Neutrino Masses and Oscillations

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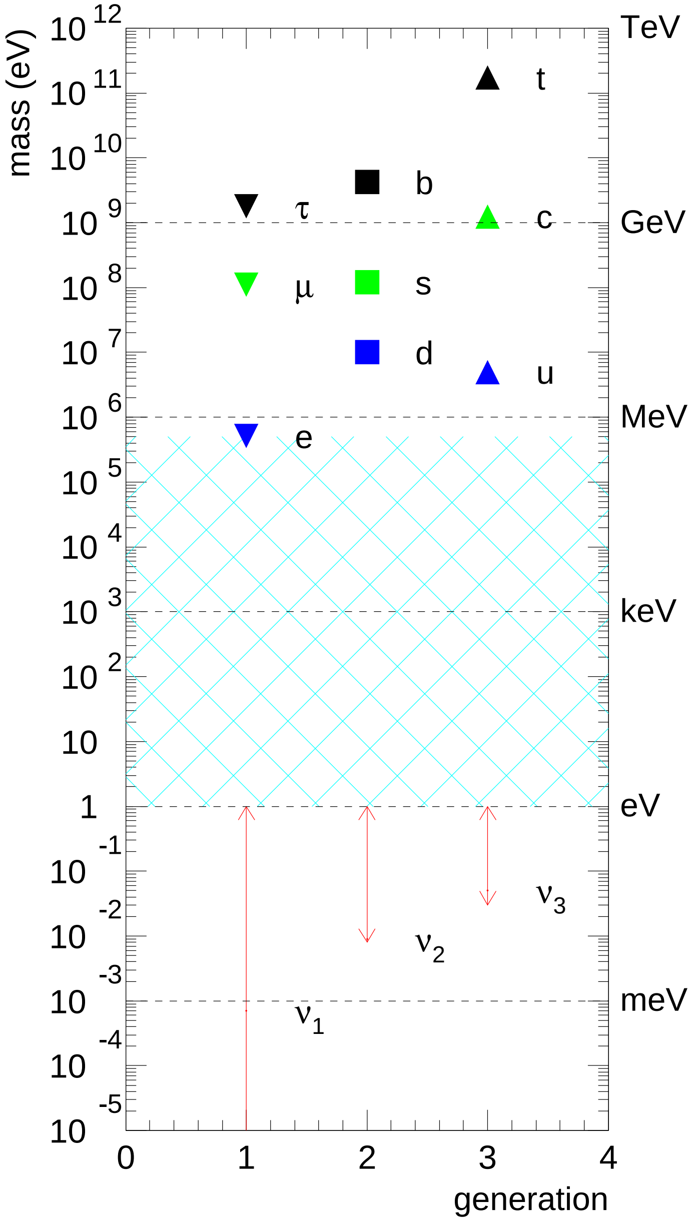
Outline

1. What We Have Learned About Neutrinos;
2. What We Know We Don't Know;
3. Neutrino Masses As Physics Beyond the Standard Model;
4. Ideas for Neutrino Masses, with Consequences;
5. Conclusions.



NEUTRINOS HAVE MASS

[albeit very tiny ones...]

We don't know why that is, but we have a "gut feeling" it means something important.

Are neutrinos fundamentally different?

Are neutrino masses generated by a distinct dynamical mechanism?

How Did We Find Out: Flavor Oscillations!

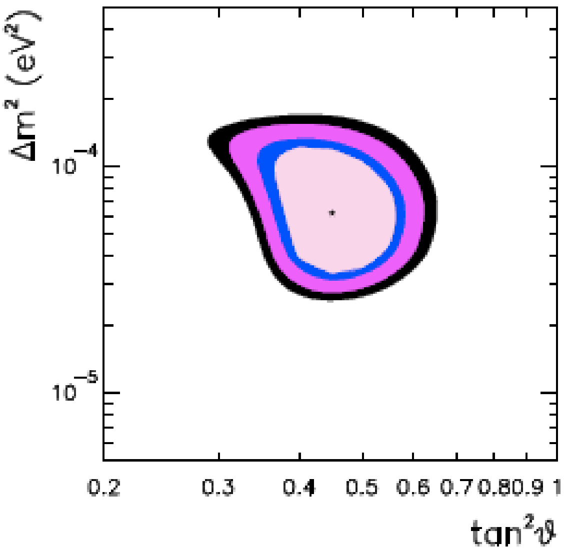
Neutrino oscillation experiments have revealed that neutrinos change flavor after propagating a finite distance. The rate of change depends on the neutrino energy E_ν and the baseline L.

- $\nu_\mu \to \nu_\tau$ and $\bar{\nu}_\mu \to \bar{\nu}_\tau$ — atmospheric experiments ["indisputable"];
- $\nu_e \to \nu_{\mu,\tau}$ — solar experiments ["indisputable"];
- $\bar{\nu}_e \to \bar{\nu}_{\rm other}$ — reactor neutrinos ["indisputable"];
- $\nu_\mu \to \nu_{\rm other}$ from accelerator experiments ["really strong"].

The simplest and **only satisfactory** explanation of **all** this data is that neutrinos have distinct masses, and mix. → $\boxed{P_{\rm osc} = \sin^2 2\theta \sin^2\left(\frac{\Delta m^2 L}{4E_\nu}\right)}$

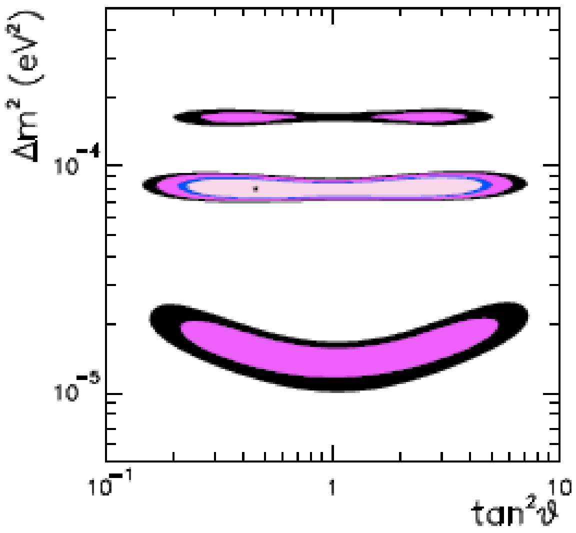
Solar

$\nu_e \to \nu_{\rm active}$

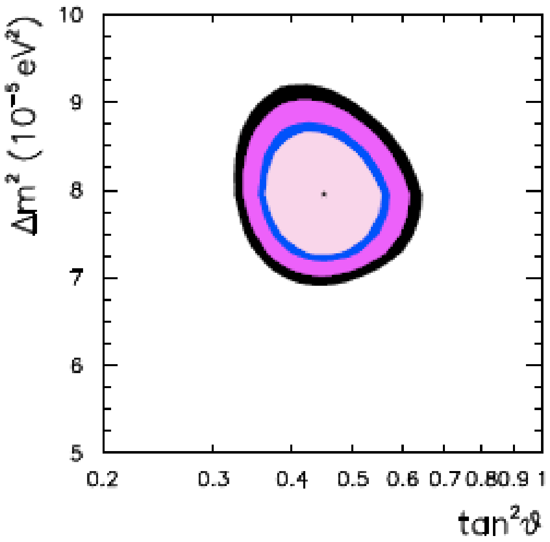


+ KamLAND

$\bar\nu_e \not\to \bar\nu_e$



ν_e oscillation parameters compatible with $\overline{\nu}_e$: *Sensible* to assume CPT: $P_{ee} = P_{\bar e\bar e}$



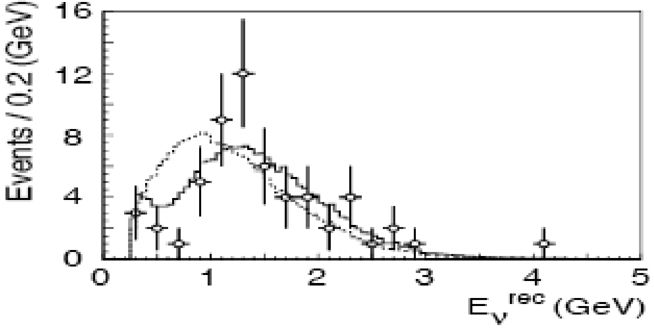
$$\Delta m^2_\odot = (8^{+0.4}_{-0.5}) \times 10^{-5}\ \mathrm{eV}^2 \quad (1\sigma)$$

$$\tan^2\theta_\odot = 0.45^{+0.05}_{-0.05}$$

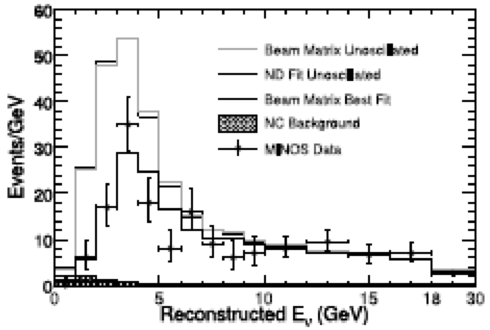
[Gonzalez-Garcia, PASI 2006]

K2K	ν_μ at KEK	SK	L=250 km
MINOS	ν_μ at Fermilab	Soundan	L=735 km
Opera/Icarus	ν_μ at CERN	Gran Sasso	L=740 km

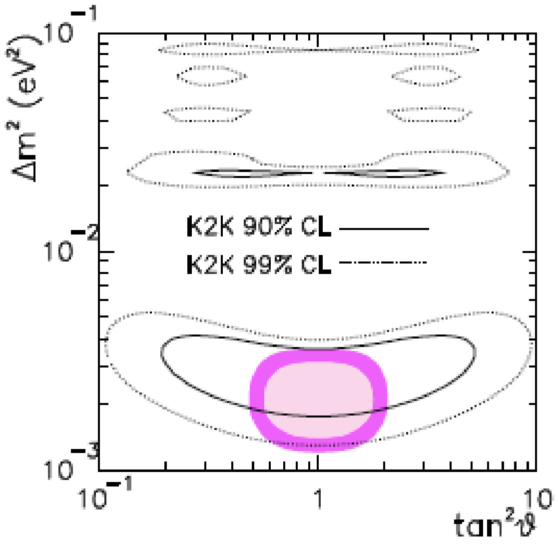
K2K 2004: spectral distortion



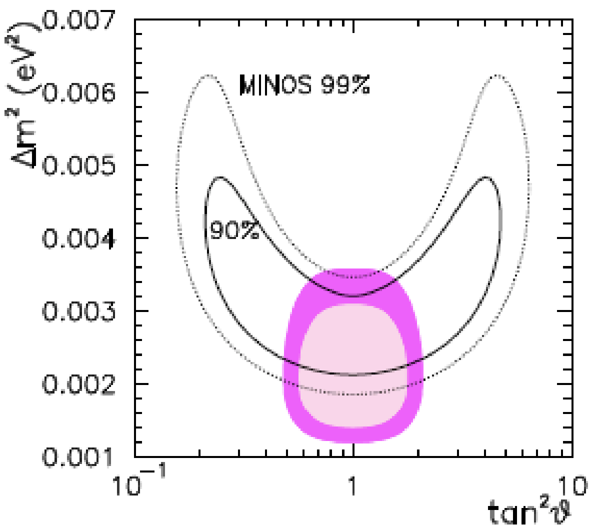
MINOS 2006: spectral distortion



Confirmation of ATM oscillations



Confirmation of ATM oscillations



[Gonzalez-Garcia, PASI 2006]

Phenomenological Understanding of Neutrino Masses & Mixing

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{e\tau2} & U_{\tau3} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

Definition of neutrino mass eigenstates (who are ν_1, ν_2, ν_3?):

