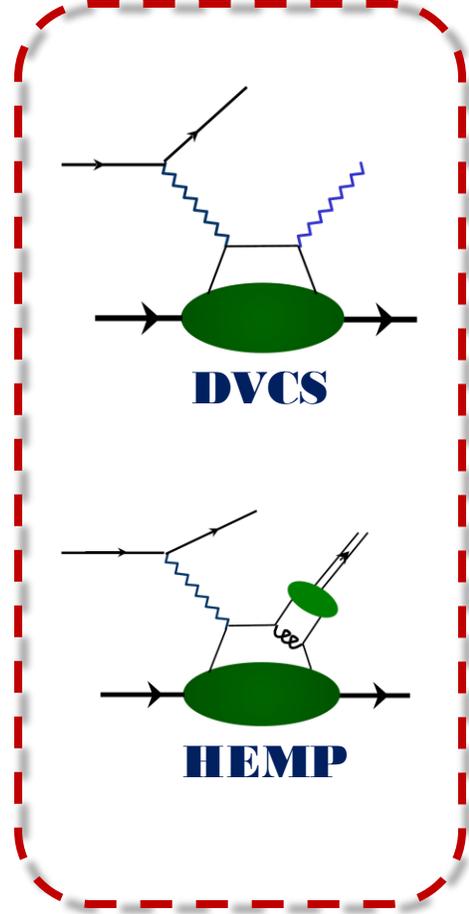
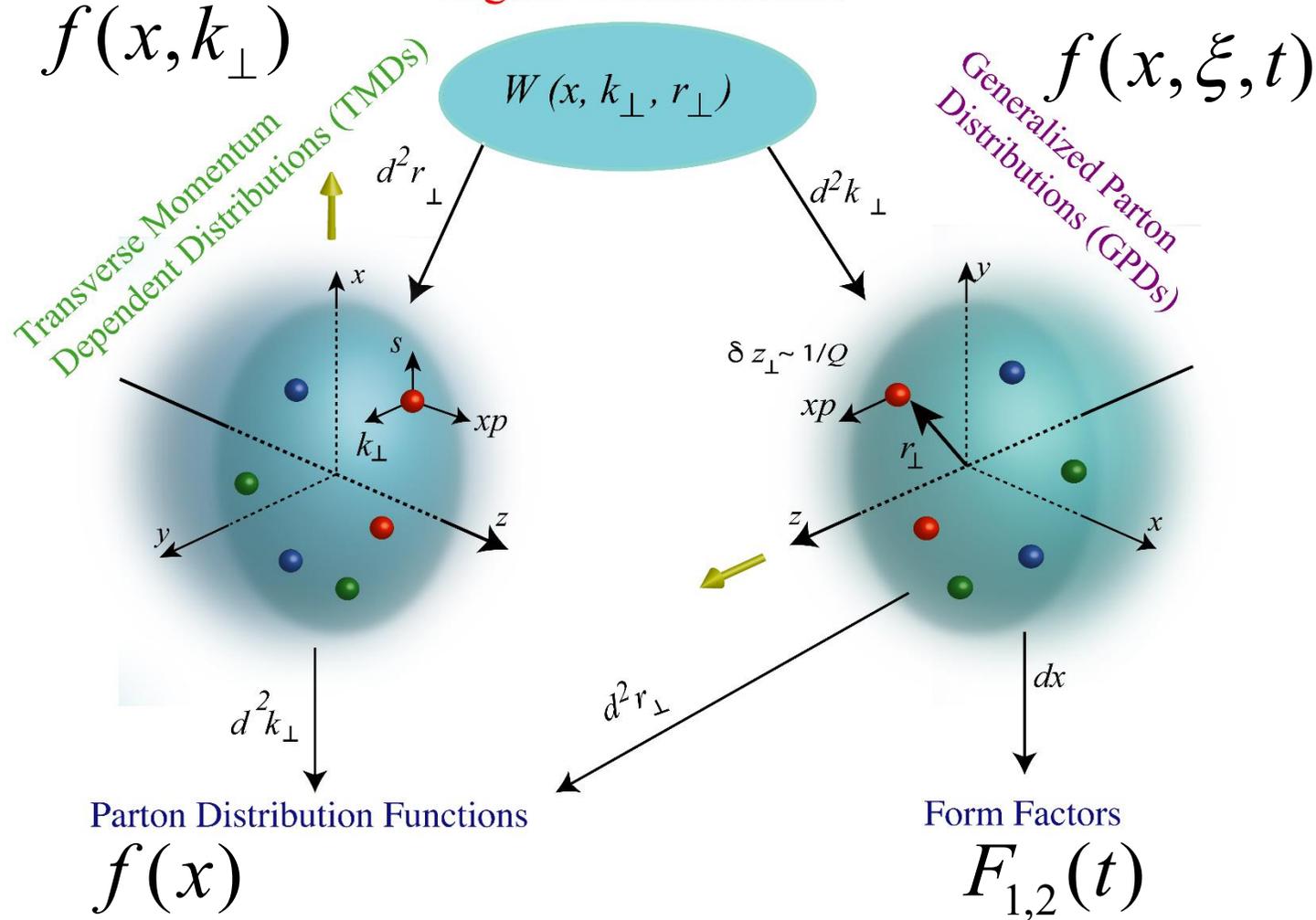
On behalf of the COMPASS Collaboration

# Multi-dimensional Partonic Structures

<http://www.int.washington.edu/PROGRAMS/17-3/>



## Wigner Distributions

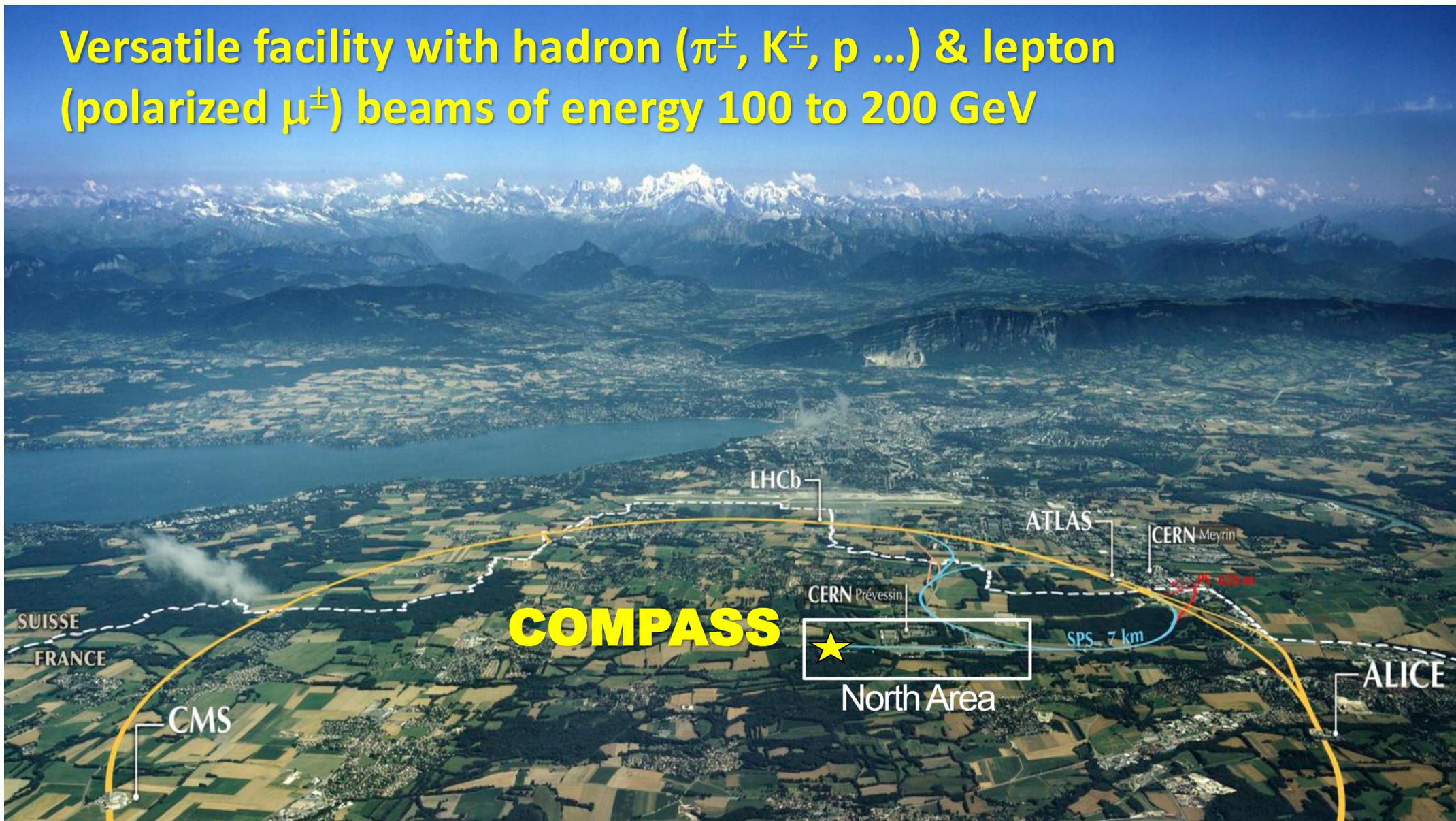


**COMPASS investigates the multi-dimensional structure of nucleon via various processes**

# COMPASS Experiment

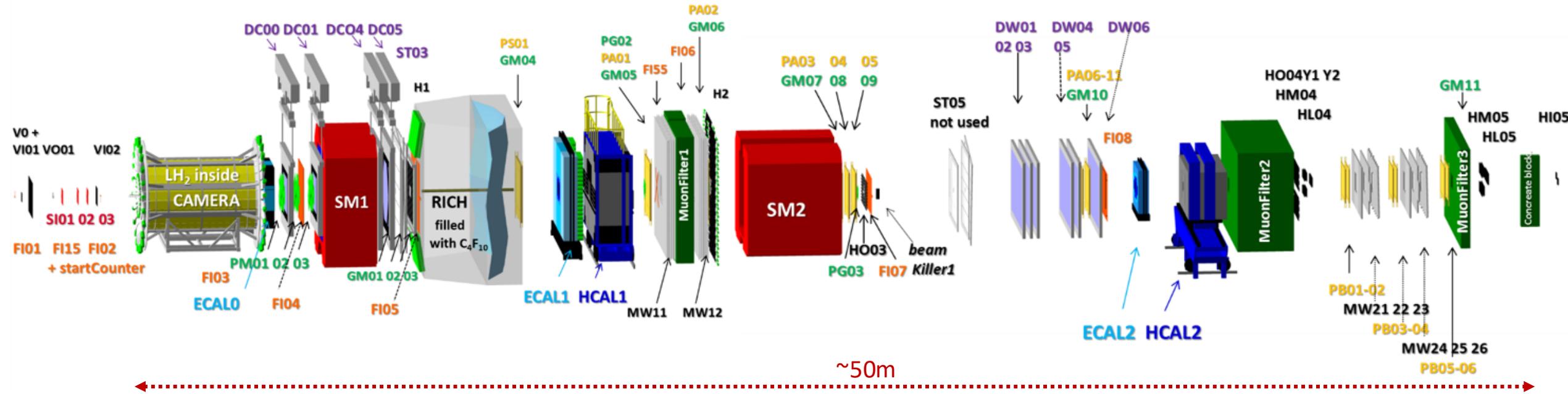


Versatile facility with hadron ( $\pi^\pm$ ,  $K^\pm$ ,  $p$  ...) & lepton (polarized  $\mu^\pm$ ) beams of energy 100 to 200 GeV



**CO**mmun  
**MU**on and  
**PR**oton  
**AP**paratus for  
**ST**ructure and  
**SP**ectroscopy

# COMPASS Experimental Setup



- Primary beam – 400 GeV p from SPS
  - Impinging on Be production target
- 190 GeV secondary hadron beams
  - $h^-$  beam: 97%  $\pi^-$ , 2%  $K^-$ , 1%  $p$
  - $h^+$  beam: 75%  $\pi^+$ , 24%  $p$ , 1%  $K^+$
- 160 GeV tertiary muon beams
  - $\mu^\pm$  longitudinally polarized

Large-acceptance forward spectrometer

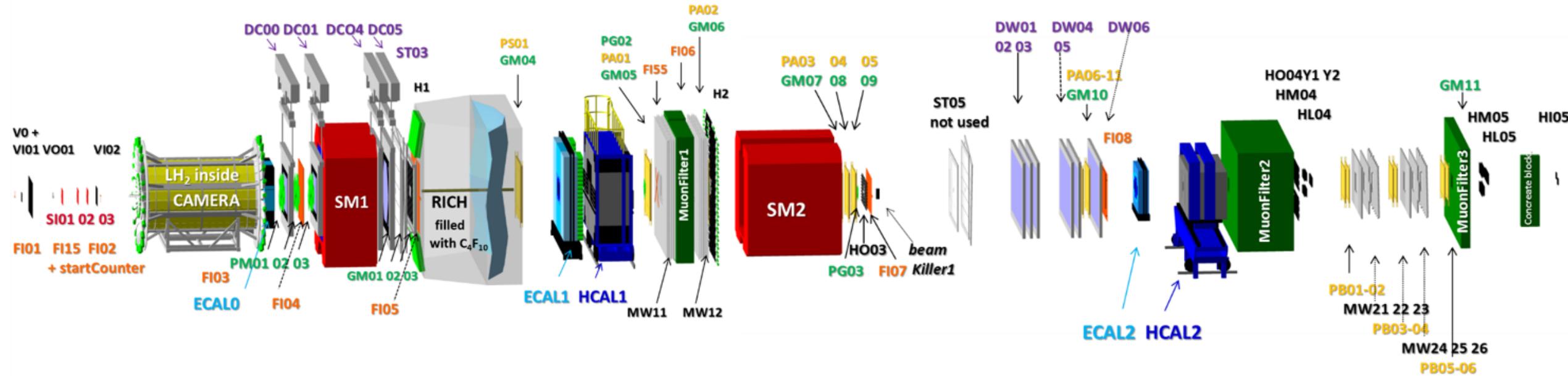
- Precise tracking (350 planes)  
SciFi, Silicon, MicroMegas, GEM, MWPC, DC, straw
- PID – CEDARs, RICH, calorimeters, Muon Walls

Various targets:

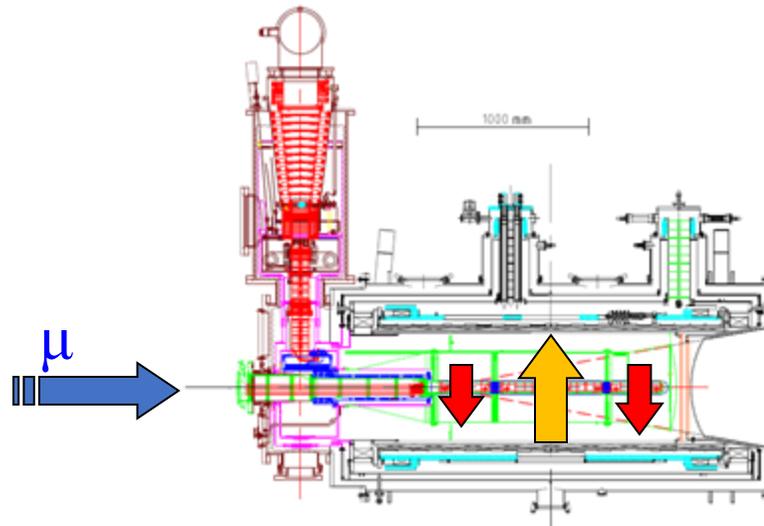
- Polarized solid-state  $\text{NH}_3$  or  $^6\text{LiD}$
- Liquid  $\text{H}_2$
- Solid-state nuclear targets

❖ NIM A 577 (2007) & NIM A 779 (2015) 69

# COMPASS Experimental Setup

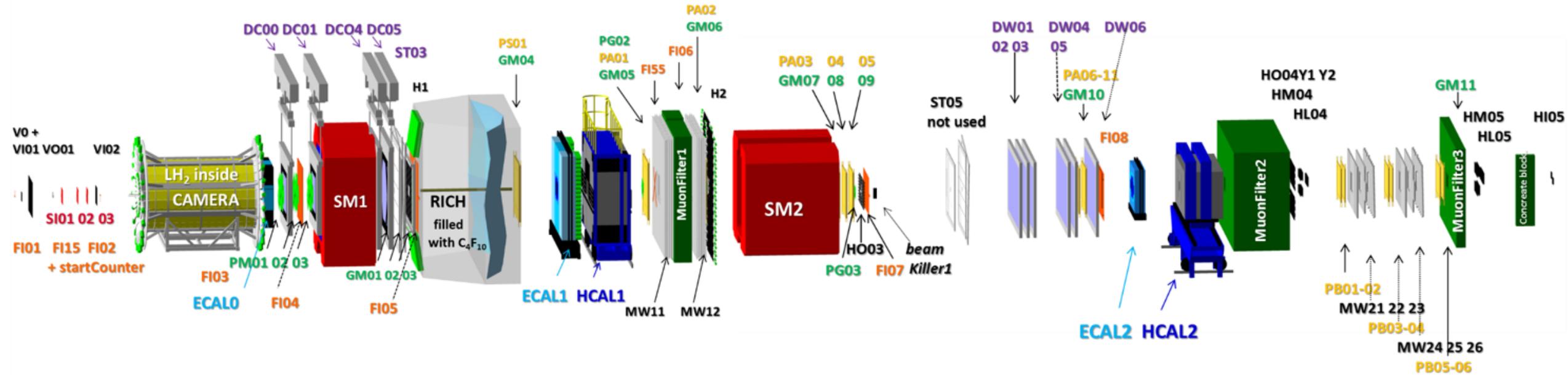


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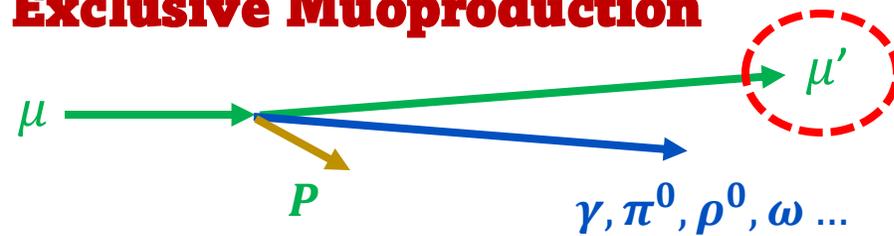
- In early GPD studies, transversely polarized target was used.
- Polarization reversal by magnetic field rotation
- 2.5m unpolarized  $LH_2$  target used in GPD dedicated runs

# COMPASS Setup for Exclusive Processes

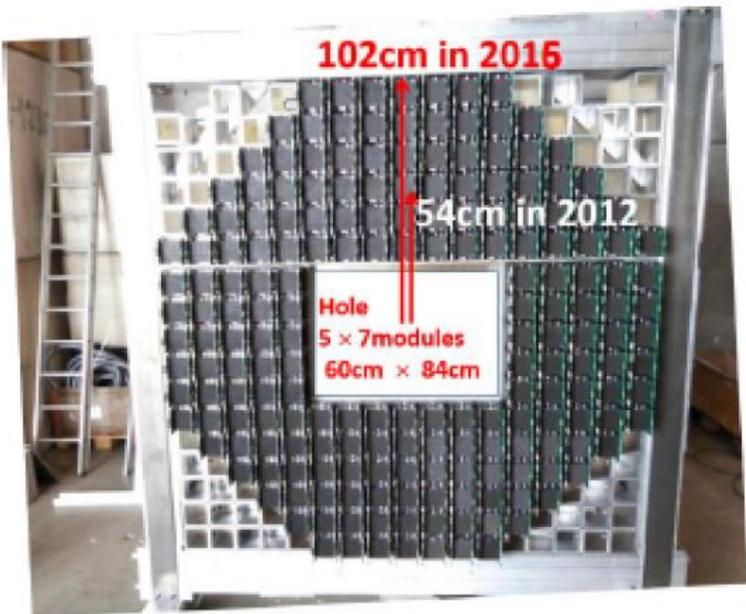
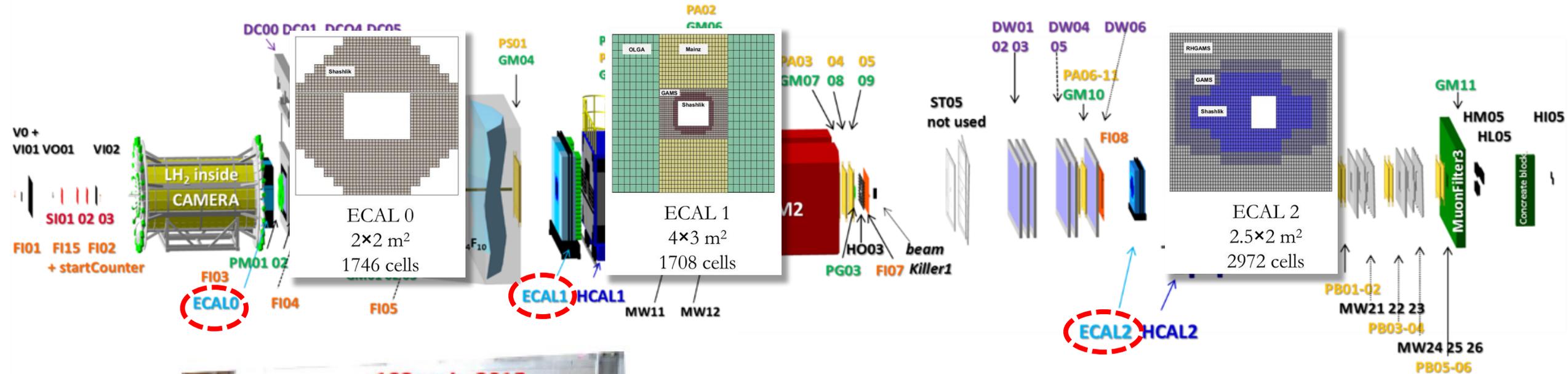


