

# Theoretical predictions on the QCD critical point

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of Particle and Nuclear Physics

  
**KENT STATE**  
UNIVERSITY



# Outline

- ① The QCD phase diagram
- ② Effective models for QCD
- ③ Lattice QCD constraints
- ④ Theory and Experiment
- ⑤ Summary

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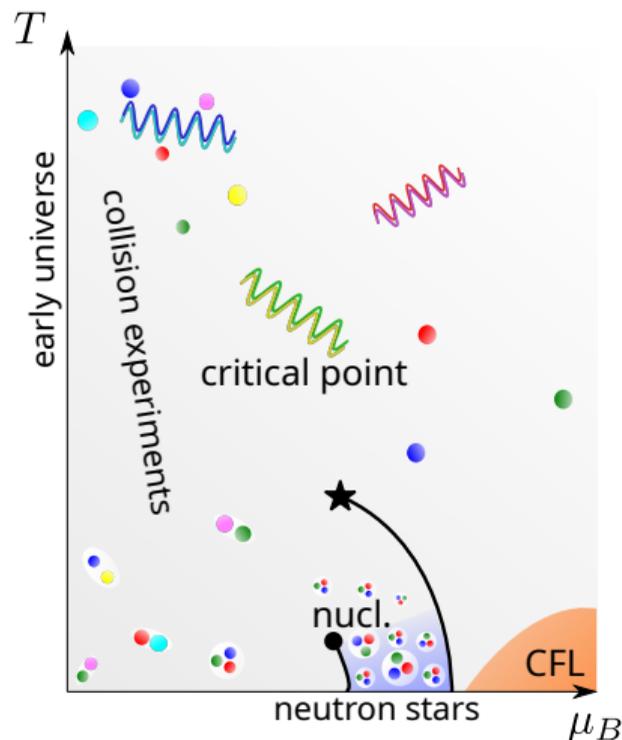
- 1 The QCD phase diagram
- 2 Effective models for QCD
- 3 Lattice QCD constraints
- 4 Theory and Experiment
- 5 Summary

# QCD Phase Diagram

We can explore the QCD phase diagram by changing  $\sqrt{s}$  in relativistic heavy ion collisions

Models predict a first order phase transition line with a critical point

Lattice QCD is the most reliable theoretical tool to study the QCD phase diagram.



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We can explore the QCD phase diagram by changing  $\sqrt{s}$  in relativistic heavy ion collisions

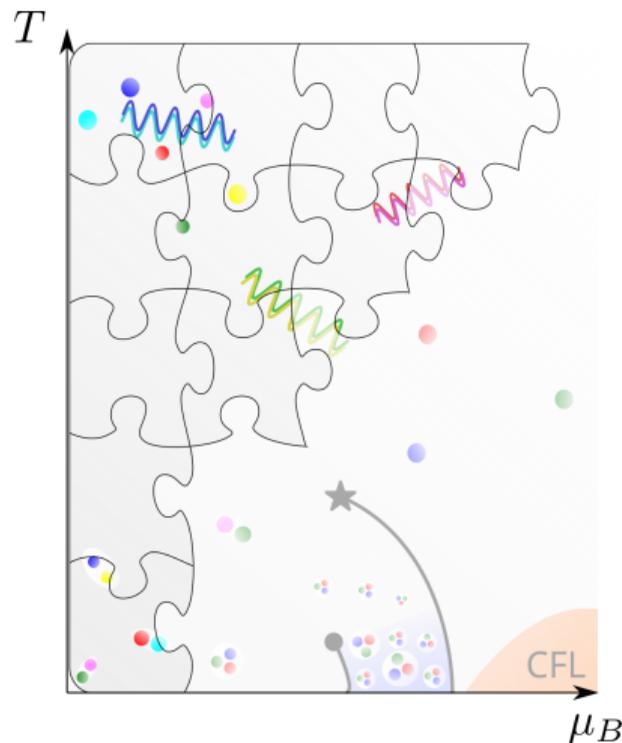
Models predict a first order phase transition line with a critical point

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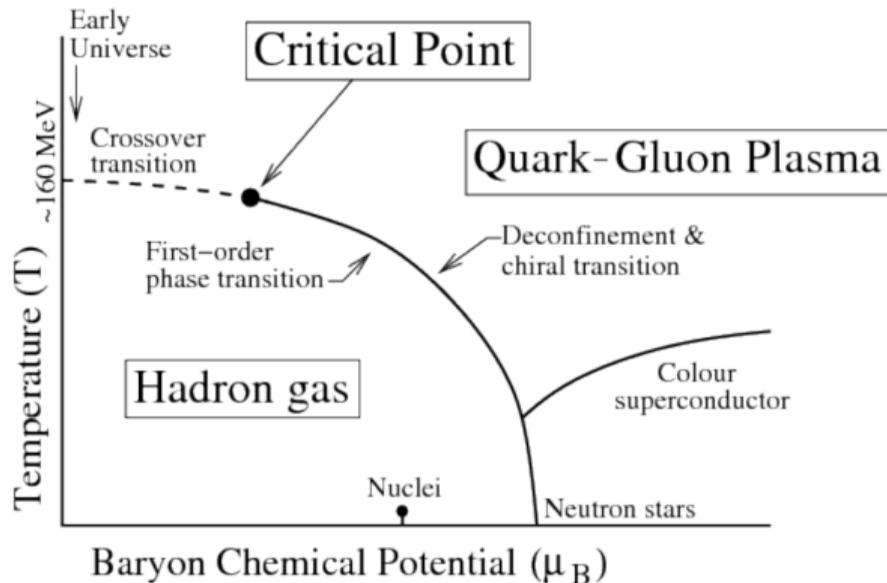
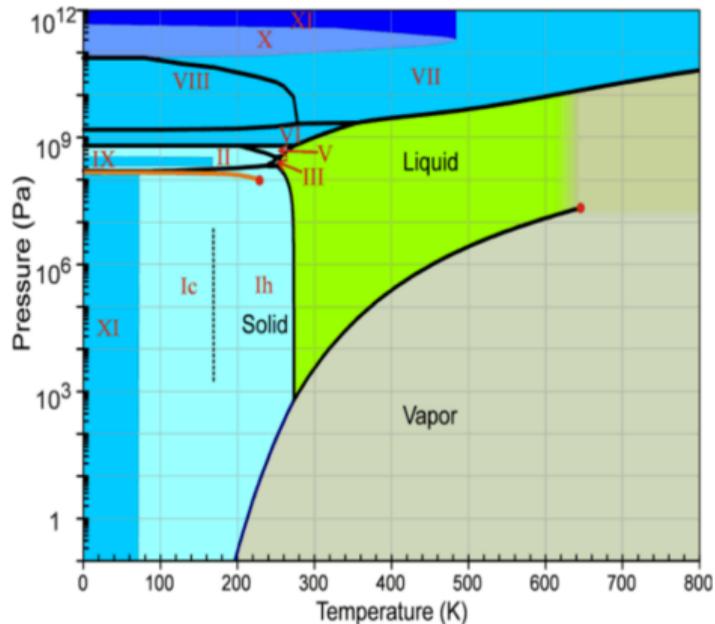
Sign problem:

Equation of state for low to moderate  $\mu_B/T$ .

