## Predictive model of inverted neutrino hierarchy and resonant leptogenesis

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Neutrinos have tiny masses & their flavor mixing involves two large angles and one small angle:

#### Best fit values:

$$\Delta m_s^2 = 8.0 \times 10^{-5} \text{eV}^2 \quad \Delta m_a^2 = 2.5 \times 10^{-3} \text{eV}^2$$
  

$$\theta_{12} = 34^{\circ}$$
  

$$\theta_{23} = 45^{\circ}$$
  

$$\theta_{13} = 0^{\circ}$$

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- Why are the two mixing angles large?
- What is the sign of  $\Delta m_{atm}^2$  ?
- The sign is directly linked to neutrino mass hierarchy
- The value of  $\theta_{13}$ ?

(if  $\theta_{13} \neq 0 \rightarrow \delta$  of *CP* is physical)

⇒ A reasonable extension of the SM
 should provide an understanding of
 this pattern

#### Inverted hierarchy and *L*- Symmetry $L=L_e-L_{\mu}-L_{\tau}$

Starting from the following minimal seesaw

$$\mathfrak{I} = \overline{\nu}_L m_D N - \frac{1}{2} N^T M_R N$$

to ensure maximal  $v_{\mu} - v_{\tau}$  mixing, impose an interchange  $S_2$ symmetry between  $v_{\mu} \leftrightarrow v_{\tau}$  ( $\theta_{23} \simeq 45^{\circ}$ )

The most general couplings (that respect  $S_2$ ) is

$$Y_{\nu}^{D} = -1 \begin{pmatrix} \alpha & 0 \\ \beta' & \beta \\ -1 \begin{pmatrix} \beta' & \beta \\ \beta' & \beta \end{pmatrix}, \qquad M_{R} = -1 \begin{pmatrix} -\delta_{N} & 1 \\ 1 & -\delta'_{N} \end{pmatrix} M^{L=L_{e}-L_{\mu}-L_{\tau}}$$

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using seesaw;

$$m_{\nu} = \begin{pmatrix} 2\delta_{\nu}' & \sqrt{2} & \sqrt{2} \\ \sqrt{2} & \delta_{\nu} & \delta_{\nu} \\ \sqrt{2} & \delta_{\nu} & \delta_{\nu} \end{pmatrix} \frac{m}{2}$$

$$m = \frac{\langle h_u^0 \rangle^2}{M(1 - \delta_N \delta'_N)} \sqrt{2} \alpha \left(\beta + \beta' \delta'_N\right)$$

$$\delta_{\nu} = \frac{\sqrt{2}}{\alpha} \frac{2\beta\beta' + \beta^2 \delta_N + (\beta')^2 \delta'_N}{\beta + \beta' \delta'_N} ,$$

$$\delta_{\nu}^{'} = \frac{\alpha}{\sqrt{2}} \frac{\delta_{N}^{'}}{\beta + \beta' \delta_{N}^{'}}$$

(since they are proportional to *L*- symmetry breaking term,  $|\delta_{\nu}|, |\delta_{\nu}^{'}| \ll 1$  will also be responsible for the solar splitting) If L violating term were not present, you will only generate atmospheric mass splitting

Assuming the charged lepton sector to be diagonal, all the mixing will be contained in  $U_v = U_{23}U_{12}$  such that  $Diag(m_1, m_2, 0) = U_v^T m_v U_v$ 

$$U_{23} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} \\ 0 & \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \end{pmatrix}, \quad U_{12} \simeq \begin{pmatrix} \bar{c} & -\bar{s}e^{i\rho} & 0 \\ \bar{s}e^{-i\rho} & \bar{c} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$
$$Sin^{2}(\theta_{12}) \simeq \frac{1}{2}$$

corresponding to maximal mixing in the solar angle; imcompatible with experiments!!