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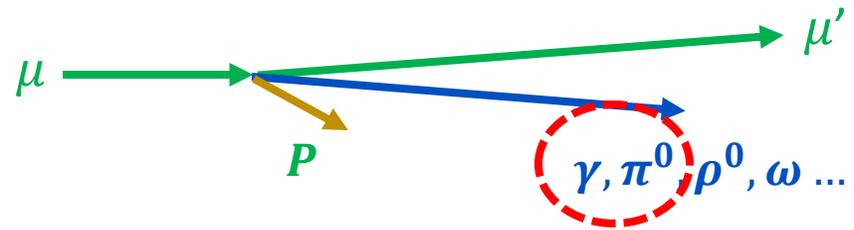
## Exclusive Muoproduction



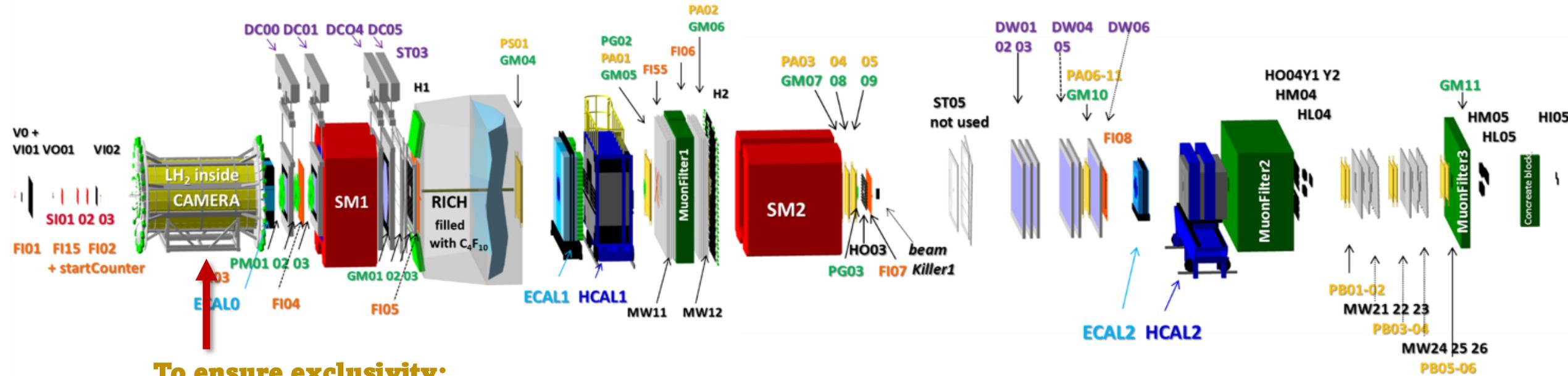
# COMPASS Setup for Exclusive Processes



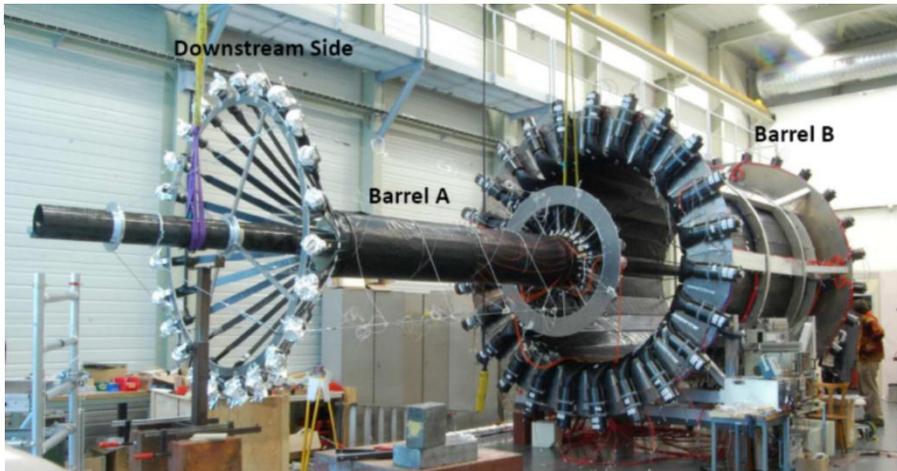
## Exclusive Muoproduction



# COMPASS Setup for Exclusive Processes

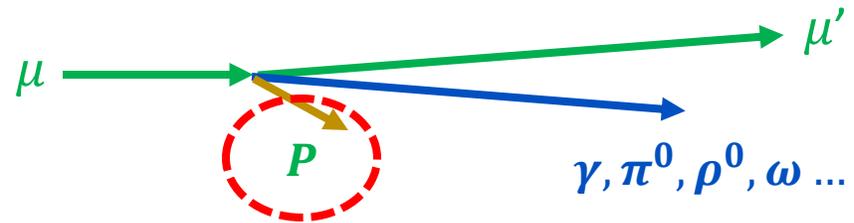


**To ensure exclusivity:**



CAMERA recoil proton detector

## Exclusive Muoproduction



# COMPASS Experiment



2002-2022 COMPASS data taking

2002-2004	DIS & SIDIS, $\mu^+$ - <b>d</b> , 160 GeV, <b>L &amp; T polarized target</b>
2005	<i>CERN accelerator shutdown, increase of COMPASS acceptance</i>
2006	DIS & SIDIS, $\mu^+$ - <b>d</b> , 160 GeV, <b>L polarized target</b>
2007	DIS & SIDIS, $\mu^+$ - <b>p</b> , 160 GeV, <b>L &amp; T polarized target</b>
2008-2009	Hadron spectroscopy & Primakoff reaction, $\pi/K/p$ beam
2010	SIDIS, $\mu^+$ - <b>p</b> , 160 GeV, <b>T polarized target</b>
2011	DIS & SIDIS, $\mu^+$ - <b>p</b> , 200 GeV, <b>L polarized target</b>
2012	Primakoff reaction, $\pi/K/p$ beam
2012 pilot run	DVCS/HEMP/SIDIS, $\mu^+$ & $\mu^-$ - <b>p</b> , 160 GeV, <b>unpolarized target</b>
2013	<i>CERN accelerator shutdown, LS1</i>
2014-2015	Drell-Yan, $\pi^-$ - <b>p</b> , <b>T polarized target</b>
2016-2017	DVCS/HEMP/SIDIS, $\mu^+$ & $\mu^-$ - <b>p</b> , 160 GeV, <b>unpolarized target</b>
2018	Drell-Yan, $\pi^-$ - <b>p</b> , <b>T polarized target</b>
2019-2020	<i>CERN accelerator shutdown, LS2</i>
2021-2022	SIDIS, $\mu^+$ - <b>d</b> , 160 GeV, <b>T polarized target</b>

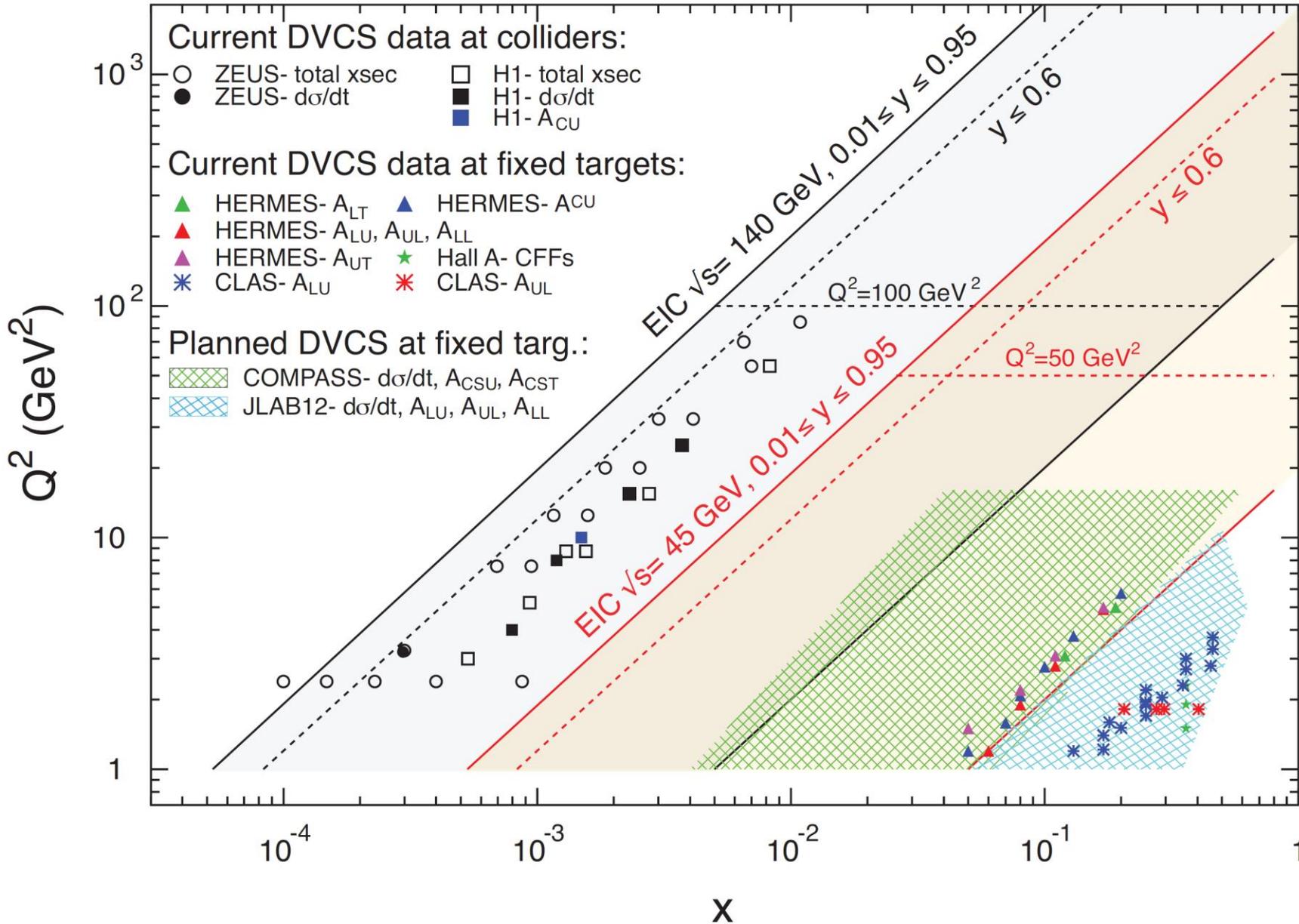
**Study hadron structure with complementary tools:**

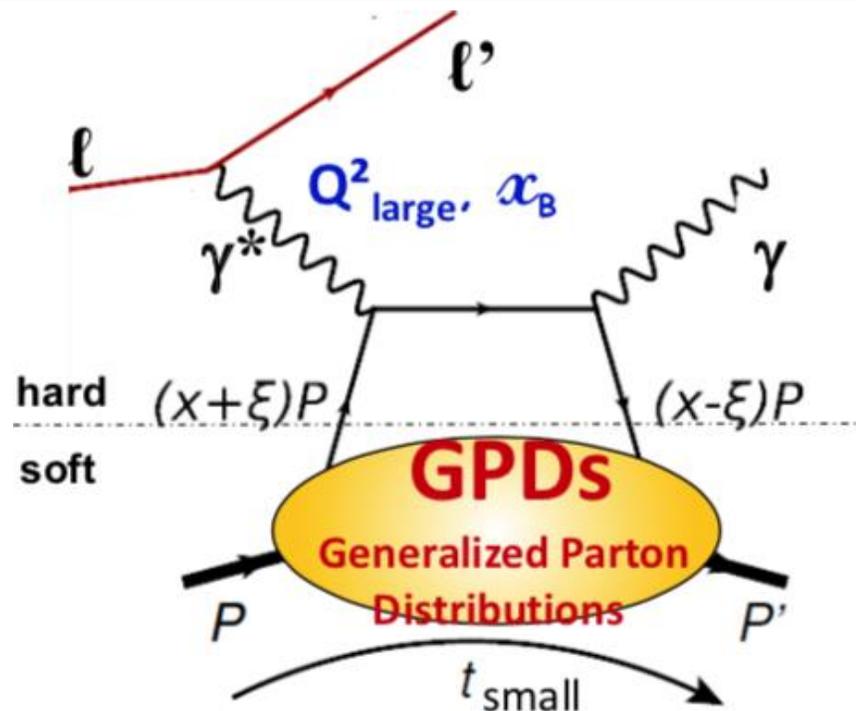
➤ **COMPASS holds the record for the longest-running CERN experiment**



- 2012 pilot run with 4-week data taking
- 2016-17 dedicated run. 2 x 6 months.

# Landscape – Global Programs of DVCS



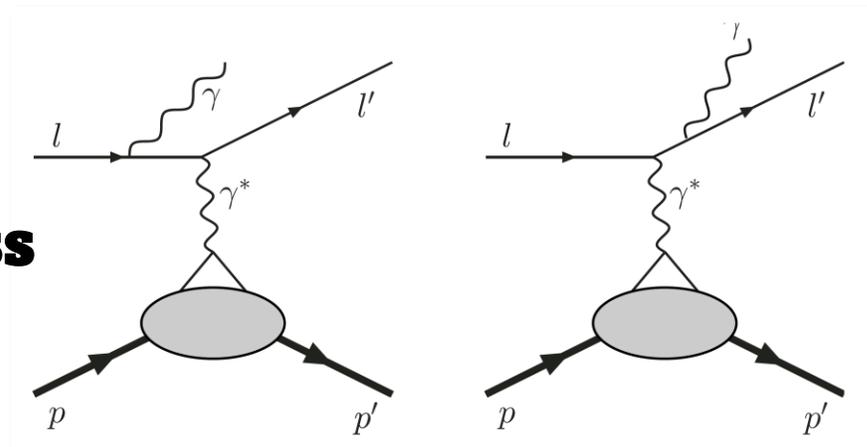


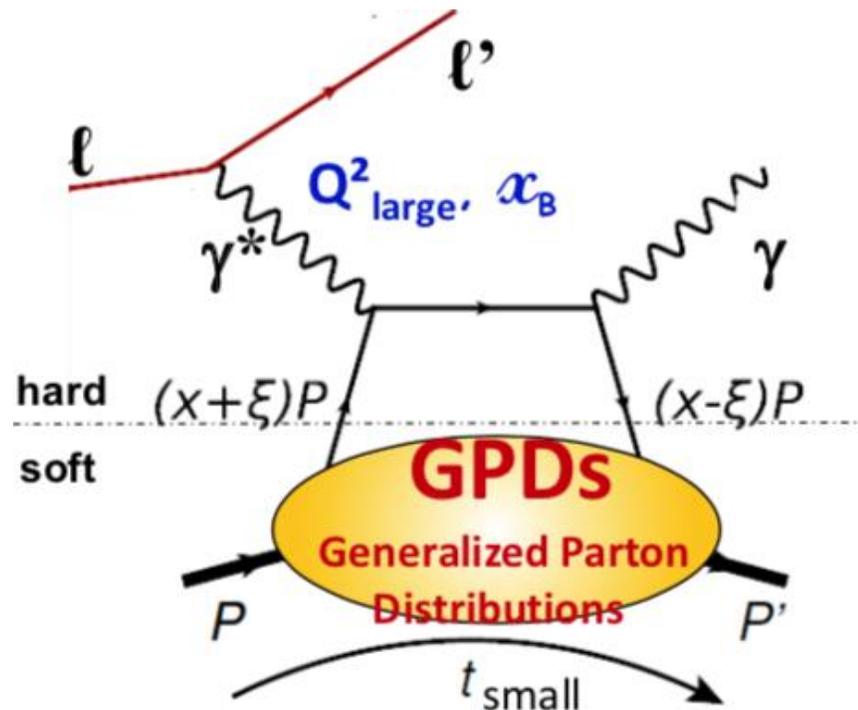
$$\text{DVCS: } l + p \rightarrow l' + p' + \gamma$$

- As the golden channel to access GPDs, DVCS has been the workhorse for GPD Extraction.
- Its interference with the well-understood Bethe-Heitler process gives access to more info.

- The GPDs depend on the following variables:
  - $x$ : average longitudinal momentum frac.
  - $\xi$ : longitudinal momentum diff.
  - $t$ : four momentum transfer  
(correlated to  $b_{\perp}$  via Fourier transform)
  - $Q^2$ : virtuality of  $\gamma^*$

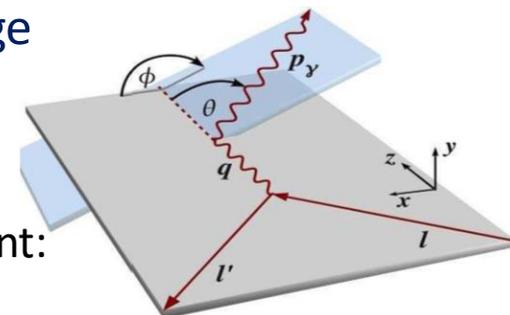
**BH  
Process**





## DVCS: $l + p \rightarrow l' + p' + \gamma$

- With LH<sub>2</sub> target and small  $x_B$  coverage  
→ focuses on H at COMPASS



- The variables measured in the experiment:  
 $E_\ell$ ,  $Q^2$ ,  $x_{Bj} \sim 2\xi / (1+\xi)$ ,  
 $t$  (or  $\theta_{\gamma^*\gamma}$ ) and  $\phi$  ( $l\ell'$  plane/ $\gamma\gamma^*$  plane)

CFP →  $\mathcal{H}(\xi, t) = \int_{-1}^{+1} dx \frac{\mathbf{H}(x, \xi, t)}{x - \xi + i\epsilon} + \dots = \mathcal{P} \int_{-1}^{+1} dx \frac{\mathbf{H}(x, \xi, t)}{x - \xi} - i\pi \mathbf{H}(x = \pm\xi, \xi, t) + \dots$

GPD →

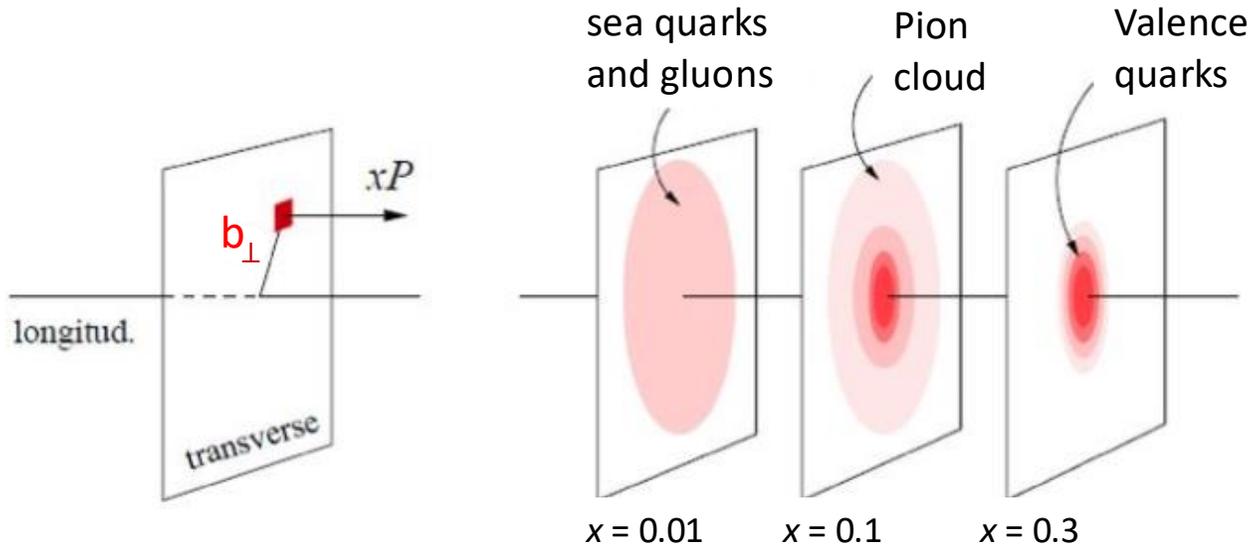
REAL part

Imaginary part

$$\text{Re } \mathcal{H}(\xi, t) = \mathcal{P} \int dx \frac{\text{Im } \mathcal{H}(x, t)}{x - \xi} + \Delta(t)$$