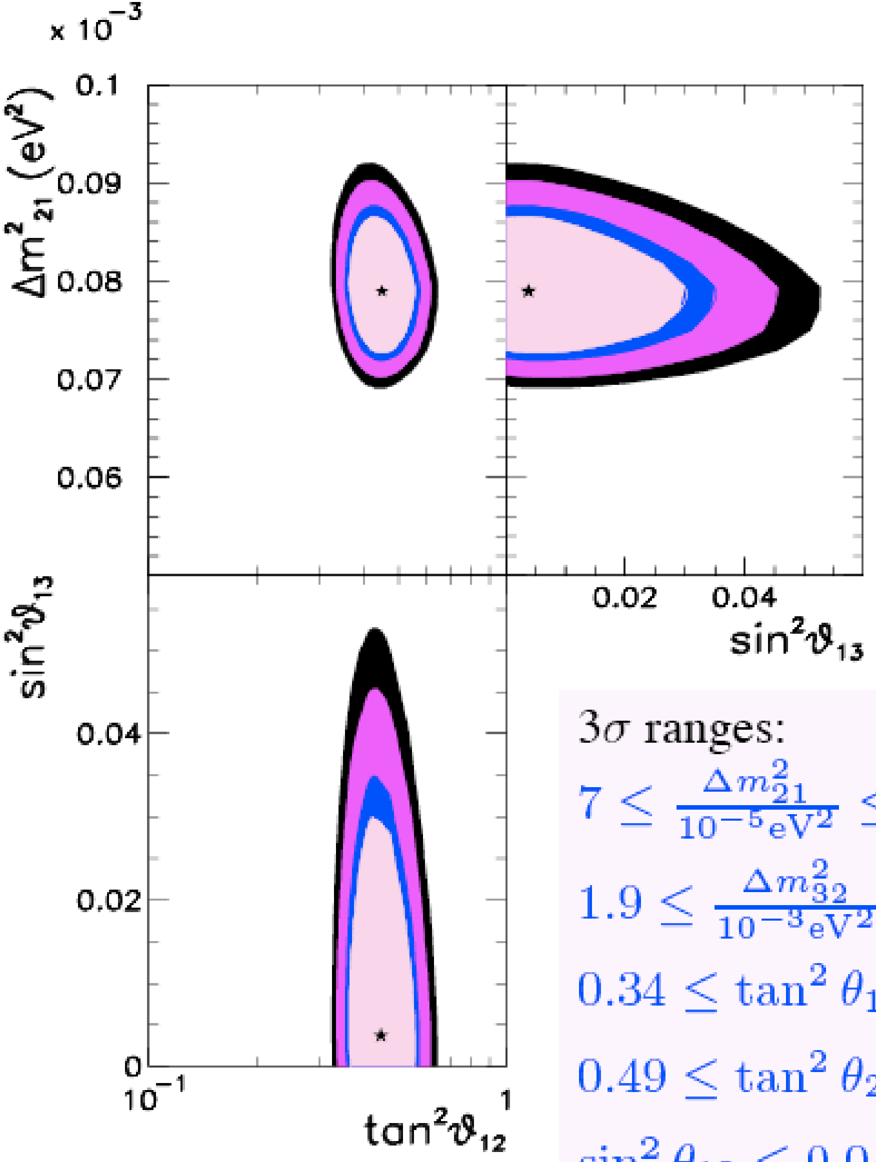
- $m_1^2 < m_2^2$
- $m_2^2 - m_1^2 \ll |m_3^2 - m_{1,2}^2|$

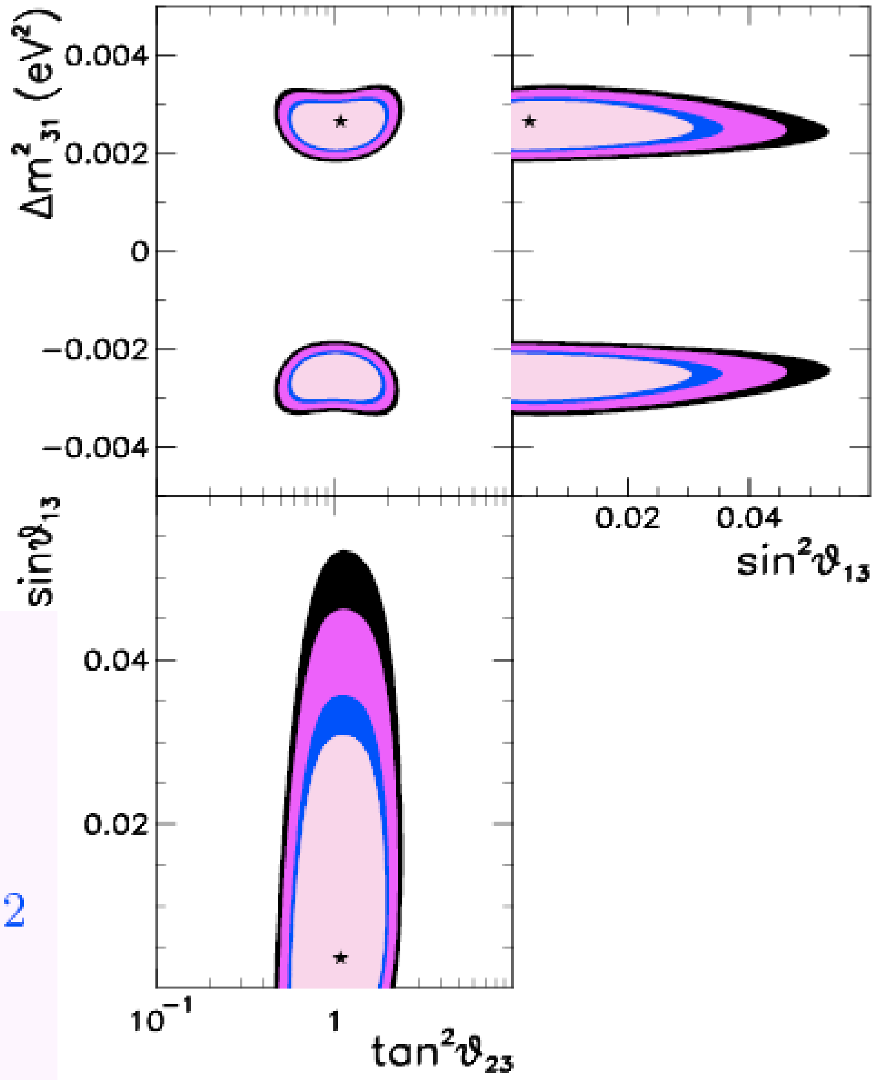
$\Delta m^2_{13} < 0$ – Inverted Mass Hierarchy

$\Delta m^2_{13} > 0$ – Normal Mass Hierarchy

$$\tan^2\theta_{12} \equiv \frac{|U_{e2}|^2}{|U_{e1}|^2}; \quad \tan^2\theta_{23} \equiv \frac{|U_{\mu3}|^2}{|U_{\tau3}|^2}; \quad U_{e3} \equiv \sin\theta_{13} e^{-i\delta}$$

Putting It All Together:





3σ ranges:

$$7 \le \frac{\Delta m^2_{21}}{10^{-5}\text{eV}^2} \le 9.1$$

$$1.9 \le \frac{\Delta m^2_{32}}{10^{-3}\text{eV}^2} \le 3.25$$

$$0.34 \le \tan^2\theta_{12} \le 0.62$$

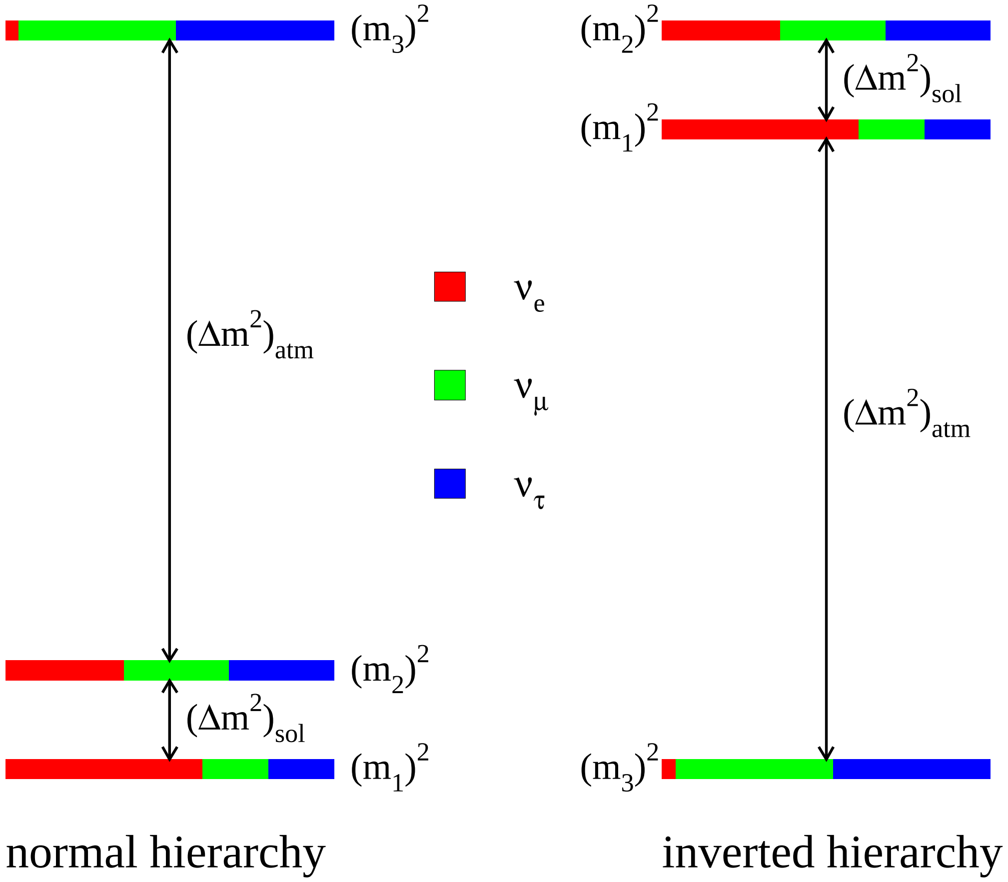
$$0.49 \le \tan^2\theta_{23} \le 2.2$$

$$\sin^2\theta_{13} \le 0.045$$

$$-\pi \le \delta \le \pi$$

[Gonzalez-Garcia, PASI 2006]

What We Know We Don't Know (1)



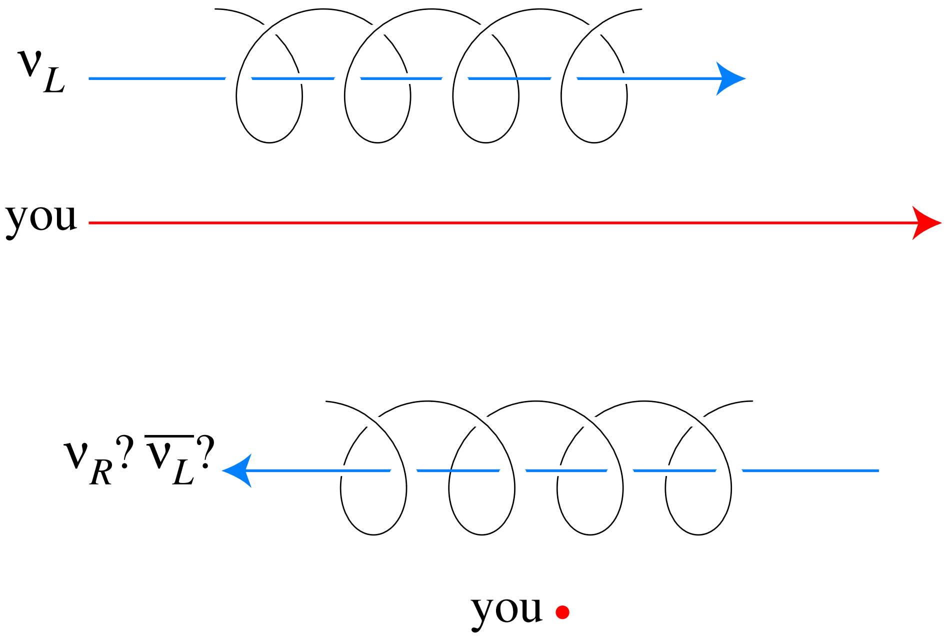
- What is the ν_e component of ν_3? ($\theta_{13} \neq 0$?)
- Is CP-invariance violated in neutrino oscillations? ($\delta \neq 0, \pi$?)
- Is ν_3 mostly ν_μ or ν_τ? ($\theta_{23} > \pi/4$, $\theta_{23} < \pi/4$, or $\theta_{23} = \pi/4$?)
- What is the neutrino mass hierarchy? ($\Delta m^2_{13} > 0$?)