Borsányi, Fodor, Guenther et al., PRL **126** (2021)



# Why a critical point?





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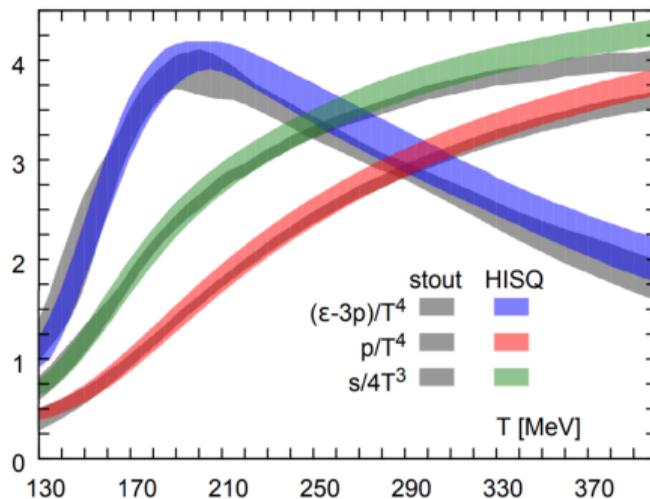
# Model Requirements

We need a simplified theoretical framework that describes QCD in the desired energy range.

Interpret data  $\iff$  make predictions

## Requirements:

- QCD symmetries, degrees of freedom, thermodynamics, and/or interactions.
- Agreement with Lattice EoS at  $\mu_B = 0$
- Agreement with lattice susceptibilities at  $\mu = 0$



stout: S. Borsanyi et al., 1309.5258, PBL (2014)  
HISQ: A. Bazavov et al., 1407.6387, PRD (2014)

# Critical point predictions as of some years ago

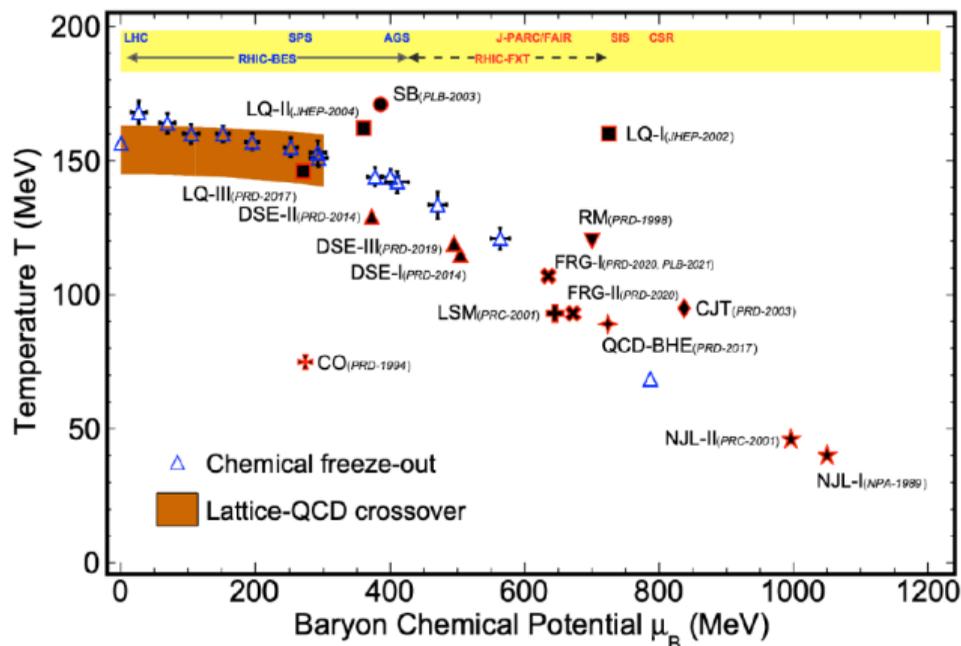


Figure adapted from A. Pandav, D. Mallick, B. Mohanty, Prog. Part. Nucl. Phys. 125 (2022)

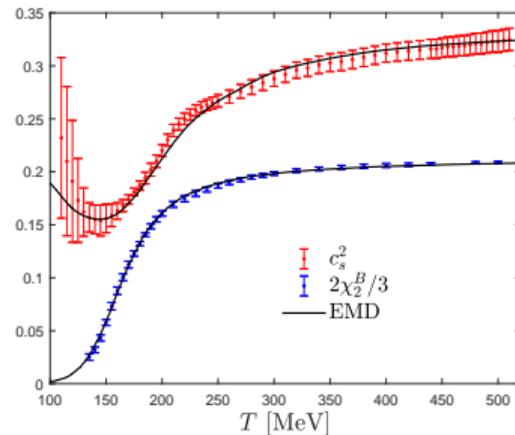
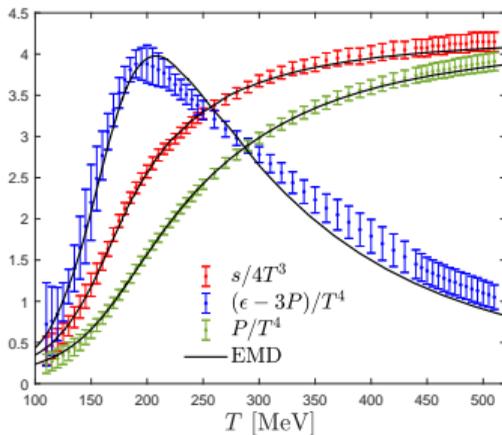
- Including the scenario of no critical point at all.

de Forcrand, Philipsen, JHEP 01, 077 (2007); VV, Steinheimer, Philipsen, Stoecker, PRD 97, 114030 (2018)

# Holography (Black Hole engineering - EMD model - gauge/gravity duality)

O DeWolfe et al. Phys.Rev.D 83, (2011). R Rougemont et al. JHEP(2016)102. R. Critelli et al., Phys.Rev.D96(2017).

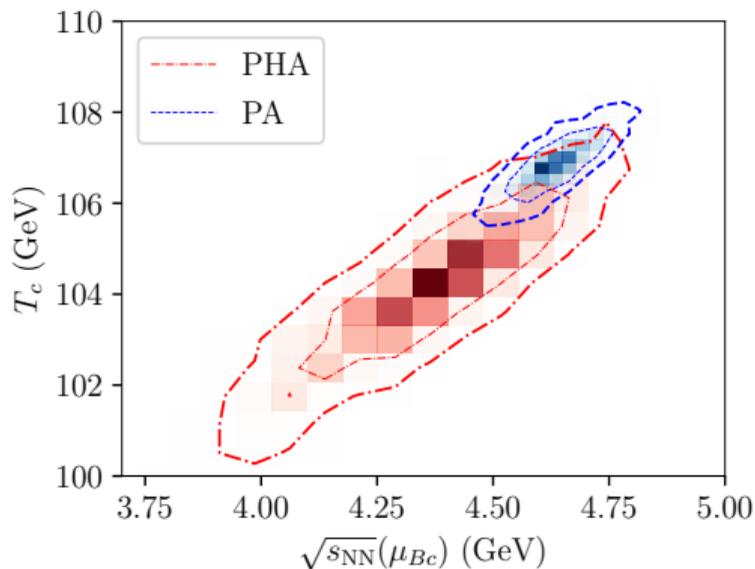
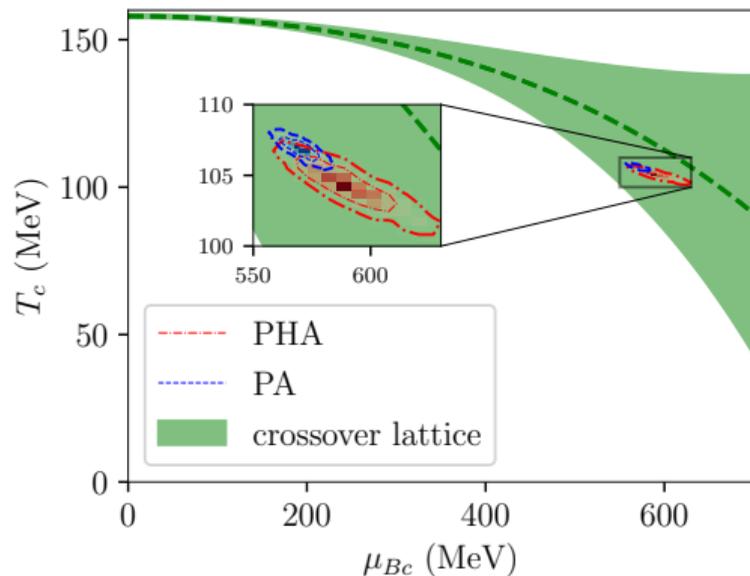
$$S = \frac{1}{2\kappa_5^2} \int_{M_5} d^5x \sqrt{-g} \left[ R - \frac{(\partial_\mu \phi)^2}{2} - \underbrace{V(\phi)}_{\text{nonconformal}} - \underbrace{\frac{f(\phi)F_{\mu\nu}^2}{4}}_{\mu_B \neq 0} \right]$$