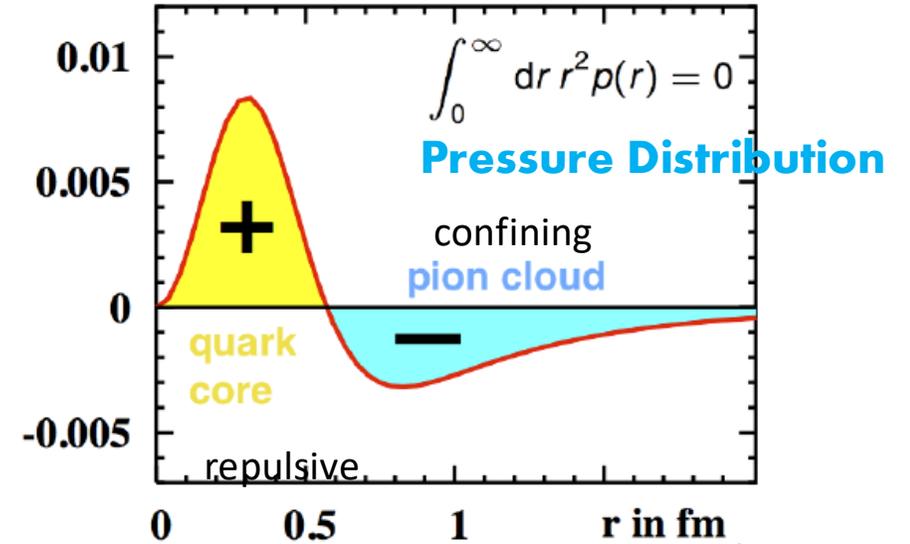
# Transverse Imaging and Pressure Distribution

## Mapping in the transverse plane



M. Polyakov, P. Schweitzer, Int.J.Mod.Phys. A33 (2018)

$r^2 p(r)$  in  $\text{GeV fm}^{-1}$



$\overset{\text{GPD}}{\mathcal{H}(\xi, t)} = \int_{-1}^{+1} dx \frac{\mathbf{H}(x, \xi, t)}{x - \xi + i\epsilon} + \dots = \mathcal{P} \int_{-1}^{+1} dx \frac{\mathbf{H}(x, \xi, t)}{x - \xi} - i\pi \mathbf{H}(x = \pm \xi, \xi, t) + \dots$

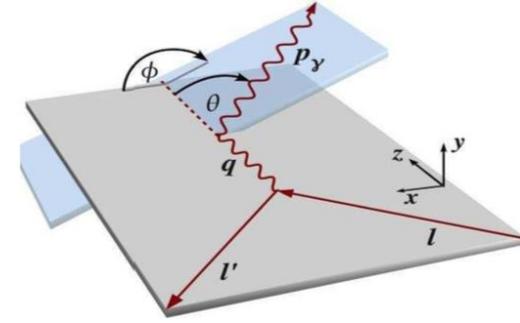
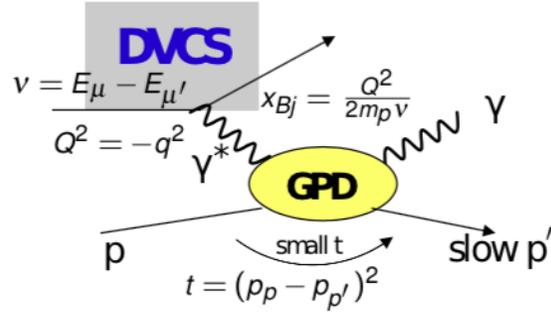
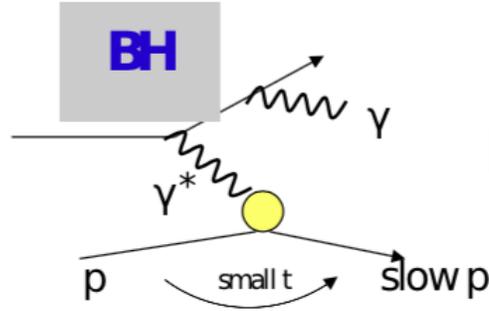
$\text{REAL part} \quad \text{Imaginary part}$

$\text{Re } \mathcal{H}(\xi, t) = \mathcal{P} \int dx \frac{\text{Im } \mathcal{H}(x, t)}{x - \xi} + \Delta(t)$

$\Delta(t) \equiv d_1(t)$  (D-term)

$\leftarrow \text{FT of } H(x, \xi=0, t)$

# Azimuthal Dependence of BH & DVCS



$$\frac{d^4\sigma(\ell p \rightarrow \ell p \gamma)}{dx_B dQ^2 d|t| d\phi} = \underbrace{d\sigma^{BH}}_{\text{Well known}} + \left( d\sigma_{unpol}^{DVCS} + P_\ell d\sigma_{pol}^{DVCS} \right) + (e_\ell \text{Re } I + e_\ell P_\ell \text{Im } I)$$

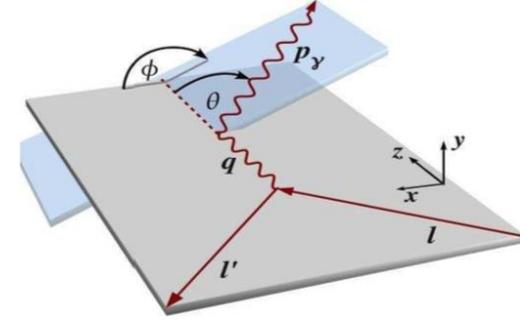
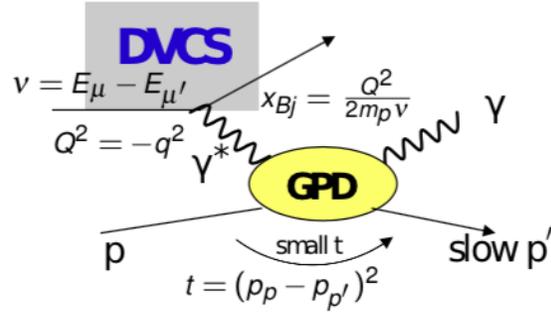
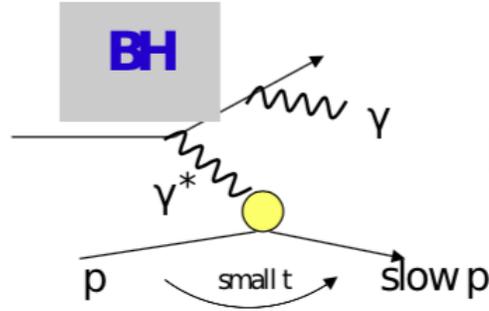
Beam Charge-spin difference & sum

$$\mathcal{D}_{CS,U}(\phi) \equiv d\sigma(\mu^{+\leftarrow}) - d\sigma(\mu^{--\rightarrow})$$

$$\mathcal{S}_{CS,U}(\phi) \equiv d\sigma(\mu^{+\leftarrow}) + d\sigma(\mu^{--\rightarrow})$$

$$\begin{aligned} d\sigma^{BH} &\propto c_0^{BH} + c_1^{BH} \cos \phi + c_2^{BH} \cos 2\phi \\ d\sigma_{unpol}^{DVCS} &\propto c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi \\ d\sigma_{pol}^{DVCS} &\propto s_1^{DVCS} \sin \phi \\ \text{Re } I &\propto c_0^I + c_1^I \cos \phi + c_2^I \cos 2\phi + c_3^I \cos 3\phi \\ \text{Im } I &\propto s_1^I \sin \phi + s_2^I \sin 2\phi \end{aligned}$$

# Azimuthal Dependence of BH & DVCS



$$\frac{d^4\sigma(\ell p \rightarrow \ell p \gamma)}{dx_B dQ^2 d|t| d\phi} = d\sigma^{BH}_{\text{Well known}} + \left( d\sigma^{DVCS}_{\text{unpol}} + P_\ell d\sigma^{DVCS}_{\text{pol}} \right) + (e_\ell \text{Re } I + e_\ell P_\ell \text{Im } I)$$

Beam Charge-spin difference & sum

$$\mathcal{D}_{CS,U}(\phi) \equiv d\sigma(\mu^{+\leftarrow}) - d\sigma(\mu^{-\rightarrow})$$

$$\mathcal{S}_{CS,U}(\phi) \equiv d\sigma(\mu^{+\leftarrow}) + d\sigma(\mu^{-\rightarrow})$$

$$d\sigma^{BH} \propto c_0^{BH} + c_1^{BH} \cos \phi + c_2^{BH} \cos 2\phi$$

$$d\sigma^{DVCS}_{\text{unpol}} \propto c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi$$

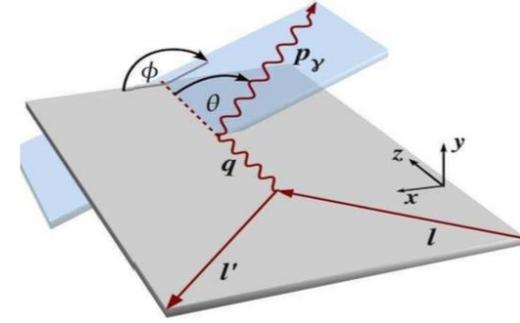
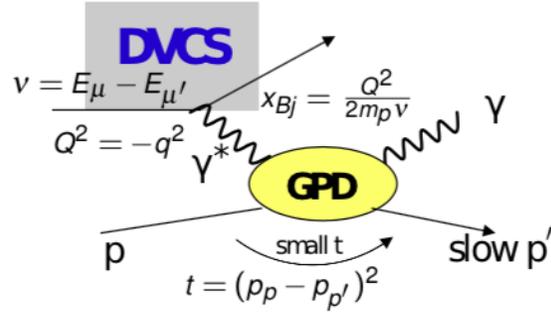
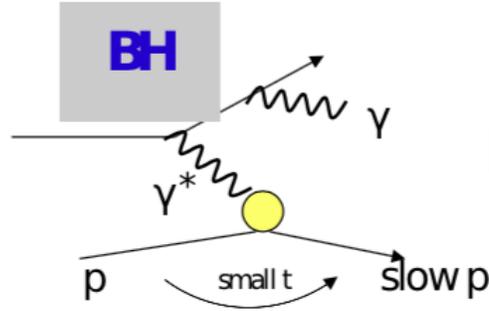
$$d\sigma^{DVCS}_{\text{pol}} \propto s_1^{DVCS} \sin \phi$$

$$\text{Re } I \propto c_0^I + c_1^I \cos \phi + c_2^I \cos 2\phi + c_3^I \cos 3\phi$$

$$\text{Im } I \propto s_1^I \sin \phi + s_2^I \sin 2\phi$$

$\mathcal{D}_{CS,U}(\phi)$

# Azimuthal Dependence of BH & DVCS



$$\frac{d^4\sigma(\ell p \rightarrow \ell p \gamma)}{dx_B dQ^2 d|t| d\phi} = d\sigma^{BH} + \left( d\sigma_{unpol}^{DVCS} + P_\ell d\sigma_{pol}^{DVCS} \right) + (e_\ell \text{Re } I + e_\ell P_\ell \text{Im } I)$$

Well known

Beam Charge-spin difference & sum

$$\mathcal{D}_{CS,U}(\phi) \equiv d\sigma(\mu^{+\leftarrow}) - d\sigma(\mu^{-\rightarrow})$$

$$\mathcal{S}_{CS,U}(\phi) \equiv d\sigma(\mu^{+\leftarrow}) + d\sigma(\mu^{-\rightarrow})$$

$$\begin{aligned} d\sigma^{BH} &\propto c_0^{BH} + c_1^{BH} \cos \phi + c_2^{BH} \cos 2\phi \\ d\sigma_{unpol}^{DVCS} &\propto c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi \end{aligned}$$

$$d\sigma_{pol}^{DVCS} \propto s_1^{DVCS} \sin \phi$$

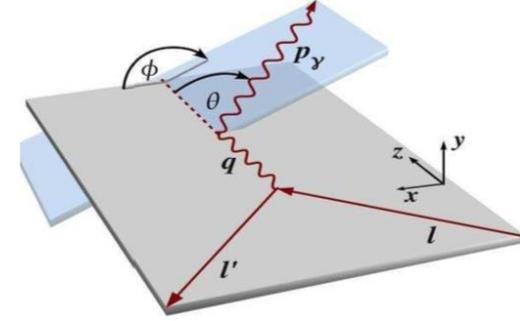
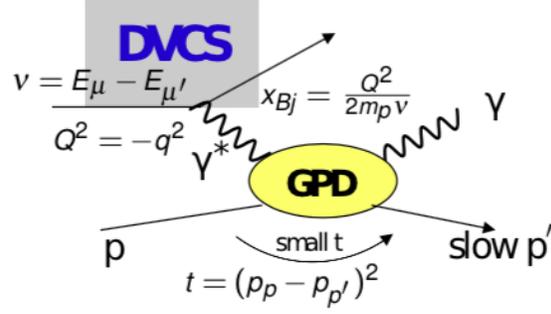
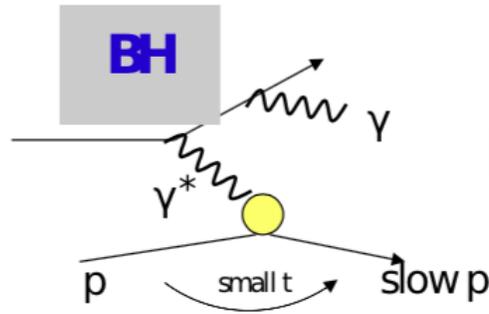
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$\mathcal{S}_{CS,U}(\phi)$

$\mathcal{S}_{CS,U}(\phi)$

# Azimuthal Dependence of BH & DVCS



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Beam Charge-spin difference & sum

$$\mathcal{D}_{CS,U}(\phi) \equiv d\sigma(\mu^{+\leftarrow}) - d\sigma(\mu^{-\rightarrow})$$

$$\rightarrow \underline{C_1^I \propto \text{Re } \mathcal{F}}$$

More challenging

$$\mathcal{S}_{CS,U}(\phi) \equiv d\sigma(\mu^{+\leftarrow}) + d\sigma(\mu^{-\rightarrow})$$

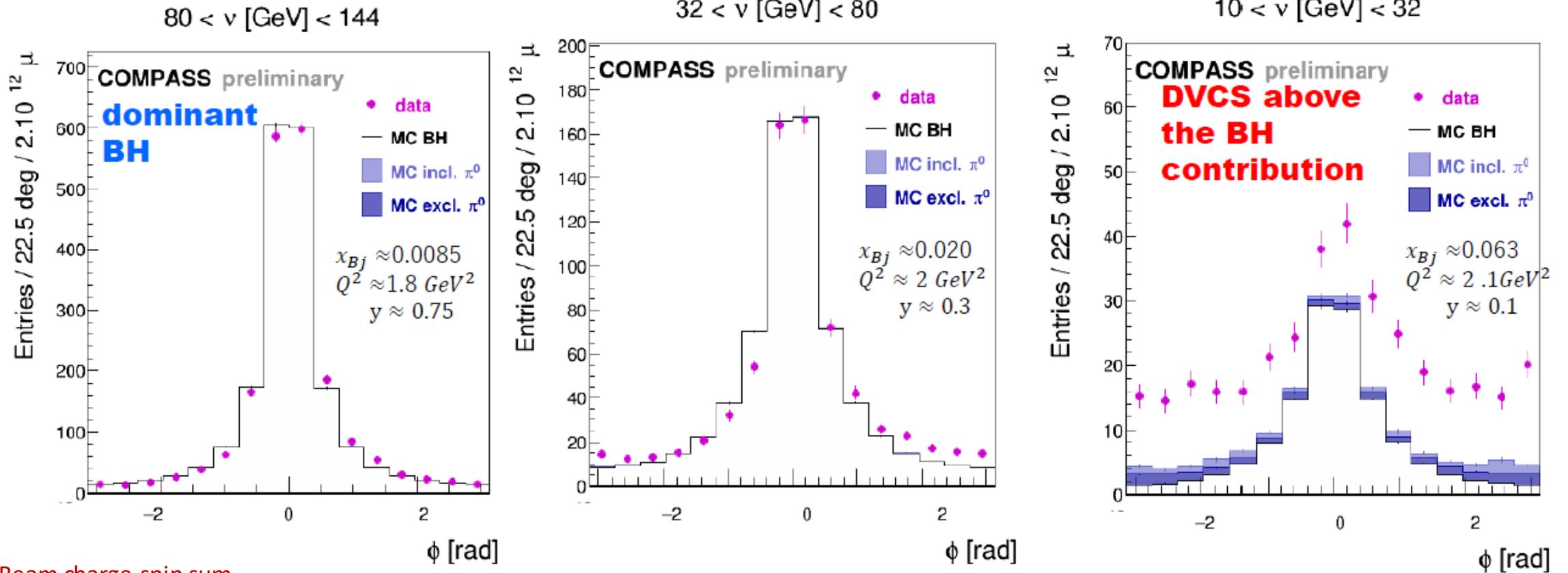
$$\rightarrow \underline{C_0^{DVCS} \propto (\text{Im } \mathcal{H})^2} \text{ and } \underline{S_1^I \propto \text{Im } \mathcal{F}}$$

Easier to measure

$$\mathcal{F} = F_1 \mathcal{H} + \xi(F_1 + F_2) \mathcal{H} + t/4m^2 F_2 \mathcal{E} \xrightarrow[\text{Small } x_B \text{ at COMPASS}]{\text{Proton Target}} F_1 \mathcal{H}$$

Compton Form factor linked to GPD  $\mathcal{H}$

# COMPASS 2016 Preliminary Results



➤ Beam charge-spin sum

$$\begin{aligned}
 S_{CS,u}(\phi) &\equiv d\sigma(\mu^{+\leftarrow}) + d\sigma(\mu^{-\rightarrow}) = 2[ d\sigma^{BH} + d\sigma_{unpol}^{DVCS} + \text{Im } I ] \\
 &= 2[ d\sigma^{BH} + c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi + s_1^I \sin \phi + s_2^I \sin 2\phi ]
 \end{aligned}$$

$$c_0^{DVCS} \underset{\text{small } x_{Bj}}{\propto} 4(\mathcal{H}\mathcal{H}^* + \tilde{\mathcal{H}}\tilde{\mathcal{H}}^*) + \frac{t}{M^2} \mathcal{E}\mathcal{E}^* \rightarrow 4 (\text{Im } \mathcal{H})^2 \underset{\text{model dependent}}{}$$