$\Rightarrow$ All of the above can be addressed in neutrino oscillation experiments if we get lucky, that is if θ_{13} is large enough†

- What is the smallest neutrino mass?

(†talk by Bonnie Fleming)

What We Know We Don't Know (2): Are Neutrinos Majorana Fermions?



The neutrino is the only neutral elementary fermion. There is a left-handed one and a right-handed one.

as far as we can tell (experiments) . . .
the left-handed has lepton number $L = +1$, while the right-handed one has $L = -1$:

$(\nu_\ell)_L + X \to \ell^- + X'$, while
$(\nu_\ell)_R + X \to \ell^+ + X'$, so we call $(\nu_\ell)_R \equiv \bar{\nu}_\ell$

However:

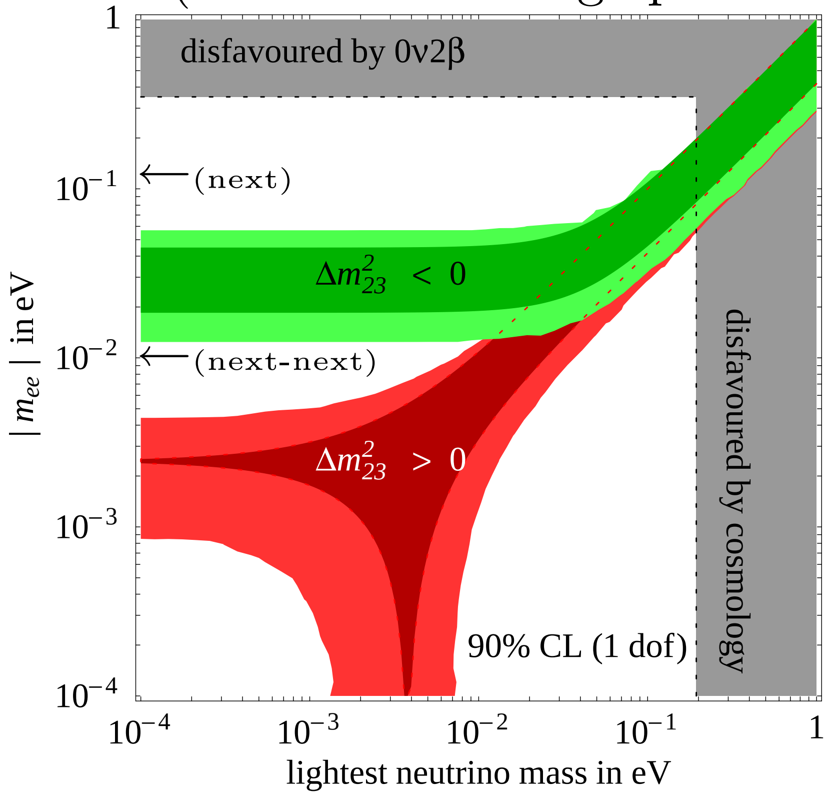
If the neutrino is its own antiparticle (Majorana fermion), then the lepton number conservation law must not be exact → look for L-violation.

Search for the Violation of Lepton Number (or $B-L$)

In order to make significant theoretical progress, we need to decide whether the neutrinos are Dirac of Majorana fermions

Best Bet: search for Neutrinoless Double-Beta decay: $\boxed{Z \to (Z+2)e^-e^-}$

(neutrino exchange picture: $2n \to 2p + 2e^- + \bar{\nu}_e + \bar{\nu}_e \to 2p + 2e^-$)



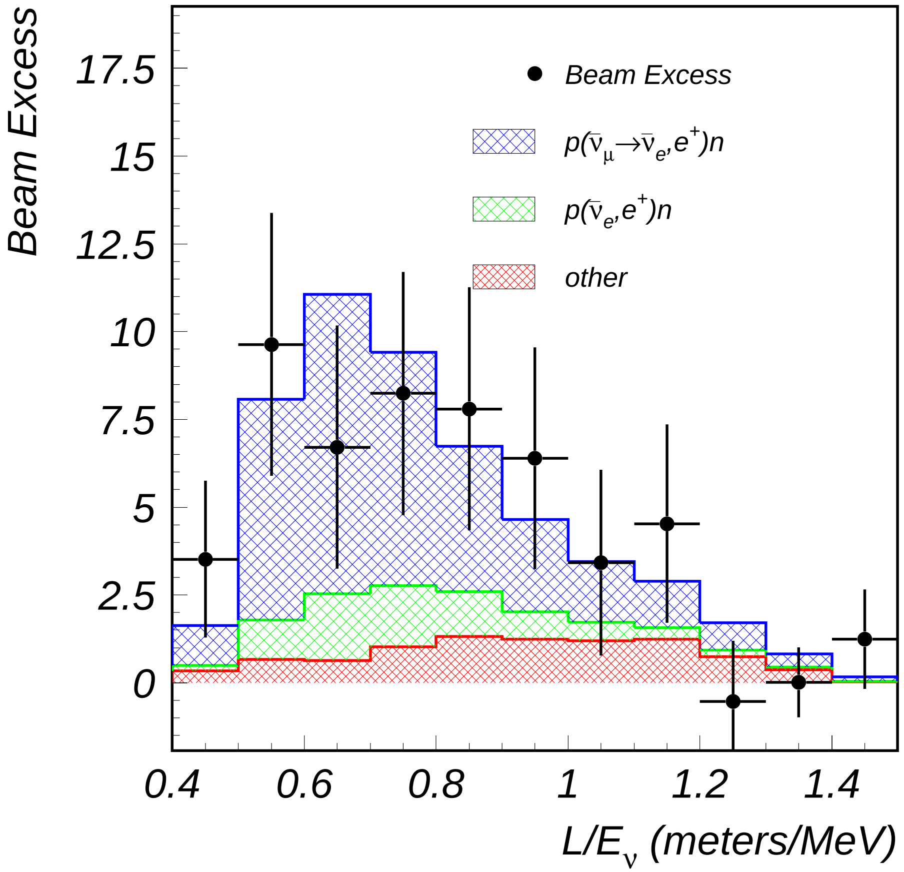
Helicity Suppressed Amplitude $\propto \frac{m_{ee}}{E}$

Observable: $m_{ee} \equiv \sum_i U_{ei}^2 m_i$

[Are there any other competitive probes?]

Detour: the LSND Anomaly

LSND: strong evidence for $\bar{\nu}_\mu \to \bar{\nu}_e$



If oscillations (??) $\Rightarrow \Delta m^2 \sim 1$ eV^2

× does not fit into 3 ν picture;

× $2+2$ scheme ruled out (solar, atm);

× $3+1$ scheme ruled out;

× 3 ν's CPTV ruled out (KamLAND, atm);

× $\mu \to e\nu_e\bar{\nu}_e$ ruled out (KARMEN, TWIST);

×? $3+1+1$ scheme;

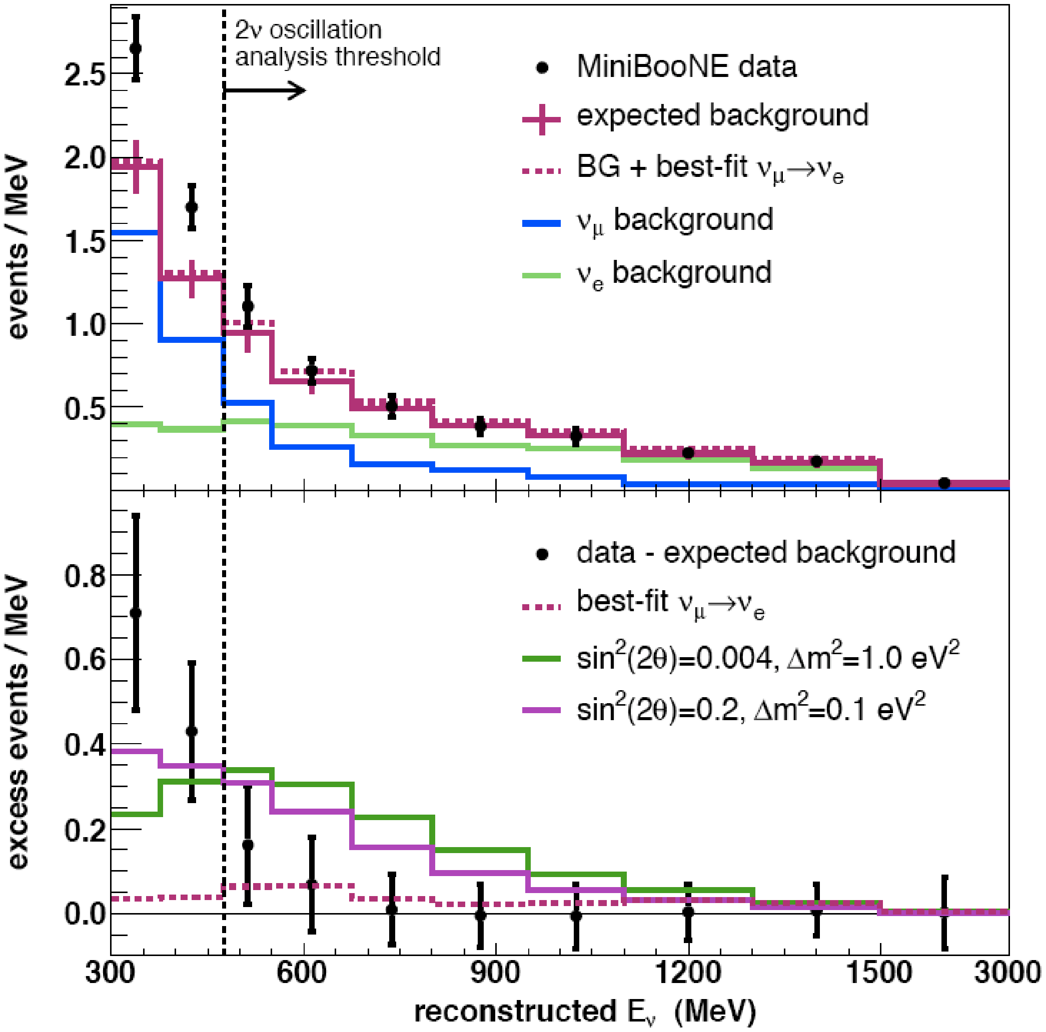
○ 4 ν's CPTV

×? "heavy" decaying sterile neutrinos;

○ 3 νs and Lorentz-invariance violation;

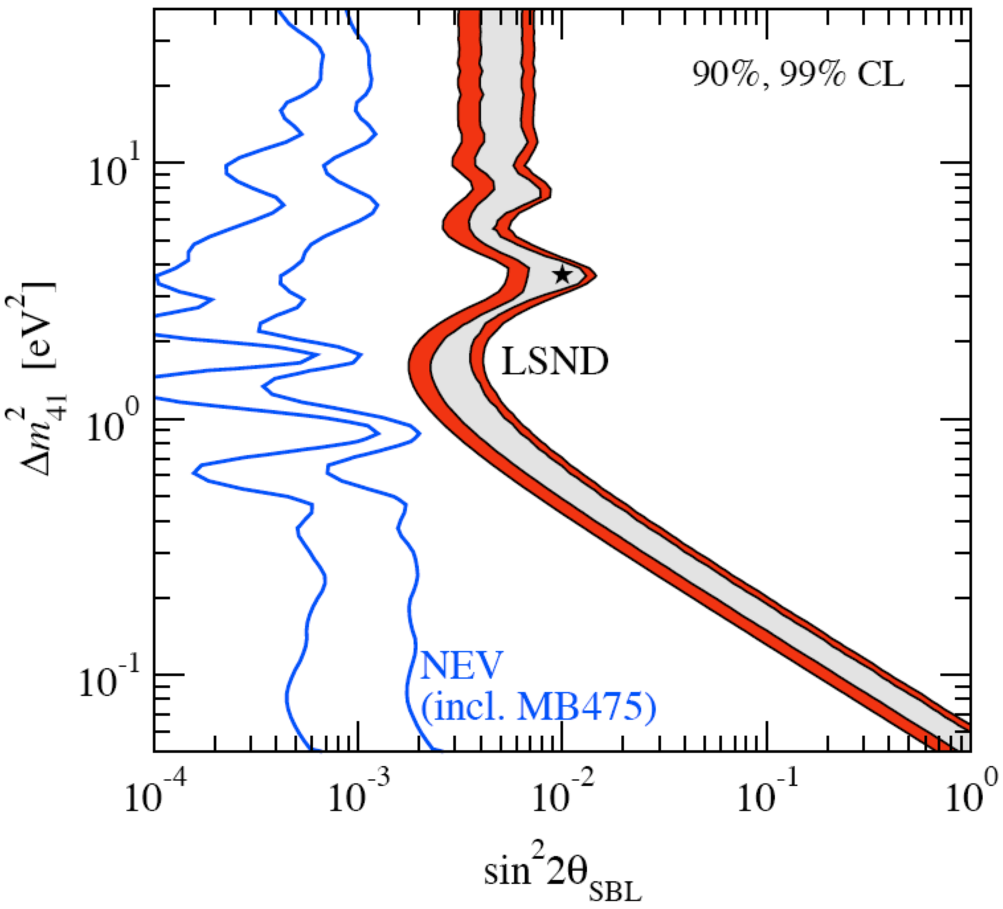
○ something completely different.

