

# Path from the LHC to the New Standard Model via On-Shell Effective Theories

Nima Arkani-Hamed

Harvard - Philip Schuster, Natalia Toro

Berkeley - Jesse Thaler    Princeton - Lian-Tao Wang

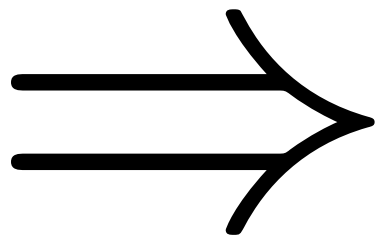
MIT - Bruce Knuteson    FNAL - Steve Mrenna

SLAC - Johan Alwall    CERN - Michelangelo Mangano

[www.marmoset-mc.net](http://www.marmoset-mc.net)

[hep-ph/0703088](https://arxiv.org/abs/hep-ph/0703088)

LHC  
Data



$\mathcal{L}$  ?

# Collective Hadron Collider Discovery Experience

$W/Z$

No undetermined  
parameters

$t$

One unknown  
parameter  
( $m_t$ )

In both cases, knew **exactly** what to look for..

# Characterizing New Physics

Only a few dynamical variables control the essential phenomenology of new physics at hadron colliders.

Characterizing new physics directly in terms of these variables permits a simple, accurate parametrization of almost **any** model of new physics —

an On-Shell Effective Theory.

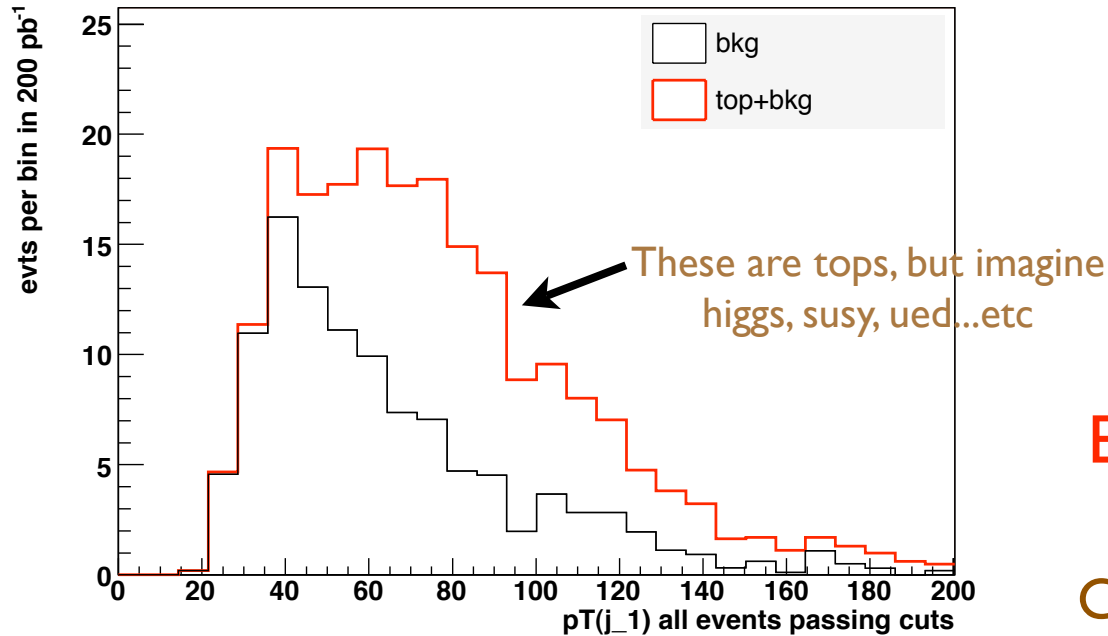
Only with this information can we tackle  
the fundamental questions:

What new physical principles are  
being revealed at the TeV scale?

Is nature supersymmetric?

Is the electroweak scale natural?

Suppose we have “interesting” events clearly inconsistent with the SM



How do we characterize “interesting” events?

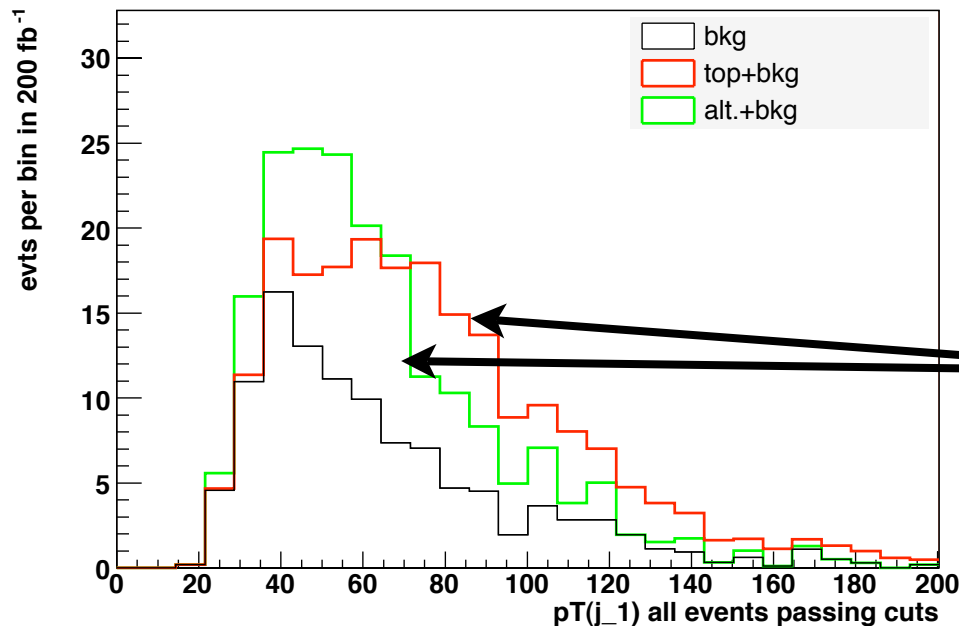
The “composite sketch” approach

Exploit an effective description:

Throw out irrelevant info

Characterize topology and kinematics

Close connection to Lagrangian, but far simpler



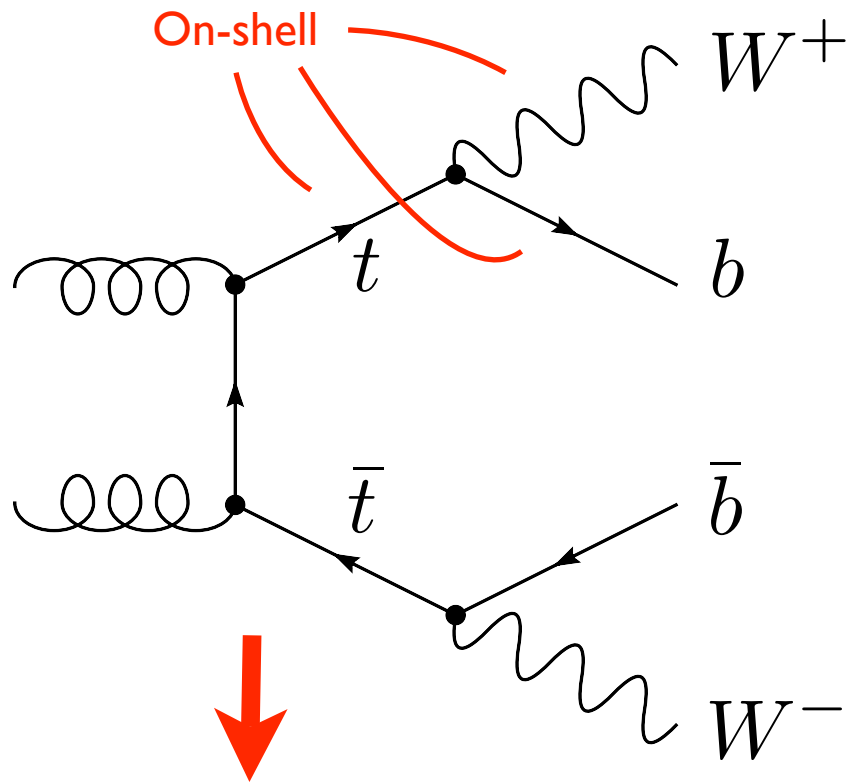
Use On-Shell Effective Theories (OSET)

Fast, accurate, transparent characterization and discrimination among possibilities

Ask good questions about the data early in the analysis process

# What's an OSET?

# The Basic Idea:



**Example: Top Quark**  
Masses, Rates, and Topology  
vs. **Amplitudes**

**Dominant Top Properties:**

$$\sigma(gg \rightarrow t\bar{t})$$

$$\text{Br}(t \rightarrow bW)$$

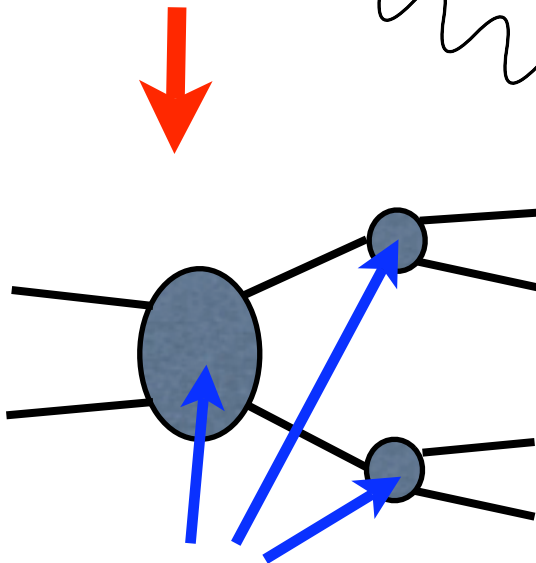
$$m_t, m_W, m_b$$

**Detailed Top Properties:**

$$d\sigma/d\hat{t}$$

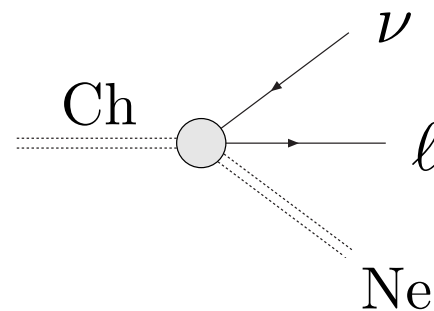
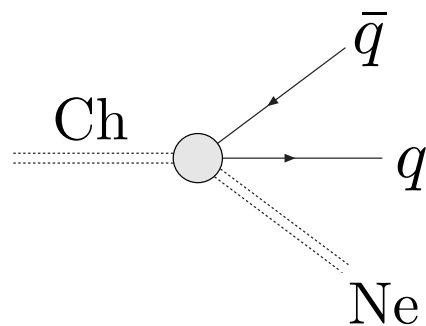
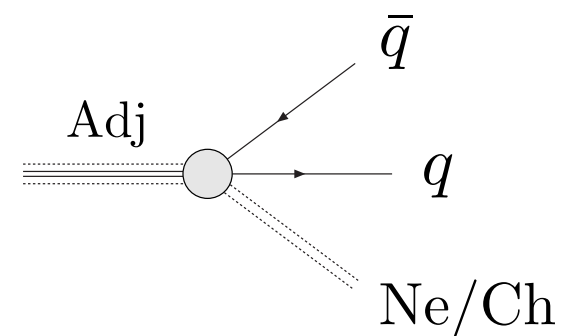
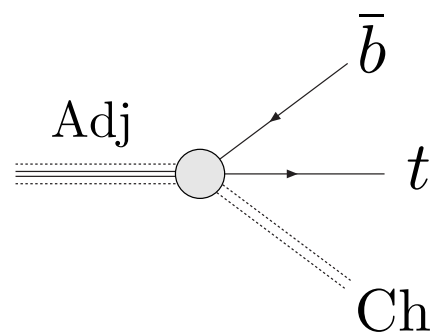
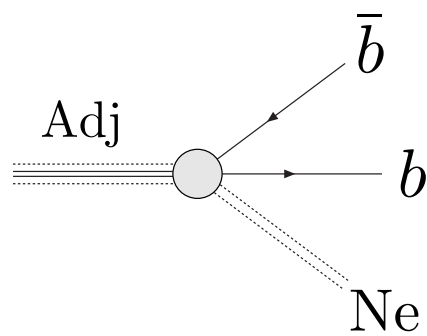
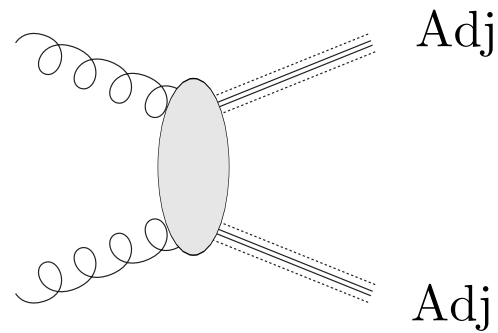
$W$  helicity

$t$  charge



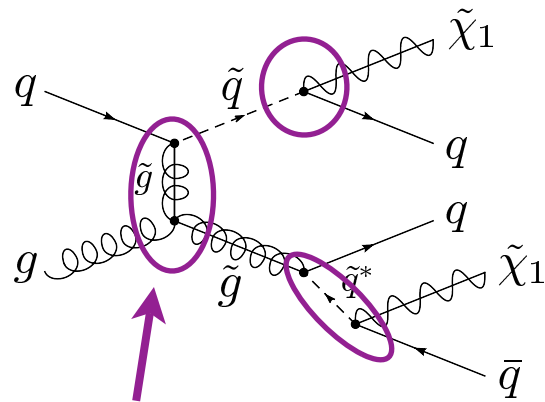
Simple rules given for these parts

# What's an OSET? Complicated Case



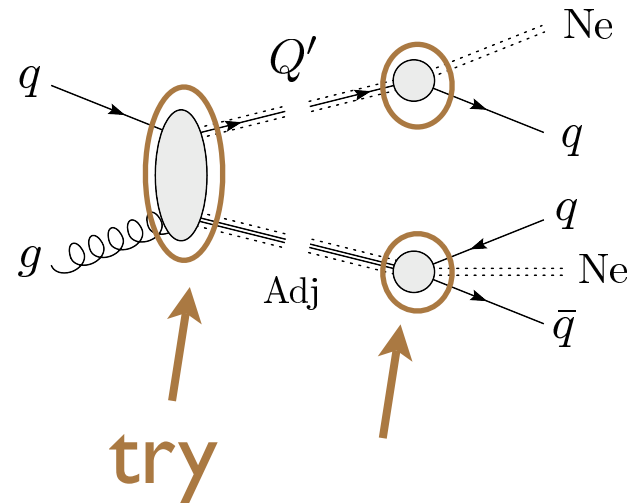


# OSET Rules: Modeling Production



$$|\mathcal{M}|^2 = \frac{\pi\alpha_s^2}{s^2} \left( \frac{4m_g^2 - t}{9s} + \frac{[(m_g^2 - t)s + 2m_g^2(m_q^2 - t)]}{(t - m_g^2)^2} \right. \\ \left. + u + st + su + tu \text{ channels} \right)$$

- Monte Carlo
- compare to data

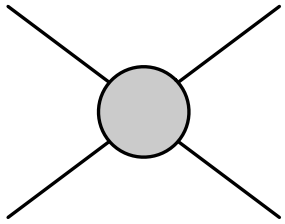


try

$$|\mathcal{M}|^2 = \text{const.}$$

(do kinematics right!)

# 2→2 Folk Theorem



$$|\mathcal{M}|^2 = f(s, \xi) \quad \xi = \beta_{34} \cos \theta = \frac{\hat{t} - \hat{u}}{\hat{s}} = \frac{p_z}{\sqrt{s}}$$

$$\frac{d\sigma}{d\hat{t}} = \int \text{Parton Luminosity} \times \text{Phase Space (Threshold)} \times |\mathcal{M}|^2$$

Well Approximated by Constant!