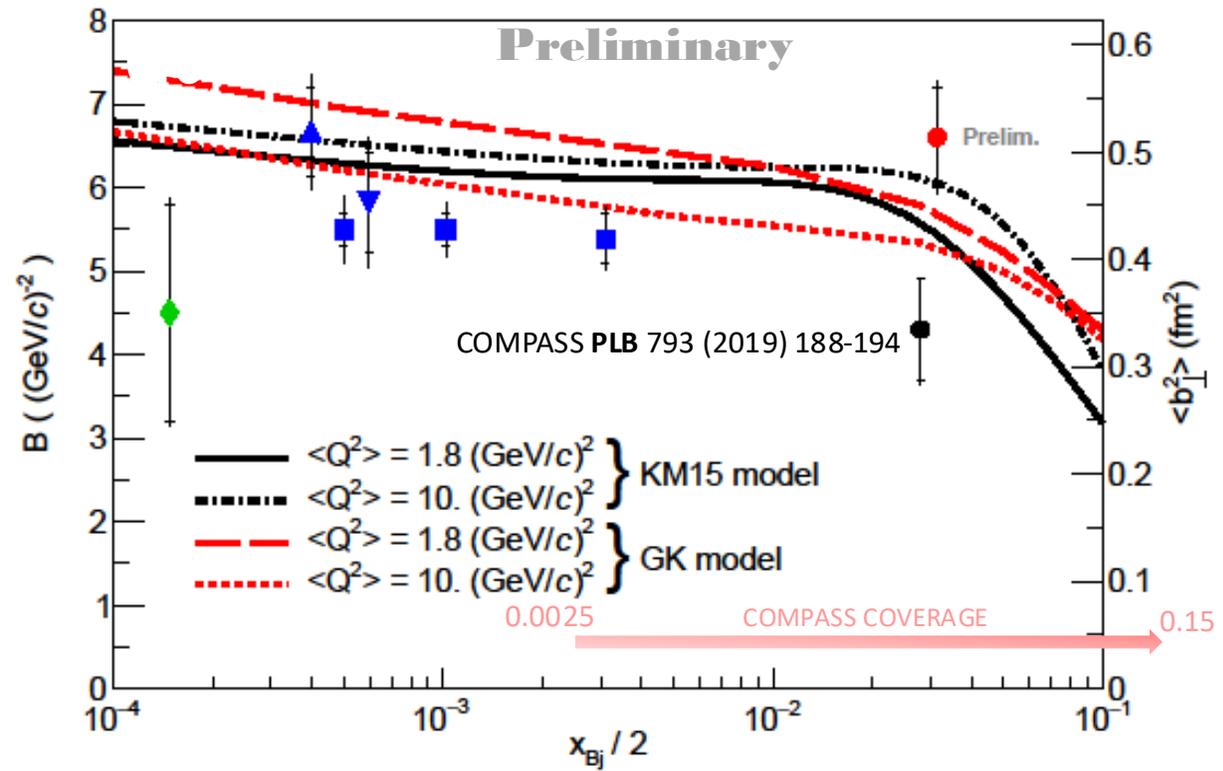
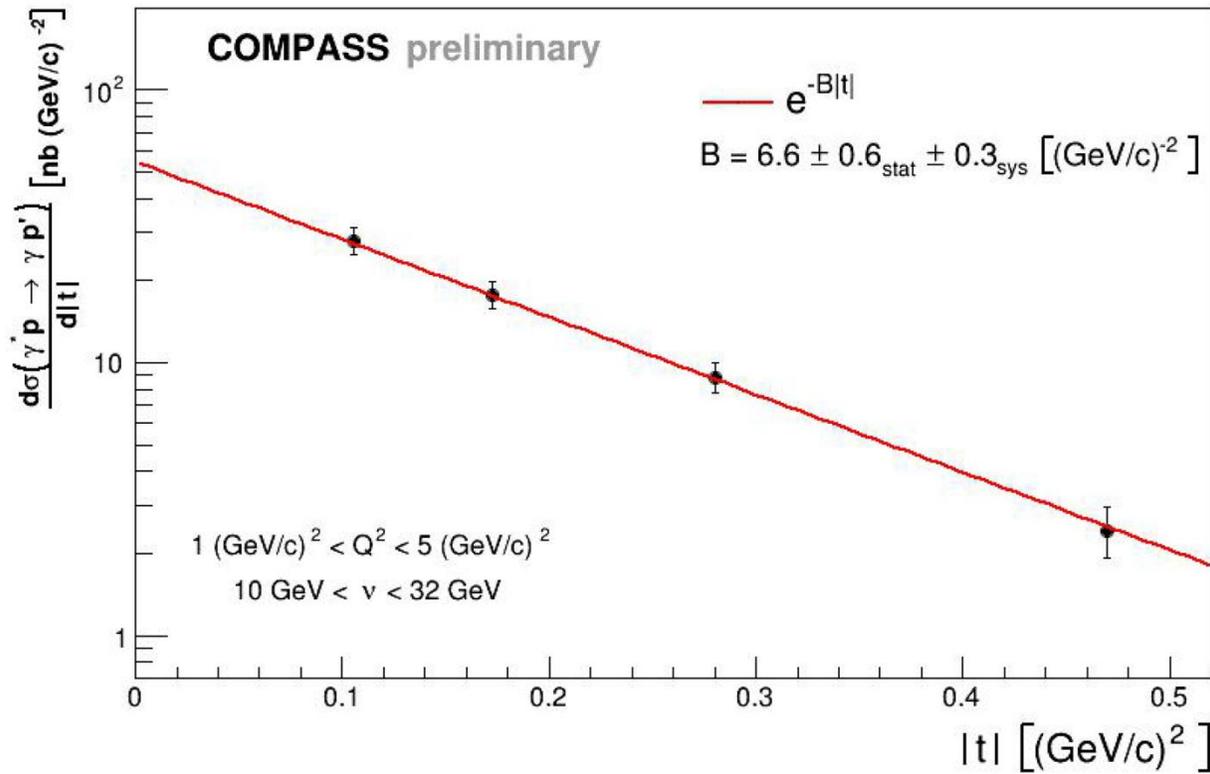
# Transverse extension of partons – 2016 data



$$d\sigma^{DVCS}/d|t| \propto e^{-B|t|}$$

$$\langle r_{\perp}^2(x_B) \rangle \approx 2B(x_B) \quad \text{At small } x_B$$

- COMPASS:  $\langle Q^2 \rangle = 1.8 \text{ (GeV/c)}^2$
- ◆ ZEUS:  $\langle Q^2 \rangle = 3.2 \text{ (GeV/c)}^2$
- ▲ H1:  $\langle Q^2 \rangle = 4.0 \text{ (GeV/c)}^2$
- ▼ H1:  $\langle Q^2 \rangle = 8.0 \text{ (GeV/c)}^2$
- H1:  $\langle Q^2 \rangle = 10. \text{ (GeV/c)}^2$

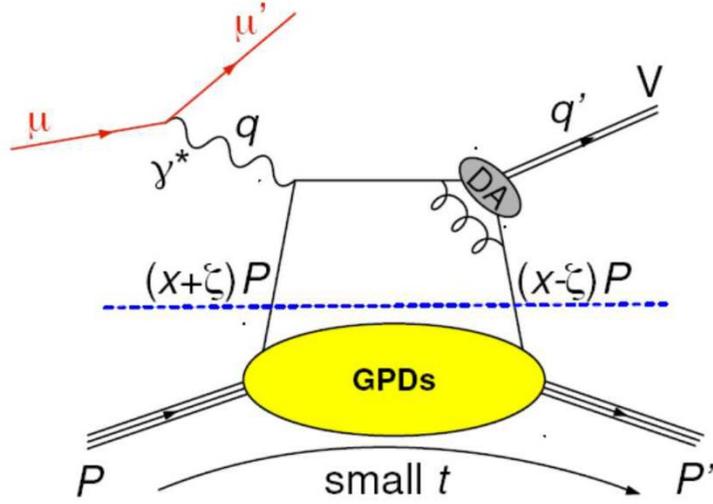


➤ The transverse-size evolution as a function of  $x_{Bj}$  → Expect at least 3  $x_{Bj}$  bins from 2016-17 data

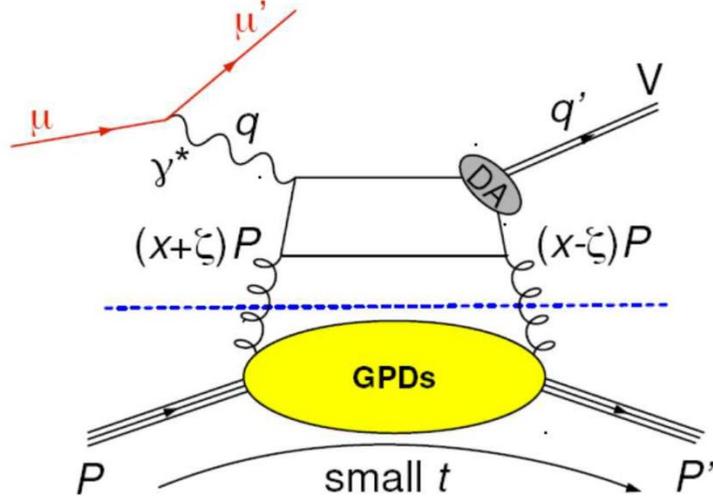
# GPDs in Hard Exclusive Meson Production



quark contribution



gluon contribution



4 chiral-even GPDs: helicity of parton unchanged

$$\mathbf{H}^q(x, \xi, t) \quad \mathbf{E}^q(x, \xi, t) \quad \rightarrow \text{Vector Meson}$$

$$\tilde{\mathbf{H}}^q(x, \xi, t) \quad \tilde{\mathbf{E}}^q(x, \xi, t) \quad \rightarrow \text{Pseudo-Scalar Meson}$$

+ 4 chiral-odd (transversity) GPDs: helicity of parton changed  
(not possible in DVCS)

$$\mathbf{H}_T^q(x, \xi, t)$$

$$\mathbf{E}_T^q(x, \xi, t)$$

$$\tilde{\mathbf{H}}_T^q(x, \xi, t)$$

$$\tilde{\mathbf{E}}_T^q(x, \xi, t)$$

$$\overline{\mathbf{E}}_T^q = 2 \tilde{\mathbf{H}}_T^q + \mathbf{E}_T^q$$

- Ability to probe the chiral-odd GPDs.
- Universality of GPDs, quark flavor filter
- In addition to nuclear structure, provide insights into reaction mechanism.
- Additional non-perturbative term from meson wave function.

# Exclusive $\pi^0$ Production on Unpolarized Proton



$\mu p \rightarrow \mu \pi^0 p$

$$\frac{d^2\sigma}{dt d\phi_\pi} = \frac{1}{2\pi} \left[ \left( \frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} \right) + \epsilon \cos 2\phi_\pi \frac{d\sigma_{TT}}{dt} + \sqrt{2\epsilon(1+\epsilon)} \cos \phi_\pi \frac{d\sigma_{LT}}{dt} \right]$$

$\epsilon$  : degree of longitudinal polarization

$$\frac{d\sigma_L}{dt} = \frac{4\pi\alpha}{k'} \frac{1}{Q^6} \left\{ (1-\xi^2) |\langle \tilde{H} \rangle|^2 - 2\xi^2 \text{Re} [\langle \tilde{H} \rangle^* \langle \tilde{E} \rangle] - \frac{t'}{4m^2} \xi^2 |\langle \tilde{E} \rangle|^2 \right\}$$

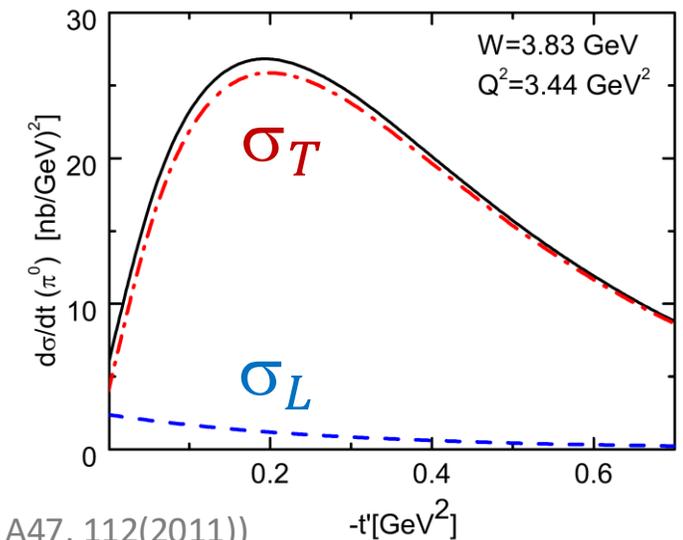
Leading twist expected be dominant  
But measured as  $\approx$  only a few % of  $\frac{d\sigma_T}{dt}$

The other contributions arise from coupling between chiral-odd (quark helicity flip) GPDs to the **twist-3** pion amplitude

$$\frac{d\sigma_T}{dt} = \frac{4\pi\alpha}{2k'} \frac{\mu_\pi^2}{Q^8} \left[ (1-\xi^2) |\langle H_T \rangle|^2 - \frac{t'}{8m^2} |\langle \bar{E}_T \rangle|^2 \right]$$

$$\frac{\sigma_{LT}}{dt} = \frac{4\pi\alpha}{\sqrt{2}k'} \frac{\mu_\pi}{Q^7} \xi \sqrt{1-\xi^2} \frac{\sqrt{-t'}}{2m} \text{Re} [\langle H_T \rangle^* \langle \tilde{E} \rangle]$$

$$\frac{\sigma_{TT}}{dt} = \frac{4\pi\alpha}{k'} \frac{\mu_\pi^2}{Q^8} \frac{t'}{16m^2} |\langle \bar{E}_T \rangle|^2$$

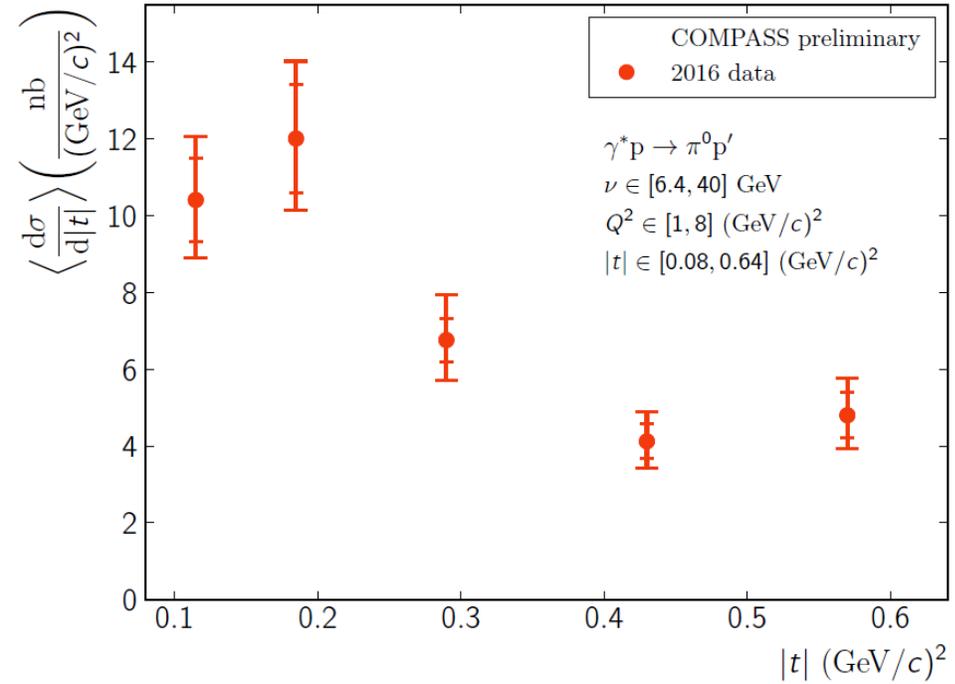
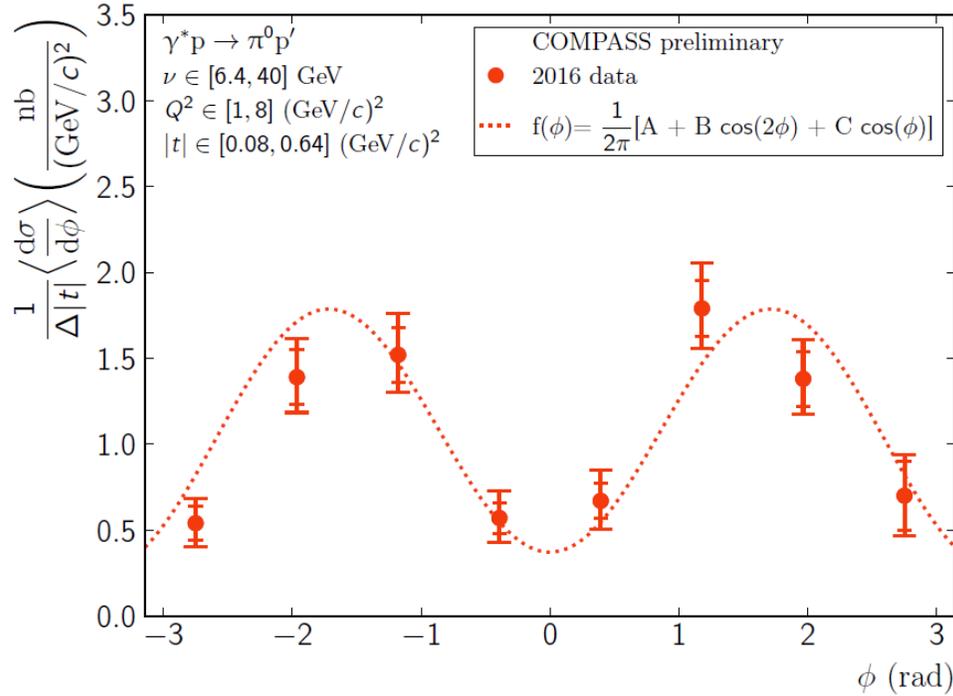




# New 2016 Exclusive $\pi^0$ Prod. on Unpolarized Proton



- Kinematic domain:  $\nu \in [6.4, 40]$  GeV and  $Q^2 \in [1, 8]$   $\text{GeV}^2/c^2$ ,  $\langle x_B \rangle = 0.134$



$$\left\langle \frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} \right\rangle = \left( 6.6 \pm 0.3_{stat} \begin{matrix} +0.9 \\ -0.8 \end{matrix} \Big|_{sys} \right) \frac{nb}{(\text{GeV}/c)^2}$$

$$\left\langle \frac{d\sigma_{TT}}{dt} \right\rangle = \left( -4.6 \pm 0.5_{stat} \begin{matrix} +0.3 \\ -0.3 \end{matrix} \Big|_{sys} \right) \frac{nb}{(\text{GeV}/c)^2}$$

$$\left\langle \frac{d\sigma_{LT}}{dt} \right\rangle = \left( 0.2 \pm 0.2_{stat} \begin{matrix} +0.2 \\ -0.2 \end{matrix} \Big|_{sys} \right) \frac{nb}{(\text{GeV}/c)^2}$$

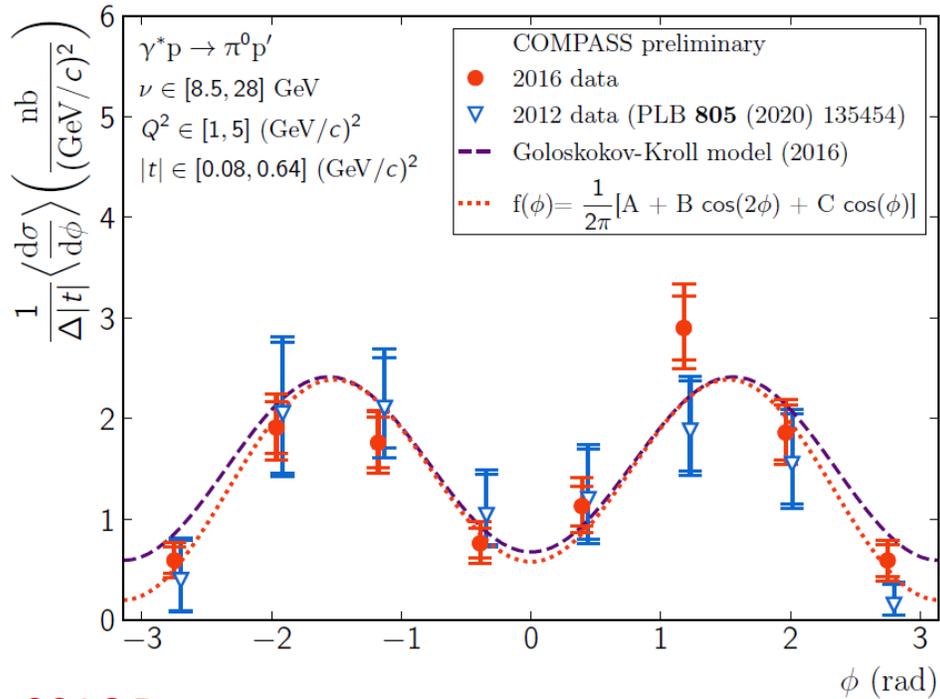
$$\langle \epsilon \rangle = 0.997$$

- Cross section extracted in a larger ( $\nu, Q^2$ ) domain, compared with the 2012 result. (COMPASS, PLB 805 (2020) 135454)
- Comparable  $|\sigma_{TT}|$  and  $\sigma_T + \epsilon\sigma_L$

# 2012-16 Exclusive $\pi^0$ Prod. Comparison



➤ Kinematic domain:  $\nu \in [8.5, 28]$  GeV and  $Q^2 \in [1, 5]$   $\text{GeV}^2/c^2$ ,  $\langle x_B \rangle = 0.10$