3+1 scheme “ruled out”



(talk by Chris Polly)

[Maltoni, Schwetz, arXiv:0705.0107 [hep-ph]]



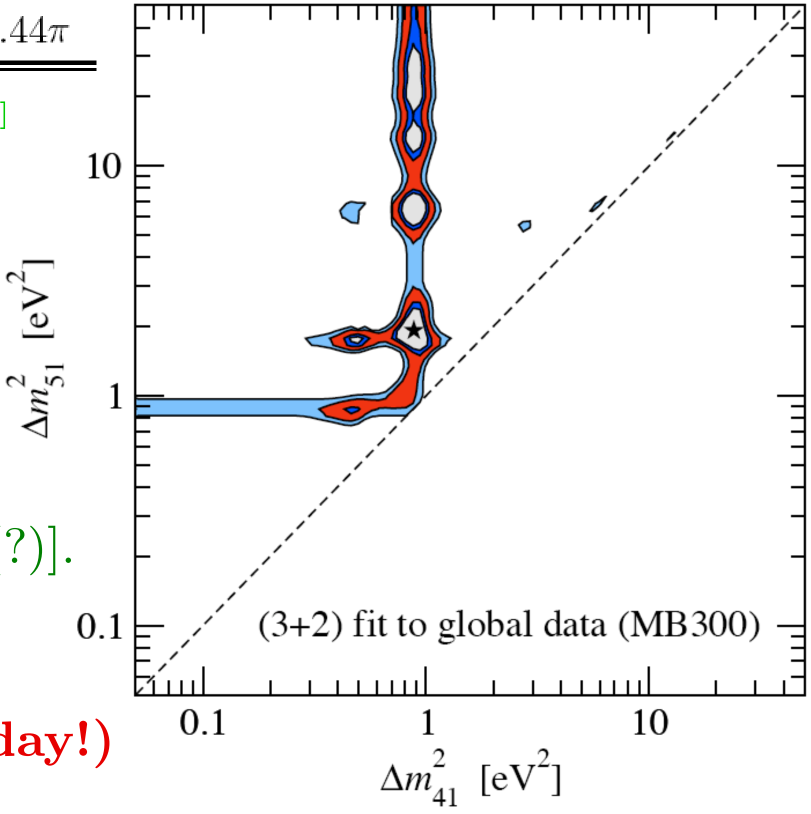
3+1+1 Fits Introduce an Extra Δm^2 and New Mixing Parameters

data set	$\lvert U_{e4}U_{\mu4}\rvert$		Δm^2_{41}	$\lvert U_{e5}U_{\mu5}\rvert$		Δm^2_{51}	δ – CP-violating phase
appearance (MB475)	0.044		0.66	0.022		1.44	1.12π
appearance (MB300)	0.31		0.66	0.27		0.76	1.01π
	$\lvert U_{e4}\rvert$	$\lvert U_{\mu4}\rvert$		$\lvert U_{e5}\rvert$	$\lvert U_{\mu5}\rvert$		
global data (MB475)	0.11	0.16	0.89	0.12	0.12	6.49	1.64π
global data (MB300)	0.12	0.18	0.87	0.11	0.089	1.91	1.44π

[Maltoni, Schwetz, arXiv:0705.0107 [hep-ph]]

Mini-BooNE and LSND fit "perfectly," including low-energy excess (MB300).

However, severely disfavored by disappearance data, especially if MB300 is included [$3\sigma - 4\sigma$ (?)].



(talk by Georgia Karagiorgi on Tuesday!)

Who Cares About Neutrino Masses: Only* "Palpable" Evidence of Physics Beyond the Standard Model

The SM we all learned in school predicts that neutrinos are strictly massless. Massive neutrinos imply that the the SM is incomplete and needs to be replaced/modified.

Furthermore, the SM has to be replaced by something qualitatively different.

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* There is only a handful of questions our model for fundamental physics cannot explain properly. These are, in order of "palpability" (my opinion!):

- What is the physics behind electroweak symmetry breaking? (Higgs *or* not in SM).
- What is the dark matter? (not in SM).
- Why does the Universe appear to be accelerating? Why does it appear that the Universe underwent rapid acceleration in the past? (not in SM – is this "particle physics?").

What I Mean By the Standard Model

The SM is a quantum field theory with the following defining characteristics:

- Gauge Group ($SU(3)_c \times SU(2)_{\rm L} \times U(1)_{\rm Y}$);
- Particle Content (fermions: Q, u, d, L, e, scalars: H).

Once this is specified, the SM is unambiguously determined:

- Most General Renormalizable Lagrangian;
- Measure All Free Parameters, and You Are Done! (after several decades of hard experimental work…)

If you follow these rules, neutrinos have no mass. Something has to give.

What is the New Standard Model? [νSM]

The short answer is – WE DON'T KNOW. Not enough available info!

$$\Updownarrow$$

Equivalently, there are several completely different ways of addressing neutrino masses. The key issue is to understand what else the νSM candidates can do. [are they falsifiable?, are they "simple"?, do they address other outstanding problems in physics?, etc]

We need more experimental input, and it looks like it may be coming in the near/intermediate future!

Options include:

- modify SM Higgs sector (e.g. Higgs triplet) and/or
- modify SM particle content (e.g. $SU(2)_L$ Triplet or Singlet) and/or
- modify SM gauge structure and/or
- supersymmetrize the SM and add R-parity violation and/or
- augment the number of space-time dimensions and/or
- etc

Massive Neutrinos and the Seesaw Mechanism

A simple[a], renormalizable Lagrangian that allows for neutrino masses is

$$\mathcal{L}_\nu = \mathcal{L}_{\rm old} - \lambda_{\alpha i} L^\alpha H N^i - \sum_{i=1}^{3} \frac{M_i}{2} N^i N^i + H.c.,$$

where N_i ($i = 1, 2, 3$, for concreteness) are SM gauge singlet fermions. $\mathcal{L}_\nu$ is the most general, renormalizable Lagrangian consistent with the SM gauge group and particle content, plus the addition of the N_i fields.

After electroweak symmetry breaking, $\mathcal{L}_\nu$ describes, besides all other SM degrees of freedom, six Majorana fermions: **six neutrinos.**

---

[a] Only requires the introduction of three fermionic degrees of freedom, no new interactions or symmetries.

To be determined from data: λ and M.

The data can be summarized as follows: there is evidence for three neutrinos, mostly "active" (linear combinations of ν_e, ν_μ, and ν_τ). At least two of them are massive and, if there are other neutrinos, they have to be "sterile."

This provides very little information concerning the magnitude of M_i (assume $M_1 \sim M_2 \sim M_3$)

Theoretically, there is prejudice in favor of very large M: $M \gg v$. Popular examples include $M \sim M_{\rm GUT}$ (GUT scale), or $M \sim 1$ TeV (EWSB scale).

Furthermore, $\lambda \sim 1$ translates into $M \sim 10^{14}$ GeV, while thermal leptogenesis requires the lightest M_i to be around 10^{10} GeV.

we can impose very, very few experimental constraints on M

What We Know About M:

- $M = 0$: the six neutrinos "fuse" into three Dirac states. Neutrino mass matrix given by $\mu_{\alpha i} \equiv \lambda_{\alpha i} v$.

  The symmetry of $\mathcal{L}_\nu$ is enhanced: $U(1)_{B-L}$ is an exact global symmetry of the Lagrangian if all M_i vanish. Small M_i values are 'tHooft natural.

- $M \gg \mu$: the six neutrinos split up into three mostly active, light ones, and three, mostly sterile, heavy ones. The light neutrino mass matrix is given by $m_{\alpha\beta} = \sum_i \mu_{\alpha i} M_i^{-1} \mu_{\beta i}$ $[m = 1/\Lambda \;\Rightarrow\; \Lambda = M/\mu^2]$.

  This the **seesaw mechanism.** Neutrinos are Majorana fermions. Lepton number is not a good symmetry of $\mathcal{L}_\nu$, even though L-violating effects are hard to come by.

- $M \sim \mu$: six states have similar masses. Active–sterile mixing is very large. This scenario is (generically) ruled out by active neutrino data (atmospheric, solar, KamLAND, K2K, etc).

High-energy seesaw has no observable consequence other than non-zero neutrino masses, except, perhaps,

Baryogenesis via Leptogenesis

One of the most basic questions we are allowed to ask (with any real hope of getting an answer) is whether the observed baryon asymmetry of the Universe can be obtained from a baryon–antibaryon symmetric initial condition plus well understood dynamics. [Baryogenesis]

This isn't just for aesthetic reasons. If the early Universe undergoes a period of inflation, baryogenesis is required, as inflation would wipe out any pre-existing baryon asymmetry.

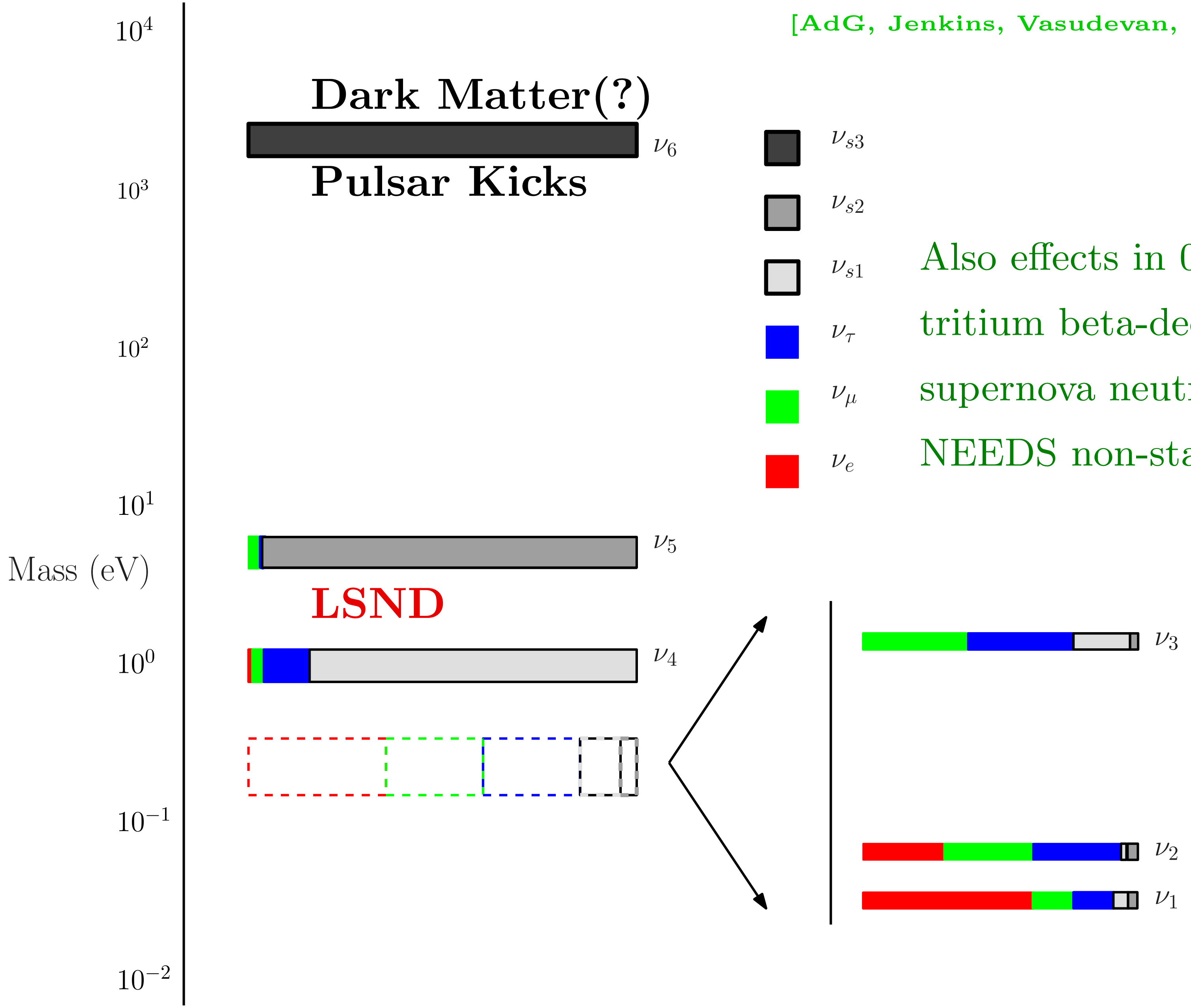
It turns out the seesaw mechanism contains all necessary ingredients to explain the baryon asymmetry of the Universe as long as the right-handed neutrinos are heavy enough – $M > 10^9$ GeV (with some exceptions that I won't have time to mention).