• This picture leads to Bimaximal mixing  $(\theta_{12} = \pi/4)$ however,  $\theta_{12}$  is significantly away from maximal mixing (Kamland & Solar)

We take a step forward to correct the situation: The  $S_2$  – Symmetry is exact in neutrino mass matrix but broken in the charged lepton mass matrix (such that some predictions are preserved)

The charged lepton matrix will now be assumed non-diagonal,

- $\Rightarrow$  Contribute to leptonic mixing
- $\Rightarrow$  Use this contribution to fix the value of  $\theta_{12}$
- ⇒ Find interesting correlation between solar & reactor angles

Assume the following texture for charged lepton mass matrix:

$$Y_E = \begin{pmatrix} 0 & a' & 0 \\ a & \lambda_{\mu} & 0 \\ 0 & 0 & \lambda_{\tau} \end{pmatrix}, \quad a' = \lambda_{\mu} \theta_e e^{i\omega}$$

Can be diagonalized by the rotation :

$$U_e = \begin{pmatrix} c & se^{i\omega} & 0\\ -se^{-i\omega} & c & 0\\ 0 & 0 & 1 \end{pmatrix} \quad c \equiv \cos t, \ s \equiv \sin t \text{ and } \tan t = -\theta_e$$

The total mixing matrix:

$$U^l = U^*_e U_\nu \; ,$$

### **Prediction:**

$$\sin^2 \theta_{12} = \frac{1}{2} - \sqrt{1 - \tan^2 \theta_{13}} \tan \theta_{13} \cos \delta ,$$
$$\sin^2 \theta_{23} = \frac{1}{2} \left( 1 - \tan^2 \theta_{13} \right)$$

 $\Rightarrow$  The deviation of the solar angle from its maximal value

of 
$$\frac{\pi}{4}$$
 is controlled by  $\theta_{13}$ .

Implying several constraints:

 $0.31 \le Sin^2(\theta_{12}) \le 0.7, \rightarrow 0 \le Cos(\delta) \le 0.7, \theta_{13} \ge 0.13$ 



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### Resonant Leptogenesis

Our model of inverted hierarchy involves two quasi-degenerate RHN Resonant enhancement of CP Violation, successfully leading to Baryon Asymmetry

$$\epsilon_1 = \frac{\mathrm{Im}(\hat{Y}_{\nu}^{\dagger}\hat{Y}_{\nu})_{21}^2}{(\hat{Y}_{\nu}^{\dagger}\hat{Y}_{\nu})_{11}} \frac{|\delta_N^* + \delta_N'|}{8\pi |\delta_N^* + \delta_N'|^2 + (\hat{Y}_{\nu}^{\dagger}\hat{Y}_{\nu})_{22}} , \quad \text{Pilaftsis, 1997}$$

Allowing for lower RHN masses, in way that avoid Gravitino problem

 $M_R \leq 10^8 GeV$ 

Two cases:  $8\pi |\delta_N^* + \delta_N'|^2 \gg (\hat{Y}_{\nu}^{\dagger} \hat{Y}_{\nu})_{22}$ Extreme resonant case  $\longleftarrow 8\pi |\delta_N^* + \delta_N'|^2 \sim (\hat{Y}_{\nu}^{\dagger} \hat{Y}_{\nu})_{22}$  Both heavy RHN equally contribute to CP asymmetry:

$$\epsilon_1 \simeq \epsilon_2 \simeq -\frac{1}{32\sqrt{2}\pi} \frac{(2-x^2)^2}{x(2+x^2)} \frac{\sqrt{\Delta m_{\rm atm}^2} M}{\langle h_u^0 \rangle^2 \sin^2 \beta} \frac{\sin 2r}{|\delta_N^* + \delta_N'|}$$
$$\frac{n_B}{s}\Big|_{x=1} \simeq 1.34 \cdot 10^{-15} \left(\frac{M}{10^8 \text{GeV}}\right) \frac{\sin 2r}{|\delta_N^* + \delta_N'|}$$

Gives the right order for;

$$\left|\delta_N^* + \delta_N'\right| \simeq 1.5 \cdot 10^{-5} \sin 2r \left(\frac{M}{10^8 \text{GeV}}\right)$$

Can be explained via symmetry...

# **Conclusion:**

- Predictive model of inverted hierarchical neutrinos was presented
- New symmetries play central role for predictions (Can be tested in experiments)

- $0.31 \le Sin^2(\theta_{12}) \le 0.7$ ,  $0 \le Cos(\delta) \le 0.7$  and  $\theta_{13} \ge 0.13$ , all which testable in forthcoming experiments.
- Model involves two quasi-degenerate RHNs → succesful Baryon Asymmetry via Resonant Leptogenesis