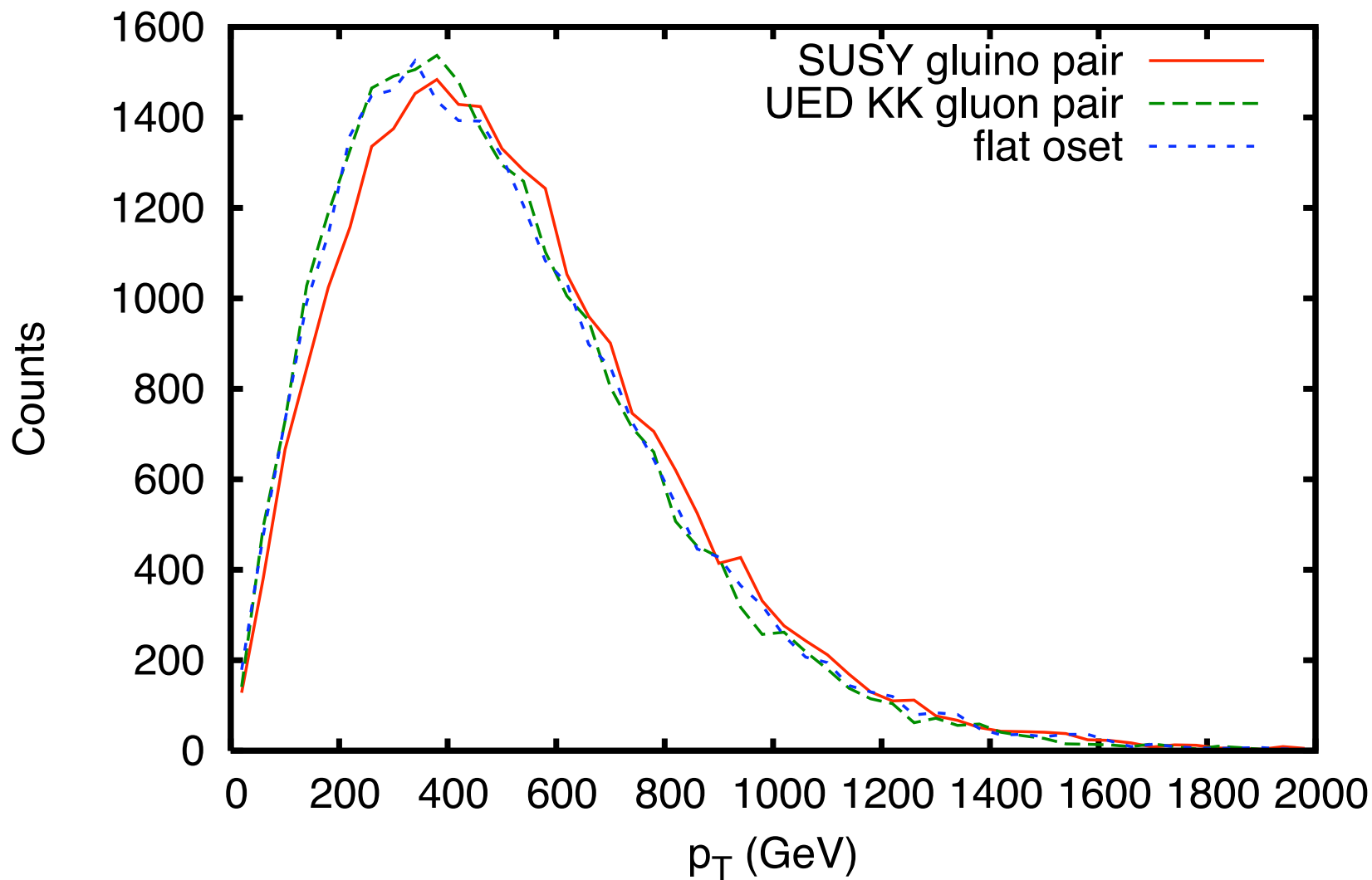
Cross Sections **Dominated by Thresholds!**

# Modeling Production

Just an example.

Not measurable in this case

Glauino  $p_T$



# Modeling Production

Rapidly falling PDFs

S-wave often dominates



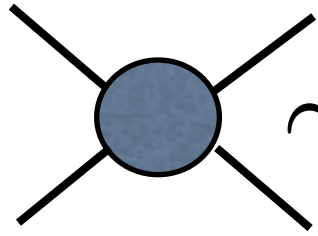
Transverse momentum and rapidity shapes well modeled  
by constant Matrix Element in most cases!

This is not very surprising...

What about extreme kinematics?

What about corrections?

# Adding Corrections:



Full ME

$$\sim |M|^2 = \sum_{p,q} C_{pq} X^p \xi^q$$

Contains CM  
threshold behavior

Contains CM  
angular and threshold  
structure

General Expansion

$$X \equiv \frac{s}{s_{th}}$$

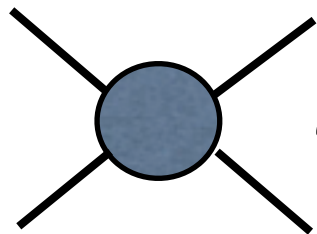
$$\xi \equiv \beta_{34} \cos(\theta)$$

Parton Luminosity

CM rapidity integrated

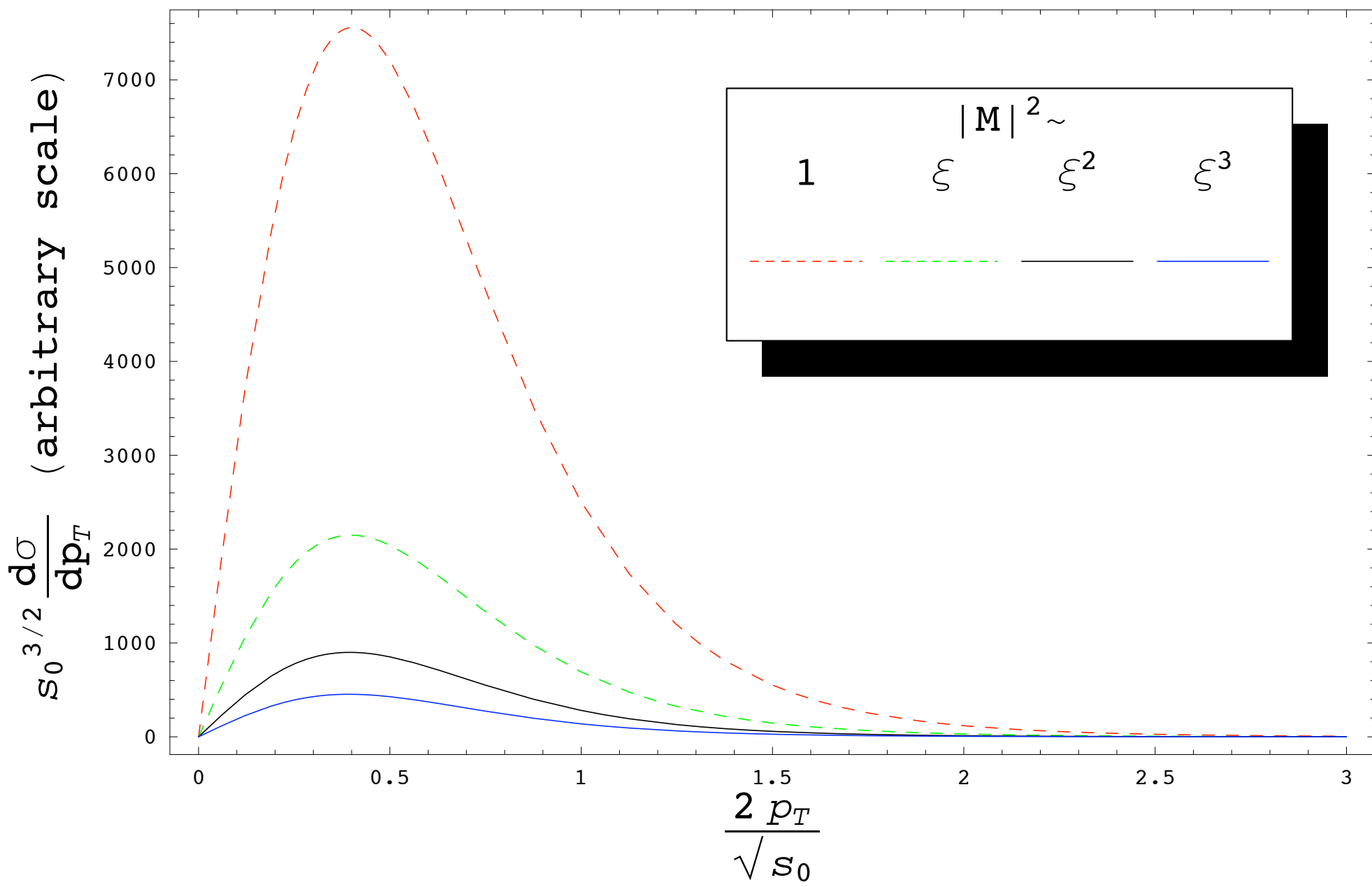
$$\rho_{PDF} \sim \tau^a \log\left(\frac{1}{\tau}\right)$$

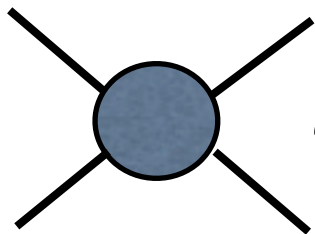
How much matrix element  
structure survives  
PDF convolution?