# Holographic Bayesian Analysis: posterior critical points

$$(T_c, \mu_{Bc})_{PHA} = (104 \pm 3, 589^{+36}_{-26}) \text{ MeV},$$

$$(T_c, \mu_{Bc})_{PA} = (107 \pm 1, 571 \pm 11) \text{ MeV}.$$



- Both Ansätze overlap at  $1\sigma$ . **Robust results!**

M. Hippert, J.G., T.A. Manning, J. Noronha, J. Noronha-Hostler, I. Portillo, C. Ratti, R. Rougemont, M. Trujillo, [arXiv:2309.00579](https://arxiv.org/abs/2309.00579).

Based on the truncated expansion of the QCD functional, and requires control of errors to compute thermodynamics.

## Dyson–Schwinger Equations (DSE)

- integral equations derived from the QCD action that relate the propagators (Green's functions) of quarks and gluons to each other.
- are solved approximately using truncation schemes and modeling of the interaction vertices.

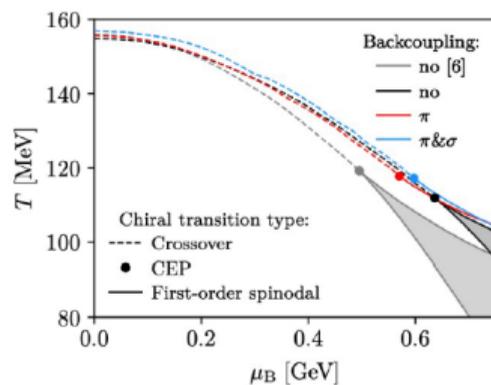
## Functional renormalization group (FRG)

- describing how the effective QCD action changes as one varies the energy scale.
- allows for a continuous evolution from microscopic physics to macroscopic phenomena, capturing quantum, thermal, and density fluctuations along the way.

# Effective QCD theories prediction

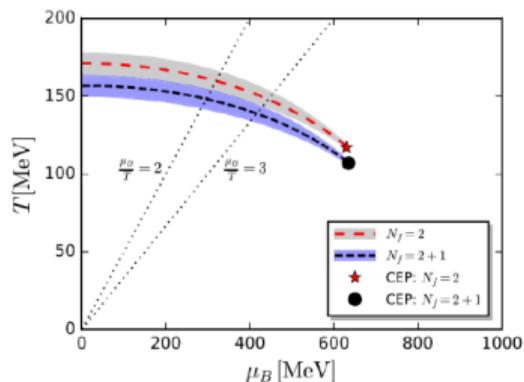
- Different effective approaches, all in excellent agreement with lattice QCD at  $\mu_B = 0$  (and  $\mu_B/T \sim 3.5$ ), predict the location of the critical point in a similar region.
- If true, reachable in heavy ion collisions at  $\sqrt{s_{NN}} \sim 3 - 5$  GeV.

## Dyson-Schwinger equations



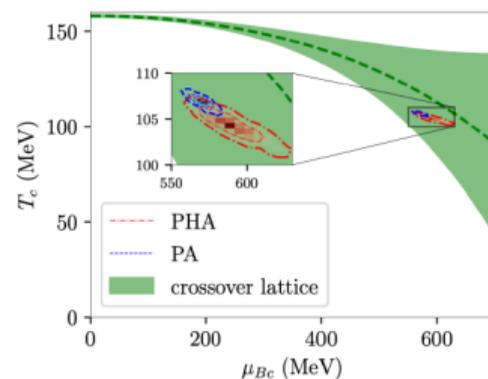
Gunkel, Fischer, PRD 104, 054202 (2021)

## Functional renormalization group



Fu, Pawłowski, Rennecke, PRD 101, 053032 (2020)

## Black-hole engineering



Hippert, J.G., et al., PRD 110, 094006 (2024)

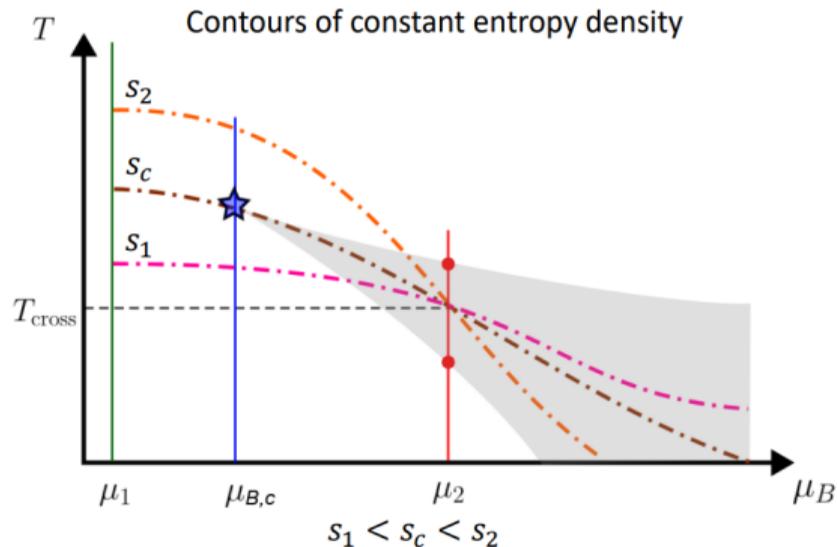
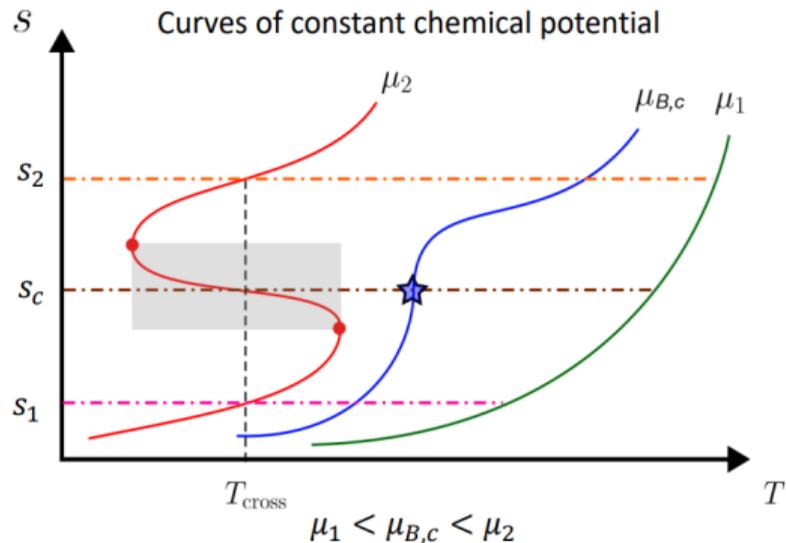
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# Entropy density contours