Low-Energy Seesaw [AdG PRD72,033005)]

Lets peek in the other end of the M spectrum. What do we get?

- Neutrino masses are small because the Yukawa couplings are very small $\lambda \in [10^{-6}, 10^{-11}]$;
- No standard thermal leptogenesis – right-handed neutrinos way too light;
- No obvious connection with other energy scales (EWSB, GUTs, etc);
- Right-handed neutrinos are propagating degrees of freedom. They look like sterile neutrinos ⇒ sterile neutrinos associated with the fact that the active neutrinos have mass;
- sterile–active mixing can be predicted – hypothesis is falsifiable!
- Small values of M are natural (in the 'tHooft sense). In fact, theoretically, no value of M should be discriminated against!

[AdG, Jenkins, Vasudevan, PRD75, 013003 (2007)]



Also effects in $0\nu\beta\beta$,

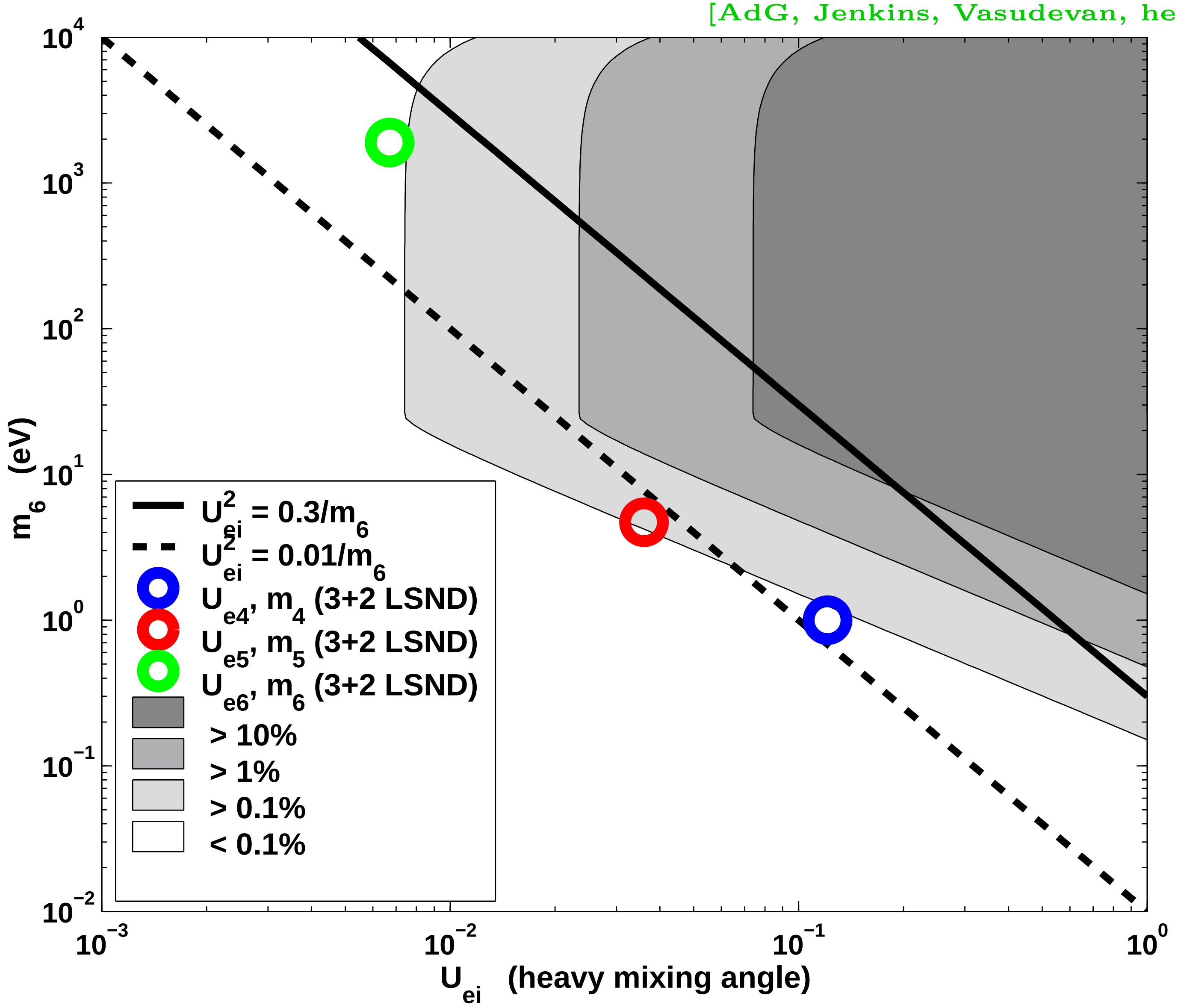
tritium beta-decay,

supernova neutrino oscillations,

NEEDS non-standard cosmology.

sensitivity of tritium beta decay to seesaw sterile neutrinos

[AdG, Jenkins, Vasudevan, hep-ph/0608147]



Other predictions: **Neutrinoless Double-Beta Decay**

The exchange of Majorana neutrinos mediates lepton-number violating neutrinoless double-beta decay, $0\nu\beta\beta$: $Z \to (Z+2)e^-e^-$.

For light enough neutrinos, the amplitude for $0\nu\beta\beta$ is proportional to the effective neutrino mass

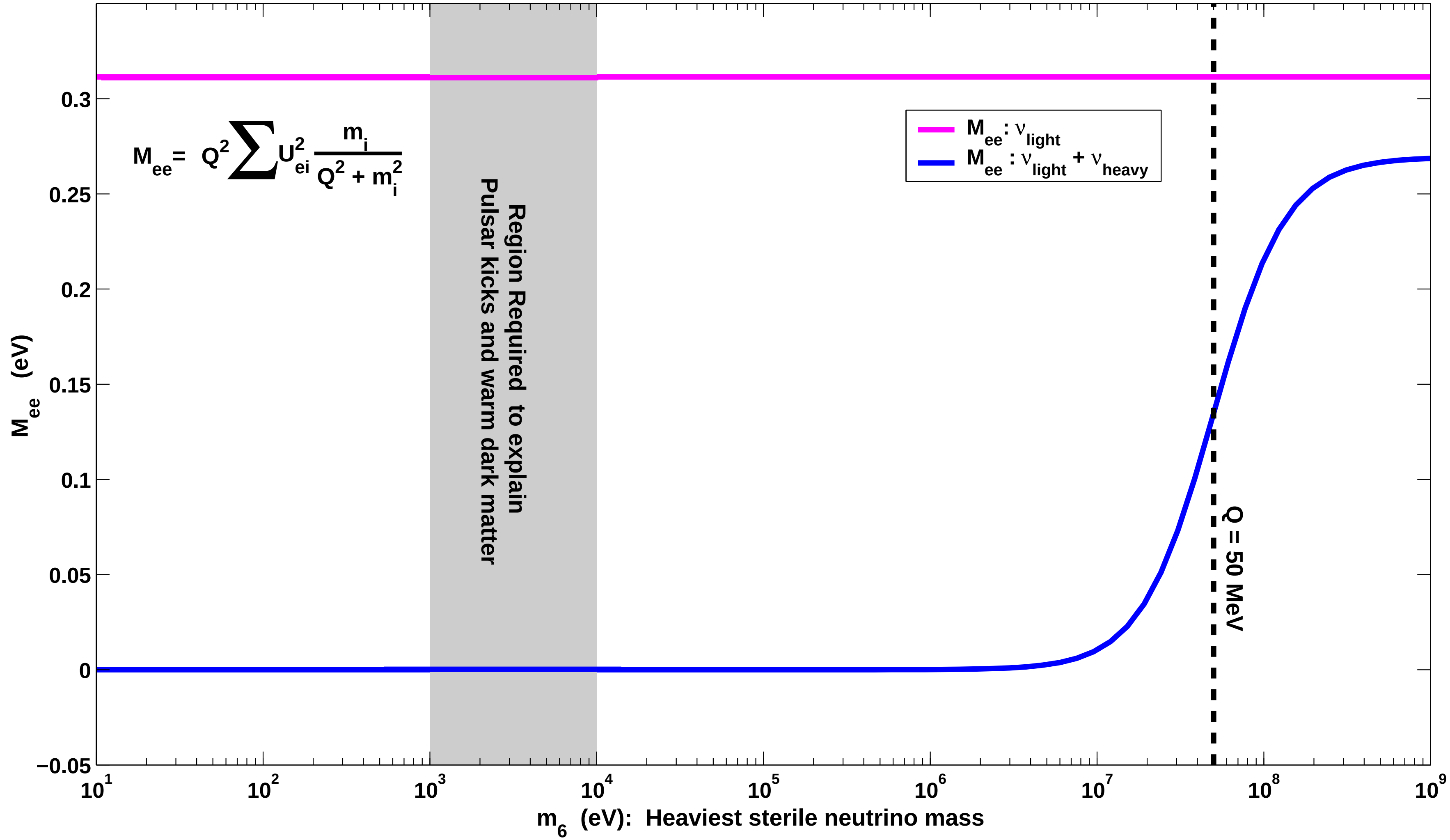
$$m_{ee} = \left|\sum_{i=1}^{6} U_{ei}^2 m_i\right| \sim \left|\sum_{i=1}^{3} U_{ei}^2 m_i + \sum_{i=1}^{3} \vartheta_{ei}^2 M_i\right|.$$

However, upon further examination, $m_{ee} = 0$ in the eV-seesaw. **The contribution of light and heavy neutrinos exactly cancels!** This seems to remain true to a good approximation as long as $M_i \ll 1$ MeV.

$$\left[\ \mathcal{M} = \begin{pmatrix} 0 & \mu^{\rm T} \\ \mu & M \end{pmatrix} \quad \to \quad m_{ee} \text{ is identically zero!}\ \right]$$

(lack of) sensitivity in $0\nu\beta\beta$ due to seesaw sterile neutrinos

[AdG, Jenkins, Vasudevan, hep-ph/0608147]



Why are Neutrino Masses Small?