$$\sim |M|^2 \sim X^q \xi^p$$

Threshold Suppression

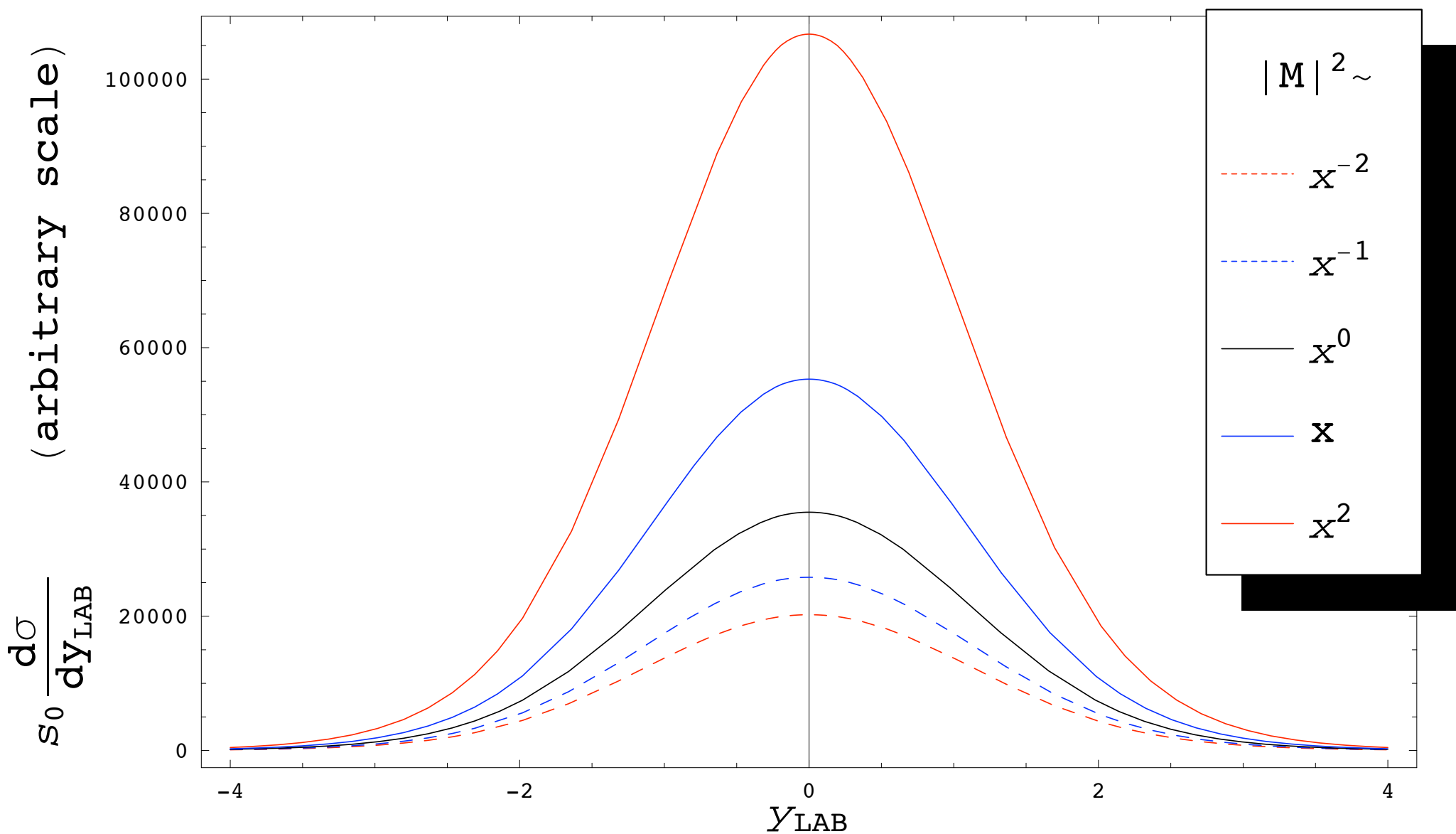


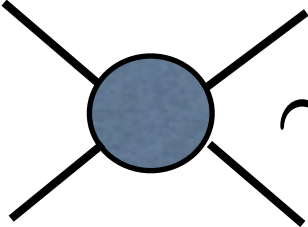


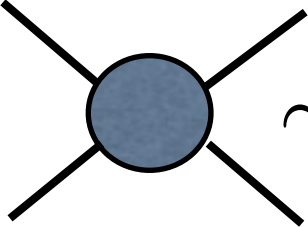

$$\sim |M|^2 \sim X^q \xi^p$$

Threshold Suppression

gg Initial States

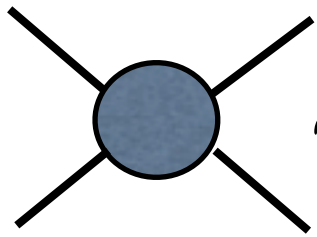



$$\sim |M|^2 = \sum_{p,q} C_{pq} X^p \xi^q$$


$$\sim |M|^2 \sim X^q \xi^p$$

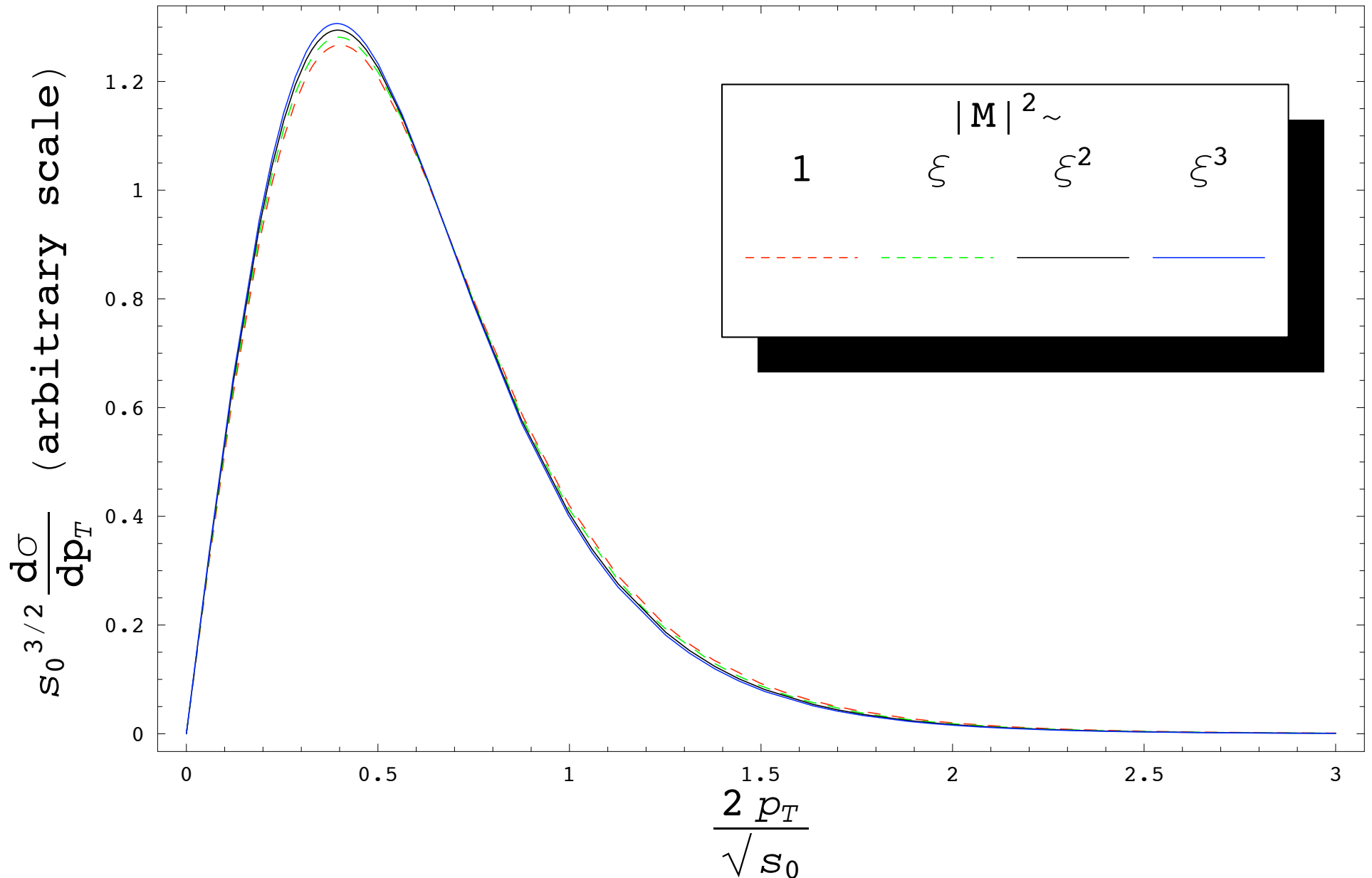
Only near-threshold behavior survives

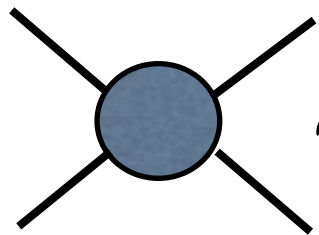




$$\sim |M|^2 \sim X^q \xi^p$$

$\xi$ -Independence of  
Transverse Shape!

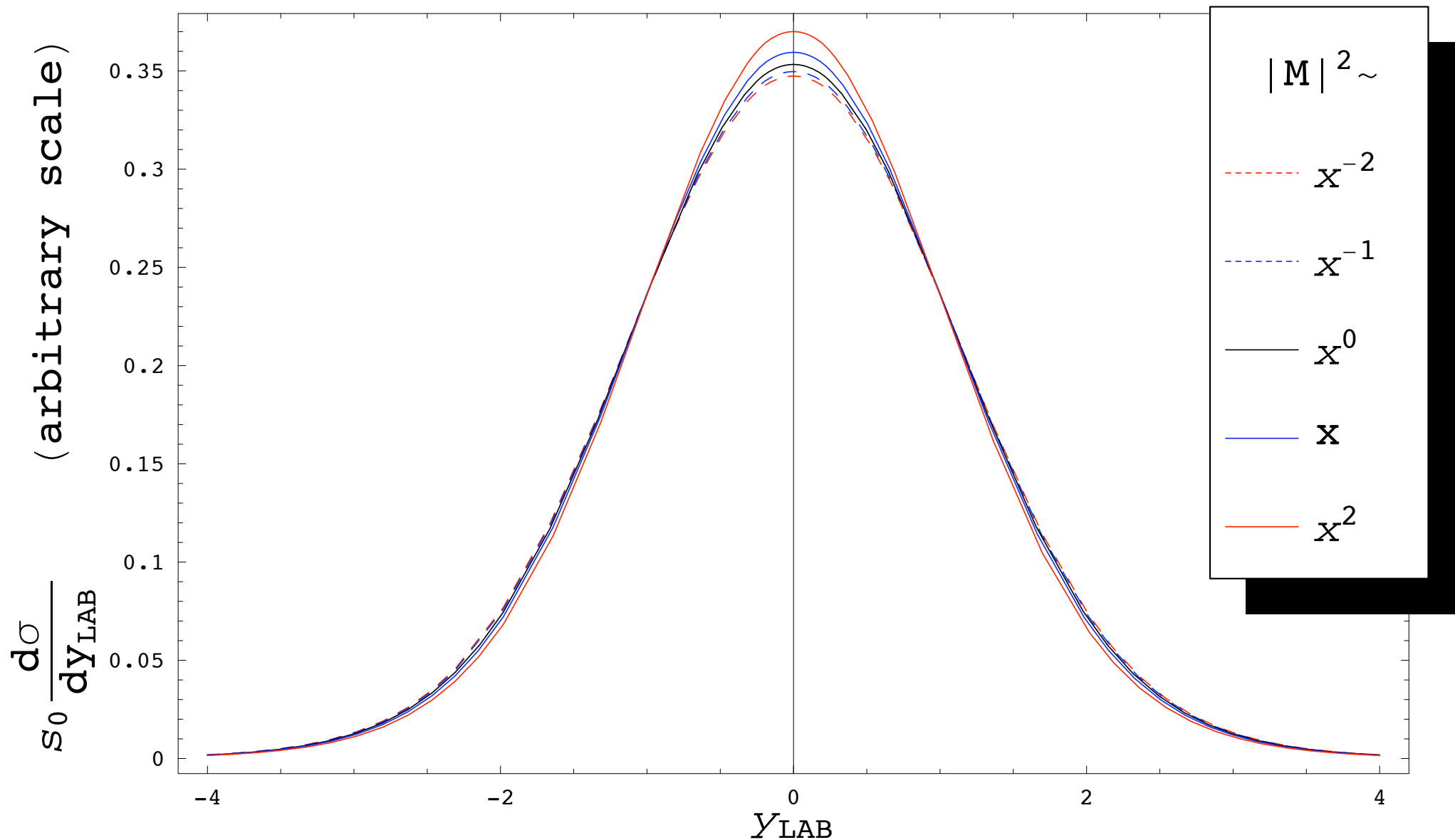




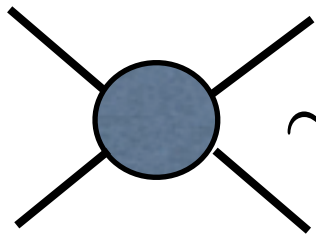
$$\sim |M|^2 \sim X_{S_0}^{qgp}$$

**X-Independence of rapidity Shape!**

gg Initial States



# Shape Invariance



$$\sim |M|^2 \sim X^q \xi^p$$

PDF  $E_{cm}$  and  $y_{cm}$   
homogeneity  
properties

Inclusive  $p_T$  shape invariant under:

$$|M|^2 \rightarrow |M|^2 \xi^m$$

Inclusive  $y_{lab}$  shape invariant under:

$$|M|^2 \rightarrow |M|^2 X^n$$

Simple “Universal” corrections to constant ME!

Messy collider environment turned to our advantage

Correct PDFs necessary

Caveats: Large final state mass asymmetry requires care

Transverse momentum-rapidity correlations not included beyond phase space

# Defining an OSET

## Production:

2 → 1 Use Breit Wigner

2 → 2

$$|\mathcal{M}|^2 = A + B \left( 1 - \frac{s_{\text{thresh}}}{s} \right)$$

or

$$|\mathcal{M}|^2 = A + B \left( \frac{s}{s_{\text{thresh}}} - 1 \right)$$

Usually dominates

“Normal” Behavior

“Contact” Operator Behavior

Dominant  $\xi$  correction  
can be included  
(not usually necessary)

2 → 3 Use “standard” modes with OSET decay scheme

## Decay:

- Polynomial in  $\cos \theta$  : rank determined by spins, coefficients by masses. Spin correlations can be included.
- Single-object lab-frame distributions well approximated by phase space decays.

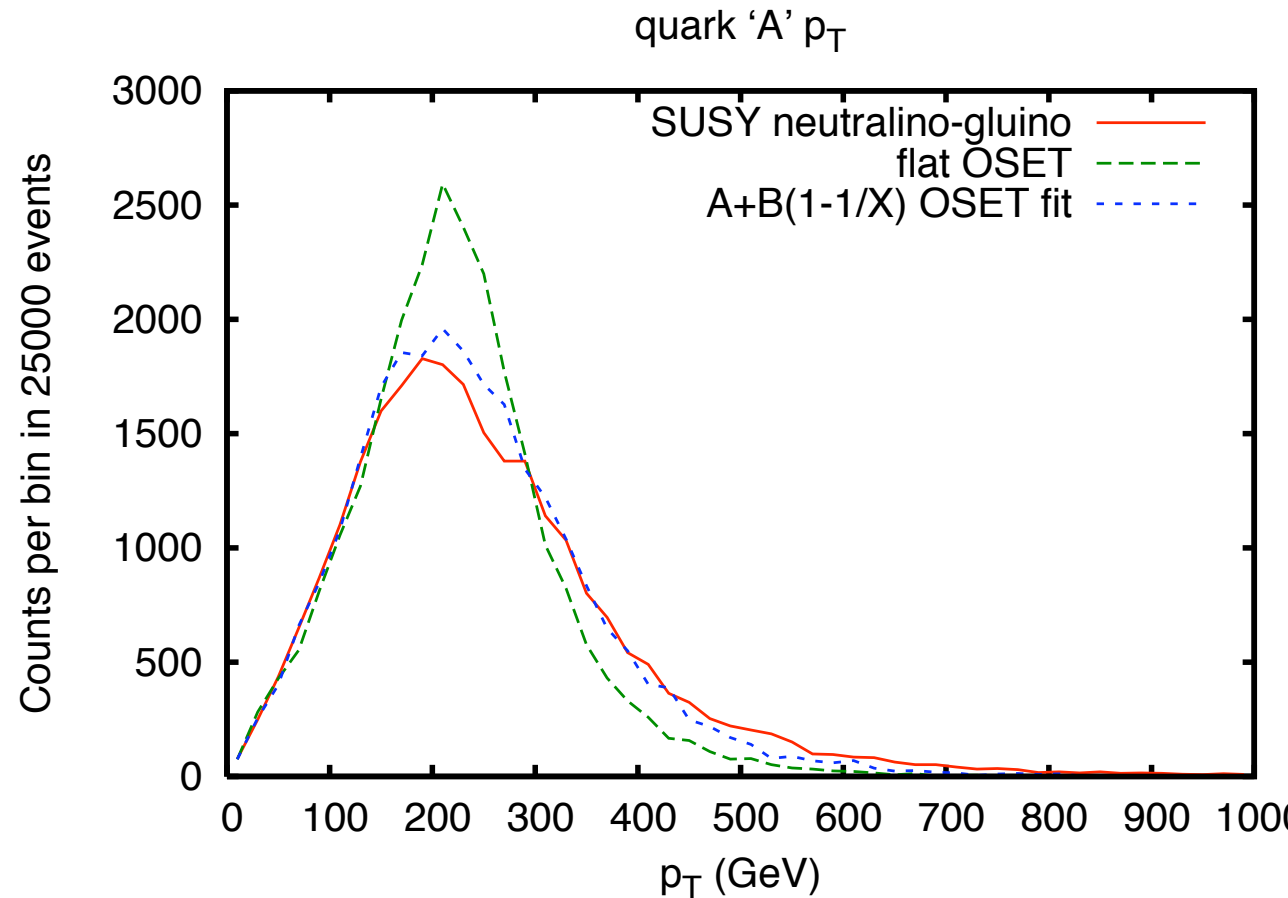
See: [hep-ph/0703088](https://arxiv.org/abs/hep-ph/0703088) for detail...

# “Normal” Behavior

$$m_{t\text{-chan}} \sim m_{\tilde{g}}$$

$p$ -wave  $\rightarrow$   
suppressed  
near threshold.

