

Unparticle Solution to Hierarchy

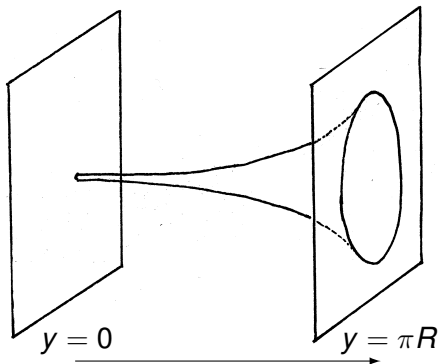
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May 10, 2010

Randall Sundrum One (RS1)

$$ds^2 = e^{-2ky} \eta_{\mu\nu} dx^\mu dx^\nu + dy^2$$

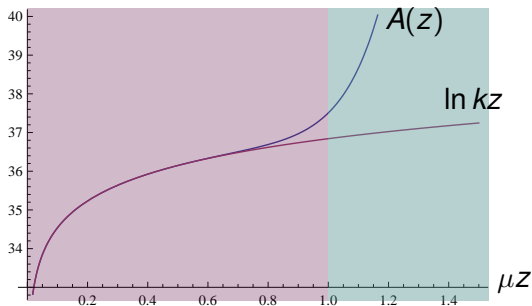


- Need to stabilize R distance (Goldberger-Wise)
- Why not use GW scalar to replace IR brane?

BG Soft-Wall Geometry

$$ds^2 = e^{-2A(z)} \left(\eta_{\mu\nu} dx^\mu dx^\nu + dz^2 \right)$$

$$A(z) = \ln kz + \frac{2}{3}(\mu z)^\nu$$



Planck Weak Hierarchy

- Need to achieve $\mu/k \sim 10^{-16}$
- What sets μ/k ?
- Consider:
 - μ sets scale where scalar back-reaction strong
 - Must fix field at one location
 - Boundary condition on UV brane fixes field
- Boundary potential:

$$\lambda_{UV} = W(\eta_0) + \partial_\eta W(\eta_0)(\eta - \eta_0) + m_{UV}(\eta - \eta_0)^2$$

- Boundary conditions require $\eta_0 = \langle \eta \rangle_0$

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Planck Weak Hierarchy

- The value of the field at the UV brane is

$$\eta_0 = \pm \sqrt{3} \left(\frac{\nu + 1}{\nu} \right) \left[\sqrt{\frac{2}{3} \frac{\nu}{\nu + 1} \left(\frac{\mu}{k} \right)^\nu + \left(\frac{2}{3} \frac{\nu}{\nu + 1} \left(\frac{\mu}{k} \right)^\nu \right)^2} + \sinh^{-1} \sqrt{\frac{2}{3} \frac{\nu}{\nu + 1} \left(\frac{\mu}{k} \right)^\nu} \right]$$

- Not easily inverted...