In the old SM, neutrino masses are zero. However, the SM content allows for neutrino masses at the non-renormalizable level (dimension five operator):

$$\mathcal{L}_5 = \frac{LHLH}{\Lambda}.$$

Neutrino masses are small if $\Lambda \gg \langle H \rangle$. In the case of the seesaw,

$$\Lambda \sim \frac{M}{\lambda^2},$$

so neutrino masses are small if either

- they are generated by physics at a high energy scale (usual seesaw);

  or

- they arise out of a very weak coupling between the SM and a new, hidden sector (low energy seesaw).

Another possibility is that the physics responsible for neutrino masses leads to higher-dimensional effective operators, like

$$\mathcal{L}_\nu \propto \left(\frac{\phi}{\Lambda}\right)^N \frac{LHLH}{\Lambda},$$

or, if there are right-handed neutrino fields N,

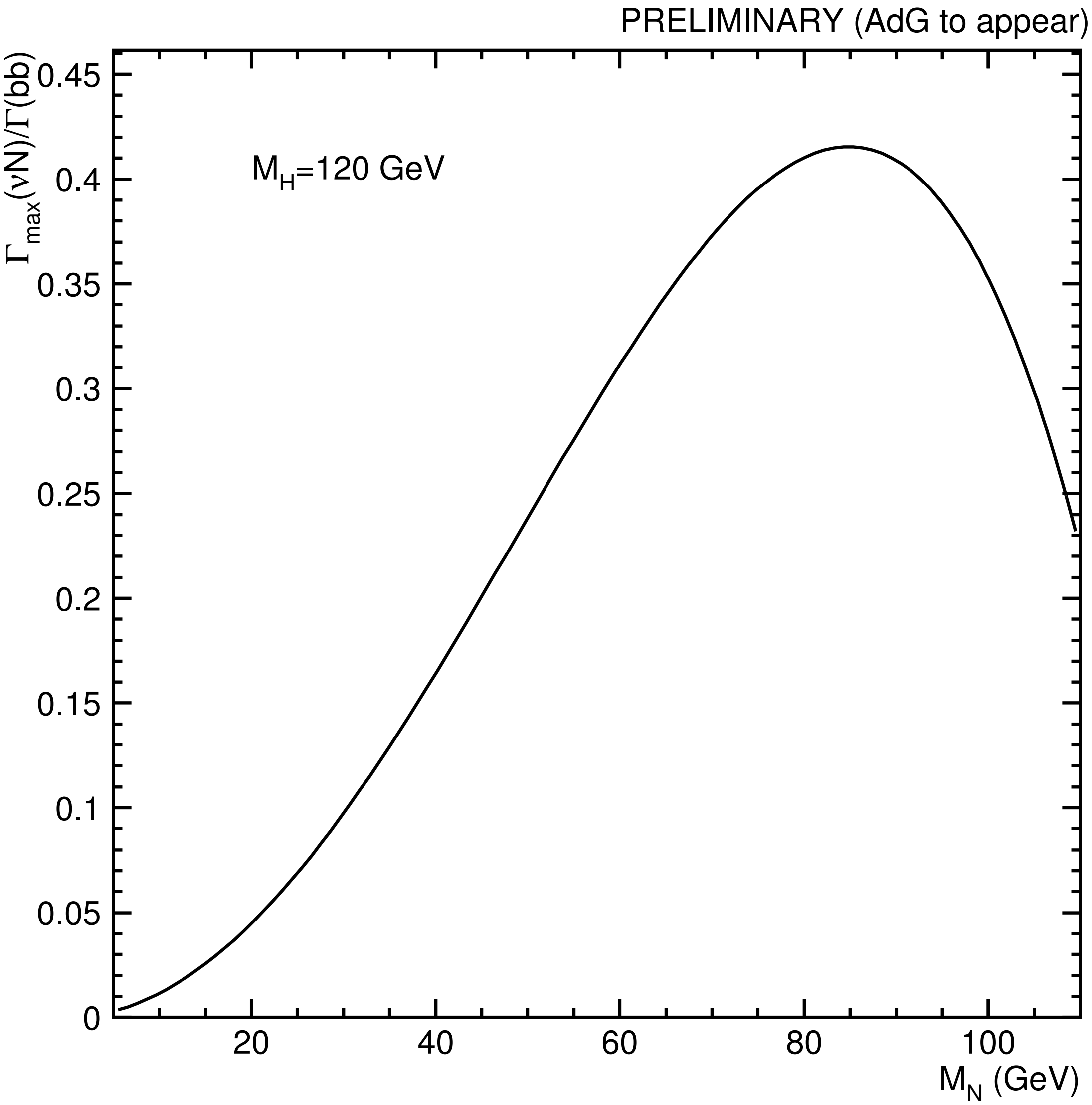
$$\mathcal{L}_\nu \propto \left(\frac{\phi}{\Lambda}\right)^N LHN.$$

In this case, the physics responsible for neutrino masses is neither very heavy, nor very weakly coupled.

⇒ potentially accessible in particle physics experiments!

Weak Scale Seesaw, and Accidentally Light Neutrino Masses

[AdG to appear]



What does the seesaw Lagrangian predict for the LHC?

Nothing much, unless. . .

- $M_N \sim 1 - 100$ GeV,
- Yukawa couplings larger than naive expectations.

$\Leftarrow H \to \nu N$ as likely as $H \to b\bar{b}$!

$\big($NOTE: $N \to \ell q'\bar{q}$ or $\ell\ell'\nu$) either prompt or with displaced vertex. “Weird” Higgs decay signature! $\big)$

ALSO: “Majorana neutrinos at the LHC,” see Han, Zhang, hep-ph/0604064

et cetera

How Do We Learn More?

In order to learn more, we need more information. Any new data and/or idea is welcome, including

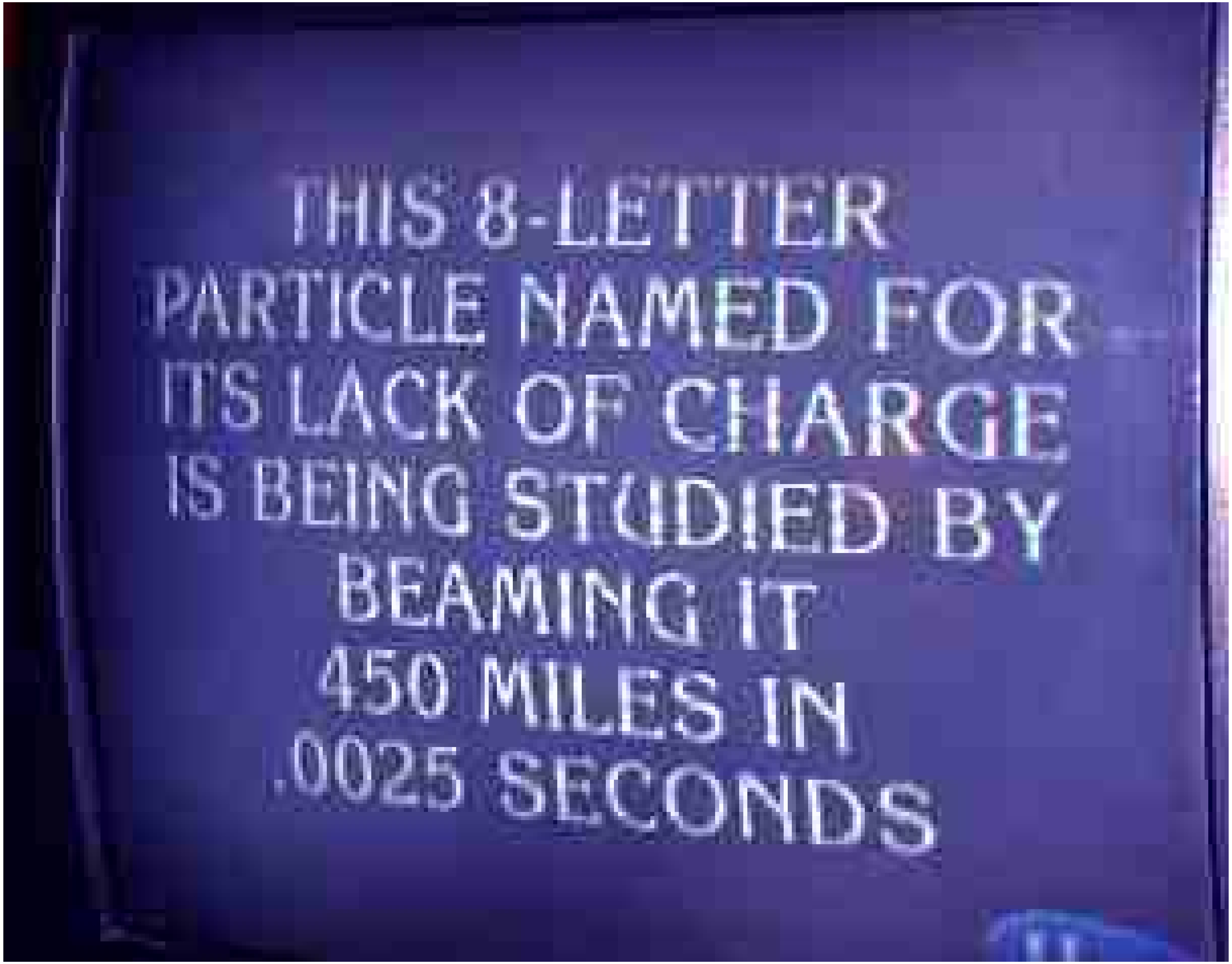
- searches for charged lepton flavor violation ($\mu \to e\gamma$, etc);
- searches for lepton number violation (neutrinoless double beta decay, etc);
- precision measurements of the neutrino oscillation parameters;
- searches for fermion electric/magnetic dipole moments (electron edm, muon $g-2$, etc);
- searches for new physics at the TeV scale – we need to understand the physics at the TeV scale before we can really understand the physics behind neutrino masses (is the low-energy SUSY?, etc).

CONCLUSIONS

The venerable Standard Model has finally sprung a leak – neutrinos are not massless!

1. we have a very successful parametrization of the neutrino sector, and we have identified what we know we don't know.
2. neutrino masses are very small – we don't know why, but we think it means something important.
3. lepton mixing is very different from quark mixing – we don't know why, but we think it means something important.
4. we need a minimal νSM Lagrangian. In order to decide which one is "correct" (required in order to attack 2. and 3. above) we must uncover the faith of baryon number minus lepton number ($0\nu\beta\beta$ is the best [only?] bet).

5. We need more experimental input – and more seems to be on the way (this is a truly data driven field right now). We only started to figure out what is going on.
6. The fact that neutrinos have mass may be intimately connected to the fact that there are more baryons than antibaryons in the Universe. How do we test whether this is correct?
7. There is plenty of room for surprises, as neutrinos are very narrow but deep probes of all sorts of physical phenomena. Remember that neutrino oscillations are "quantum interference devices" – potentially very sensitive to whatever else may be out there (e.g., $\Lambda \simeq 10^{14}$ GeV).



THIS 8-LETTER
PARTICLE NAMED FOR
ITS LACK OF CHARGE
IS BEING STUDIED BY
BEAMING IT
450 MILES IN
.0025 SECONDS

Back-up Slides:

(mostly stolen from earlier presentations)

Understanding Fermion Mixing

The other puzzling phenomenon uncovered by the neutrino data is the fact that Neutrino Mixing is Strange. What does this mean?