$$\propto \beta^2$$



$$|\mathcal{M}|^2 = A + B \left( 1 - \frac{s_{\text{thresh}}}{s} \right)$$

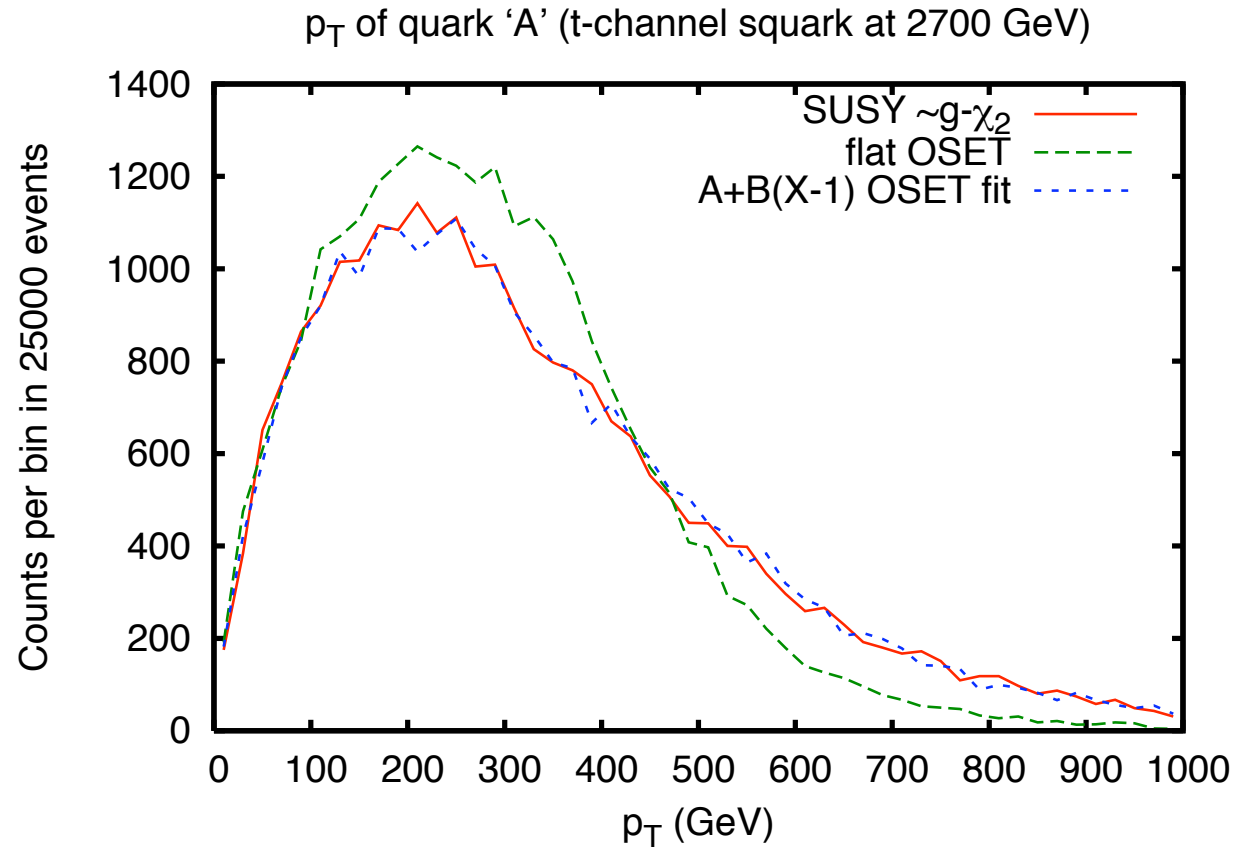
# “Contact” Operator Behavior

$$m_{t\text{-chan}} \gg m_{\tilde{g}}$$

contact interaction

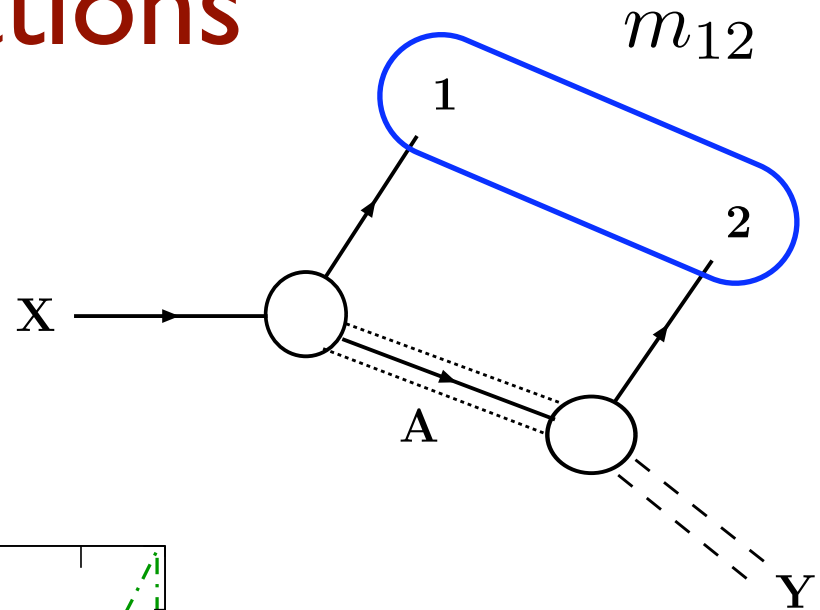
$$|\mathcal{M}|^2 \propto s^2$$

up to  $\sqrt{s} \sim m_{t\text{-chan}}$

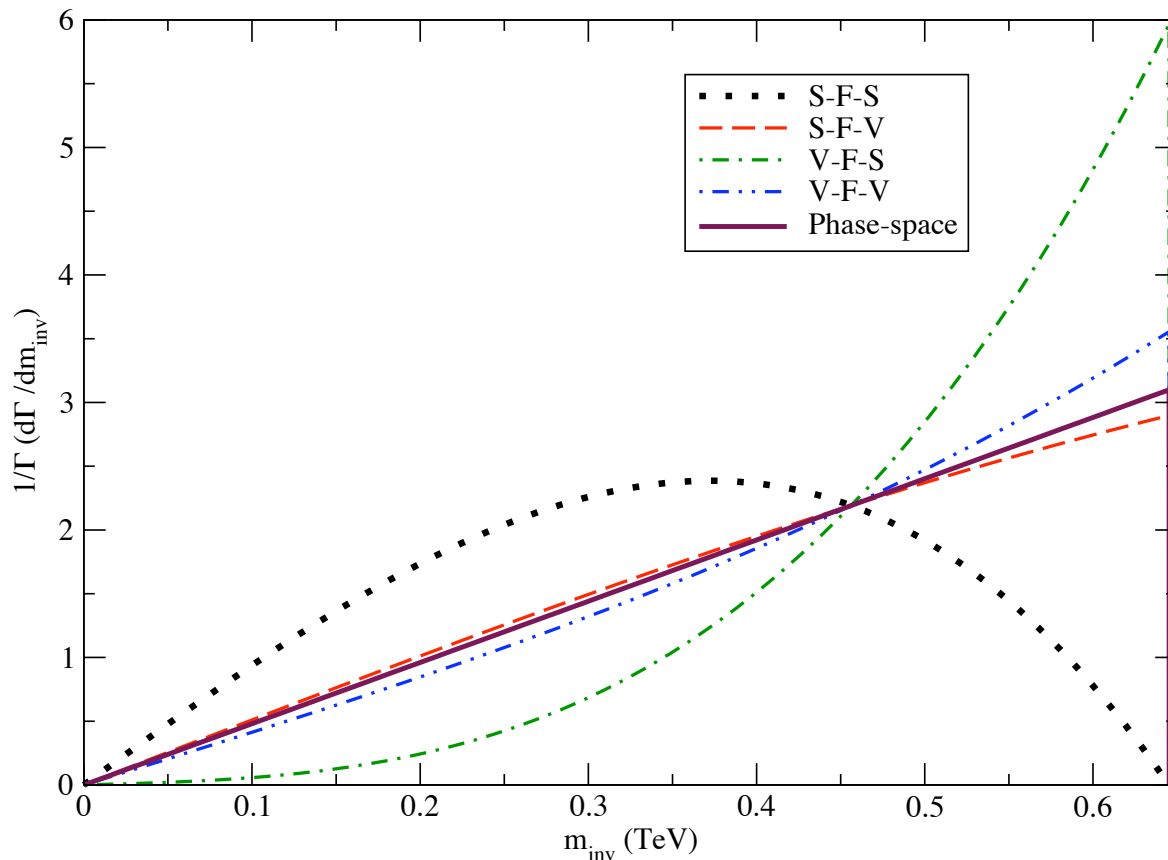


$$|\mathcal{M}|^2 = A + B \left( \frac{s}{s_{\text{thresh}}} - 1 \right)$$

# Spin/Helicity Correlations in Decay?



## Di-object Distributions



Angular correlations  
in cascade decays  
are known to be  
important!

See J. Alwall's talk about  
MadGraph/OSETs  
at Princeton MC4BSM

# Why Monte-Carlo an OSET?

- MODEL: **Fast, sufficiently accurate** event generation for your favorite model
- DATA: Check consistency of **partial interpretations** of candidate new-physics signals



# MARMOSET

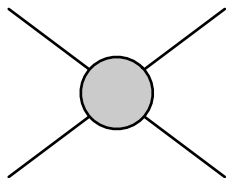
## Mass and Rate Modeling in On-Shell Effective Theories

- Masses and  $SU(3)$  and  $U(1)$  quantum numbers of new particles, their production and decay modes **fully specify model**
- An OSET implies many topologies, and Monte Carlo is generated separately for each
- Topologies can be combined by weighting according to hadronic cross-sections and decay branching fractions

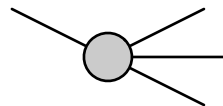
“offline” reweighting of MC  $\longrightarrow$  **fast!**

scan models by simulating coarsely spaced mass points

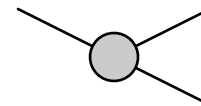
# OSET



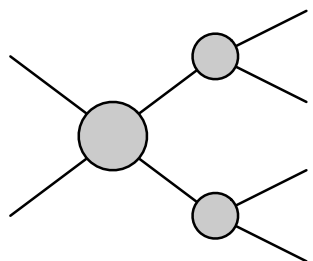
Production Mode A



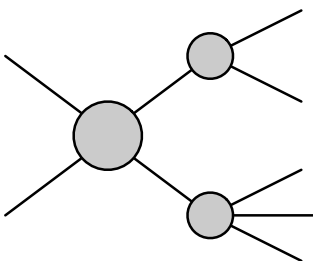
Decay Mode 1



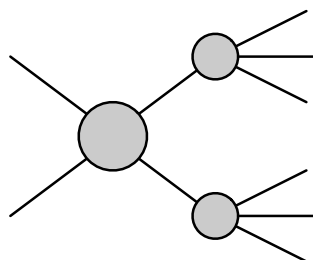
Decay Mode 2



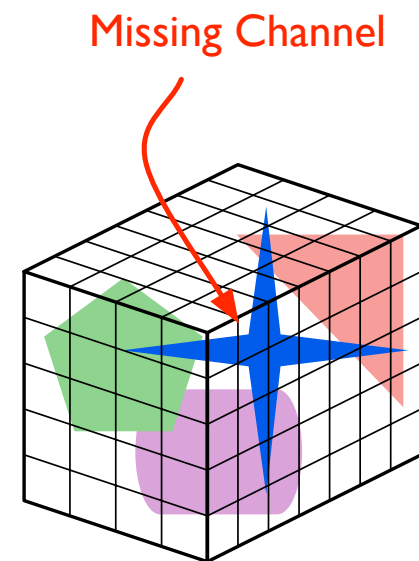
$$= \text{mc\_A11} \times \sigma_A \times \text{Br}_1 \times \text{Br}_1$$



$$= \text{mc\_A12} \times 2 \times \sigma_A \times \text{Br}_1 \times \text{Br}_2$$



$$= \text{mc\_A22} \times \sigma_A \times \text{Br}_2 \times \text{Br}_2$$



**LHC**  
**Signatures**

Correlations among final  
states

OSETs have predictive  
power!

# MARMOSET: Simple Monte Carlo for Any Model

Coarse Mass  
Scanning

Rate Scanning by  
Reweighting Existing  
Monte Carlo

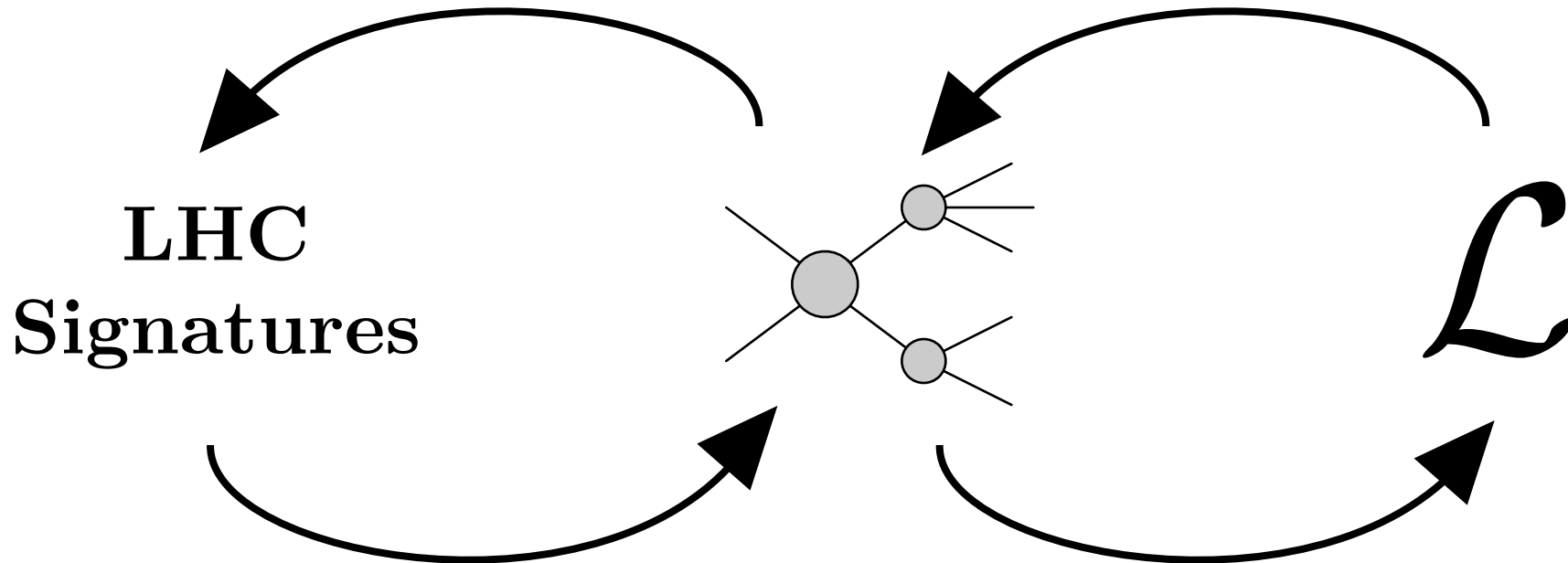
vs.

## Lagrangian Simulation

Intricate dependence of  $\sigma, Br$  on many Lagrangian parameters demands fine scanning of all of them.

OSET concisely describes  
many topologies with  
correlated rates

Lagrangian  
consistency  
constrains OSET



Simple observables and  
correlations between  
them constrain OSET

OSET constrains and  
motivates new  
physics Lagrangian

Use OSET MC like a “think-pad”

# The Michigan Black Box

1st and 2nd LHC Olympics

$\tilde{q}_{1,2}$	—————	1.8–2.1 TeV
$\tilde{b}_R$	—————	2.0 TeV
$\tilde{t}_R, \tilde{q}_{3L}$	—————	1.4 TeV
$\widetilde{W}$	—————	1.0 TeV

A SUSY model

(nothing like CMSSM)

$\widetilde{B}$	—————	480 GeV
$\tilde{g}$	—————	450 GeV

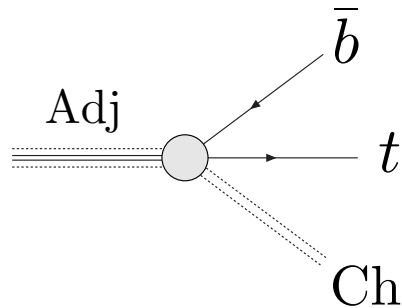
Signatures are quite simple...

visible in many searches

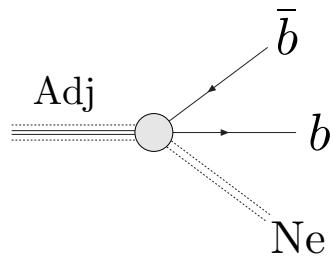
how do you combine results?  
what params reproduce them?

( $\tan \beta = 19$ )

# Where do you see it?



First discovery in 1lepton+jets+MET search?  
2l+jets+MET?

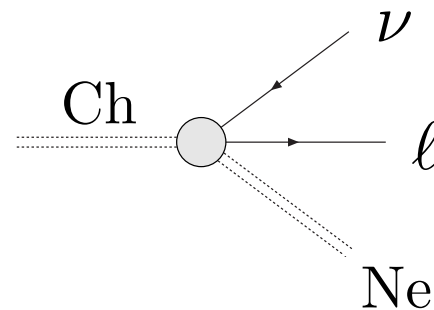
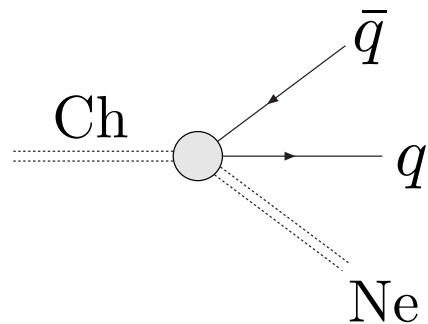
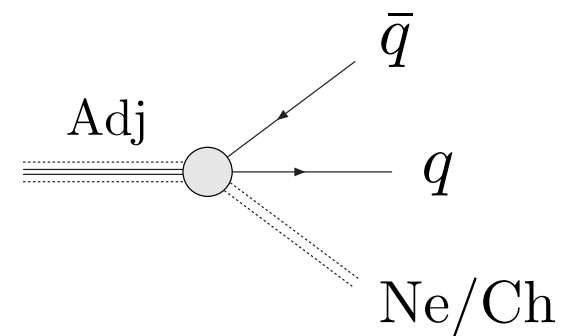
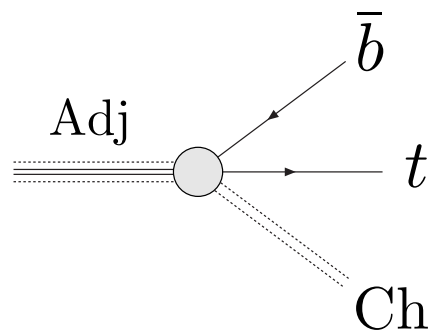
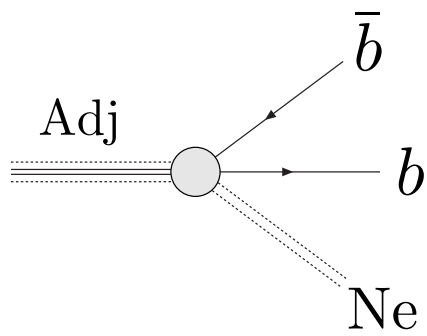
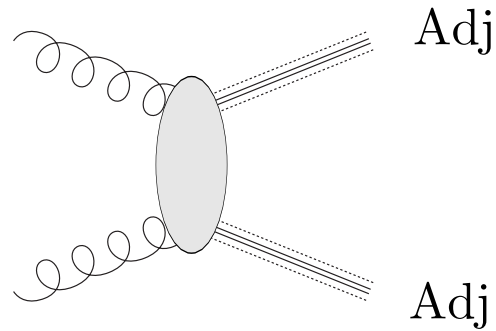