- If lines of constant  $s$  cross, it suggests a first order phase transition with a CP; otherwise, the EoS only exhibits a crossover.

$$T_s(\mu_B; T_0) = T_0 + \alpha_2(T_0) \frac{\mu_B^2}{2}; \quad \alpha_2(T_0) = -\frac{2T_0\chi_2^B(T_0) + T_0^2\chi_2^{B'}(T_0)}{s'(T_0)}; \quad \chi_2^B = \left[ \frac{\partial^2(p/T)^4}{\partial(\mu_B/T)^2} \right]_T$$



H. Shah et al. arXiv:2410.16026

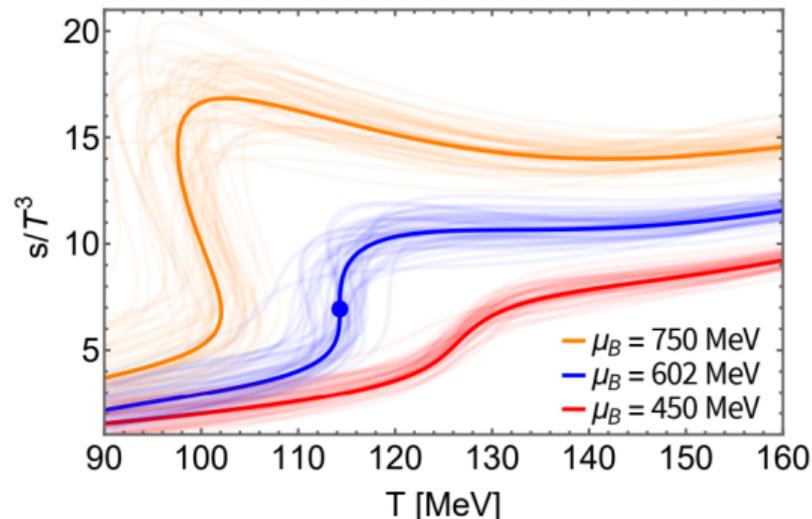
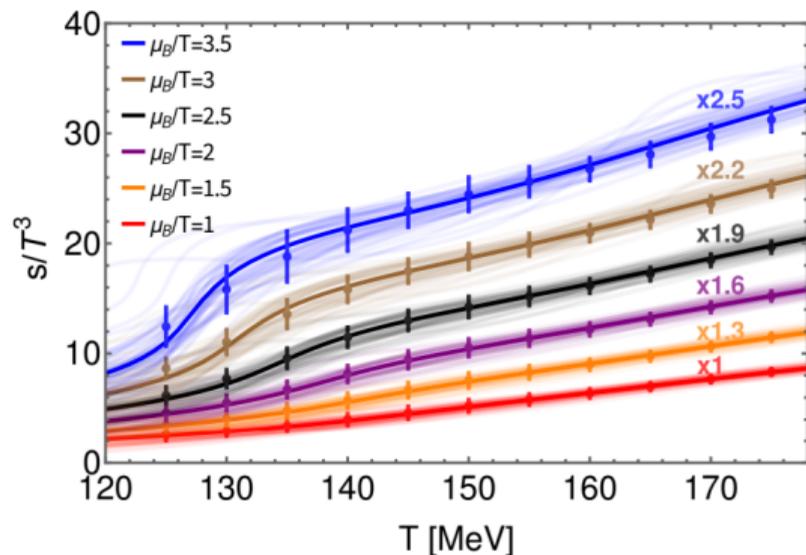
# Entropy density contours

- Excellent agreement with state-of-the-art lattice QCD data up to  $\mu_B/T = 3.5$

Borsanyi et al., PRL 126, 232001 (2021)

$$\mu_B^c = 602 \pm 62 \text{ MeV}$$

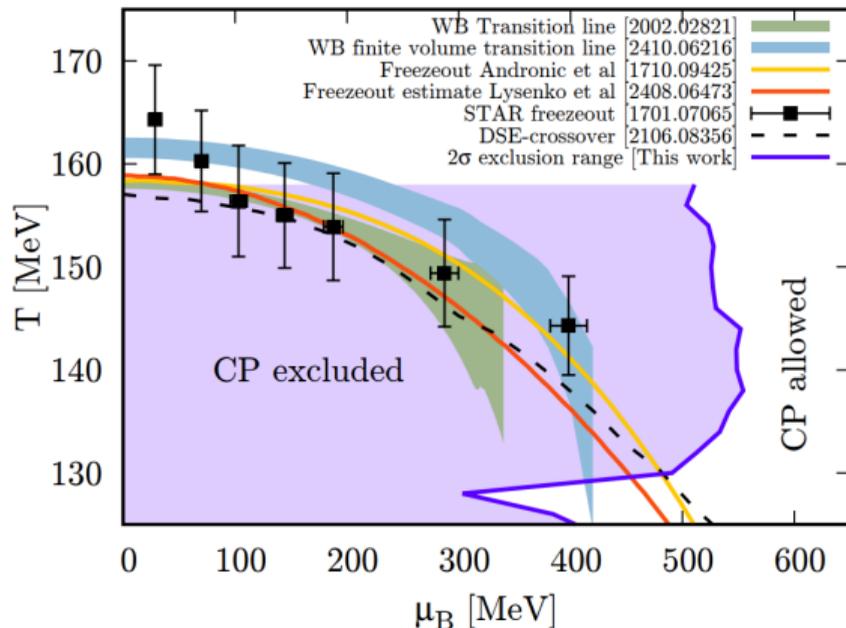
$$T^c = 114 \pm 7 \text{ MeV}$$



H. Shah et al. arXiv:2410.16026

## Constraints with improved lattice data

- New continuum extrapolated equation of state at zero density with improved precision & new data at imaginary chemical potential
- CP excluded at  $\mu_B < \sim 450$  MeV



