# **Space-Time Propagation of Neutrino Wave-Packets in the Early Universe**

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# Introduction

 $\bigstar$  Neutrino wave-packets:

(I) involved during production and detection processes in realistic experiments

(II) produced before BBN through nuclear beta decays

#### $\bigstar$ Current studies:

(I) focus on the vacuum case(II) reveal the existence of a *coherence time limit*beyond which neutrino oscillation vanishes

★ Propagation of neutrino wave-packets with dispersions at high temperature (for example, the early universe) has never been investigated

#### Neutrino Wave-Packet

★ At t = 0, a gaussian wave-packet of  $\nu_e$ and no  $\nu_{\mu}$ :  $\exp\left[-\frac{(\vec{k} - \vec{k}_0)^2}{4\sigma^2}\right]$ 

- ★ How does this neutrino wave-packet evolve and propagate (in *space-time*) through the high temperature thermal medium?
- ★ Transition probability  $\mathcal{P}_{\nu_e \to \nu_\mu}(\vec{r}, t) = ?$

### **Space-Time Evolution**

- ★ Implement finite-temperature QFT⇒ the effective Dirac equation in the medium ⇒ include self-energy corrections from the medium
- ★ Laplace transform solves the effective Dirac equation ⇒ time evolution of  $\nu_{\mu}$  for any momentum k
- ★ Fourier transform to obtain space-time evolution of  $\nu_{\mu}$ ⇒ transition probability  $\mathcal{P}_{\nu_e \to \nu_{\mu}}(\vec{r}, t)$

### Intuitive Picture # 1

- ★ Physically propagating modes are two mass eigenstates (almost degenerate):  $M_a = \overline{M} \left[ 1 + (-1)^{a-1} \frac{\delta M^2}{4\overline{M}^2} \right] ; |\delta M^2| / \overline{M}^2 \ll 1$
- ★ WMAP3 data:  $\overline{M} = \frac{1}{2} (M_1 + M_2) < \frac{1}{2} (M_1 + M_2 + M_3) \approx 0.34 \,\text{eV}$
- ★ Solar + KamLAND data: (I)  $|\delta M^2| = |M_1^2 - M_2^2| \approx 7.9 \times 10^{-5} \,(\text{eV})^2$ (II)  $\tan^2 \theta \approx 0.40$

#### Intuitive Picture #2

★ Small difference in group velocities: ⇒ two mass eigenstates separate progressively

★ Neutrino oscillations will be exponentially suppressed when two mass eigenstates cease to overlap: ⇒ medium coherence time limit  $T_{\text{medium}}^{\text{coherence}}$ 

★ Oscillatory term in  $\mathcal{P}_{\nu_e \to \nu_\mu}(\vec{r}, t)$  is proportional to  $\exp\left[-\left(\frac{t}{T_{\text{medium}}^{\text{coherence}}}\right)^2\right]\cos\left(\frac{t}{T_{\text{medium}}^{\text{oscillate}}}\right)$ 

#### What about the Dispersions?

- $\star$  Two new characteristic time scales emerge!
- ★ For  $t > T_{\perp}$ : ⇒ transverse dispersion becomes important
- ★ For  $t > T_{\parallel} \approx \left( k_0 / \overline{M} \right)^2 T_{\perp}$ : ⇒ longitudinal dispersion becomes important ⇒ competes with progressive separation between the two mass eigenstates

# Effects of Dispersions #1

- ★ Oscillatory term in  $\mathcal{P}_{\nu_e \to \nu_\mu}(\vec{r}, t \gtrsim T_{\parallel})$  is proportional to  $\exp\left[-\left(\frac{T_{\parallel}}{T_{\text{medium}}}\right)^2\right] \cos\left(\frac{t}{T_{\text{medium}}}\right)$
- ★ If  $T_{\parallel} < T_{\text{medium}}^{\text{coherence}}$ : ⇒ longitudinal dispersion will be able to catch up with progressive separation between the two mass eigenstates before the coherence time limit is exceeded
- ★ Neutrino oscillations will *never* be exponentially suppressed due to progressive separation between the two mass eigenstates  $\Rightarrow$  frozen-coherence

# Effects of Dispersions #2

- $\bigstar$  Recall that  $T_{\parallel} \gg T_{\perp}$
- $\bigstar$  Large suppression of transition probability due to enormous transverse dispersion:

$$\mathcal{P}_{\nu_e \to \nu_\mu}(\vec{r}, t \gg T_\perp) \propto \left(\frac{T_\perp}{t}\right)^2 \left( \text{oscillatory term} + \dots \right)$$

# Resonance Regime in the Early Universe

★ Assume B - L = 0 is a good symmetry: ⇒ charged-lepton and neutrino asymmetries are of the same order as  $\mathcal{L}_B \sim 10^{-9}$ 

★ Medium mixing angle:  $\sin 2\theta_m = \frac{\sin 2\theta}{\left[\left(\cos 2\theta - \frac{\text{PHENO}}{\delta M^2}\right)^2 + \sin^2 2\theta\right]^{\frac{1}{2}}}$ 

★ Resonance  $(\sin 2\theta_m \approx 1)$  occurs at  $T \sim a$  few MeV ⇒ right before BBN

## **Comparison of Time Scales**

★ Take  $T \sim a$  few MeV as the typical energy scale

★ Comparison:  $T_{\perp} << T_{\text{medium}}^{\text{oscillate}} \ll T_{\text{collision}} \ll T_{\text{Hubble}} \ll T_{\text{medium}}^{\text{coherence}} \sim T_{\parallel}$ ★  $\mathcal{P}_{\nu_e \to \nu_{\mu}}(\vec{r}, t \gg T_{\perp}) \propto \left(\frac{T_{\perp}}{t}\right)^2 \left(\text{oscillatory term} + ...\right)$ is highly suppressed on the time scale of  $T_{\text{medium}}^{\text{oscillate}}$  !!!

### Conclusions

- ★ characteristic time scales  $T_{\perp}$  and  $T_{\parallel}$  arise due to dispersions
- $\star$  frozen-coherence can occur if  $T_{\parallel} < T_{\rm medium}^{\rm coherence}$
- ★ large suppression of transition probability due to enormous transverse dispersion ⇒ distortion of  $\nu_e$  abundance due to neutrino oscillations right before BBN is negligible