Signal in  $b\bar{b}$ +jets+MET,  
jets+MET

(with unrealistic detector, we've seen it after TDR low-mass SUSY analysis cuts in inclusive jets+MET — lots to see after  $1\text{fb}^{-1}$ )

# The Michigan Black Box

## Process/Kinematic Characterization



# Parameterizing Rates in MARMOSSET

Process	fit rate	Actual rate ( $1fb^{-1}$ )
$N(gg \rightarrow Adj Adj)$	$30600 \pm 1300$	27950
$Br(Adj \rightarrow t\bar{b}Ch^- \text{ or } c.c.)$	$0.77 \pm 0.01$	0.77
$Br(Adj \rightarrow b\bar{b}Ne)$	$0.21 \pm 0.01$	0.22
$Br(Adj \rightarrow q\bar{q}Ne)$	$0.02 \pm 0.01$	0.01
$Br(Ch \rightarrow q\bar{q}'Ne)$	$0.44 \pm 0.11$	0.6
$Br(Ch \rightarrow e/\mu\bar{\nu}Ne)$	$0.55 \pm 0.11$	0.4

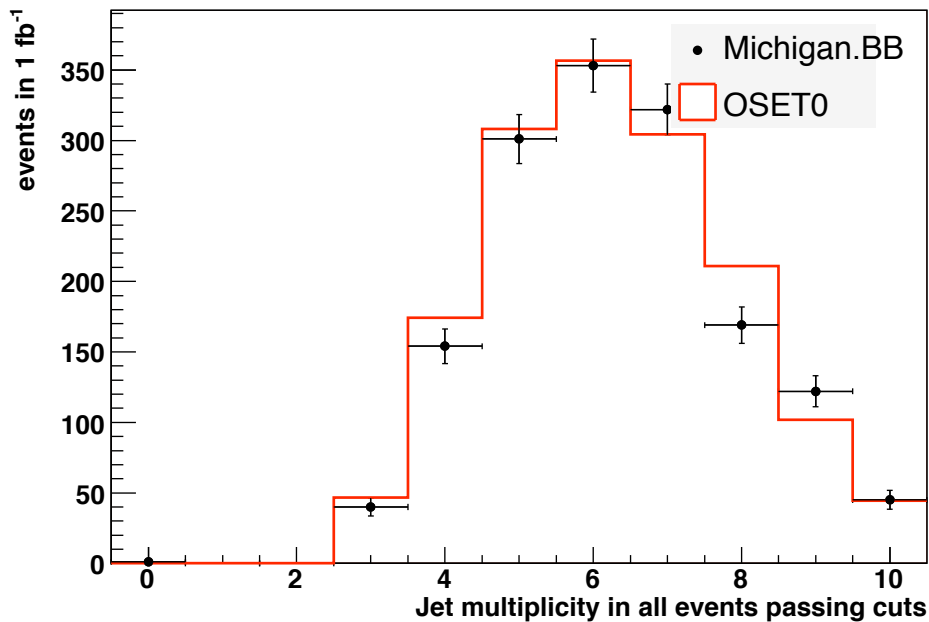


# Example Distributions (1 fb<sup>-1</sup>)

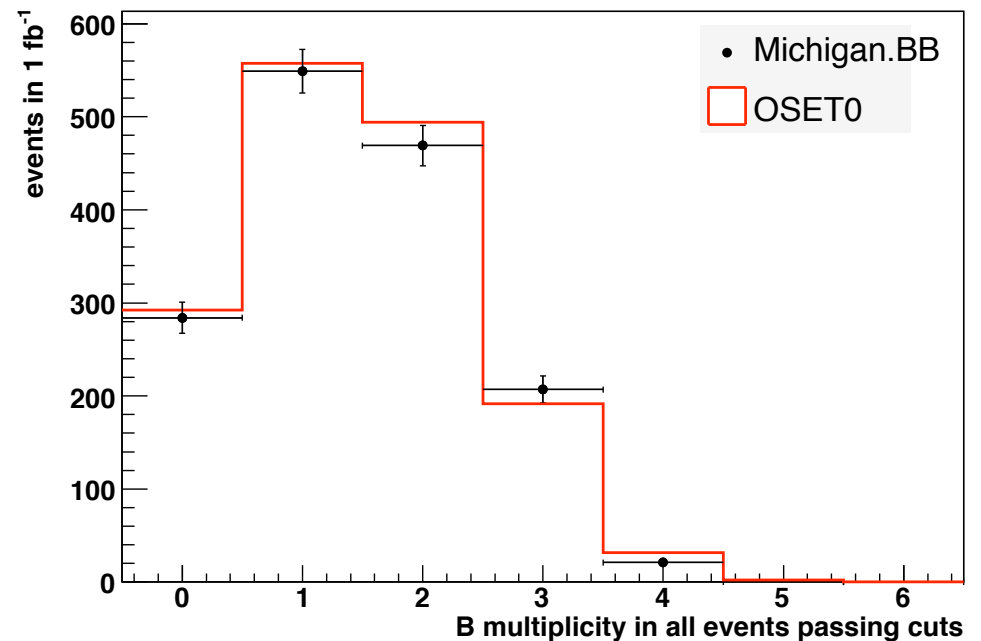
jet multiplicity

b-tag multiplicity

Jet multiplicity in all events passing cuts



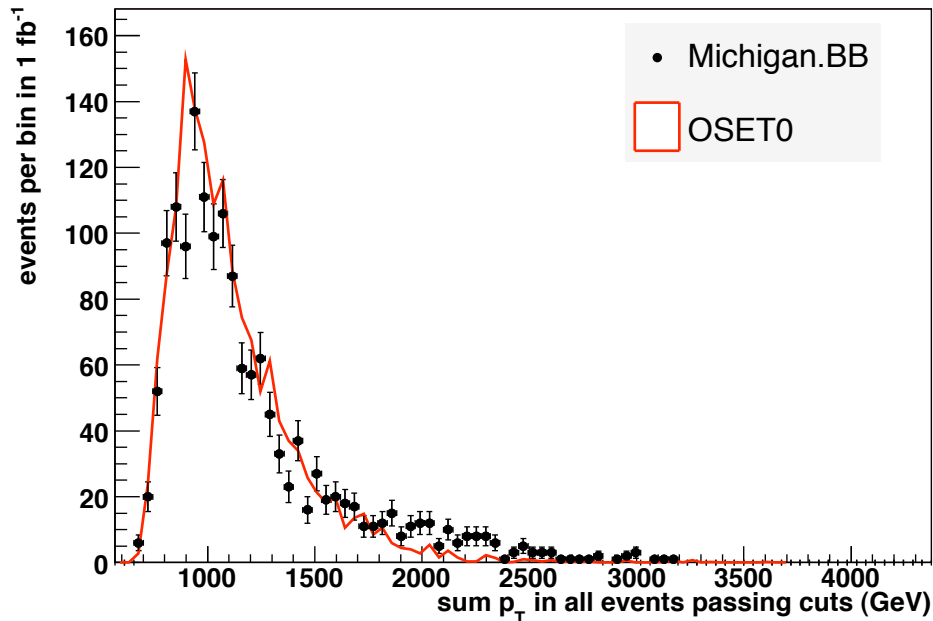
B multiplicity in all events passing cuts



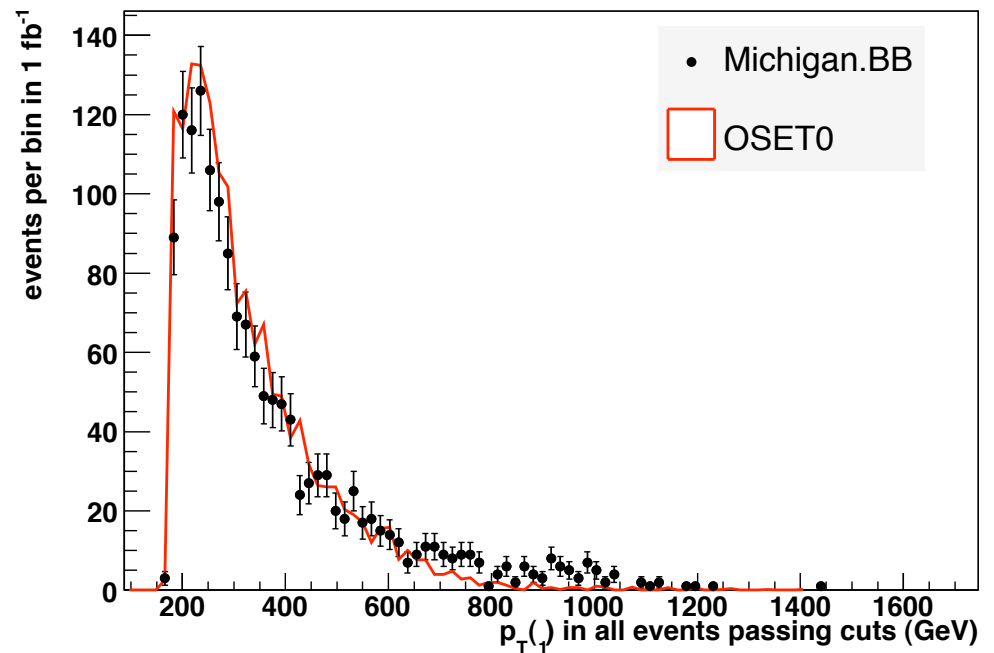
# Example Distributions (1 fb<sup>-1</sup>)

$$m_{\text{eff}} = \sum_i p_T^i \quad p_T(j_1)$$

sum  $p_T$  in all events passing cuts (GeV)

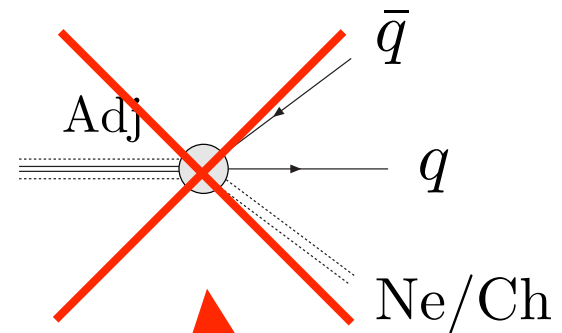
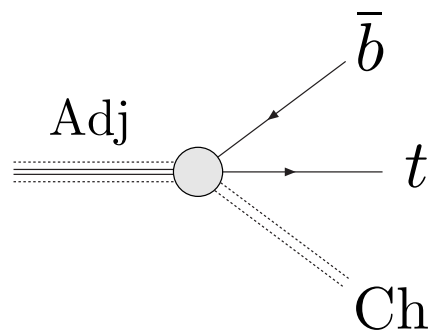
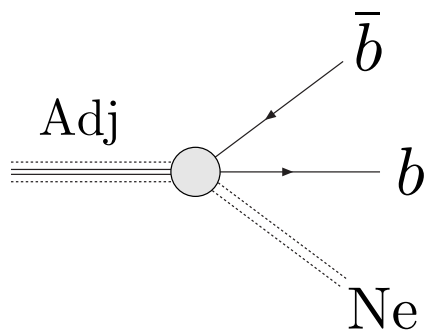
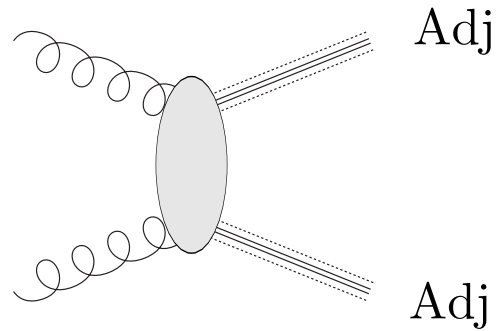


$p_{T(j_1)}$  in all events passing cuts (GeV)

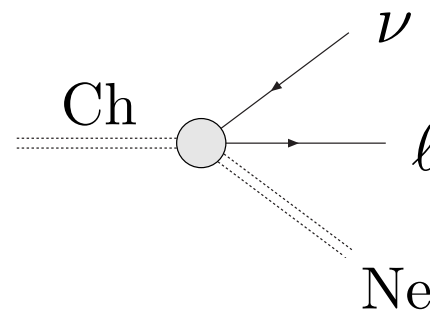
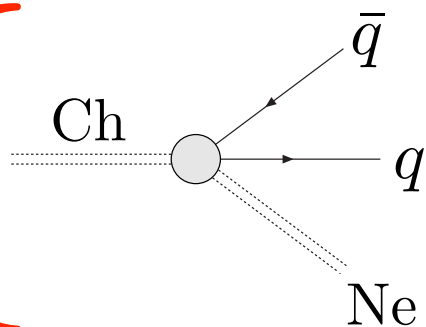