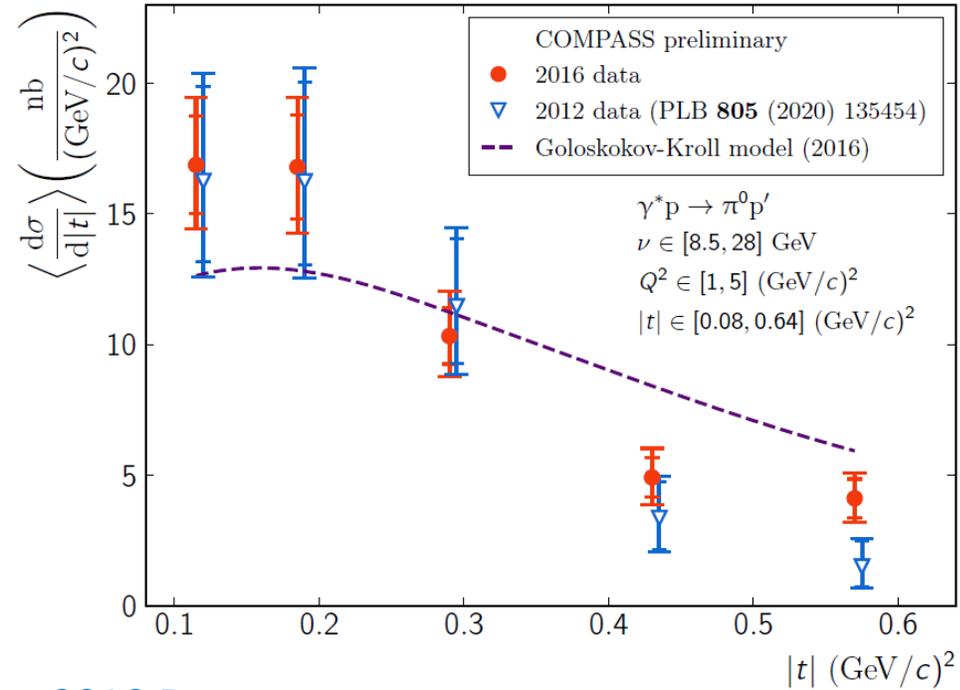
2016 Data:

$$\left\langle \frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} \right\rangle = \left( 8.7 \pm 0.5_{\text{stat}} \pm 1.0_{\text{sys}} \right) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

$$\left\langle \frac{d\sigma_{TT}}{dt} \right\rangle = \left( -6.3 \pm 0.8_{\text{stat}} \pm 0.4_{\text{sys}} \right) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

$$\left\langle \frac{d\sigma_{LT}}{dt} \right\rangle = \left( 0.6 \pm 0.3_{\text{stat}} \pm 0.3_{\text{sys}} \right) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

$$\langle \epsilon \rangle = 0.996$$



2012 Data:

$$\left\langle \frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} \right\rangle = \left( 8.1 \pm 0.9_{\text{stat}} \pm 1.1_{\text{sys}} \right) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

$$\left\langle \frac{d\sigma_{TT}}{dt} \right\rangle = \left( -6.0 \pm 1.3_{\text{stat}} \pm 0.7_{\text{sys}} \right) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

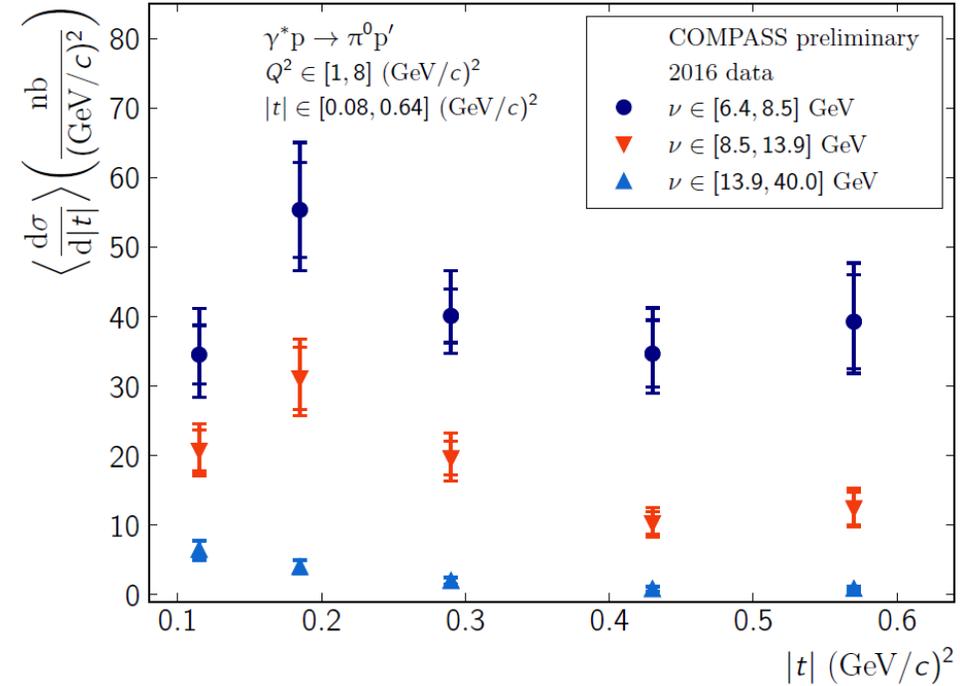
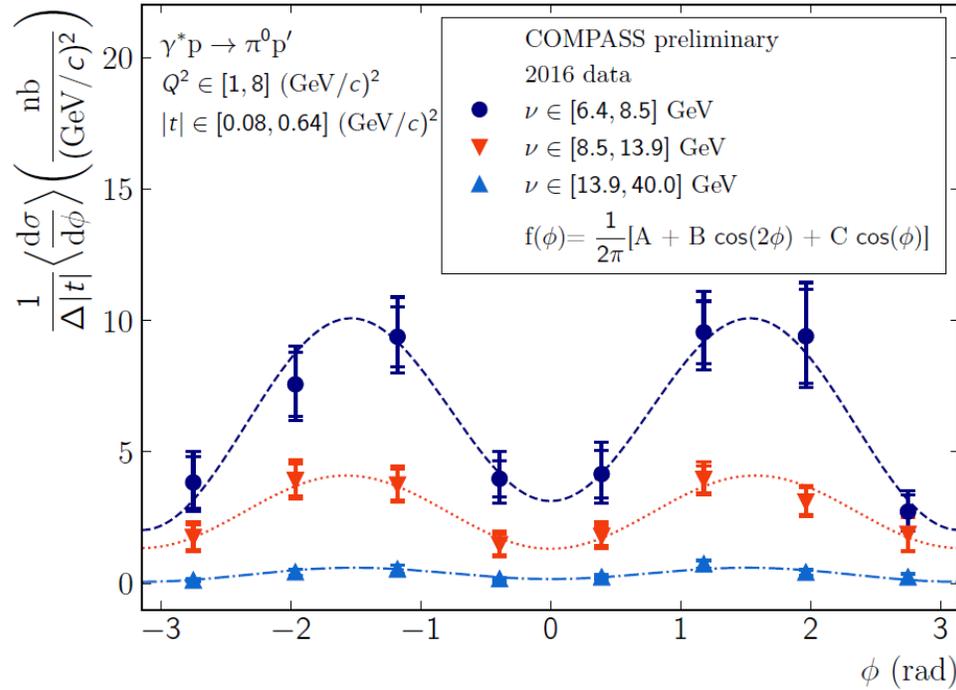
$$\left\langle \frac{d\sigma_{LT}}{dt} \right\rangle = \left( 1.4 \pm 0.5_{\text{stat}} \pm 0.3_{\text{sys}} \right) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

$$\langle \epsilon \rangle = 0.996$$

# New 2016 Exclusive $\pi^0$ Cross-section Evolution with $\nu$



➤ Cross section decreases with increasing  $\nu$

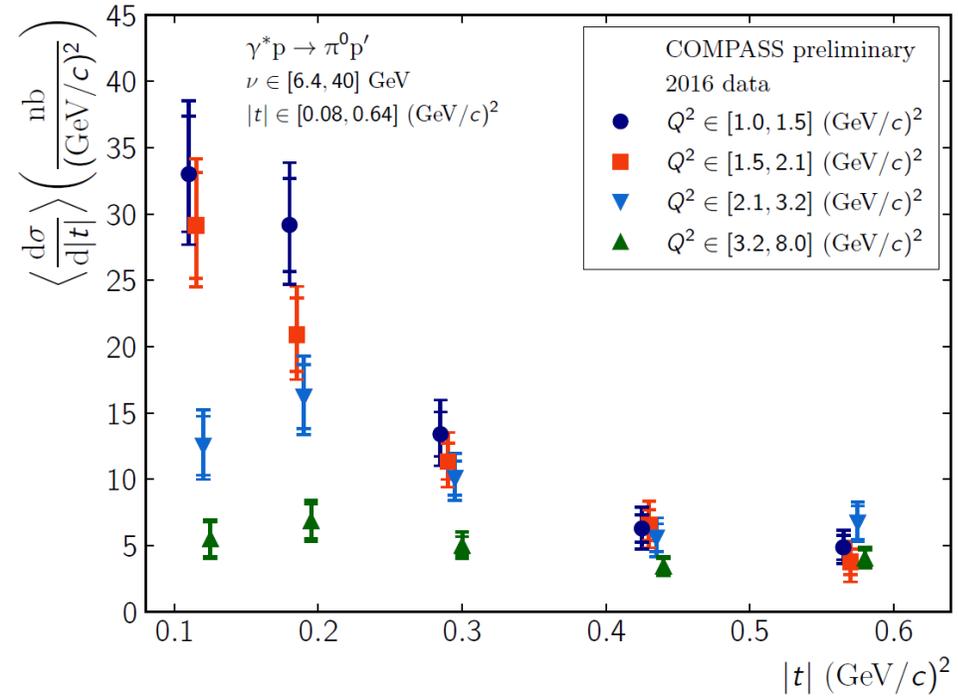
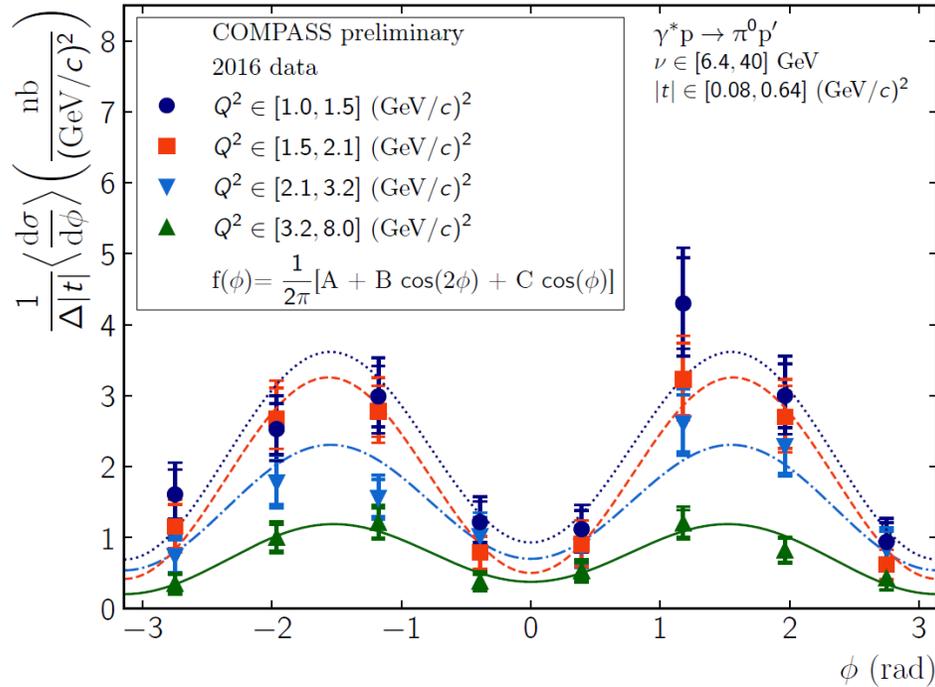


	$\langle \nu \rangle$ [GeV]	$\langle Q^2 \rangle$ [GeV <sup>2</sup> /c <sup>2</sup> ]	$\langle x_B \rangle$	$\langle \epsilon \rangle$
$\nu \in [6.4, 8.5]$	7.35	2.15	0.156	0.999
$\nu \in [8.5, 13.9]$	10.32	2.50	0.131	0.998
$\nu \in [13.9, 40.0]$	21.08	2.09	0.057	0.989

# New 2016 Exclusive $\pi^0$ Cross-section Evolution with $Q^2$

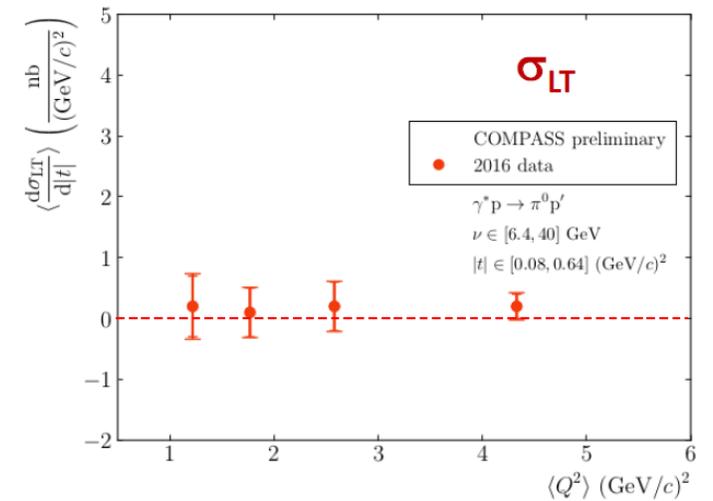
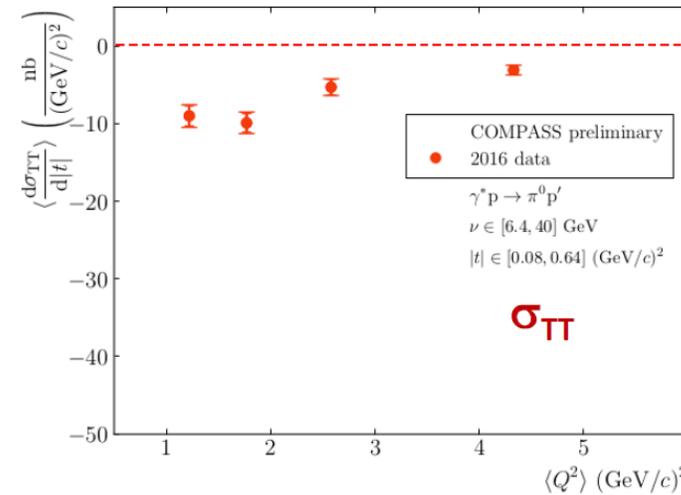
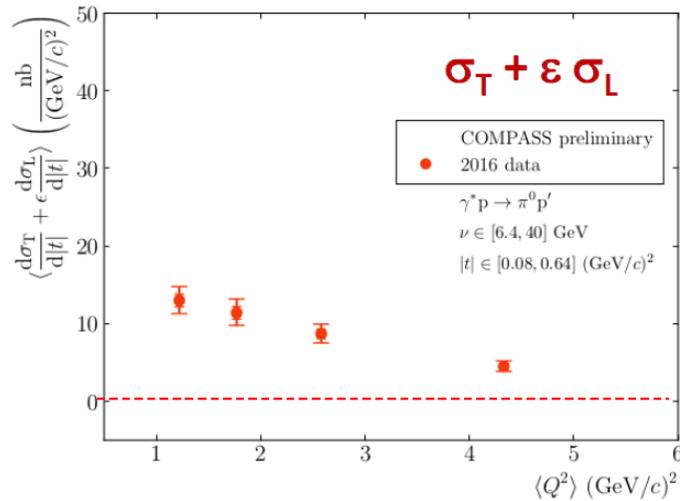
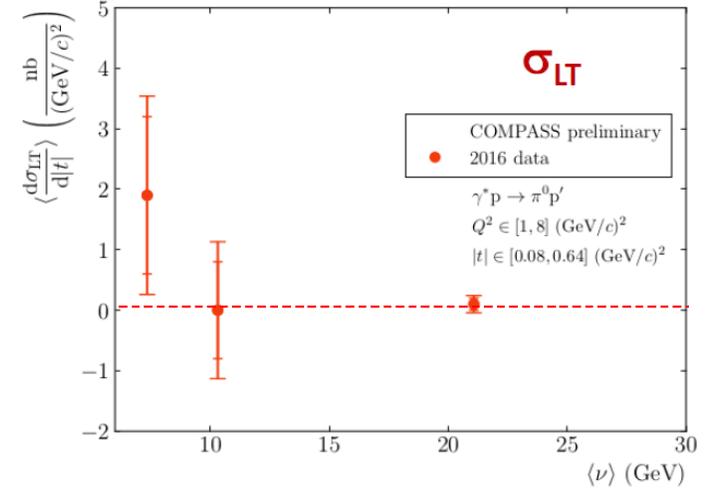
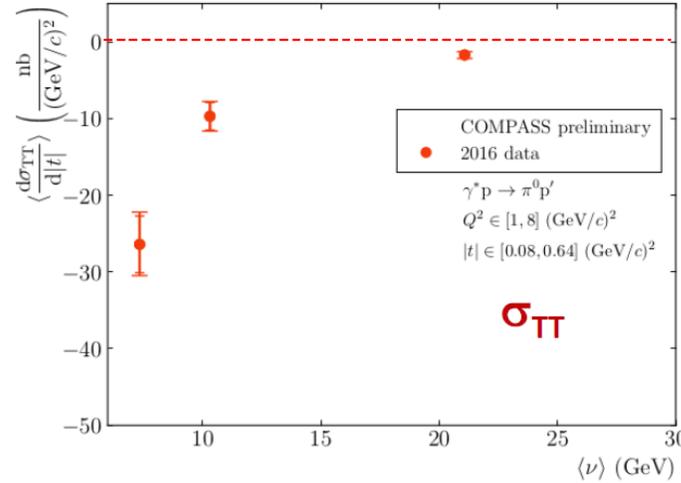
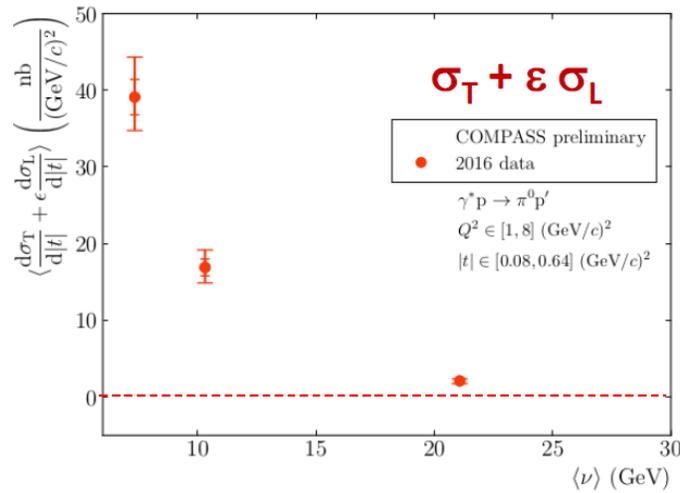


➤ Cross section decreases with increasing  $Q^2$



	$\langle Q^2 \rangle$ [GeV <sup>2</sup> /c <sup>2</sup> ]	$\langle \nu \rangle$ [GeV]	$\langle x_B \rangle$	$\langle \epsilon \rangle$
$Q^2 \in [1.0, 1.5]$	1.22	10.54	0.072	0.997
$Q^2 \in [1.5, 2.1]$	1.77	9.81	0.109	0.997
$Q^2 \in [2.1, 3.2]$	2.58	9.82	0.157	0.997
$Q^2 \in [3.2, 8.0]$	4.33	10.39	0.247	0.997

# New 2016 Evolution of the Structure Functions



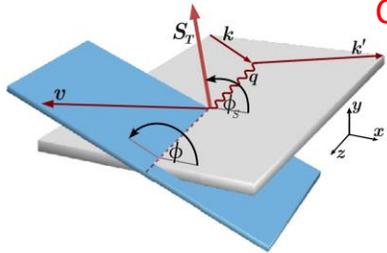
For both  $\sigma_T + \epsilon \sigma_L$  and  $\sigma_{TT}$   
 → Relatively larger evolution in  $\nu$ , smaller in  $Q^2$

$\sigma_{LT}$  consistent with 0

# 2007 & 2010 HEMP with Transversely Polarized Target



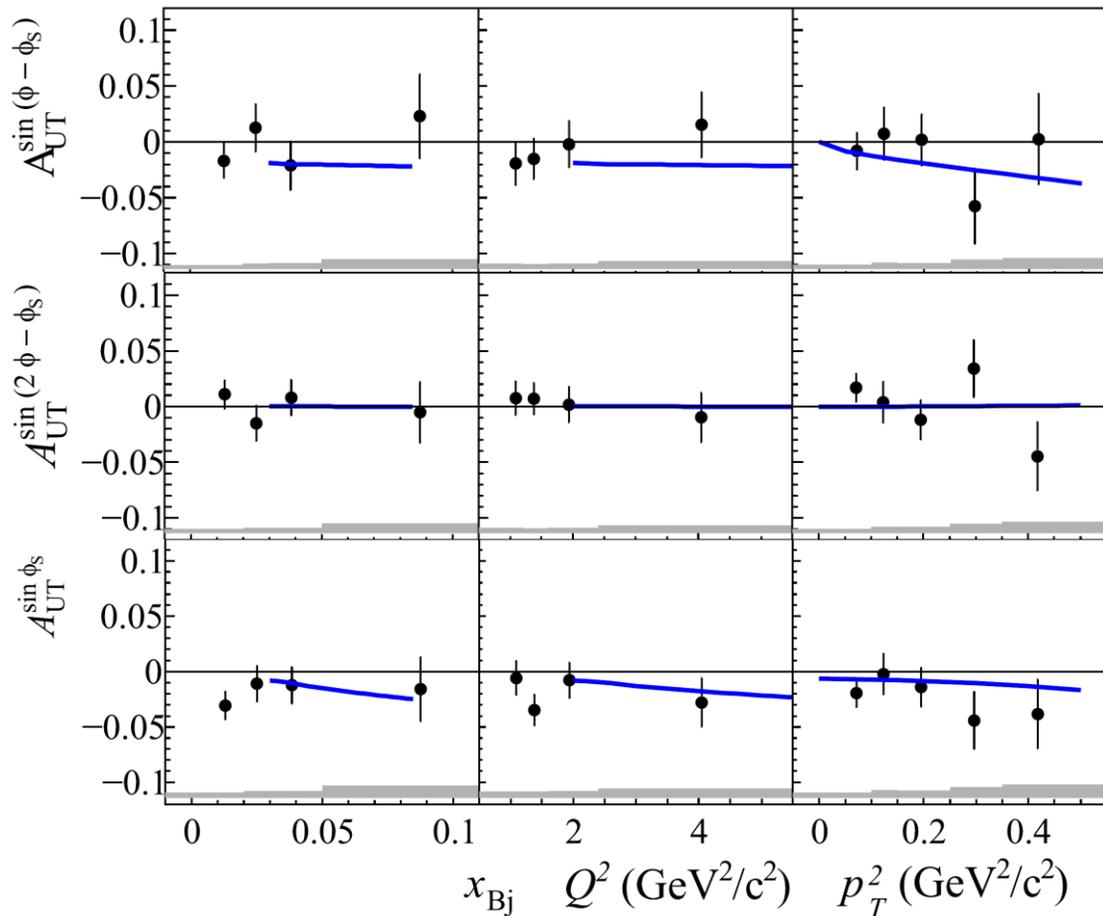
COMPASS, NPB 865 (2012) 1-20, PLB 731 (2014) 19



$$\rho^0 \rightarrow \pi^+ \pi^-$$

$$E_{\rho^0} = \frac{1}{\sqrt{2}} \left( \frac{2}{3} E^u + \frac{1}{3} E^d + \frac{3}{4} E^g / x \right)$$