Borsanyi et al. arXiv:2502.1026

# Critical point predictions as of some years ago

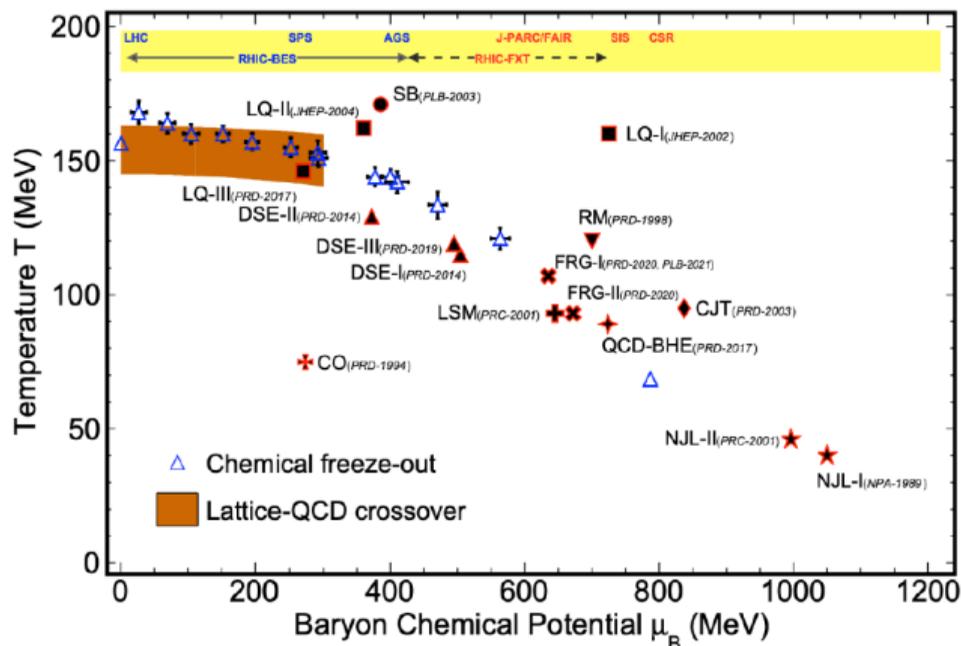


Figure adapted from A. Pandav, D. Mallick, B. Mohanty, Prog. Part. Nucl. Phys. 125 (2022)

- Including the scenario of no critical point at all.

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# Critical point with new lattice constraints: no CP at $\mu_B < 450$ MeV

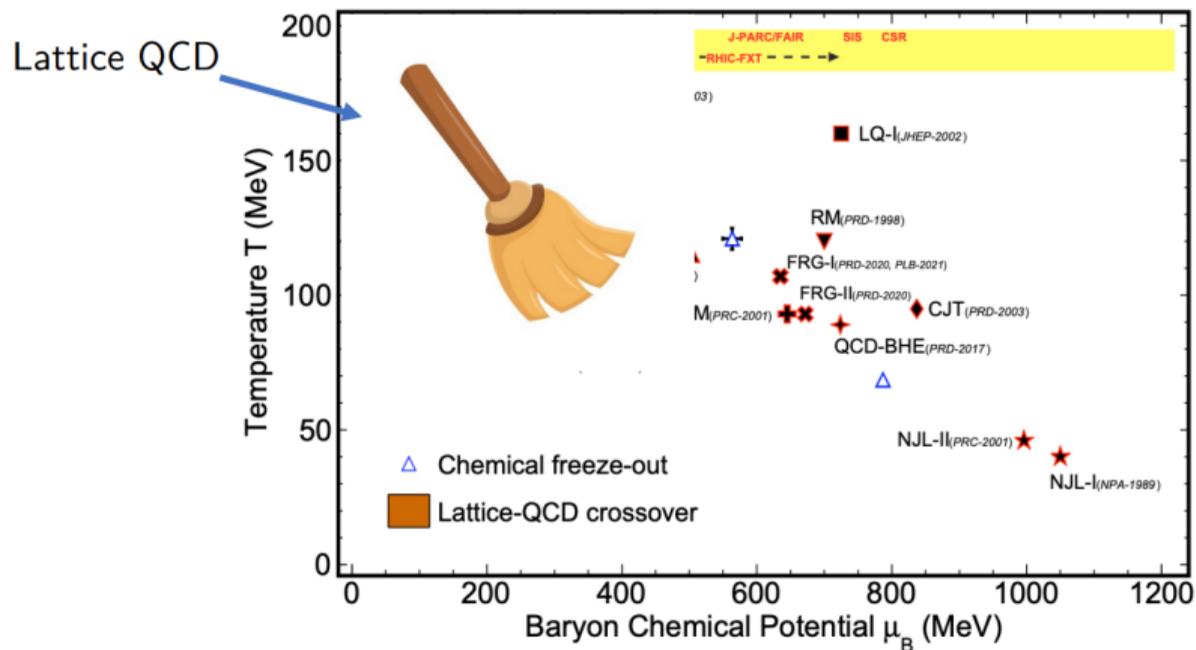


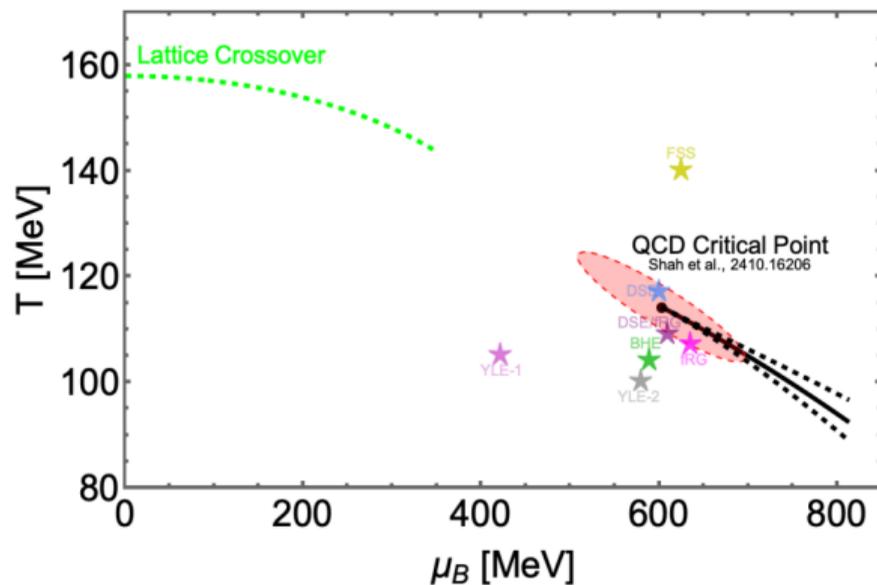
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# Current scenario

- Predictions converge to the same region...



Adapted from H. Shah et al. arXiv:2410.16026

Critical point estimate at  $O(\mu_B^2)$ :

$$T_c = 114 \pm 7 \text{ MeV}, \quad \mu_B = 602 \pm 62 \text{ MeV}$$

**Estimates from recent literature:**

YLE-1: D.A. Clarke et al. (Bielefeld-Parma), arXiv:2405.10196

YLE-2: G. Basar, PRC 110, 015203 (2024)

BHE: M. Hippert et al., arXiv:2309.00579

frg: W.-J. Fu et al., PRD 101, 054032 (2020)

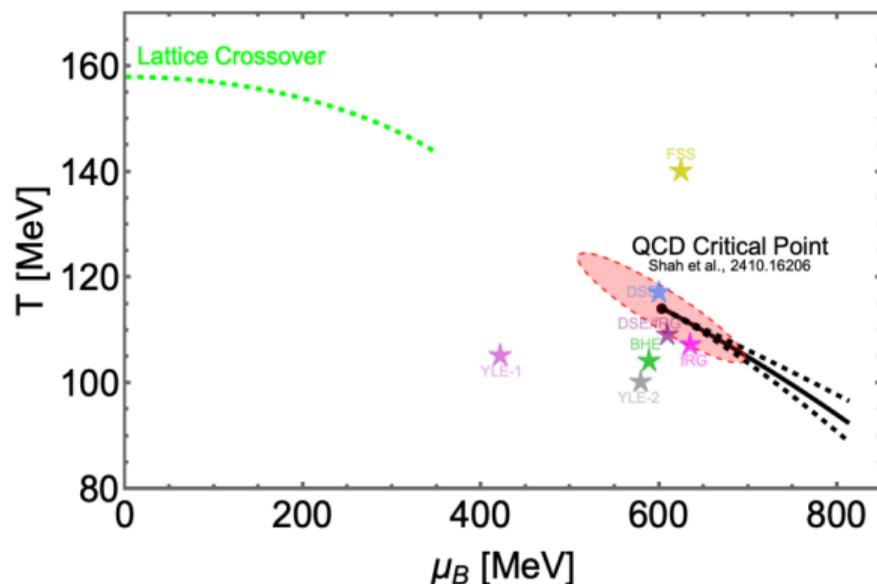
DSE/frg: Gao, Pawłowski., PLB 820, 136584 (2021)

DSE: P.J. Gunkel et al., PRD 104, 052022 (2021)

FSS: A. Sorensen et al., arXiv:2405.10278

# Current scenario

- Predictions converge to the same region...  
because lattice QCD has not ruled out that region yet?