# The Michigan Black Box

Striking features **lead to model predictions**

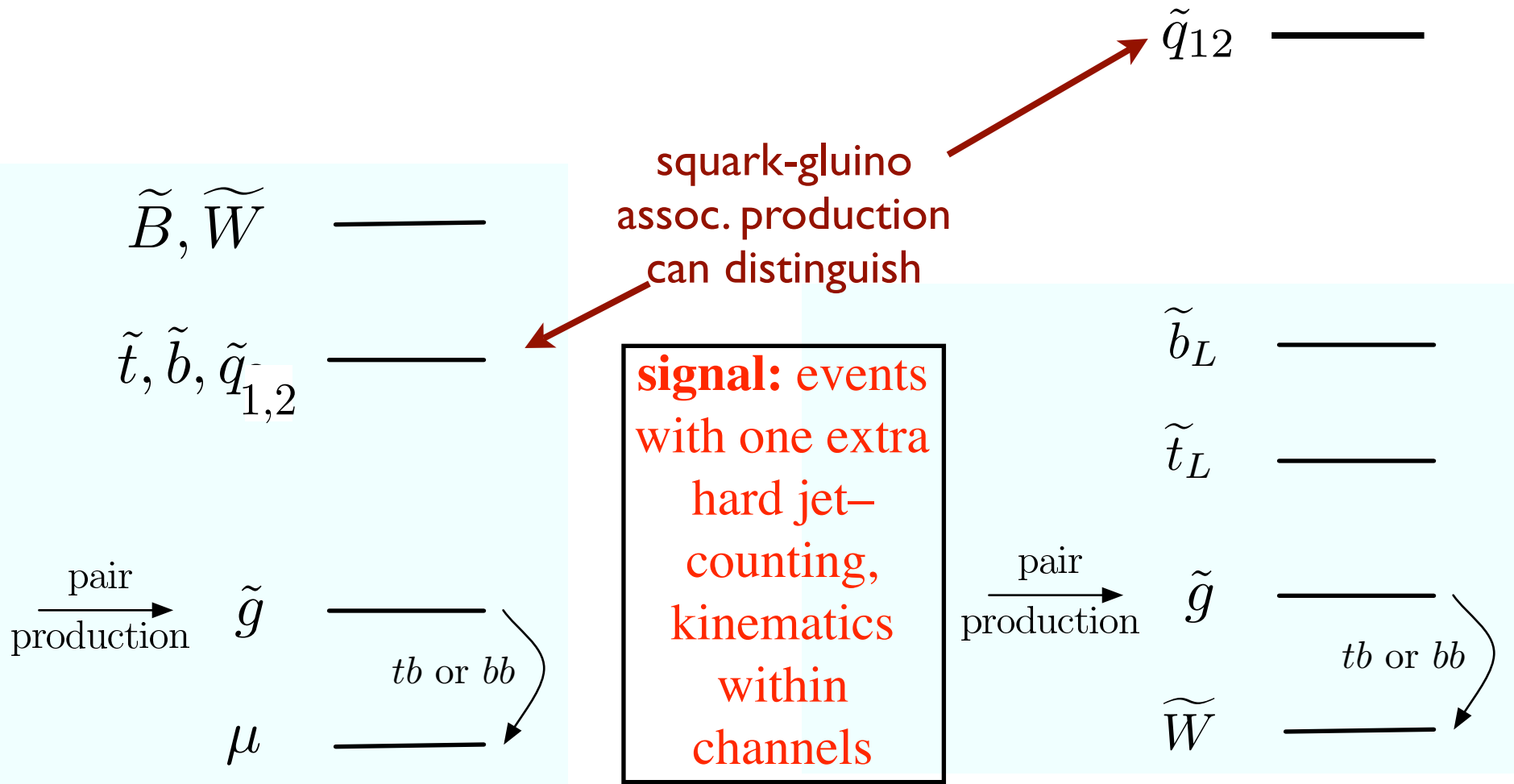


no long cascades  
↓  
**few light neutralinos (not CMSSM)**



**Model must explain its absence!**

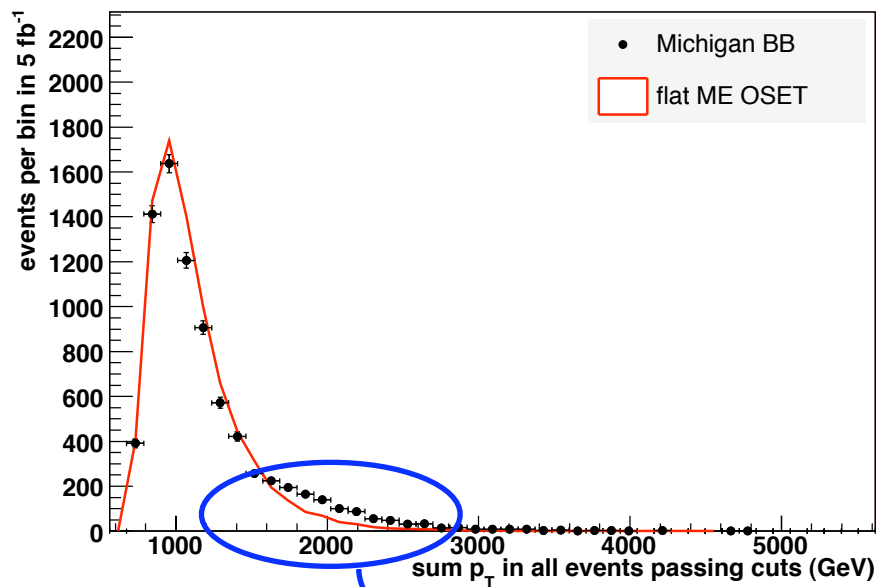
# Two SUSY Scenarios



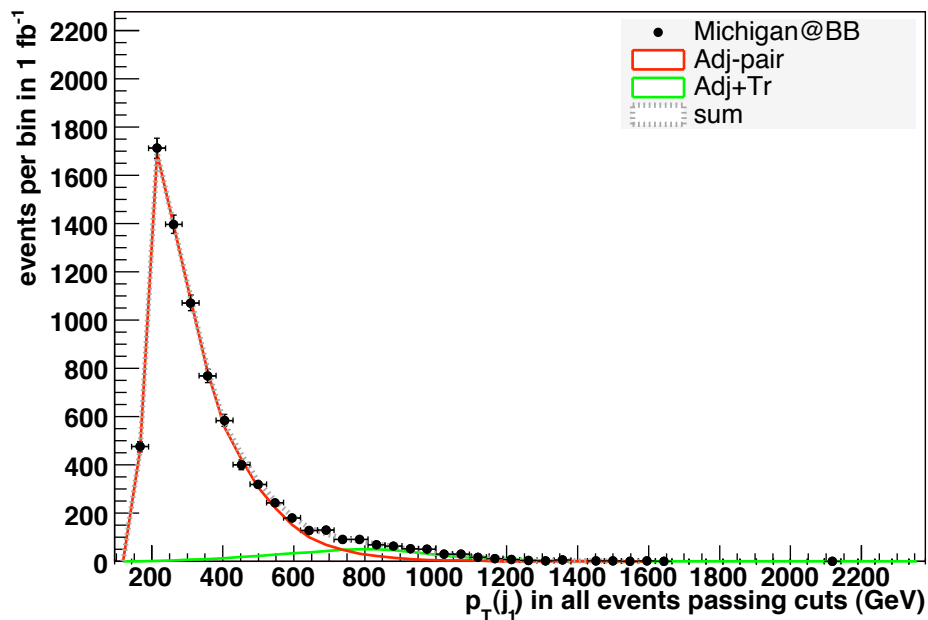
2j channel suppressed  
by small couplings to Higgsino

or

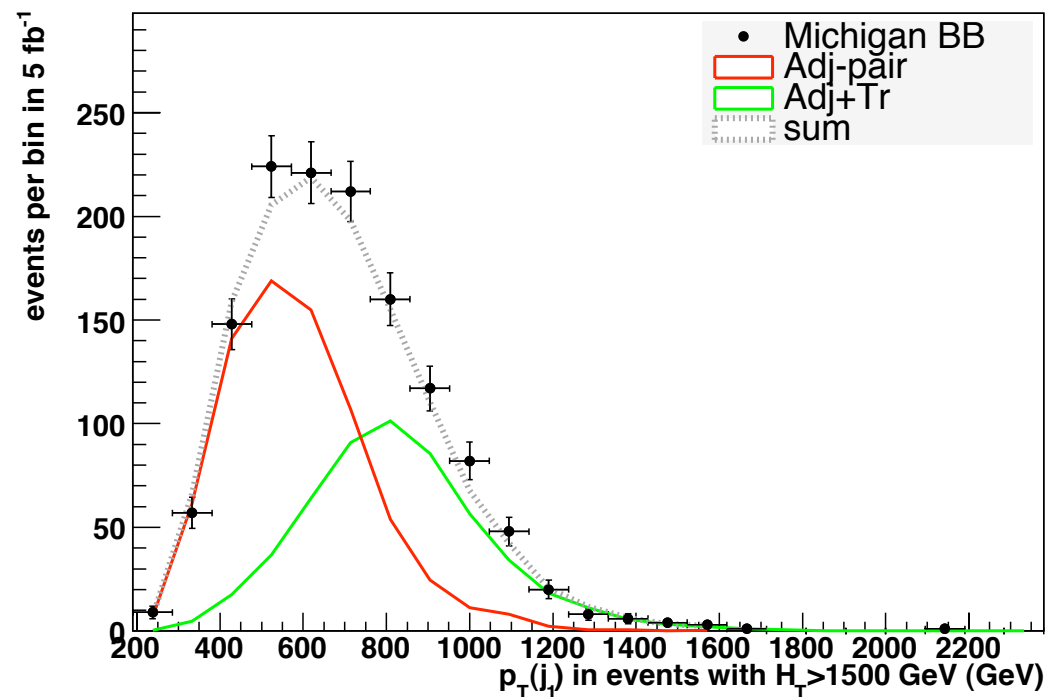
2j channel suppressed  
by high squark mass



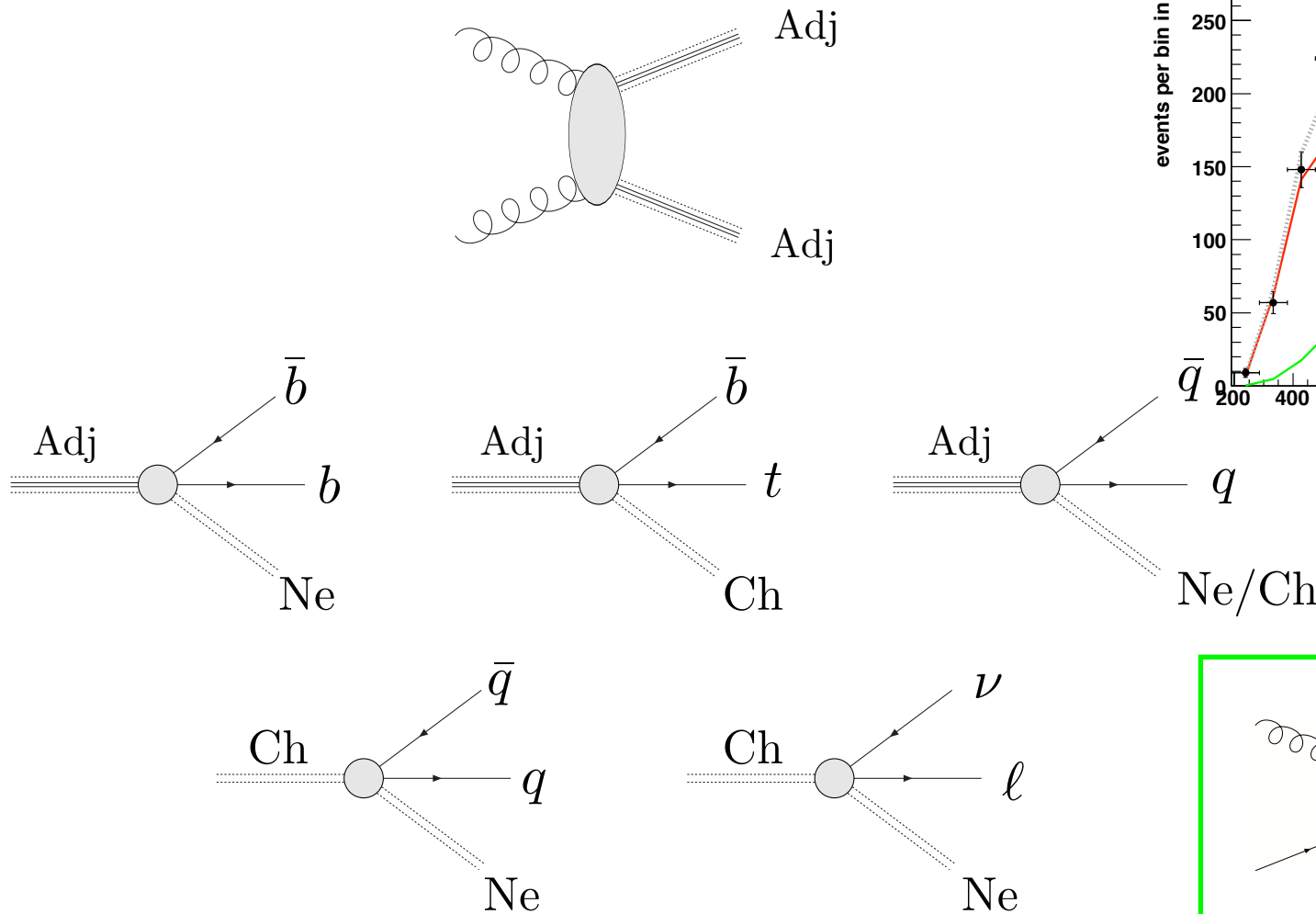
## $p_T$ of hardest jet in all events



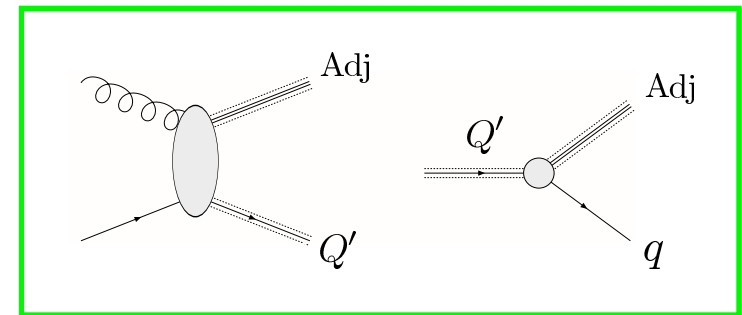
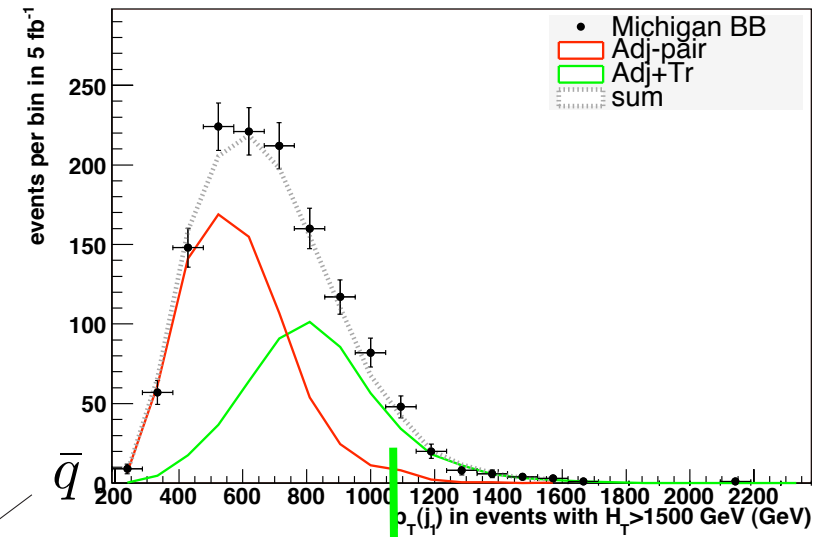
## $p_T$ of hardest jet in events w/ $H_T > 1500$ GeV