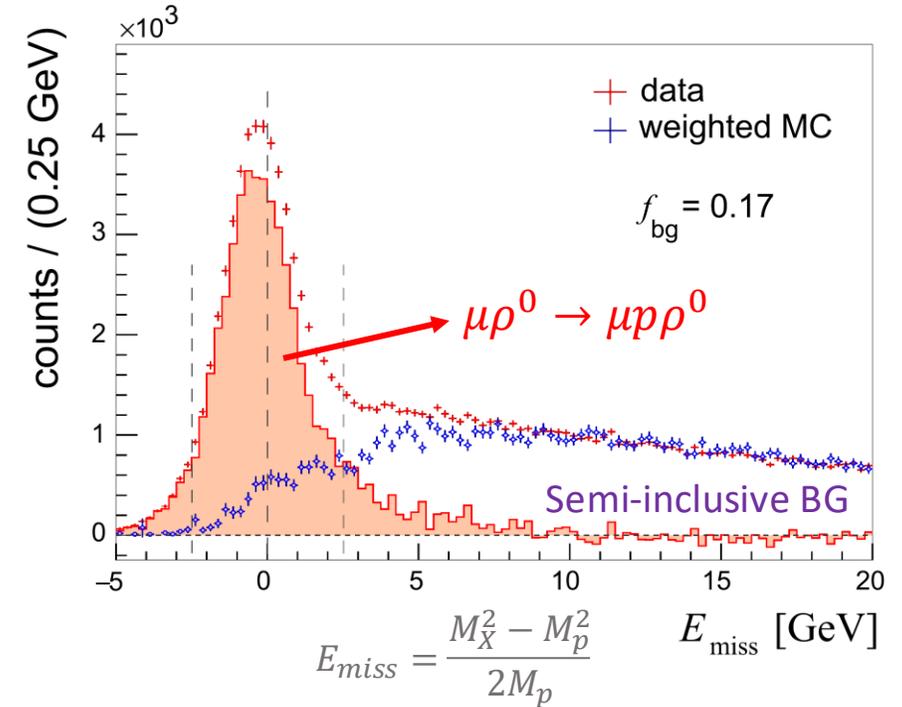
➤ Exclusivity ensured by “missing mass technique”



$$\propto \text{Im}(E^* H)$$

$$\propto \text{Im}(E^* \bar{E}_T)$$

$$\propto \text{Im}(E^* \bar{E}_T - H^* H_T)$$

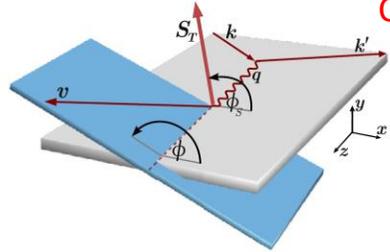


➤ Sensibility to  $E$  and  $H_T$

# 2007 & 2010 HEMP with Transversely Polarized Target

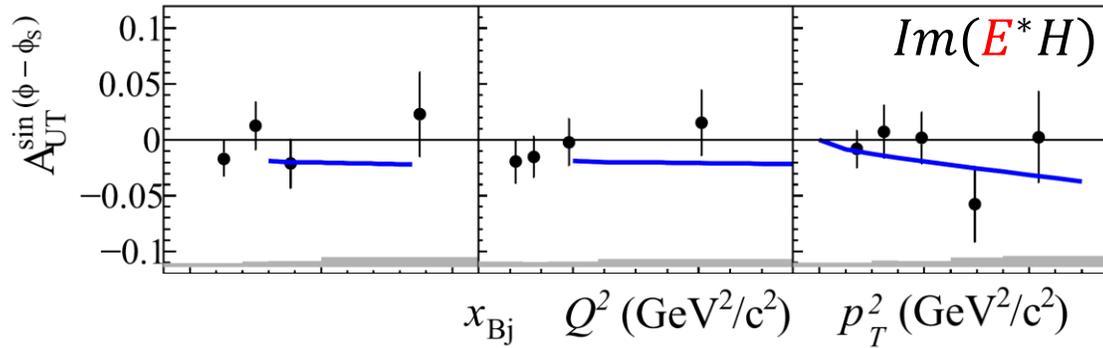


COMPASS, NPB 865 (2012) 1-20, PLB 731 (2014) 19



$$\rho^0 \rightarrow \pi^+ \pi^-$$

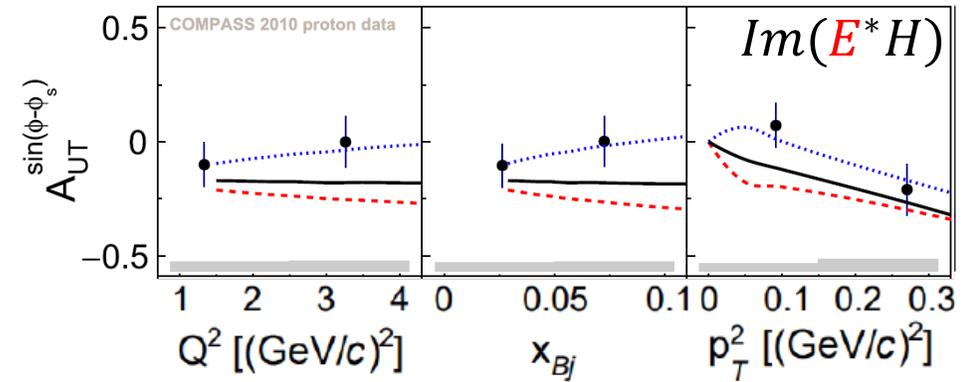
$$E_{\rho^0} = \frac{1}{\sqrt{2}} \left( \frac{2}{3} E^u + \frac{1}{3} E^d + \frac{3}{4} E^g / x \right)$$



COMPASS, NPB 865 (2012) 1-20, PLB 731 (2014) 19

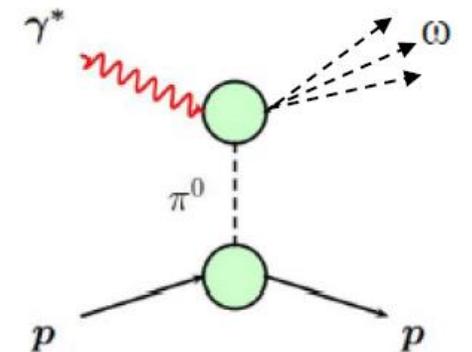
$$\omega \rightarrow \pi^+ \pi^- \pi^0$$

$$E_{\omega} = \frac{1}{\sqrt{2}} \left( \frac{2}{3} E^u - \frac{1}{3} E^d + \frac{1}{4} E^g / x \right)$$

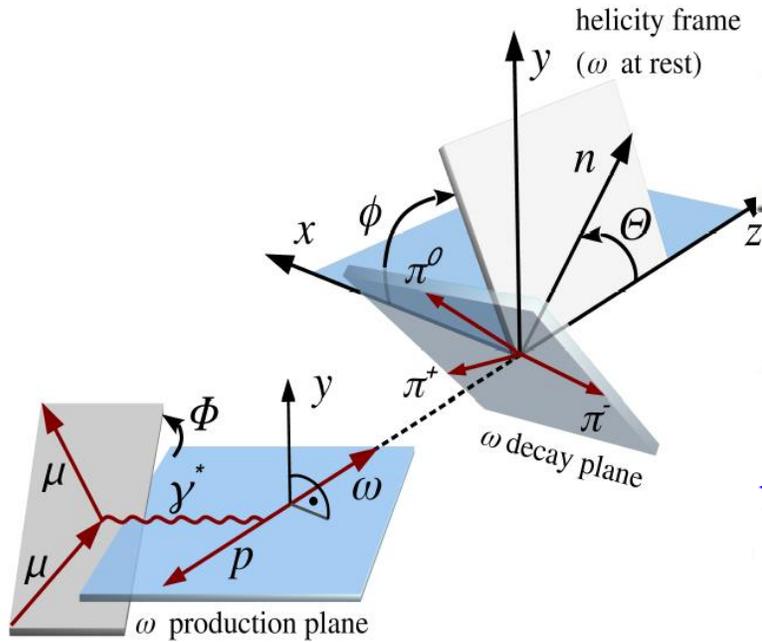


— No pion pole     
 ⋯ Negative  $\pi\omega$  form factor  
- - - Positive  $\pi\omega$  form factor

- $E^u$  and  $E^d$  are of opposite sign  $\rightarrow \omega$  is more promising for GPD study
- Nevertheless, obscured by the inherent pion pole contribution



# Exclusive $\omega$ Production on Unpolarized Proton



## Experimental angular distributions

$$\mathcal{W}^{U+L}(\Phi, \phi, \cos \Theta) = \mathcal{W}^U(\Phi, \phi, \cos \Theta) + P_b \mathcal{W}^L(\Phi, \phi, \cos \Theta)$$

## 15 unpolarized SDMEs in $\mathcal{W}^U$ and 8 polarized in $\mathcal{W}^L$

$$\begin{aligned} \mathcal{W}^U(\Phi, \phi, \cos \Theta) = & \frac{3}{8\pi^2} \left[ \frac{1}{2}(1 - r_{00}^{04}) + \frac{1}{2}(3r_{00}^{04} - 1) \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^{04}\} \sin 2\Theta \cos \phi - r_{1-1}^{04} \sin^2 \Theta \cos 2\phi \right. \\ & - \epsilon \cos 2\Phi \left( r_{11}^1 \sin^2 \Theta + r_{00}^1 \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^1\} \sin 2\Theta \cos \phi - r_{1-1}^1 \sin^2 \Theta \cos 2\phi \right) \\ & \left. - \epsilon \sin 2\Phi \left( \sqrt{2}\text{Im}\{r_{10}^2\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^2\} \sin^2 \Theta \sin 2\phi \right) \right. \\ & + \sqrt{2\epsilon(1+\epsilon)} \cos \Phi \left( r_{11}^5 \sin^2 \Theta + r_{00}^5 \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^5\} \sin 2\Theta \cos \phi - r_{1-1}^5 \sin^2 \Theta \cos 2\phi \right) \\ & \left. + \sqrt{2\epsilon(1+\epsilon)} \sin \Phi \left( \sqrt{2}\text{Im}\{r_{10}^6\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^6\} \sin^2 \Theta \sin 2\phi \right) \right], \end{aligned}$$

$$\begin{aligned} \mathcal{W}^L(\Phi, \phi, \cos \Theta) = & \frac{3}{8\pi^2} \left[ \sqrt{1-\epsilon^2} \left( \sqrt{2}\text{Im}\{r_{10}^3\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^3\} \sin^2 \Theta \sin 2\phi \right) \right. \\ & + \sqrt{2\epsilon(1-\epsilon)} \cos \Phi \left( \sqrt{2}\text{Im}\{r_{10}^7\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^7\} \sin^2 \Theta \sin 2\phi \right) \\ & \left. + \sqrt{2\epsilon(1-\epsilon)} \sin \Phi \left( r_{11}^8 \sin^2 \Theta + r_{00}^8 \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^8\} \sin 2\Theta \cos \phi - r_{1-1}^8 \sin^2 \Theta \cos 2\phi \right) \right] \end{aligned}$$

➤  $\epsilon \rightarrow 1$ , small  $\mathcal{W}^L$

# 2012 Exclusive $\omega$ Prod. on Unpolarized Proton



SCHC ( $\lambda_\gamma = \lambda_V$ )

(S-Channel Helicity Conservation)

SCHC implies:

•  $r_{1-1}^1 + \text{Im} r_{1-1}^2 = 0$

=  $-0.010 \pm 0.032 \pm 0.047$  OK

•  $\text{Re} r_{10}^5 + \text{Im} r_{10}^6 = 0$

=  $0.014 \pm 0.011 \pm 0.013$  OK

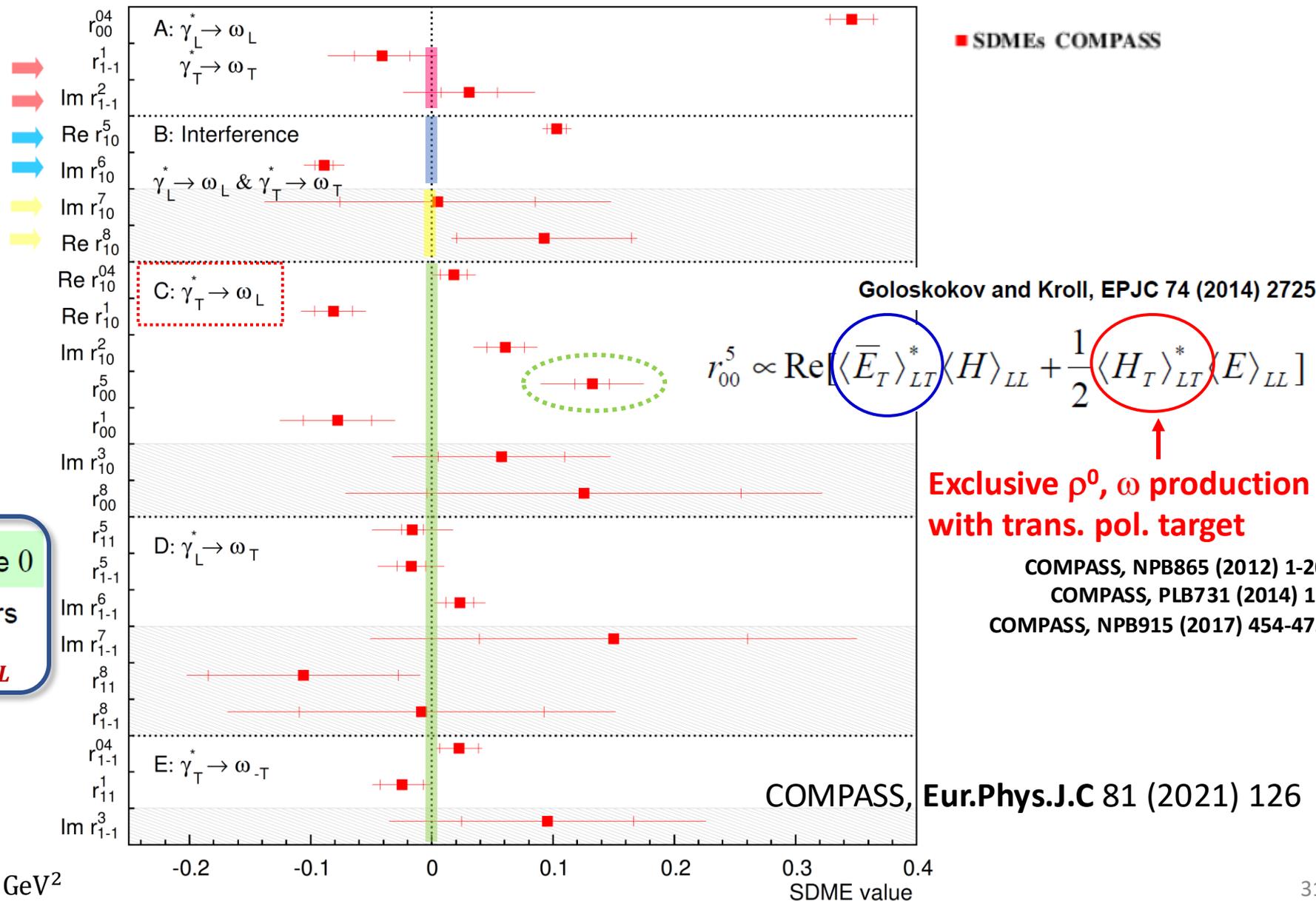
•  $\text{Im} r_{10}^7 - \text{Re} r_{10}^8 = 0$

=  $-0.088 \pm 0.110 \pm 0.196$  OK

• all elements of classes C, D, E should be 0

for  $\gamma_L^* \rightarrow \omega_T$  and  $\gamma_T^* \rightarrow \omega_T$  OK within errors

**NOT OBSERVED for transitions  $\gamma_T^* \rightarrow \omega_L$**



$\langle Q^2 \rangle = 2.1 \text{ GeV}^2, \langle W \rangle = 7.6 \text{ GeV}, \langle P_T^2 \rangle = 0.16 \text{ GeV}^2$

# 2012 Exclusive $\rho^0$ Prod. on Unpolarized Proton



SCHC ( $\lambda_\gamma = \lambda_V$ )

(S-Channel Helicity Conservation)

SCHC implies:

•  $r_{1-1}^1 + \text{Im} r_{1-1}^2 = 0$

=  $-0.000 \pm 0.006$  OK

•  $\text{Re} r_{10}^5 + \text{Im} r_{10}^6 = 0$

=  $-0.011 \pm 0.003$  Violation

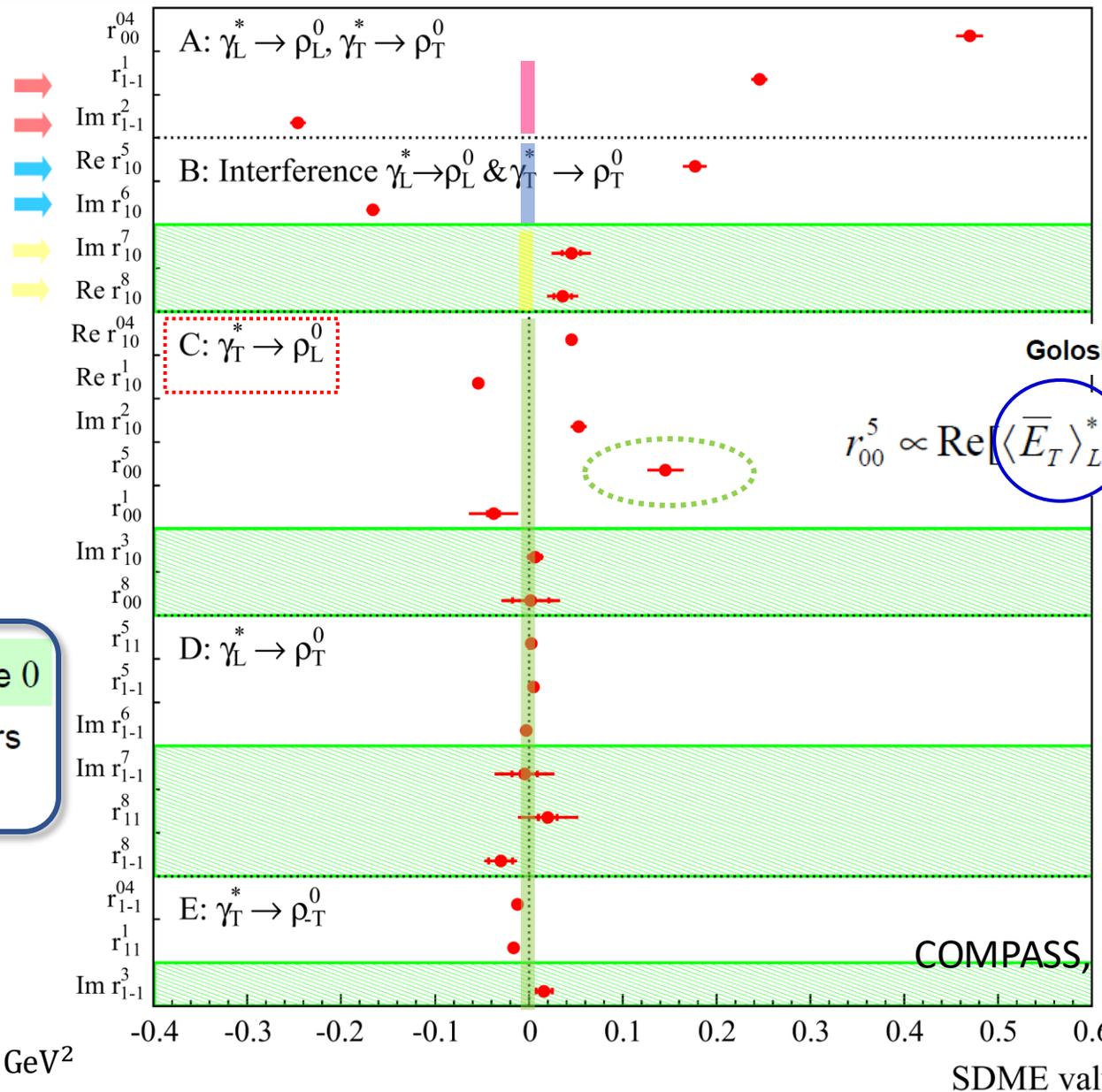
•  $\text{Im} r_{10}^7 - \text{Re} r_{10}^8 = 0$