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# Fluctuations

Cumulants measure the chemical potential derivatives of the QCD equation of state

Cumulants as moments of the particle number distribution

variance:  $\kappa_2 = \langle (\Delta N)^2 \rangle = \sigma$

skewness:  $\kappa_3 = \langle (\Delta N)^3 \rangle$

kurtosis:  $\kappa_4 = \langle (\Delta N)^4 \rangle - 3 \langle (\Delta N)^2 \rangle^2$

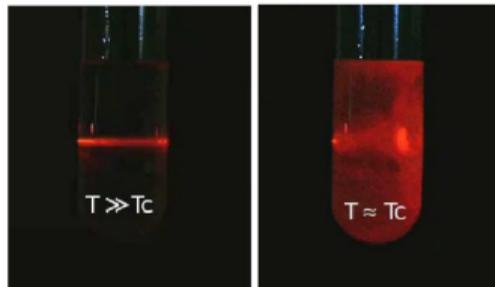
$$\kappa_2 \sim \xi^2, \quad \kappa_3 \sim \xi^{4.5}, \quad \kappa_4 \sim \xi^7$$

$$\xi \rightarrow \infty$$

Cumulants as chemical potential derivatives of the EoS

$$\ln Z(T, V, \mu) = \ln \left[ \sum_N e^{\mu N/T} Z^{ce}(T, V, N) \right]$$
$$\kappa_n \propto \frac{\partial^n (\ln Z)}{\partial \mu^n}$$

**Critical opalescence**



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Cumulants as moments of the particle number distribution

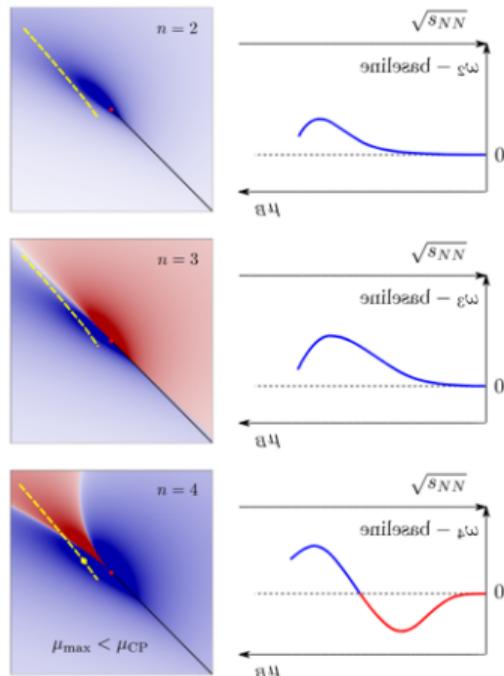
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$$\kappa_2 \sim \xi^2, \quad \kappa_3 \sim \xi^{4.5}, \quad \kappa_4 \sim \xi^7$$