It means that lepton mixing is very different from quark mixing:

$$V_{MNS} \sim \begin{pmatrix} 0.8 & 0.5 & \mathbf{0.2} \\ 0.4 & 0.6 & 0.7 \\ 0.4 & 0.6 & 0.7 \end{pmatrix} \qquad V_{CKM} \sim \begin{pmatrix} 1 & 0.2 & 0.001 \\ 0.2 & 1 & 0.01 \\ 0.001 & 0.01 & 1 \end{pmatrix}$$

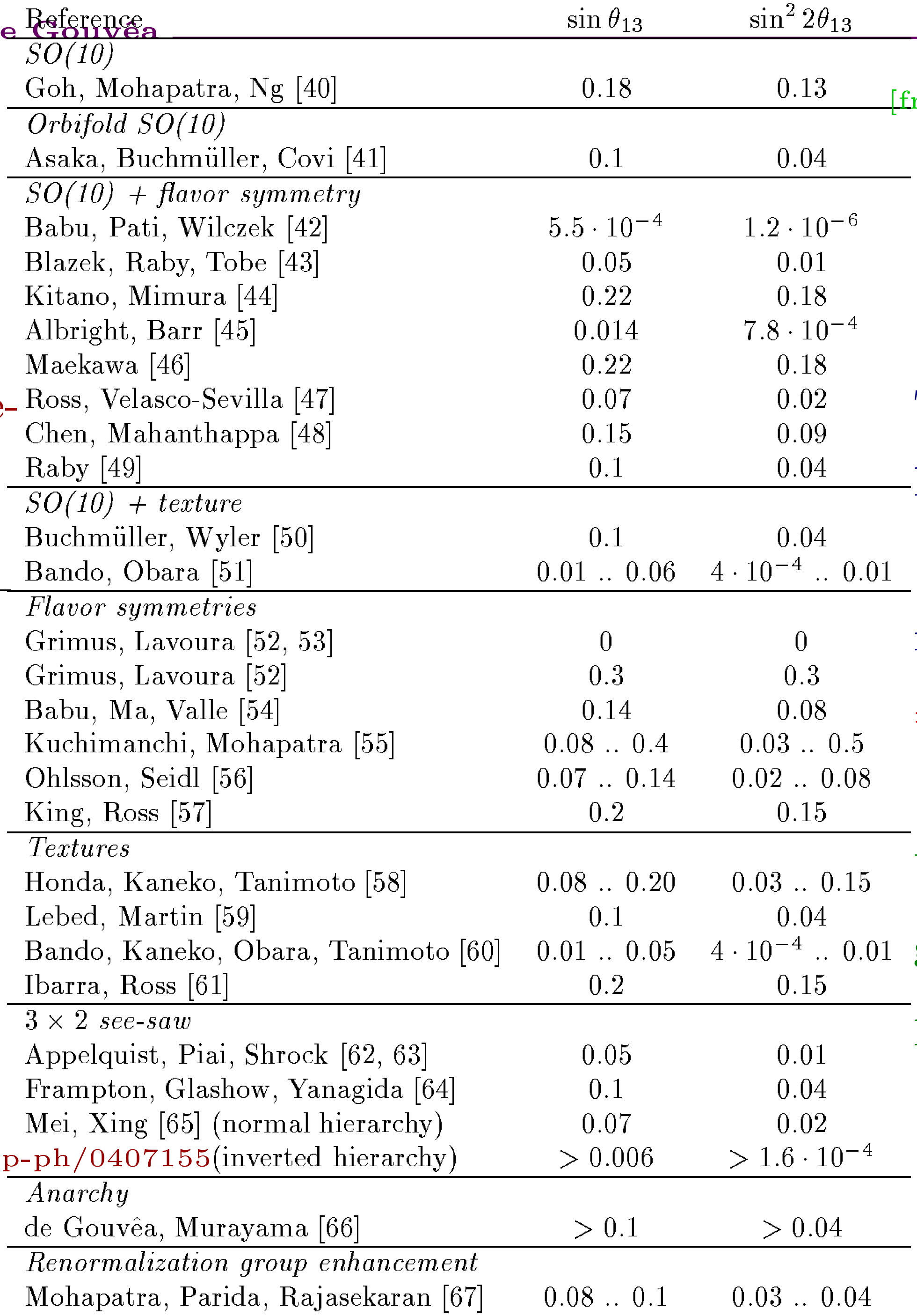
$[|(V_{MNS})_{e3}| < 0.2]$

WHY?

They certainly look VERY different, but which one would you label as “strange”?

$\Delta m^2_{13} > 0$ "typical" prediction of all* Type-I see-saw GUT models

inverted hierarchy requires* "more flavor structure"



Reference	$\sin\theta_{13}$	$\sin^2 2\theta_{13}$
SO(10)		
Goh, Mohapatra, Ng [40]	0.18	0.13
Orbifold SO(10)		
Asaka, Buchmüller, Covi [41]	0.1	0.04
SO(10) + flavor symmetry		
Babu, Pati, Wilczek [42]	$5.5 \cdot 10^{-4}$	$1.2 \cdot 10^{-6}$
Blazek, Raby, Tobe [43]	0.05	0.01
Kitano, Mimura [44]	0.22	0.18
Albright, Barr [45]	0.014	$7.8 \cdot 10^{-4}$
Maekawa [46]	0.22	0.18
Ross, Velasco-Sevilla [47]	0.07	0.02
Chen, Mahanthappa [48]	0.15	0.09
Raby [49]	0.1	0.04
SO(10) + texture		
Buchmüller, Wyler [50]	0.1	0.04
Bando, Obara [51]	0.01 .. 0.06	$4 \cdot 10^{-4}$ .. 0.01
Flavor symmetries		
Grimus, Lavoura [52, 53]	0	0
Grimus, Lavoura [52]	0.3	0.3
Babu, Ma, Valle [54]	0.14	0.08
Kuchimanchi, Mohapatra [55]	0.08 .. 0.4	0.03 .. 0.5
Ohlsson, Seidl [56]	0.07 .. 0.14	0.02 .. 0.08
King, Ross [57]	0.2	0.15
Textures		
Honda, Kaneko, Tanimoto [58]	0.08 .. 0.20	0.03 .. 0.15
Lebed, Martin [59]	0.1	0.04
Bando, Kaneko, Obara, Tanimoto [60]	0.01 .. 0.05	$4 \cdot 10^{-4}$ .. 0.01
Ibarra, Ross [61]	0.2	0.15
3 × 2 see-saw		
Appelquist, Piai, Shrock [62, 63]	0.05	0.01
Frampton, Glashow, Yanagida [64]	0.1	0.04
Mei, Xing [65] (normal hierarchy)	0.07	0.02
(inverted hierarchy)	> 0.006	$> 1.6 \cdot 10^{-4}$
Anarchy		
de Gouvêa, Murayama [66]	> 0.1	> 0.04
Renormalization group enhancement		
Mohapatra, Parida, Rajasekaran [67]	0.08 .. 0.1	0.03 .. 0.04

Table 1: Incomplete selection of predictions for θ_{13}. The numbers should be considered as order of magnitude statements.

[from reactor white paper]

Theoretical predictions:

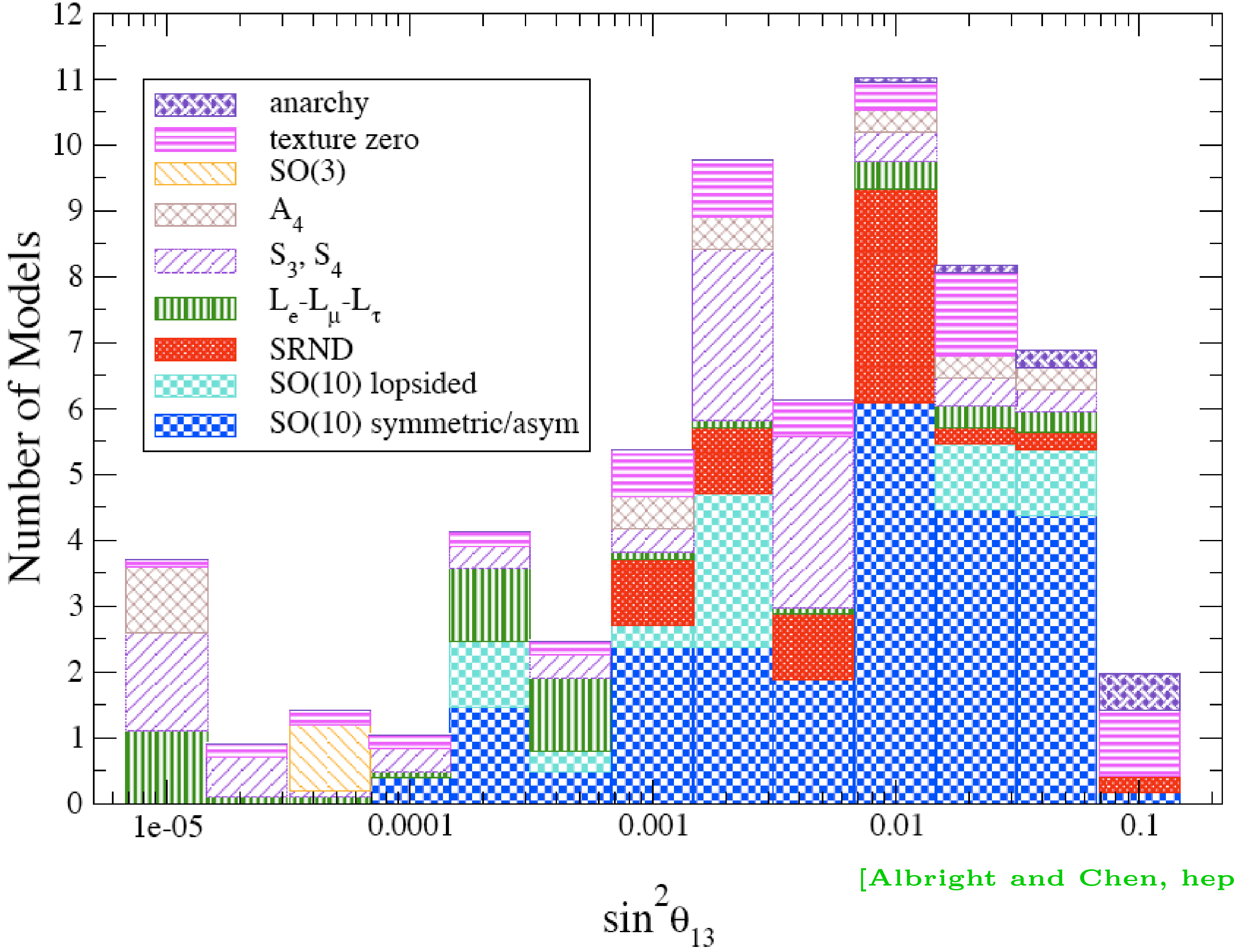
The literature on this subject is very large. The most exciting driving force (my opinion) is the fact that one can make *bona fide* predictions:

$\Rightarrow U_{e3}$, CP-violation, mass-hierarchy unknown!

Unfortunately, theorists have done too good a job, and people have successfully predicted everything. . .

More data needed to "sort things out."

* Albright,hep-ph/0407155



[Albright and Chen, hep-ph/0608137]

pessimist – “We can’t compute what $|U_{e3}|$ is – must measure it!”

(same goes for the mass hierarchy, δ)

Candidate νSM

SM as an effective field theory – non-renormalizable operators

$$\mathcal{L}_{\nu\mathrm{SM}} \supset -\lambda_{ij}\frac{L^iHL^jH}{2\Lambda} + \mathcal{O}\left(\frac{1}{\Lambda^2}\right) + H.c.$$

There is only one dimension five operator [Weinberg, 1979]. If $\Lambda \gg 1$ TeV, it leads to only one observable consequence...

$$\text{after EWSB } \mathcal{L}_{\nu\mathrm{SM}} \supset \frac{m_{ij}}{2}\nu^i\nu^j; \quad m_{ij} = \lambda_{ij}\frac{v^2}{\Lambda}.$$

- Neutrino masses are small: $\Lambda \gg v \rightarrow m_\nu \ll m_f$ ($f = e, \mu, u, d$, etc)
- Neutrinos are Majorana fermions – Lepton number is violated!
- νSM effective theory – not valid for energies above *at most* Λ.
- What is Λ? First naive guess is that M is the Planck scale – does not work. Data require $\Lambda < 10^{15}$ GeV (anything to do with the GUT scale?)

What else is this “good for”? Depends on the ultraviolet completion!