=  $-0.009 \pm 0.031$  OK

• all elements of classes C, D, E should be 0

for  $\gamma_L^* \rightarrow \omega_T$  and  $\gamma_T^* \rightarrow \omega_T$  OK within errors

**NOT OBSERVED for transitions  $\gamma_T^* \rightarrow \rho_L^0$**



■ SDMEs COMPASS

Goloskokov and Kroll, EPJC 74 (2014) 2725

$$r_{00}^5 \propto \text{Re} \left[ \langle \bar{E}_T \rangle_{LT}^* \langle H \rangle_{LL} + \frac{1}{2} \langle H_T \rangle_{LT}^* \langle E \rangle_{LL} \right]$$

First term dominates  
→ Probes  $\bar{E}_T$

COMPASS, Eur.Phys.J.C 83 (2023) 924

## DVCS cross sections with polarized $\mu^+$ and $\mu^-$

- Beam charge-spin sum  $\rightarrow \text{Im}\mathcal{H}(\xi,t) \rightarrow$  Transverse extension of partons as a function of  $x_{Bj}$
- Beam charge-spin difference  $\rightarrow \text{Re}\mathcal{H}(\xi,t) \rightarrow$  D-term, pressure distribution

## HEMP of $\pi^0$ , $\rho$ , $\omega$ , $\phi$ , $J/\psi$

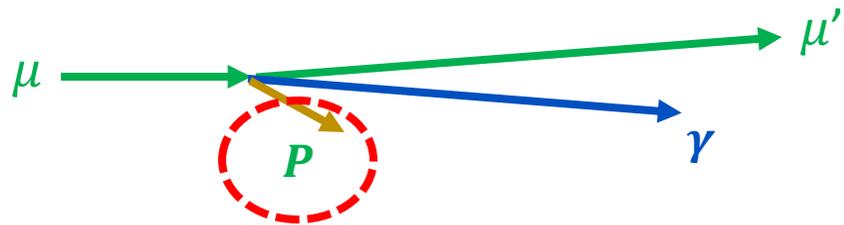
- Cross section of  $\pi^0 \rightarrow$  Submitted to Physics Letters B
- SDME of  $\rho$  &  $\omega \rightarrow$  Transversity GPDs & Flavor Decomposition
- $\phi$ ,  $J/\psi \rightarrow$  underway

➤ More results are coming!



# Backup Slides

# COMPASS 2016 Preliminary Results



$$\Delta\phi = \phi^{\text{cam.}} - \phi^{\text{spec.}}$$

Proton azimuthal angle

$$\Delta p_T = |p_T^{\text{cam.}}| - |p_T^{\text{spec.}}|$$

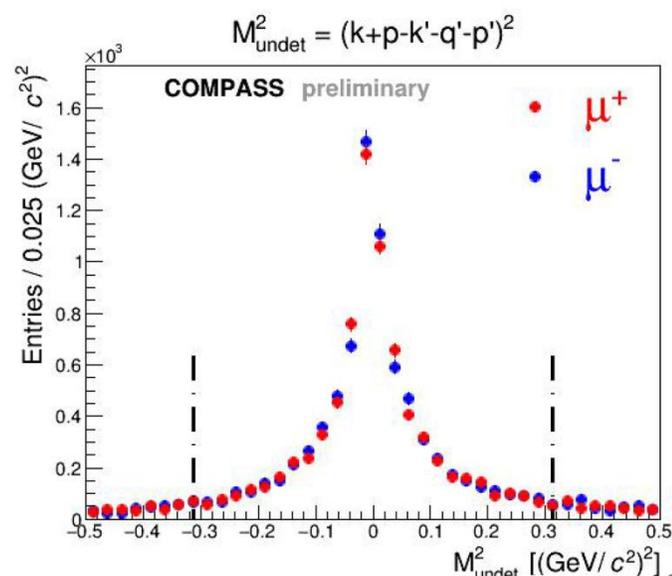
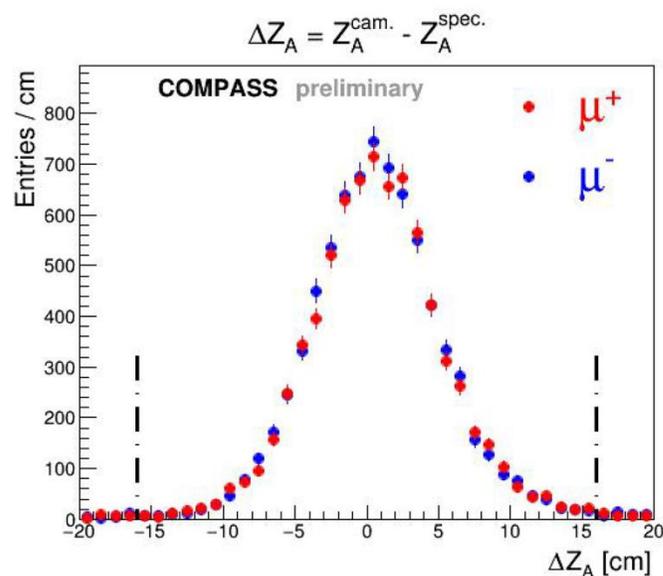
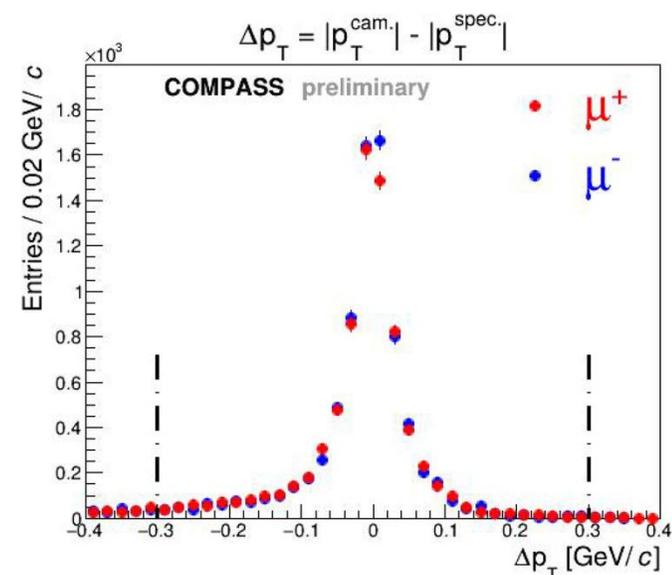
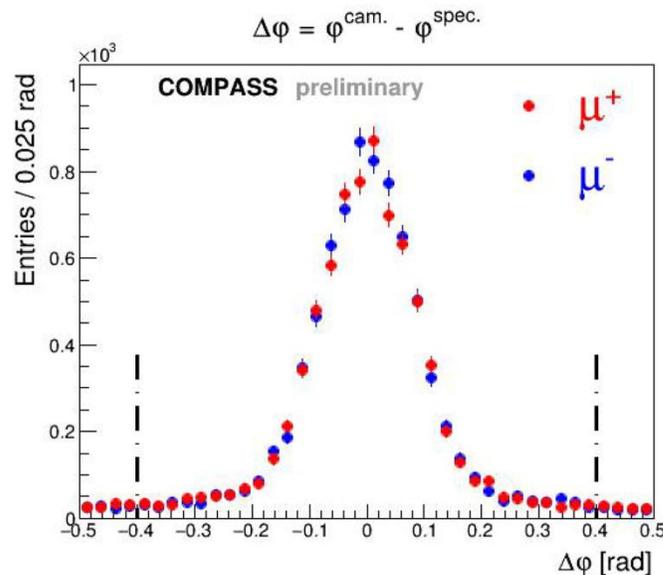
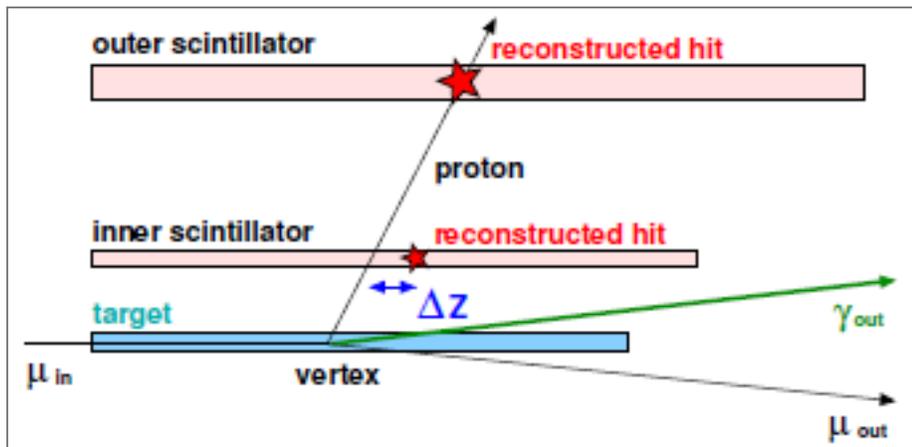
Proton momentum

$$\Delta z_A = z_A^{\text{cam.}} - z_A^{\text{spec.}}$$

Proton track position

$$M_{\text{undet}}^2 = (k + p - k' - q' - p')^2$$

Energy momentum balance



➤ **Main background of exclusive single photon events:  $\pi^0$  decay**

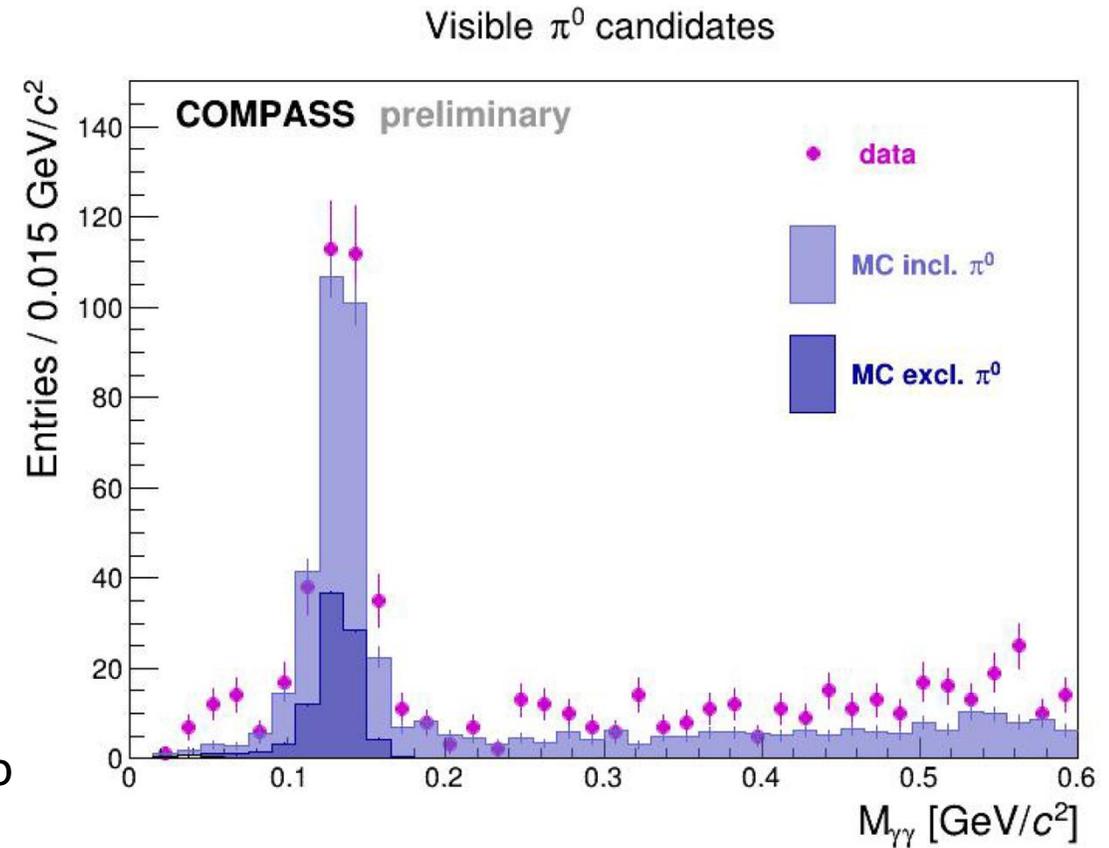
➤ **Visible (both  $\gamma$  detected) – subtracted**

A high-energy DVCS photon candidate is combined with all detected photons with energies lower than the DVCS threshold: (4,5) GeV in Ecal (0,1) respectively

➤ **Invisible (one  $\gamma$  lost) – estimated by MC**

- **Semi-inclusive LEPTO 6.1**
- **Exclusive HEPGEN  $\pi^0$  (GK model)**

The sum of LEPTO and HEPGEN contributions is normalized to the  $\pi^0$  peak in  $M_{\gamma\gamma}$  of the real data



# Beam Charge-spin Difference

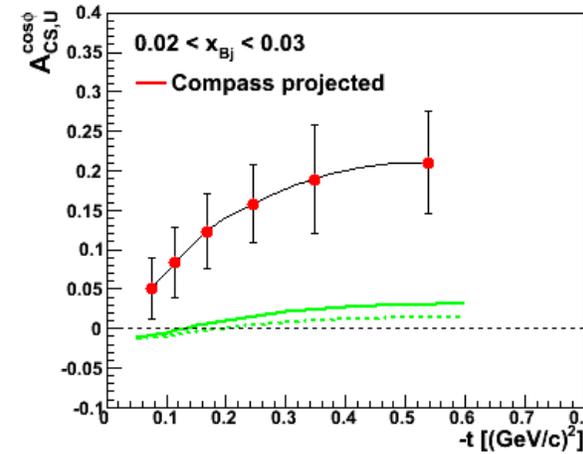
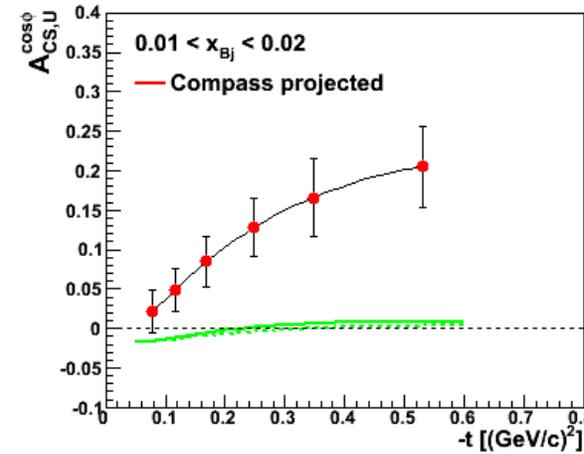
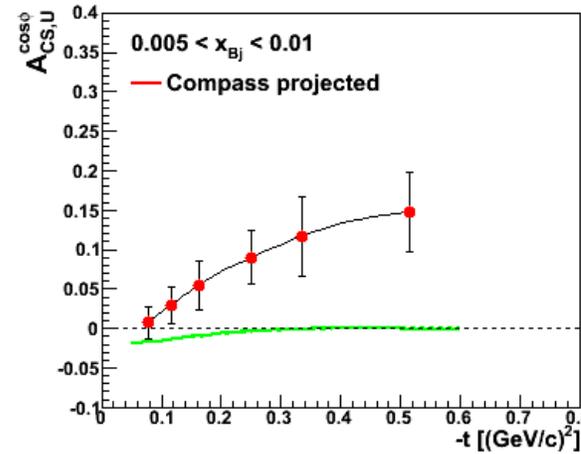
$$\mathcal{D}_{CS,U}(\phi) \equiv d\sigma(\mu^{+\leftarrow}) - d\sigma(\mu^{-\rightarrow}) \rightarrow c_0^I + c_1^I \cos \phi$$

$$BCSA = \mathcal{D}_{CS,U}/S_{CS,U} = A_0 + A_{CS,U}^{\cos\phi} \cos\phi + A_2 \cos 2\phi$$

$$c_1^I \rightarrow \text{Re } F_1 \mathcal{H}$$

**VGG**

**KM10 – fit to data**

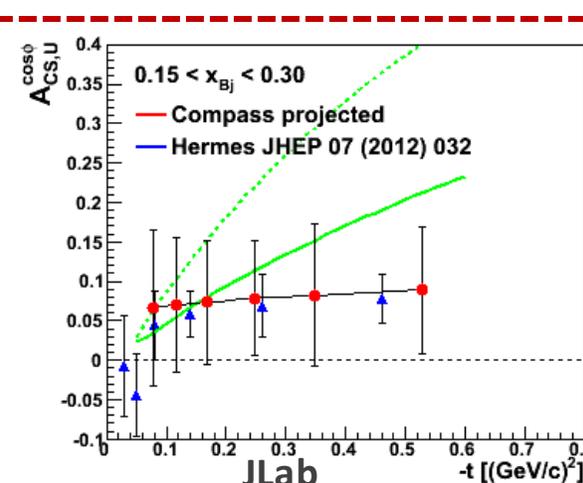
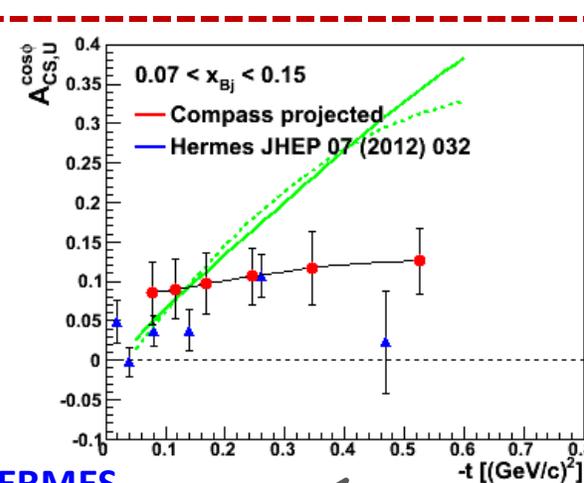
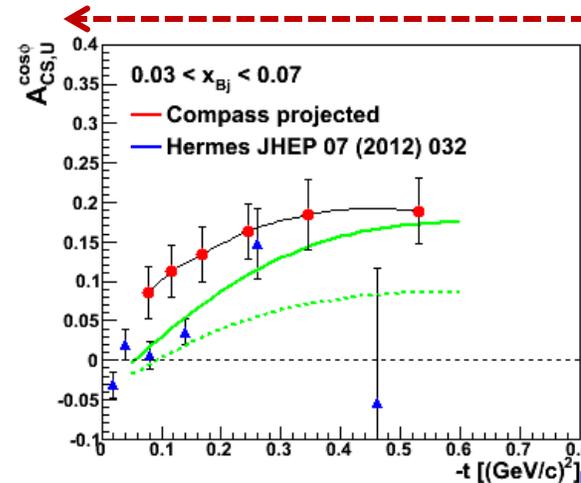


- With  $\text{Re } F_1 \mathcal{H}$  and  $\text{Im } F_1 \mathcal{H}$   
→ Extraction of **D-term**

$\text{Re } \mathcal{H} > 0$  at H1

$< 0$  at HERMES

Value of  $x_{Bj}$  for the node?