Planck Weak Hierarchy

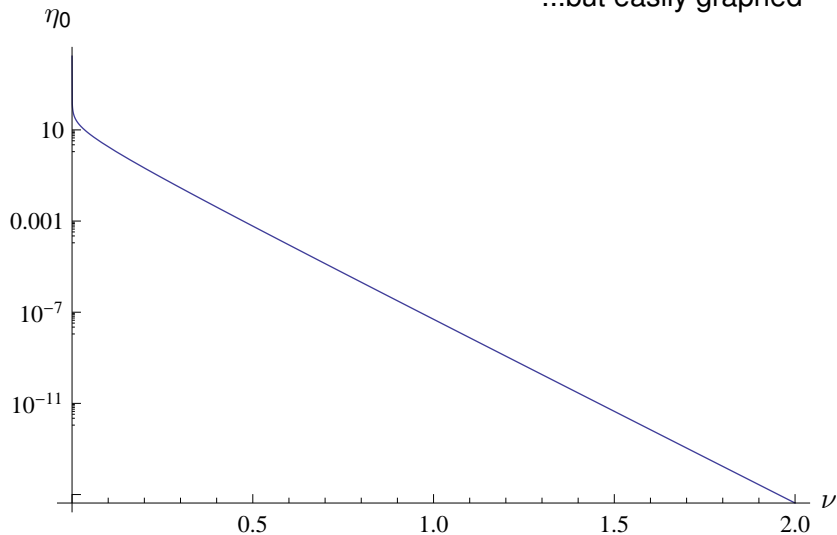
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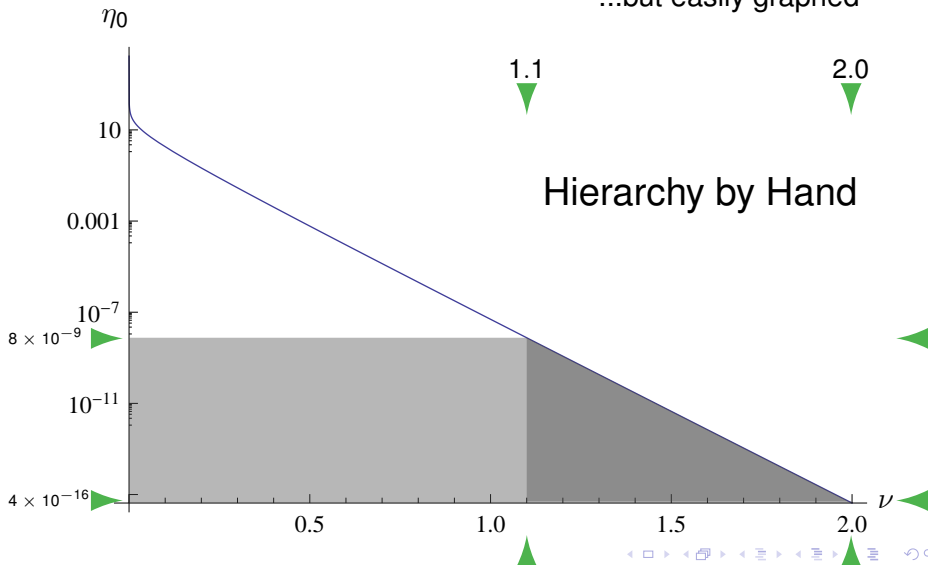
Planck Weak Hierarchy

...but easily graphed



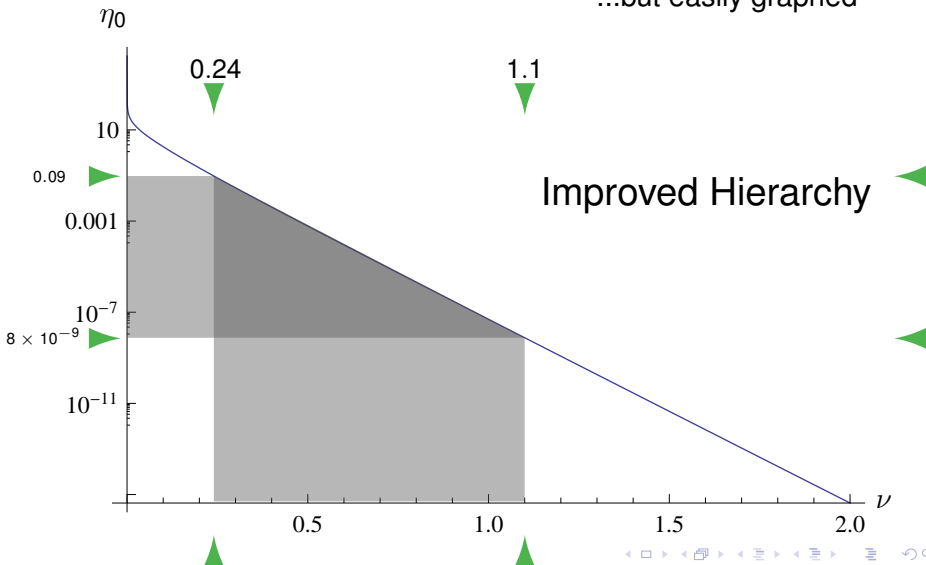
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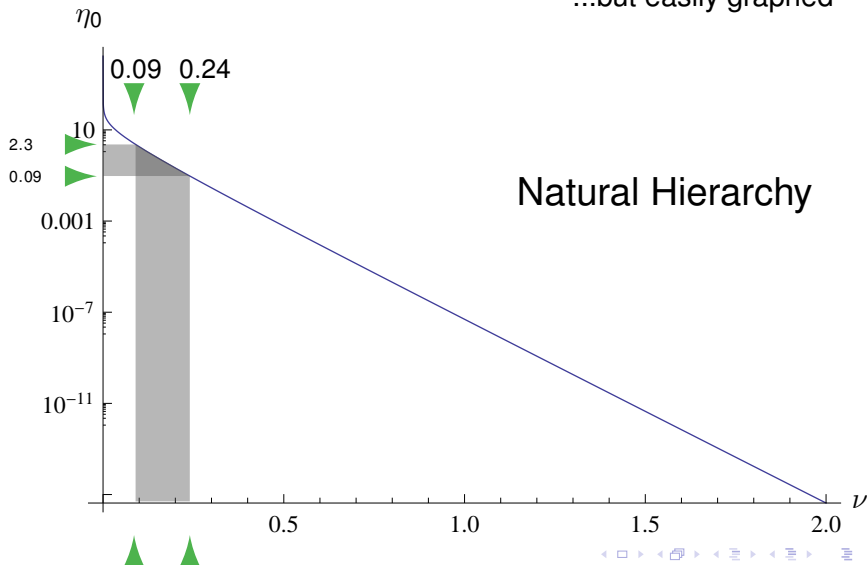
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Scalar's Potential