# The Michigan Black Box

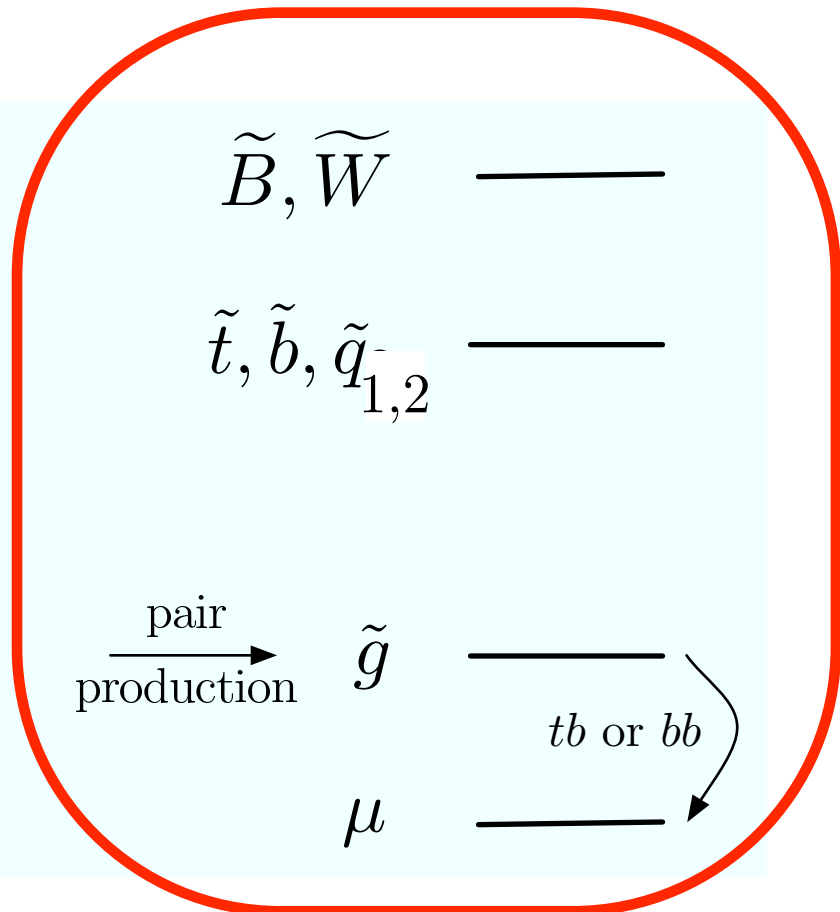


$p_T$  of hardest jet in events w/  $H_T > 1500 \text{ GeV}$

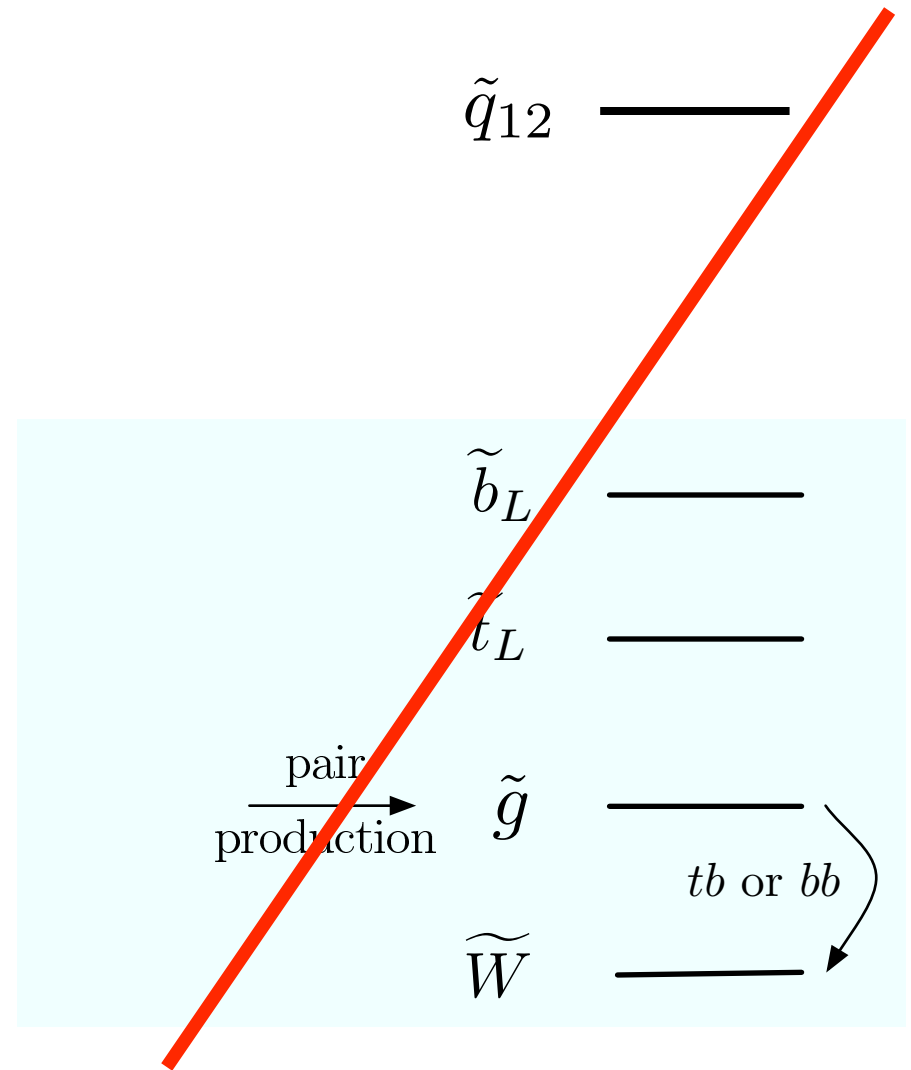


# Two SUSY Scenarios

**only this model reproduces associated production!**



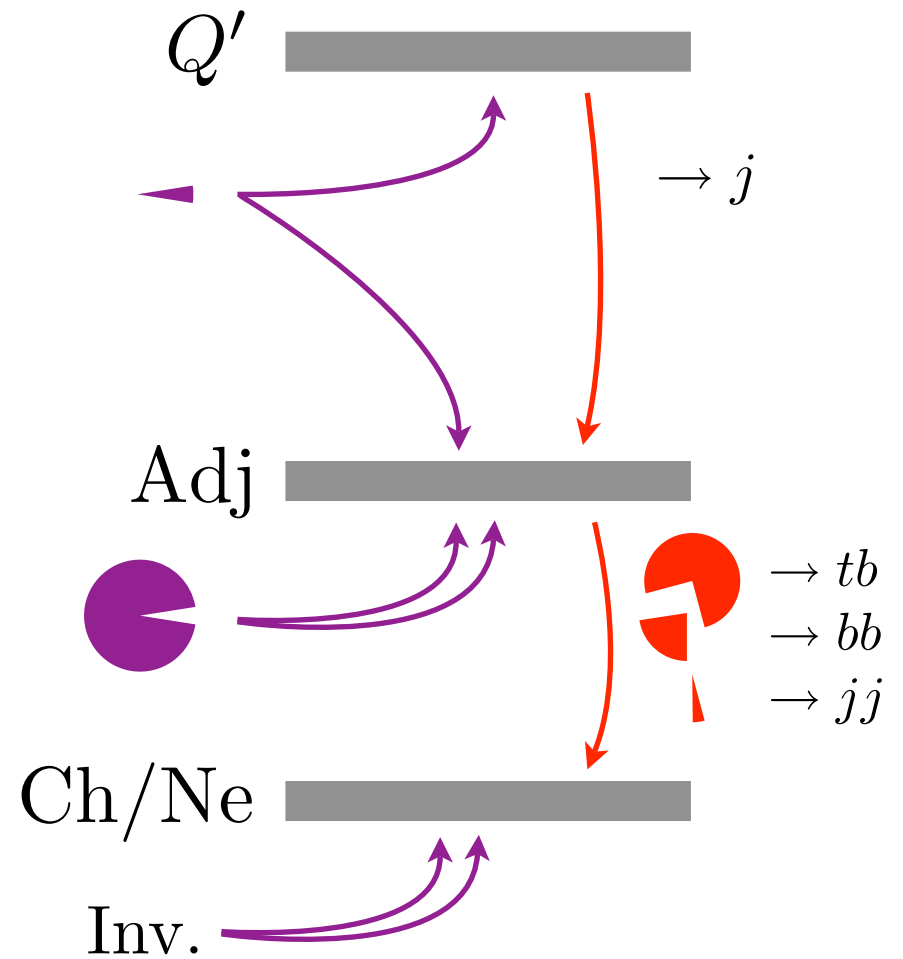
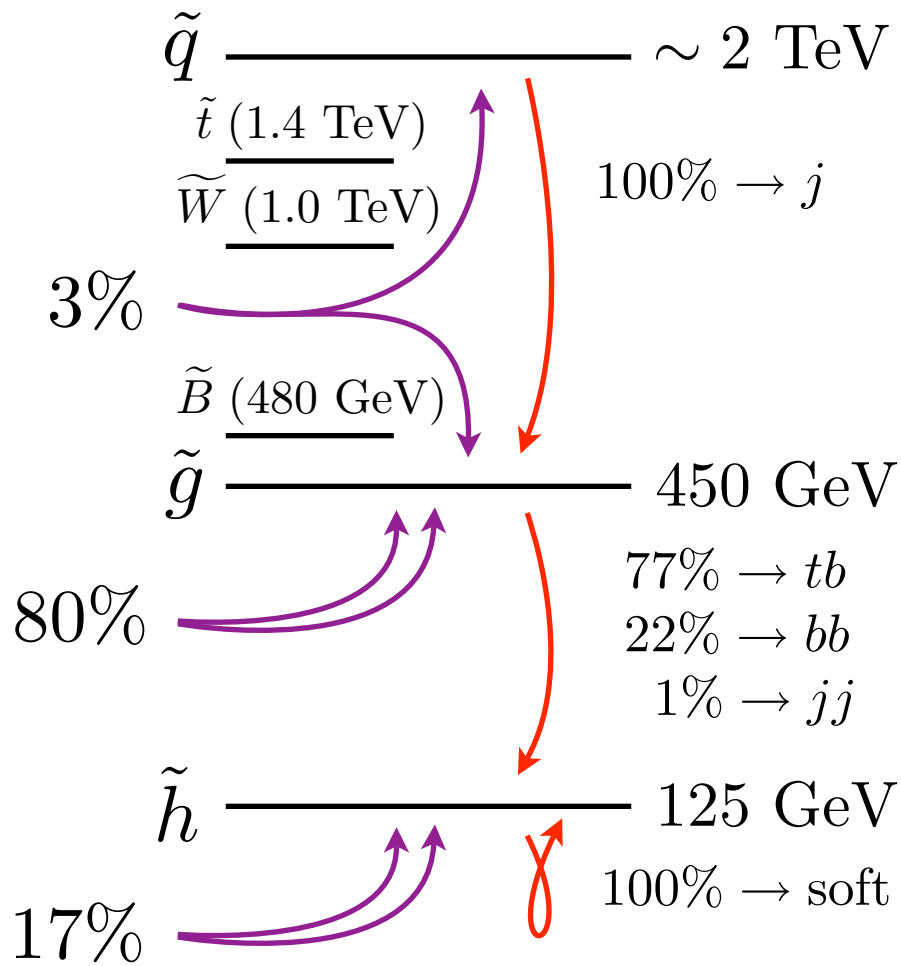
2j channel suppressed  
by small couplings to Higgsino



or

2j channel suppressed  
by high squark mass

# Michigan Box (MSSM)





# DBox

Nima & Natalia vs. Jesse & Philip  
(Marmoset Smackdown)

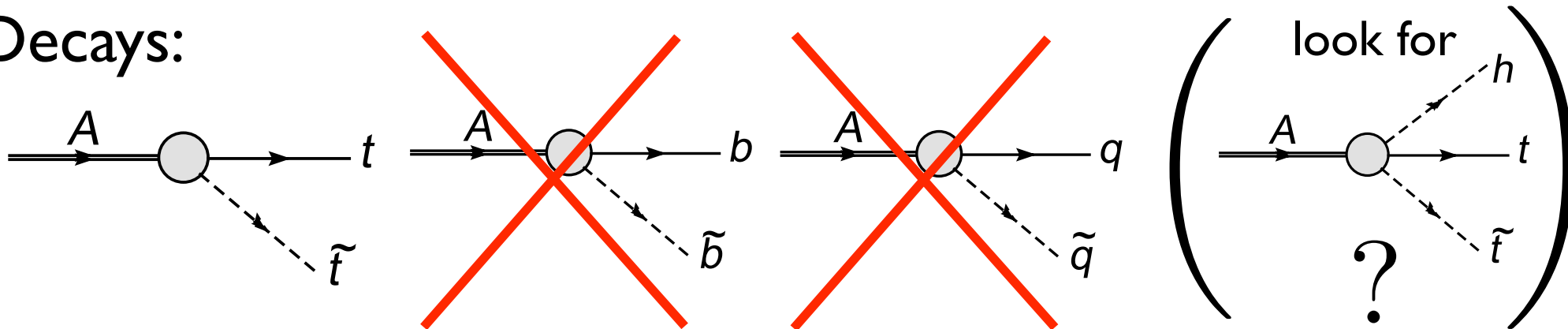
- SUSY with very heavy gauginos ( $> 5$  TeV)
- Light scalars
- **also** light colored adjoint  $\Phi$  decaying through

$$W \supset \frac{\Phi}{M} Q_3 H_u U_3^c$$

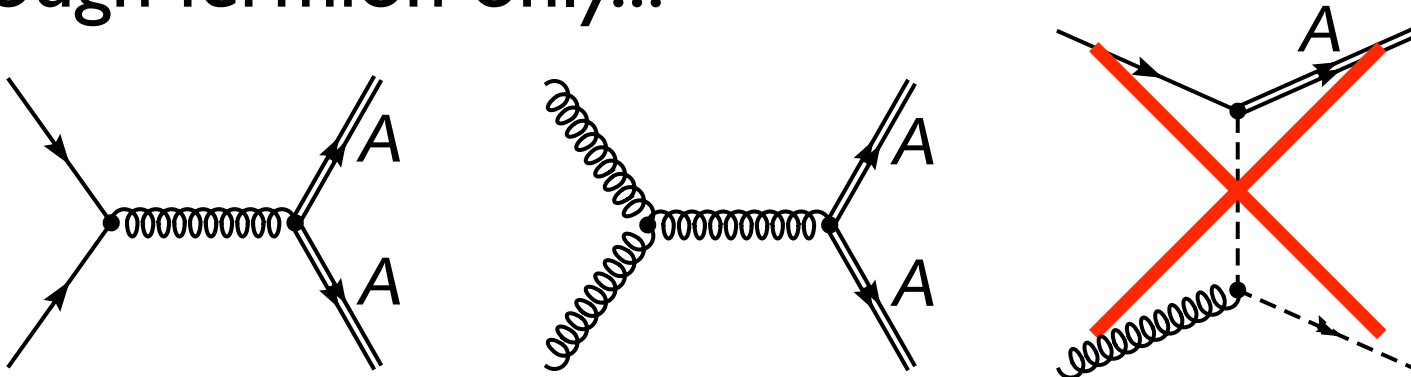
- looks like gluino with wrong decay modes!

# DBox “Gluino”?

Decays:

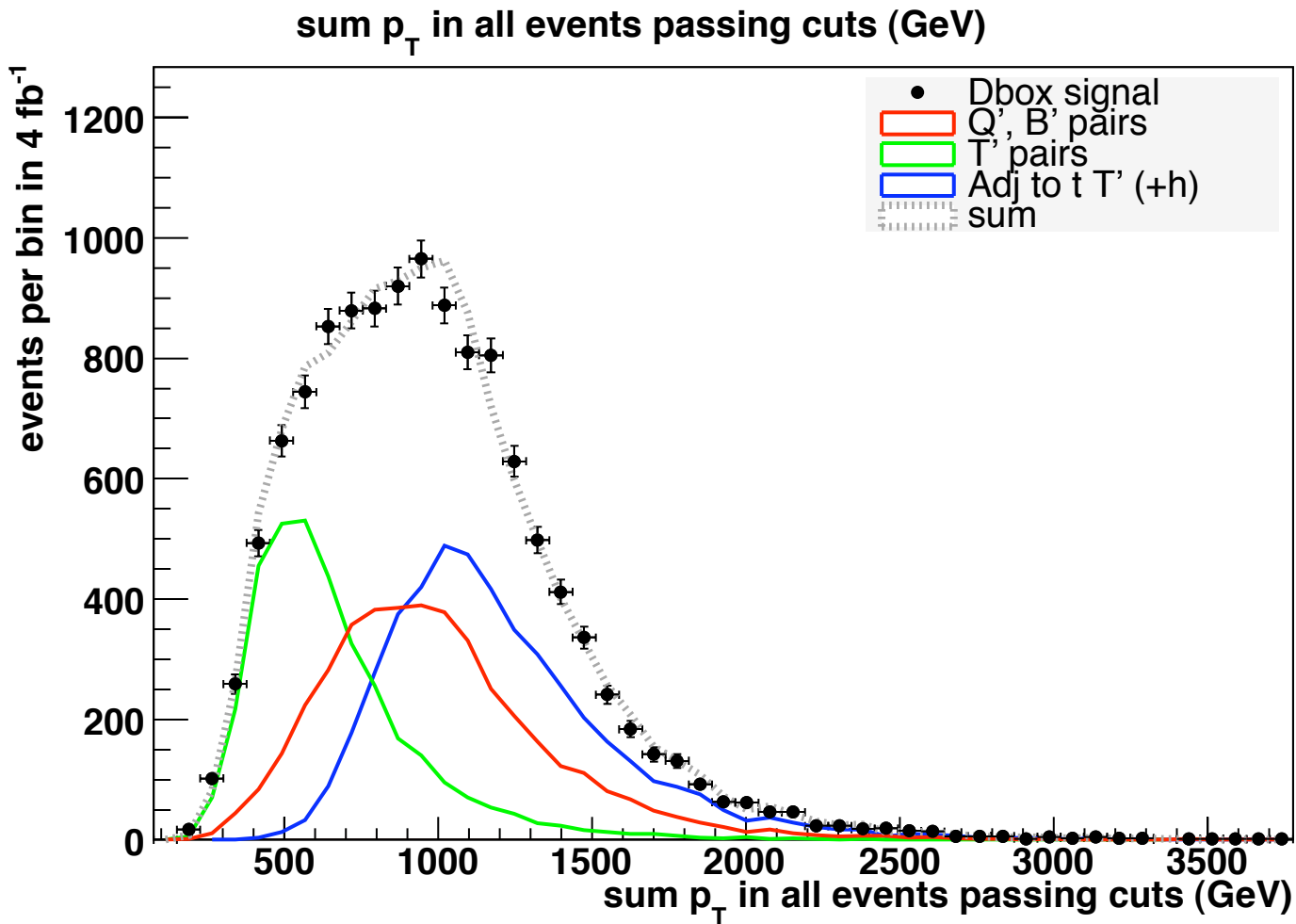


rate consistent with QCD production through fermion only...