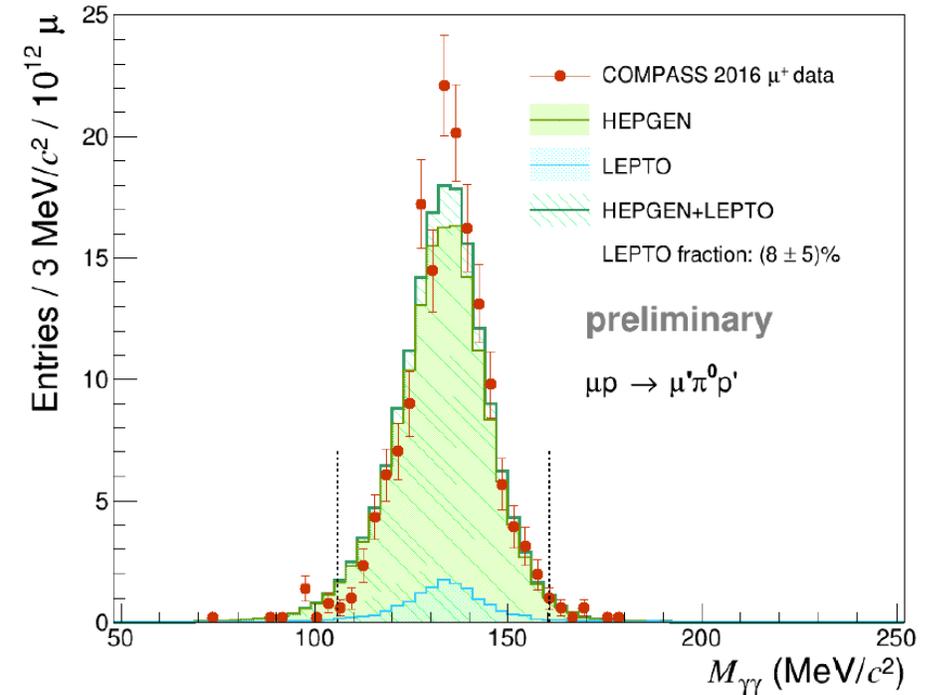
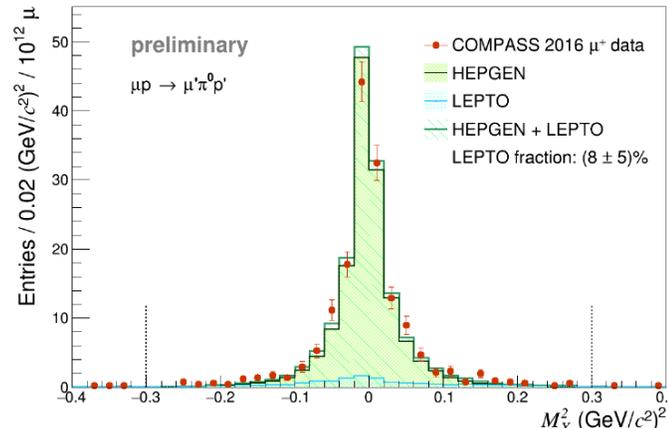
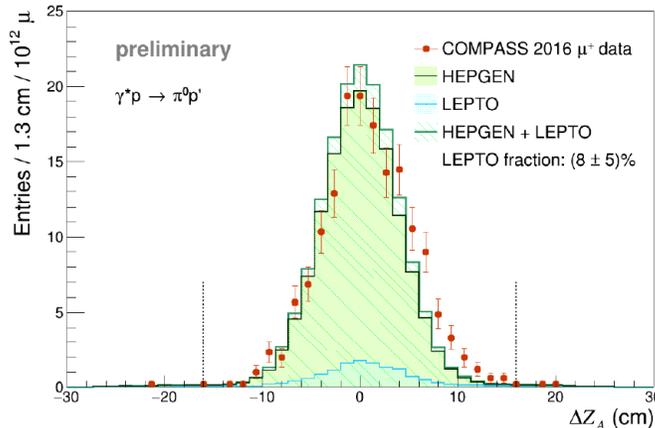
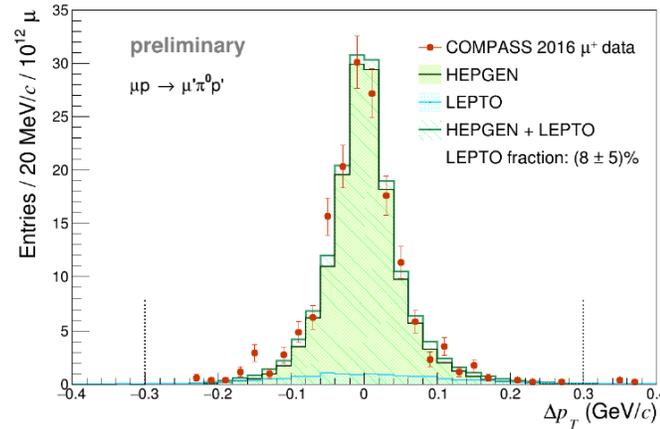
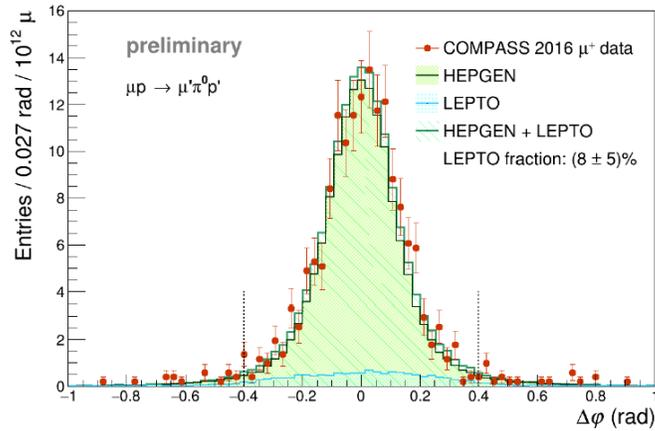


**COMPASS** 2 years of data  $E_\mu = 160 \text{ GeV}$   $1 < Q^2 < 8 \text{ GeV}^2$

# Exclusive $\pi^0$ Selection and Background Estimation



- Exclusivity ensured by cuts on *exclusivity variables*, similar to DVCS.
- Background fraction determined by fitting the exclusivity variables with Monte Carlo simulations.
  - *LEPTO* for non-exclusive background
  - *HEPGEN* of exclusive  $\pi^0$  for signal



- In 2016 data, non-exclusive background fraction in data  $\rightarrow 8 \pm 5\%$

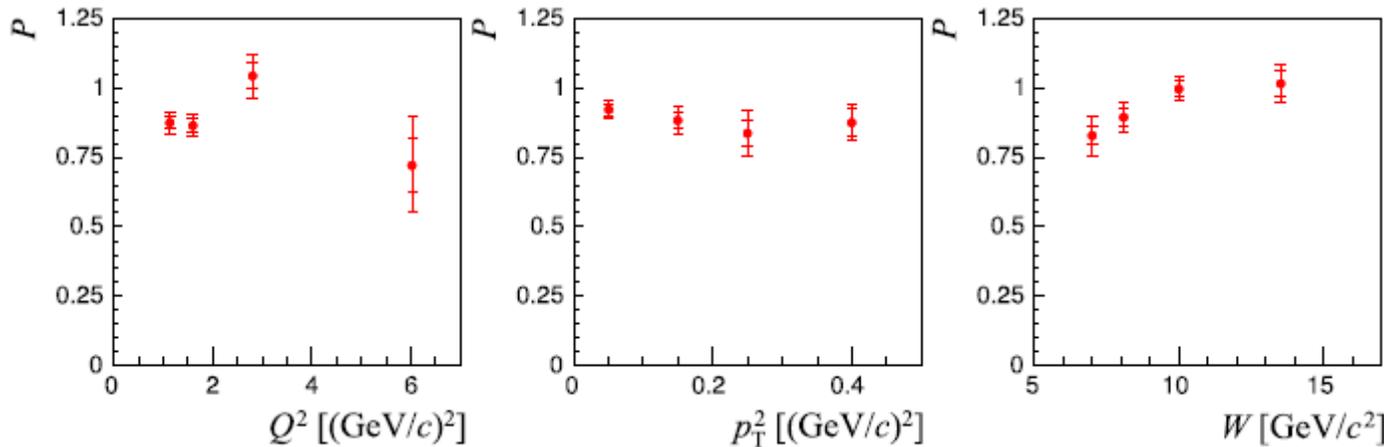
# 2012 NPE-to-UPE Asymmetry



$$P = \frac{2r_{1-1}^1}{1 - r_{00}^{04} - 2r_{1-1}^{04}} \approx \frac{d\sigma_T^N(\gamma_T^* \rightarrow V_T) - d\sigma_T^U(\gamma_T^* \rightarrow V_T)}{d\sigma_T^N(\gamma_T^* \rightarrow V_T) + d\sigma_T^U(\gamma_T^* \rightarrow V_T)}$$

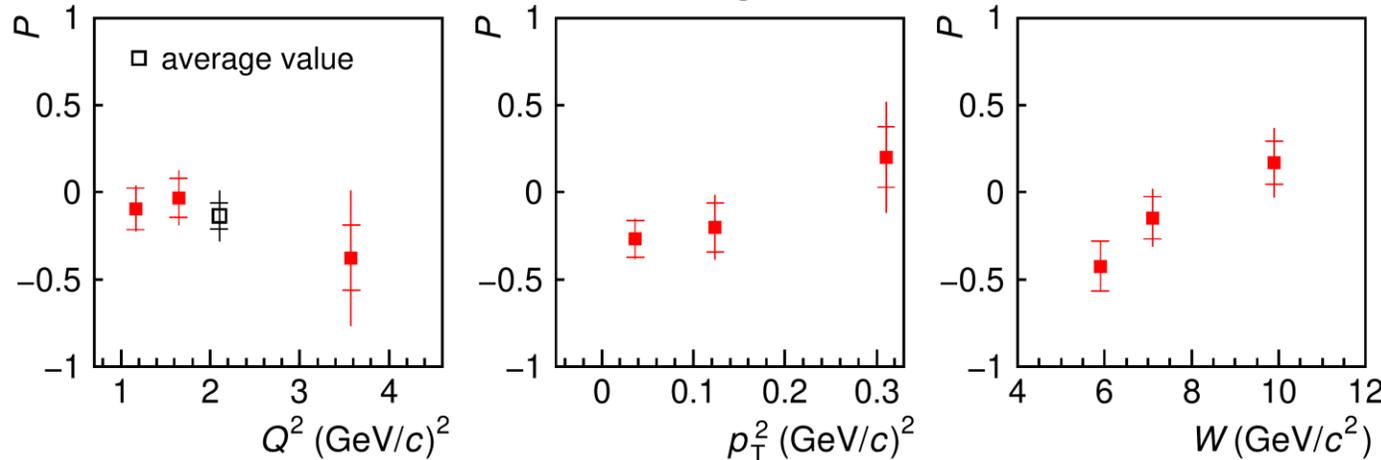
NPE-to-UPE asymmetry of cross sections for transitions  $\gamma_T^* \rightarrow V_T$

- NPE: Natural Parity Exchange
- UPE: Unnatural Parity Exchange



$\rho^0$  COMPASS, Eur.Phys.J.C 83 (2023) 924

- NPE Dominance
- NPE  $\rightarrow$  GPDs  $E, H$



$\omega$  COMPASS, Eur.Phys.J.C 81 (2021) 126

- NPE  $\approx$  UPE on average
- UPE Dominance at small  $W$  and  $p_T^2$
- UPE  $\rightarrow$  GPDs  $\tilde{E}, \tilde{H}$   
+ Pion pole (dominant)

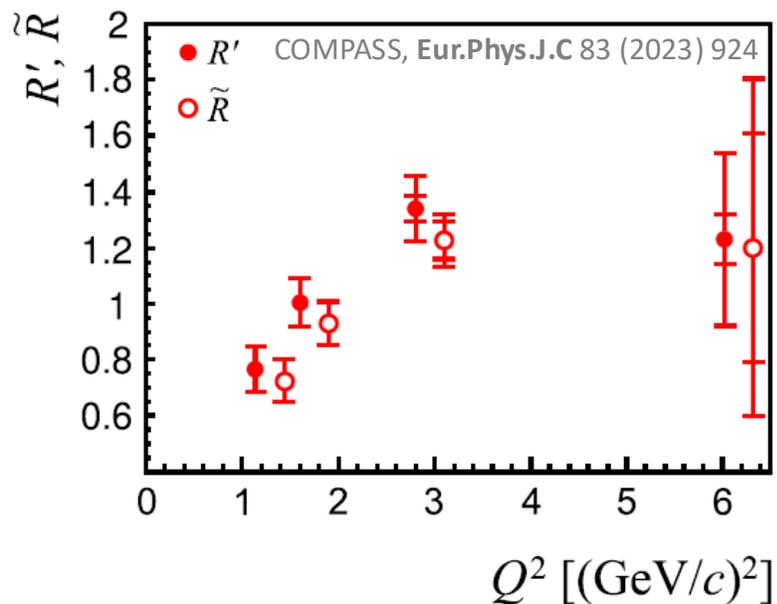
# 2012 $R = \sigma_L/\sigma_T$ for Exclusive $\rho^0$ Production



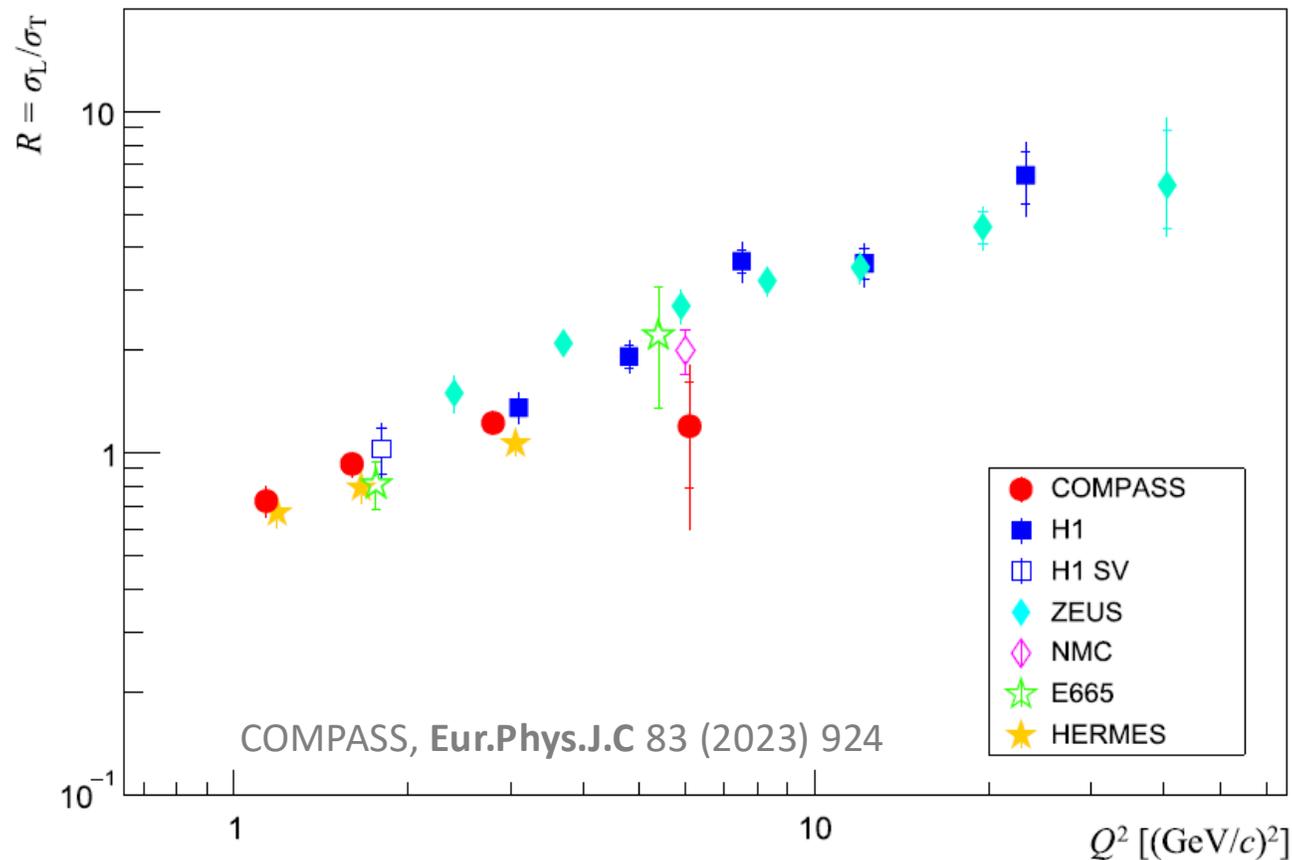
➤ Longitudinal-to-transverse  $\gamma^*$  cross section ratio:  $R = \frac{\sigma_L(\gamma_L^* \rightarrow V)}{\sigma_T(\gamma_T^* \rightarrow V)}$

• Commonly used “effective” ratio ( $R' = R$  only if SCHC):  $R' = \frac{1}{\epsilon} \frac{r_{00}^{04}}{1 - r_{00}^{04}}$

• Use of  $\tilde{R}$ , which takes SCHC violation into consideration, is preferred.



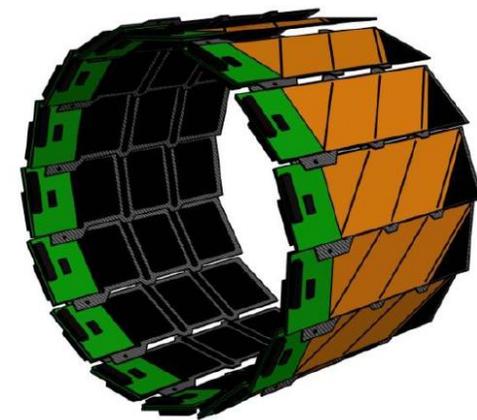
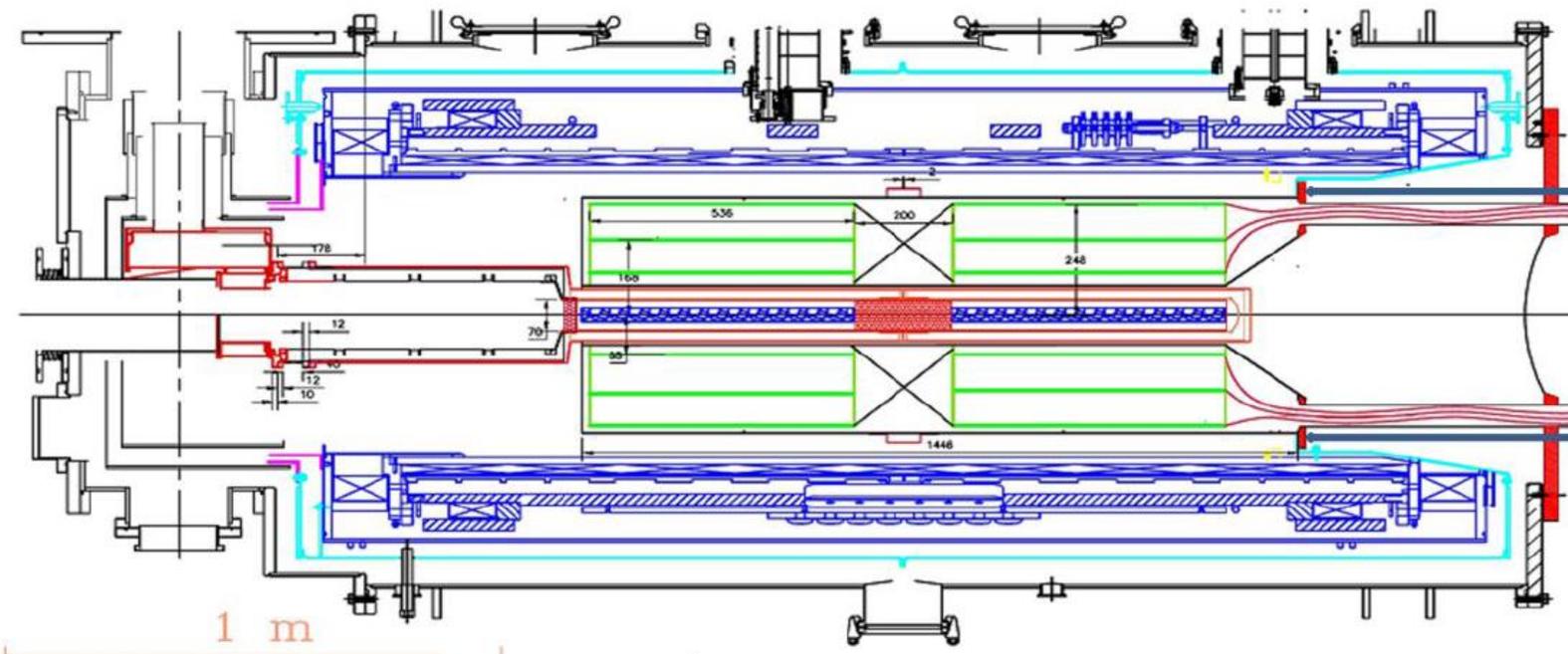
Results of all experiments with  $Q^2 > 1$  (GeV/c)<sup>2</sup>



➤ Leading-order pQCD prediction:  $Q^2/M_\rho^2 \rightarrow$  deviation due to effect of QCD evolution and  $q_T$

# Possible RPD for COMPASS++ / AMBER

A recoil proton detector (RPD) is mandatory to ensure the exclusivity. A Silicon detector is included *between* the target surrounded by the modified MW cavity *and* the polarizing magnet



A technology developed at JINR for NICA for the BM@N experiment

No possibility for ToF → PID of  $p/\pi$  with  $dE/dx$   
Momentum and trajectory measurements  
 $|t|_{\min} \sim 0.1 \text{ GeV}$