Of course, ν has other consequences...

Look at potential

$$V(\eta) = -12k^2 - k^2\nu\left(1 - \frac{\nu}{8}\right)\eta^2 + \dots$$

Gives η 's mass as

$$m_\eta^2 = -2k^2\nu\left(1 - \frac{\nu}{8}\right)$$

AdS/CFT correspondence says operator dimension is

$$\Delta = 2 + \sqrt{4 + \frac{m_\eta^2}{k^2}} = 2 + \frac{1}{2}|4 - \nu|$$

Operator Dimension

The breakdown is

Hierarchy by Hand	$\nu > 1$	$\Delta > \frac{5}{2}$
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Improved Hierarchy	$\nu \sim 1$	$\Delta \sim \frac{5}{2}$
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Natural Hierarchy	$0 < \nu < 1$	$2 < \Delta < \frac{5}{2}$
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Natural Hierarchy $0 < \nu < 1$ $2 < \Delta < \frac{5}{2}$

Fractional Dimension!

Fluctuations Parameterized

$$ds^2 = e^{2(F-A(z))} \left[((1 - 2F)\eta_{\mu\nu} + h_{\mu\nu}) dx^\mu dx^\nu + 2A_\mu dx^\mu dz + dz^2 \right]$$

$$\eta = \langle \eta \rangle + \tilde{\eta}$$

- Consider Just Scalar Modes
 - gravi-scalar, F
 - scalar tower of η
- Suitable field redefinition permits writing as Schrödinger equation

$$\left(-\partial_z^2 + V_{SE}(z) \right) \psi = m^2 \psi$$

- Start with $m = 0$ modes

No Massless Scalars

- There are no massless modes
- Consider theory without UV brane
 - Theory invariant under $A(z) \rightarrow A(z) + C$
 - Look again at parameterization of fluctuations:

$$ds^2 = e^{2(F-A(z))} \left[((1-2F)\eta_{\mu\nu} + h_{\mu\nu}) dx^\mu dx^\nu + 2A_\mu dx^\mu dz + dz^2 \right]$$

- F Goldstone boson
- But μ/k changes under shift: $\frac{\mu}{k} \rightarrow \frac{\mu}{k} e^{-C}$
- Fixing μ/k breaks symmetry
- With UV brane, μ/k determined, so no massless modes
- What about massive modes?

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Scalar Modes - Massive

- Massive modes dynamical variable

$$\nu = -\sqrt{2}e^{-3A(z)/2} \frac{\langle \eta \rangle'}{A'(z)} \left(-\frac{1}{2}F + \frac{A'(z)}{\langle \eta \rangle'} \tilde{\eta} \right)$$

- Schrödinger Potential Behavior

$$\nu > 1 \quad z \rightarrow \infty \Rightarrow V_{\text{SE}} \rightarrow \infty$$

$$\nu = 1 \quad z \rightarrow \infty \Rightarrow V_{\text{SE}} \rightarrow \mu^2$$

$$\nu < 1 \quad z \rightarrow \infty \Rightarrow V_{\text{SE}} \rightarrow 0$$

Scalar Modes - Massive

- Massive modes dynamical variable

$$v = -\sqrt{2}e^{-3A(z)/2} \frac{\langle \eta \rangle'}{A'(z)} \left(-\frac{1}{2}F + \frac{A'(z)}{\langle \eta \rangle'} \tilde{\eta} \right)$$

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Continuum
 $m^2 > 0$
 (Unparticles)

Summary

- Examined Planck weak Hierarchy for Batell-Gherghetta Soft-Wall
- Found natural hierarchy for $\nu < 1$
- $\nu < 1$ corresponds to fractional-dimension operators in dual theory
- $\nu < 1$ implies a continuum of modes without a mass gap in the $5D$ theory
- Thus, natural hierarchy implies unparticles