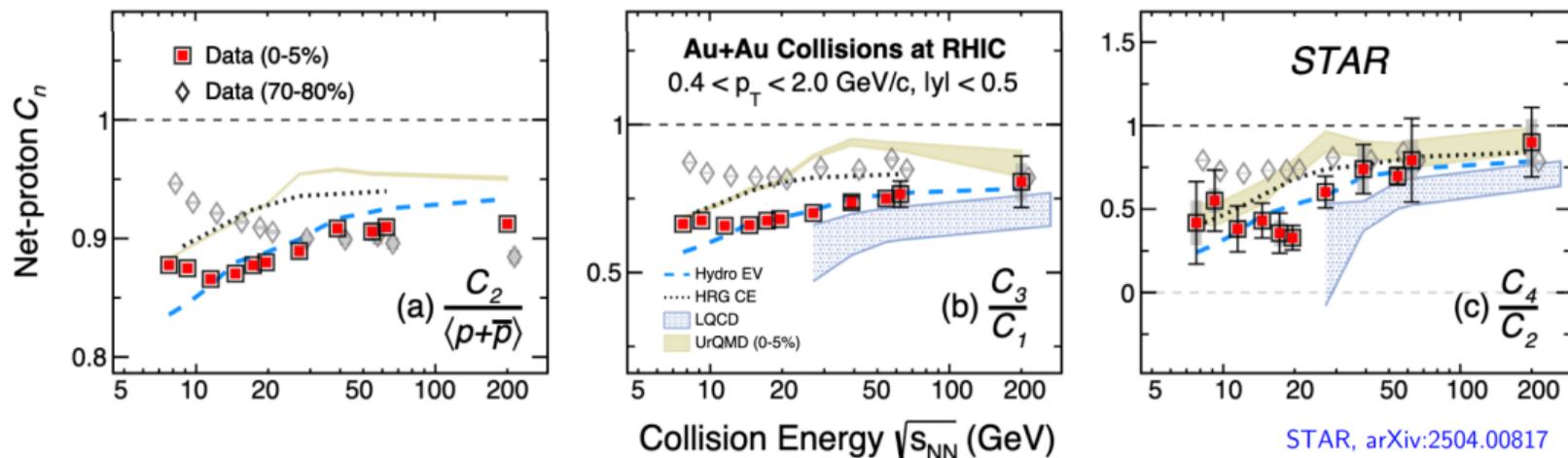
$$\xi \rightarrow \infty$$



M. Stephanov. SQM 2024

- Overall agreement with the baseline for  $\sqrt{s_{NN}} \sim 10 - 20$  GeV

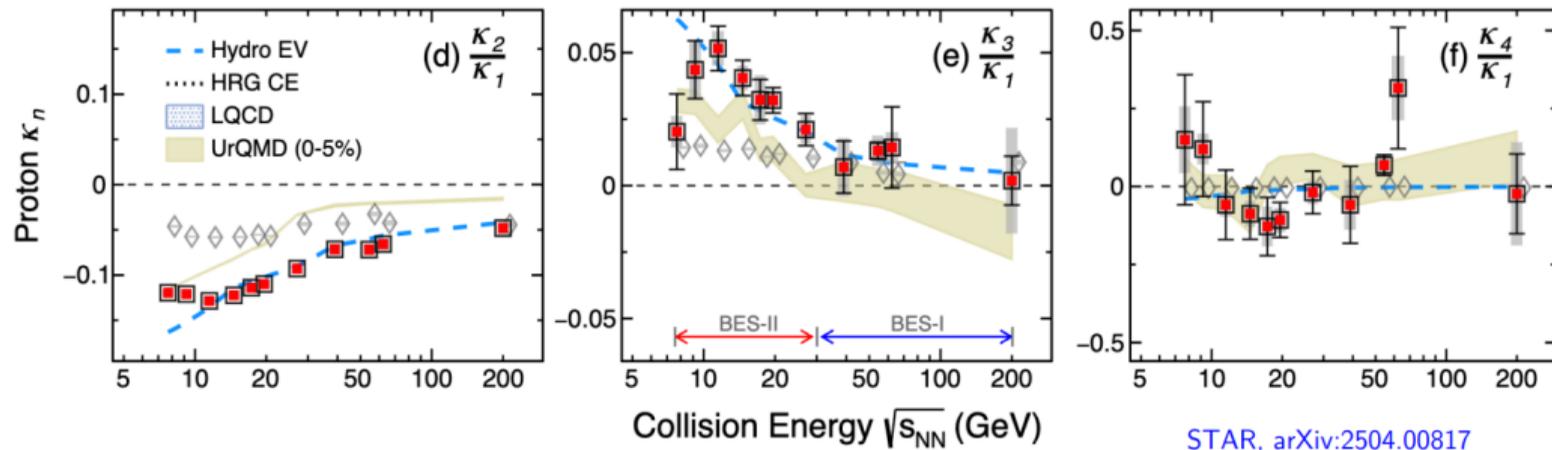
### Net-proton cumulant ratios



# Factorial cumulants!

- Exhibit more structure

## Proton factorial cumulant ratios

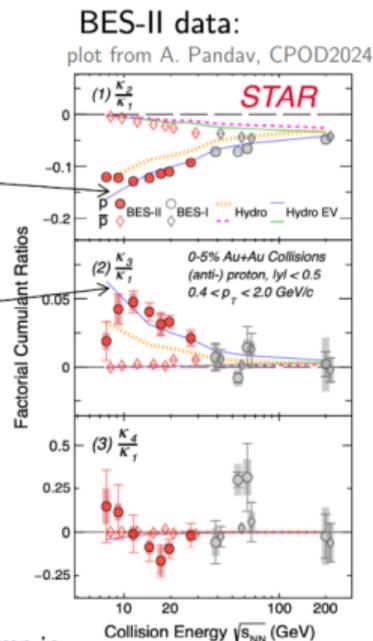
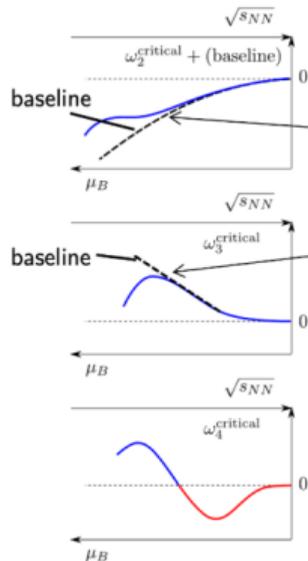
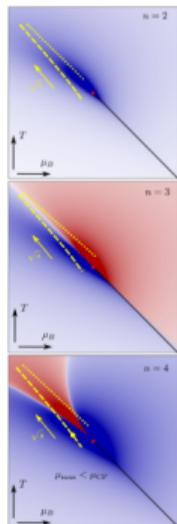


# factorial cumulants

V. Vovchenko, arXiv:2504.01368, and adapted from Stephanov, arXiv:2410.02861

$$\omega_n = \hat{C}_n / \hat{C}_1$$

(universal EOS) critical  $\chi_n$ :



Factorial cumulants: Irreducible  $n$ -particle correlations that remove Poisson contribution and probe genuine correlations

Ordinary cumulants: mix correlations of different order

$$\hat{C}_n \sim \langle N(N-1)(N-2)\dots \rangle_c$$

$$\hat{C}_1 = C_1$$

$$\hat{C}_2 = C_2 - C_1$$

$$\hat{C}_3 = C_3 - 3C_2 + 2C_1$$

$$\hat{C}_4 = C_4 - 6C_3 + 11C_2 - 6C_1$$

Expected signatures: **bump** in  $\omega_2$  and  $\omega_3$ , **dip** then **bump** in  $\omega_4$  for CP at  $\mu_B > 420$  MeV

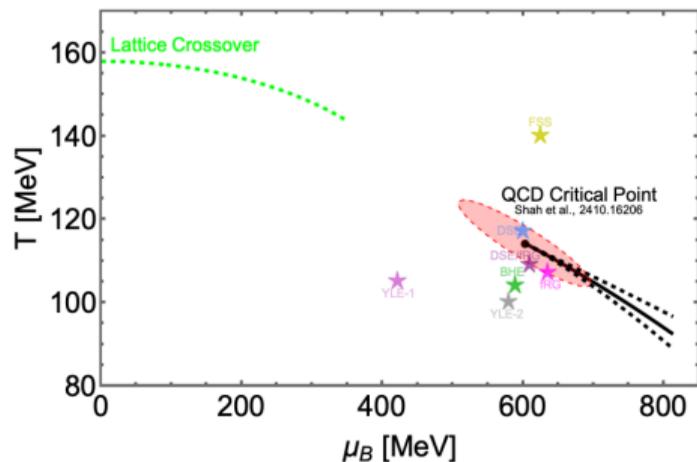
non monotonic  $\kappa_2/\kappa_1$ ,  $\kappa_3/\kappa_1$  and maybe  $\kappa_4/\kappa_1$

Bzdak, Koch, Strodthoff, PRC 95, 054906 (2017)

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# Summary



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DSE: P.J. Gunkel et al., PRD 104, 052022 (2021)

FSS: A. Sorensen et al., arXiv:2405.10278

- No indication of critical behavior from lattice QCD for  $\mu_B < 450$  MeV.
- Several effective theories predict the location of the critical point to  $T \sim 90 - 120$  MeV and  $\mu_B \sim 500 - 650$  MeV.
- No critical behavior describe proton cumulants at  $\sqrt{s_{NN}} \geq 20$  GeV.
- This trend changes around  $\sqrt{s_{NN}} \sim 10$  GeV; in particular, for factorial cumulants and the presence of the CP could be a reasonable explanation.

# Appendix

# Ordinary vs. Factorial Cumulants in Heavy-Ion Collisions

## Ordinary Cumulants

$$C_n \sim \langle \delta N^n \rangle$$

- Built from moments of distribution (mean, variance, etc).
- Sensitive to critical fluctuations.
- Connected to thermodynamic susceptibilities.
- Higher orders diverge near critical point.

## Factorial Cumulants

$$\hat{C}_n \sim \langle N(N-1)(N-2)\dots \rangle_c$$

- Built from factorial moments.
- Vanish for Poisson baseline  $\Rightarrow$  better contrast.
- More robust under detection inefficiencies.
- Useful in experimental fluctuation analyses.