Full disclosure:

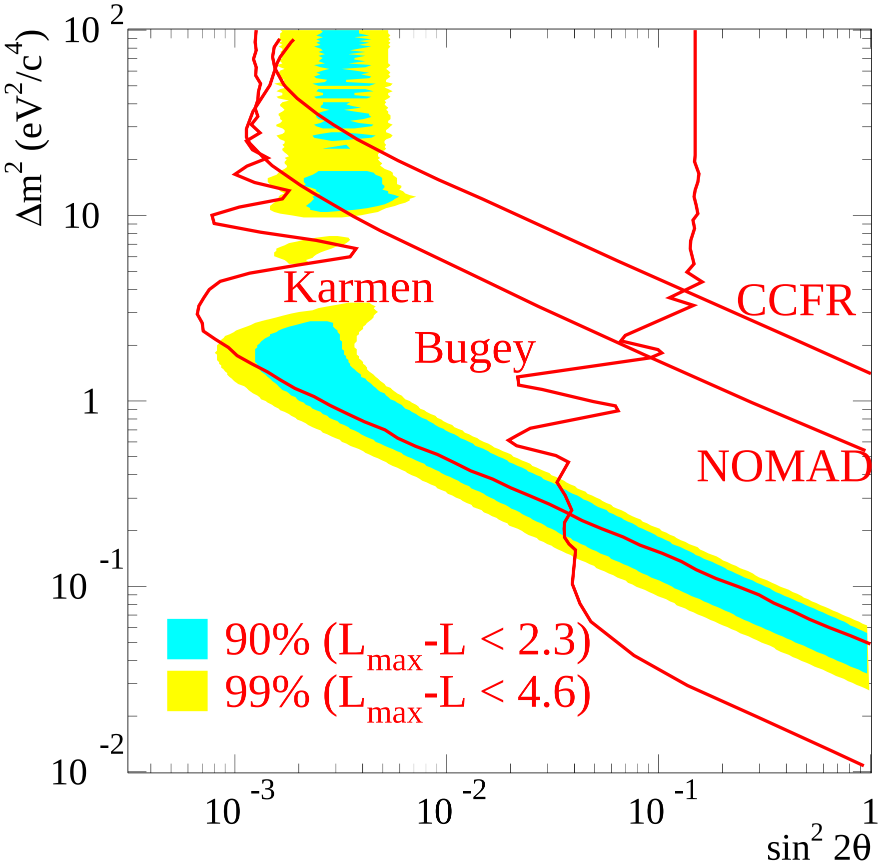
All higher dimensional operators are completely negligible, **except** those that mediate proton decay, like:

$$\frac{\lambda_B}{M^2}QQQL$$

The fact that the proton does not decay forces M/λ_B to be much larger than the energy scale required to explain neutrino masses.

Why is that? We don't know…

Is this a big deal? We don't know…



Karmen has a similar sensitivity to

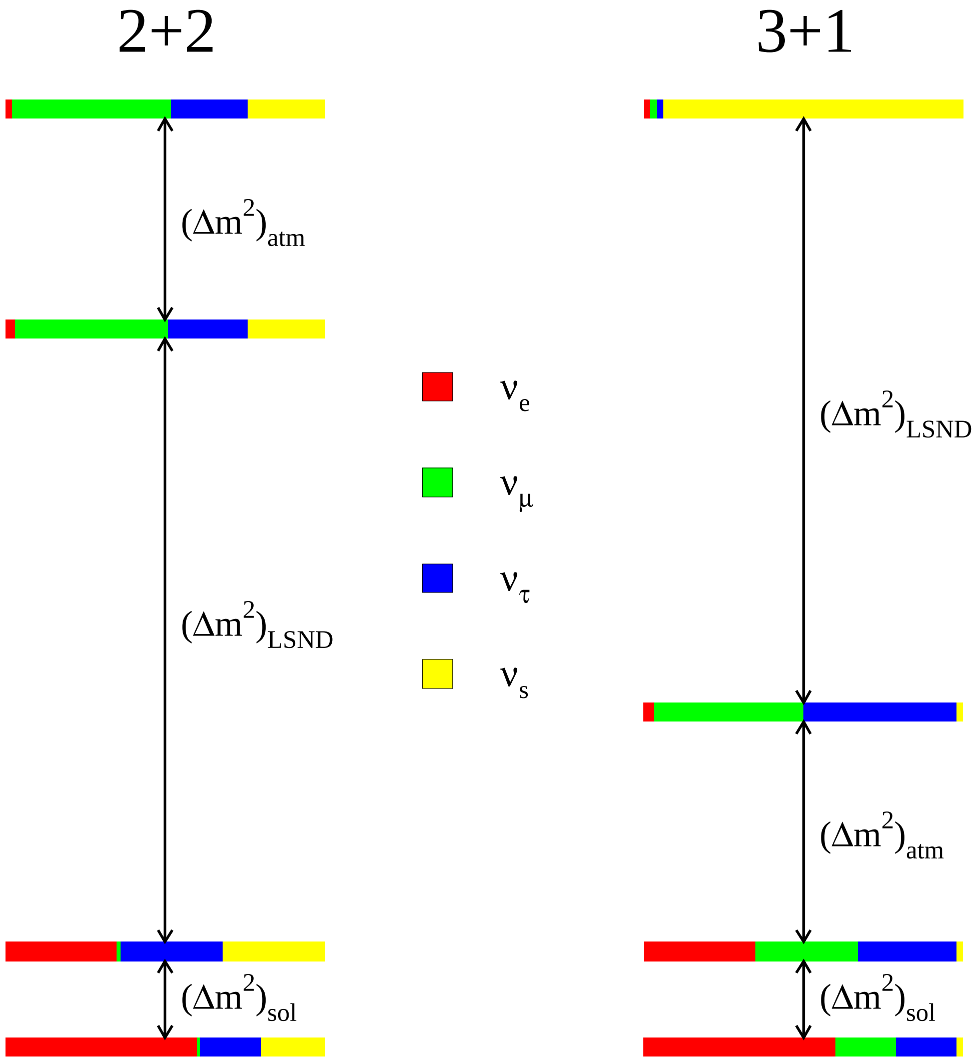
$\bar{\nu}_\mu \to \bar{\nu}_e$, but a shorter baseline ($L = 18$ m)

Other curves are failed searches for

ν_μ disappearance (CCFR),

$\bar{\nu}_e$ disapperance (Bugey), etc

Remember: $P_{\mu e} = \sin^2 2\theta \sin^2 \left[1.27 \left(\frac{\Delta m^2}{\mathrm{eV}^2}\right) \left(\frac{L}{\mathrm{m}}\right) \left(\frac{\mathrm{MeV}}{E}\right)\right]$



⇒ 2+2 requires large sterile effects in either solar or atmospheric oscillations, not observed

Another νSM

Why don't we just enhance the fermion sector of the theory?

One may argue that it is trivial and simpler to just add

$$\mathcal{L}_{\rm Yukawa} = -y_{i\alpha} L^i H N^\alpha + H.c.,$$

and neutrinos get a mass like all other fermions: $m_{i\alpha} = y_{i\alpha} v$

- Data requires $y < 10^{-12}$. Why so small?
- Neutrinos are Dirac fermions. $B - L$ exactly conserved.
- νSM is a renormalizable theory.

This proposal, however, violates the rules of the SM (as I defined them)! The operator $\frac{M_N}{2}NN$, allowed by all gauge symmetries, is absent. In order to explain this, we are forced to add a symmetry to the νSM. The simplest candidate is a global $U(1)_{B-L}$.

$U(1)_{B-L}$ is upgraded from accidental to fundamental (global) symmetry.

Old Standard Model, Encore

The SM is a quantum field theory with the following defining characteristics:

- Gauge Group $(SU(3)_c \times SU(2)_{\rm L} \times U(1)_{\rm Y})$;
- Particle Content (fermions: Q, u, d, L, e, scalars: H).

Once this is specified, the SM is unambiguously determined:

- Most General Renormalizable Lagrangian;
- Measure All Free Parameters, and You Are Done.

This model has *accidental global symmetries*. In particular, the anomaly free global symmetry is preserved: $U(1)_{B-L}$.

New Standard Model, Dirac Neutrinos

The SM is a quantum field theory with the following defining characteristics:

- Gauge Group $(SU(3)_c \times SU(2)_{\rm L} \times U(1)_{\rm Y})$;
- Particle Content (fermions: Q, u, d, L, e, N, scalars: H);
- Global Symmetry $U(1)_{B-L}$.

Once this is specified, the SM is unambiguously determined:

- Most General Renormalizable Lagrangian;
- Measure All Free Parameters, and You Are Done.

Naively not too different, but nonetheless qualitatively different → enhanced symmetry sector!

Other predictions: **Tritium beta-decay**

Heavy neutrinos participate in tritium β-decay. Their contribution can be parameterized by

$$m_\beta^2 = \sum_{i=1}^{6} |U_{ei}|^2 m_i^2 \simeq \sum_{i=1}^{3} |U_{ei}|^2 m_i^2 + \sum_{i=1}^{3} |U_{ei}|^2 m_i M_i,$$

as long as M_i is not too heavy (above tens of eV). For example, in the 3+2 scenario of the previous slide, $m_\beta^2 \simeq 0.7~\text{eV}^2 \left(\frac{|U_{e1}|^2}{0.7}\right)\left(\frac{m_1}{0.1~\text{eV}}\right)\left(\frac{M_1}{10~\text{eV}}\right)$.

NOTE: next generation experiment (KATRIN) will be sensitive to $O(10^{-1})$ eV^2.

On Early Universe Cosmology / Astrophysics

A combination of the SM of particle physics plus the "concordance cosmological model" severely constrain light, sterile neutrinos with significant active-sterile mixing. Taken at face value, not only is the eV-seesaw ruled out, but so are all oscillation solutions to the LSND anomaly.

Hence, **eV-seesaw → nonstandard particle physics and cosmology**.

On the other hand. . .

- Right-handed neutrinos may make good warm dark matter particles.

  Asaka, Blanchet, Shaposhnikov, hep-ph/0503065.

- Sterile neutrinos are known to help out with r-process nucleosynthesis in supernovae, . . .
- . . . and may help explain the peculiar peculiar velocities of pulsars.

—— ASIDE ——

On very small Yukawa couplings

We would like to believe that Yukawa couplings should naturally be of order one.

Nature, on the other hand, seems to have a funny way of showing this. Of all known fermions, only one (1) has a “natural” Yukawa coupling – the top quark!

Regardless there are several very different ways of obtaining “naturally” very small Yukawa couplings. They require more new physics.

—— END ASIDE ——

Neutrinos Masses And Colliders: Non-Anomalous, Gauged $U(1)_\nu$

And it could turn out that neutrino masses are deeply connected to physics at the electroweak symmetry breaking scale:

Add to the SM a new, non-anomalous $U(1)_\nu$ under which both SM fermions and the right-handed neutrinos transform. Charges are heavily constrained by anomaly cancellations and the fact that quarks and charged leptons have relatively large masses.

One can choose $U(1)_\nu$ charges so that all neutrino masses are forbidden by gauge invariance. This way, neutrino masses are only generated after $U(1)_\nu$ is spontaneously broken,[a] and only through higher dimensional operators, suppressed by a new ultraviolet scale Λ.

Neutrino masses might be small because they are a consequence of very high dimensional operators: $m_\nu \propto \left(\frac{\varphi}{\Lambda}\right)^{|p|}$, where p is an integer exponent.

[M.C. Chen, B. Dobrescu, hep-ph/0612017, PRD in press]

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[a] Assume $U(1)_\nu$ is spontaneous broken when SM singlet scalar Φ gets a vev, $\langle\Phi\rangle \equiv \varphi$.

After $U(1)_\nu$ breaking → see-saw Lagrangian plus "left–left" neutrino mass:

$$\mathcal{L} \supset \sum_{ik} \epsilon^{|p_{ik}|} \bar{L}_i (\lambda^\nu)^{ik} n_k \tilde{H} + \sum_{ij} \epsilon^{|q_{ij}|} \bar{L}_i^c \frac{(h^L)^{ij}}{\Lambda} L_j H H + \sum_{kk'} \epsilon^{|r_{kk'}|} \Lambda \bar{n}_k^c (h^R)^{kk'} n_{k'},$$

λ^ν – neutrino Yukawa coupling, h^L – "left–left" coupling), and h^R – "right–right" Majorana mass term). $i,j = 1,2,3$, $k,k' = 1 \ldots N$. Only allowed for integer values of p, q, and r.

Consequences for collider physics:

- Non-standard Z' – branching ratios to different fermion species can be matched to neutrino mass structure!
- Enhanced Higgs sector.

[M.C. Chen, B. Dobrescu, hep-ph/0612017, PRD in press]