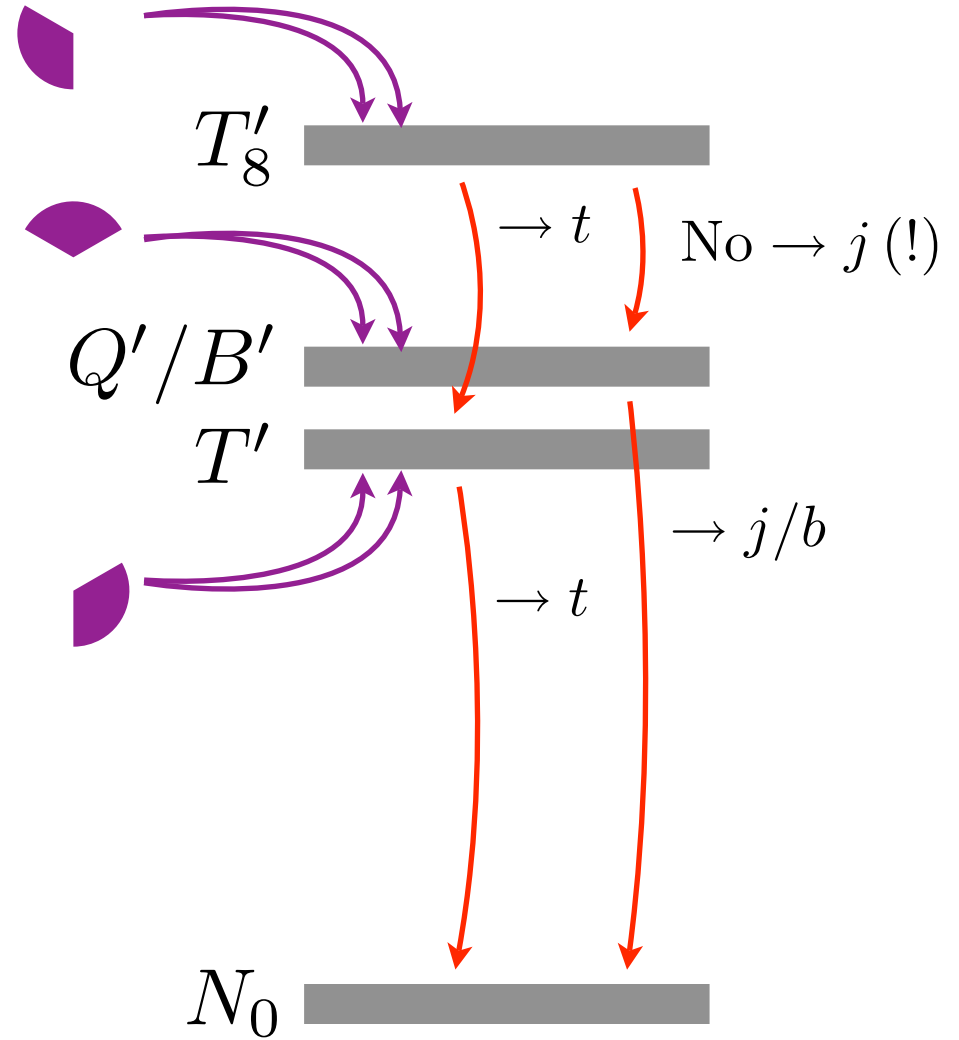
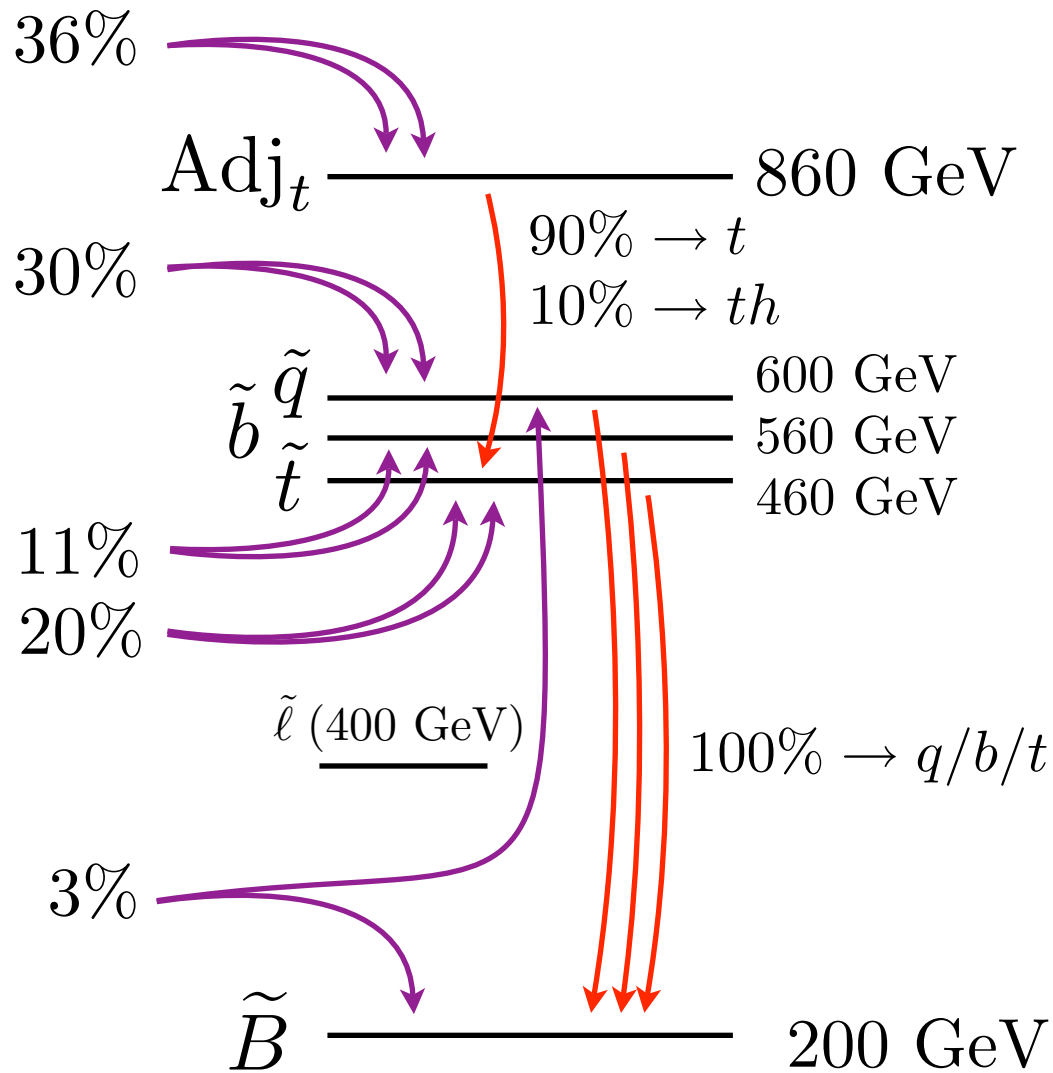
...and no associated production with squark.

**It's not a gluino!**



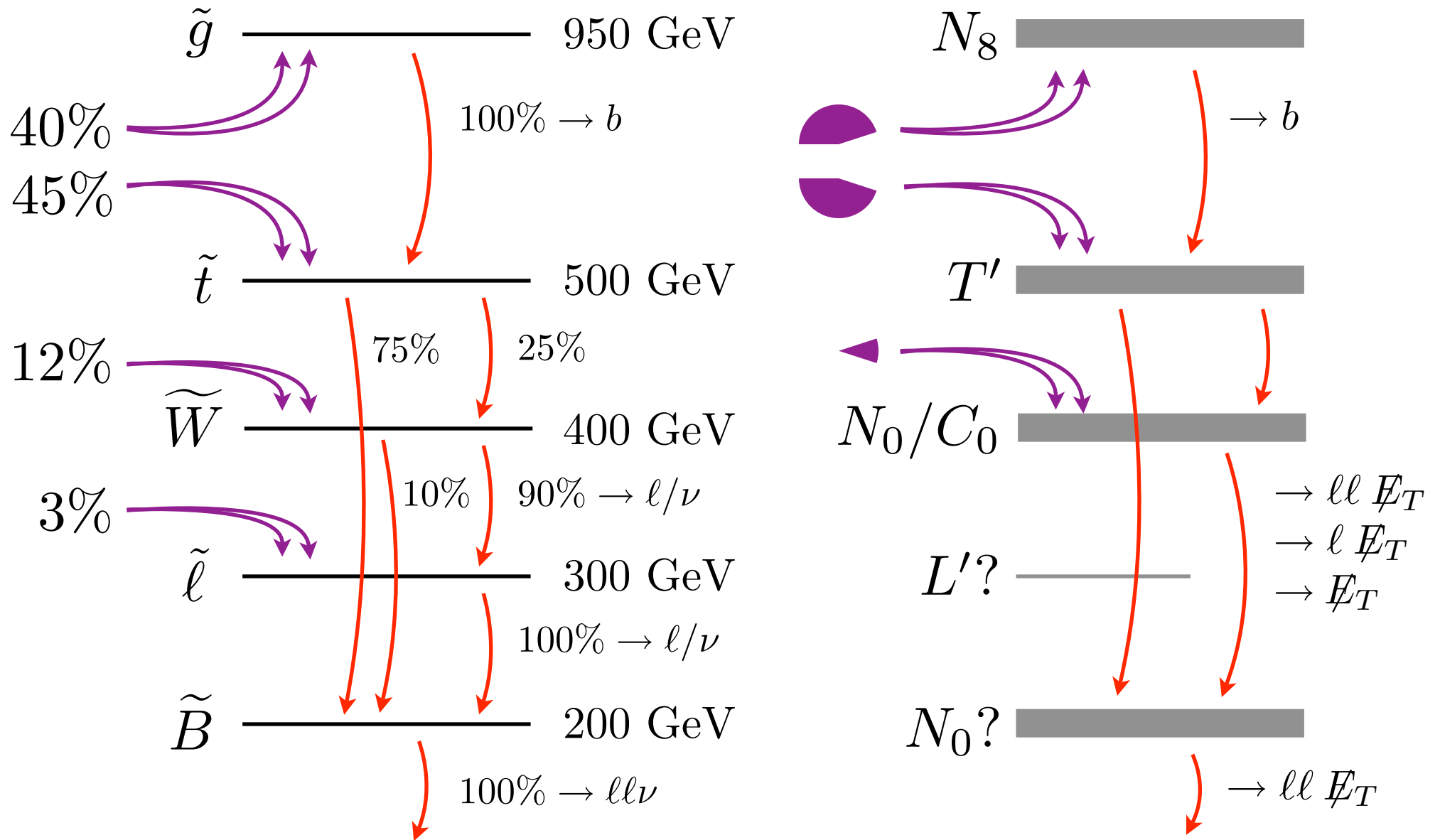
$$H_T = \sum_i E_T^i + \cancel{E}_T$$

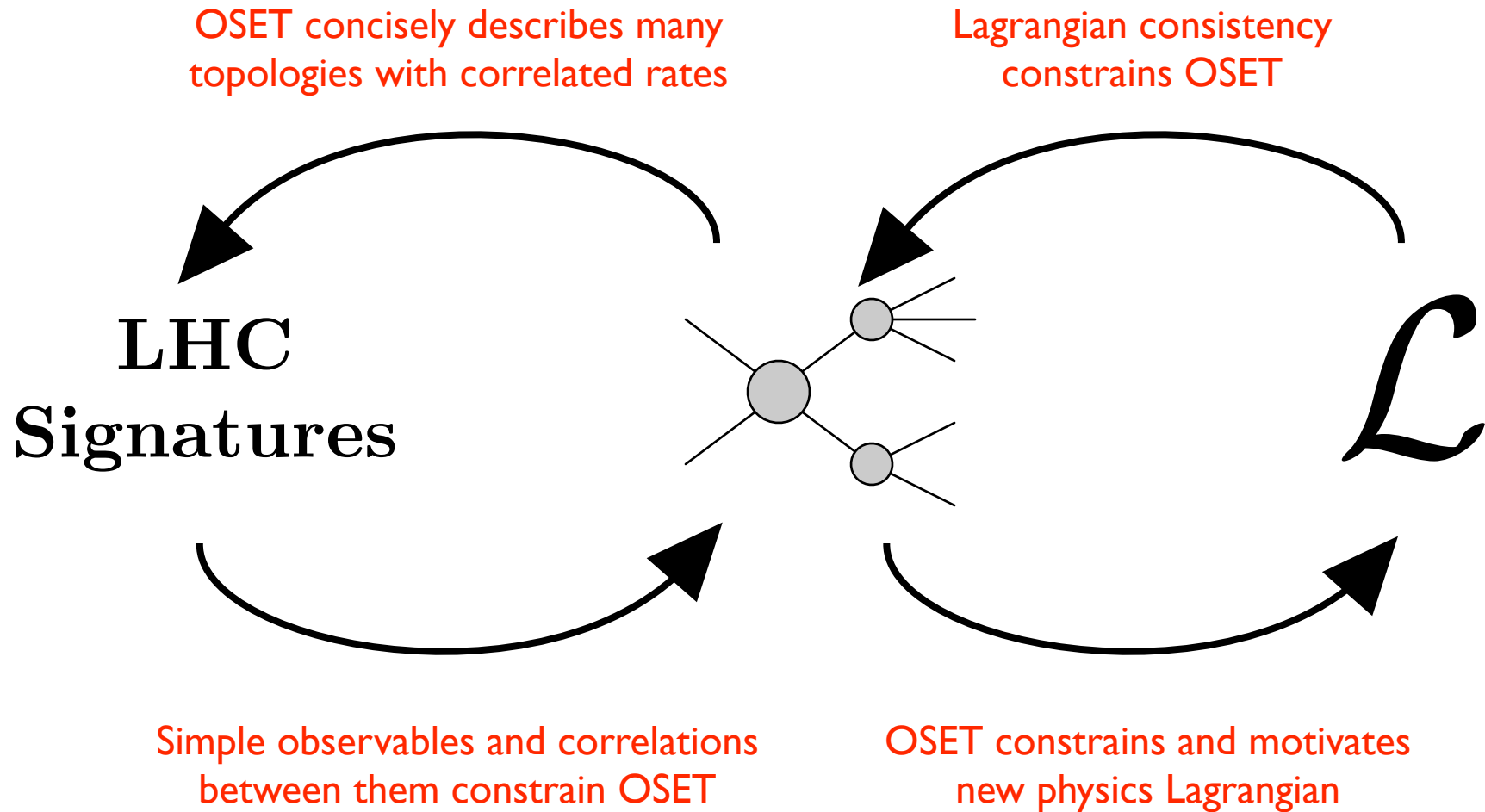
# DBox



# ABox (RPV SUSY)

Variant of Rutgers's Blackbox (3rd LHC Olympics)





**A successful Lagrangian characterization may be derived from the top down or from the bottom up, and will be most convincing when understood in both ways.**