# Which Cumulants Are Better for the QCD Critical Point?

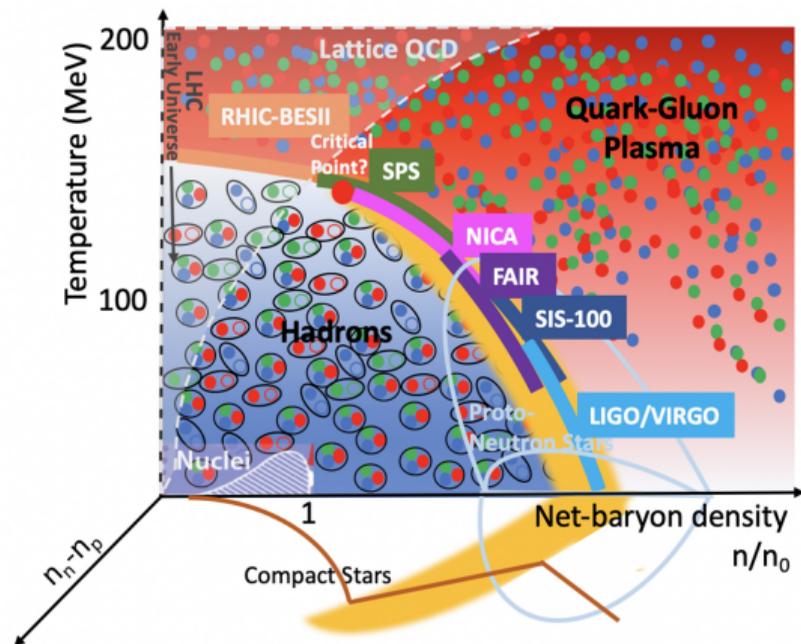
Both are useful, but serve different purposes:

- **Ordinary cumulants:**  $C_n \sim \langle \delta N^n \rangle$ 
  - Theoretically well-defined.
  - Connected to QCD susceptibilities.
- **Factorial cumulants:**  $\hat{C}_n \sim \langle N(N-1)(N-2)\dots \rangle_c$ 
  - Cleaner signals under real-world detector conditions.
  - Efficient background suppression (e.g., Poisson noise).
- *Best approach: use both and compare.*

See: Bzdak et al. Phys. Rept. 853 (2020), Kitazawa Asakawa PRC 85 (2012), STAR Collaboration (2022)

# What happens at finite/large densities?

- We need to merge the lattice QCD EoS with other effective theories.
- Study the regime of validity of each effective model.
- Constrained internal parameters to adhere know experimental and theoretical limits.
- Test models to validate/exclude them.
- EoS to guide/interpret experimental data.

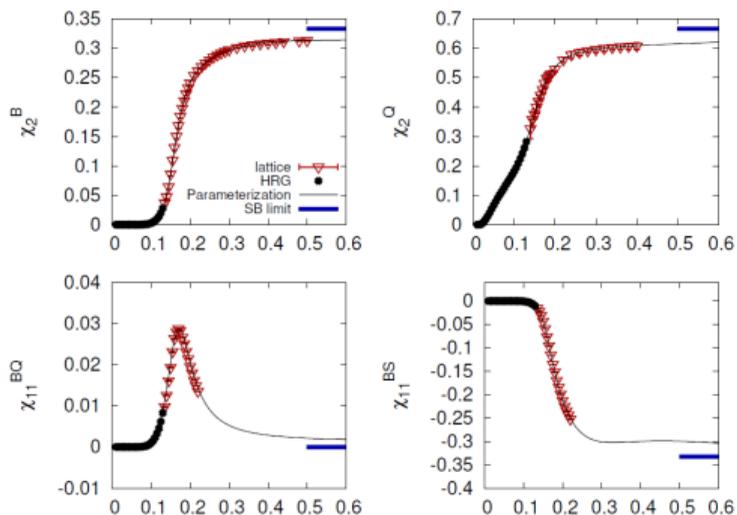


## BQS EoS

$10 < T < 600$  MeV;  $\mu_B < 450$  MeV  
J. Noronha-Hostler, et al., PRC (2019)

$$\frac{p(T, \mu_B, \mu_Q, \mu_S)}{T^4} = \sum_{i,j,k} \frac{1}{i!j!k!} \chi_{ijk}^{BQS} \left(\frac{\mu_B}{T}\right)^i \left(\frac{\mu_Q}{T}\right)^j \left(\frac{\mu_S}{T}\right)^k$$

$$\chi_{ijk}^{BQS} = \left. \frac{\partial^{i+j+k}(p/T^4)}{\partial \left(\frac{\mu_B}{T}\right)^i \partial \left(\frac{\mu_Q}{T}\right)^j \partial \left(\frac{\mu_S}{T}\right)^k} \right|_{\mu_B, \mu_Q, \mu_S=0}$$



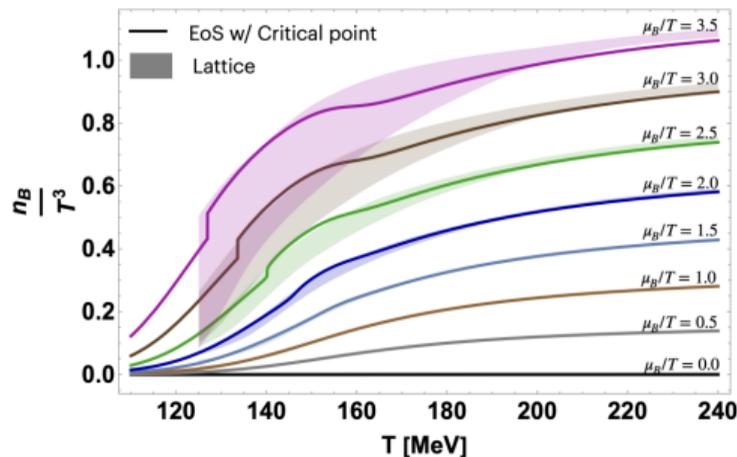
## 2D Ising T.Ex.S

$10 < T < 800$  MeV;  $\mu_B < 700$  MeV  
M. Kahangirwe, et al., PRD 109 (2024)

$$T \frac{\chi_1^B(T, \mu_B)}{\mu_B} = \chi_2^B(T', 0)$$

$$T'(T, \mu_B) = T \left[ 1 + \kappa_2^{BB}(T) \left(\frac{\mu_B}{T}\right)^2 + \kappa_4^{BB}(T) \left(\frac{\mu_B}{T}\right)^4 + \dots \right]$$

- Includes a 3D Ising model critical behavior into a lattice alternative expansion EoS.

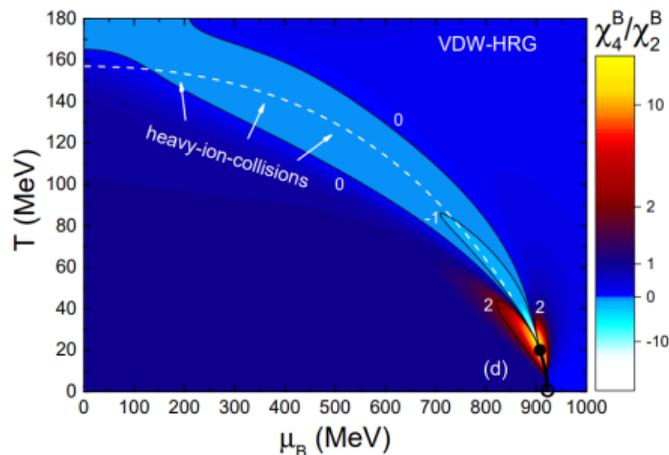


## HRG Model

$0 < T < 160$  MeV;  $\mu_B < 1000$  MeV

V. Vovchenko, CPC (2019)

- Provides a realistic hadronic EoS at low  $T$
- Interacting hadrons can be modeled by an ideal gas of resonances.
- For a realistic EoS at higher densities, Van der Waals interactions are added.
- Describes the liquid-gas phase transition.



## Holography (NumRelHolo)

$40 < T < 400$  MeV;  $\mu_B < 1200$  MeV

J. G., et al., PRD (2021), PRD (2022)

- Based on the gauge/gravity duality and constrained to reproduce lattice-QCD thermodynamics
- Large coverage of the EoS in the strongly-interacting regime.
- Predicts the location of the QCD CP.

