

# Lifting the Suppression

Electroweak Bremsstrahlung as the Dominant Dark  
Matter Annihilation Channel

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# Outline

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Suppression of dark matter annihilations



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### Types of suppression

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# Varieties of Suppression

$$x\bar{x} \rightarrow f\bar{f}$$



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In the non-relativistic limit, partial wave expansion gives that the  $L^{\text{th}}$  partial wave is suppressed by  $v^{2L}$ .

$L = 0$  is an s-wave,  $L = 1$  is a p-wave, etc... P-wave suppression is considerable in the galactic halo.

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## Helicity Suppression

For some fermionic final state currents there arises an additional suppression of  $m_f/m_\chi$  in the amplitude, leading to a suppression proportional to  $(m_f/m_\chi)^2$  in the rate

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${}^3P_0(0^{++})$	Scalar	$\bar{\Psi}\Psi$
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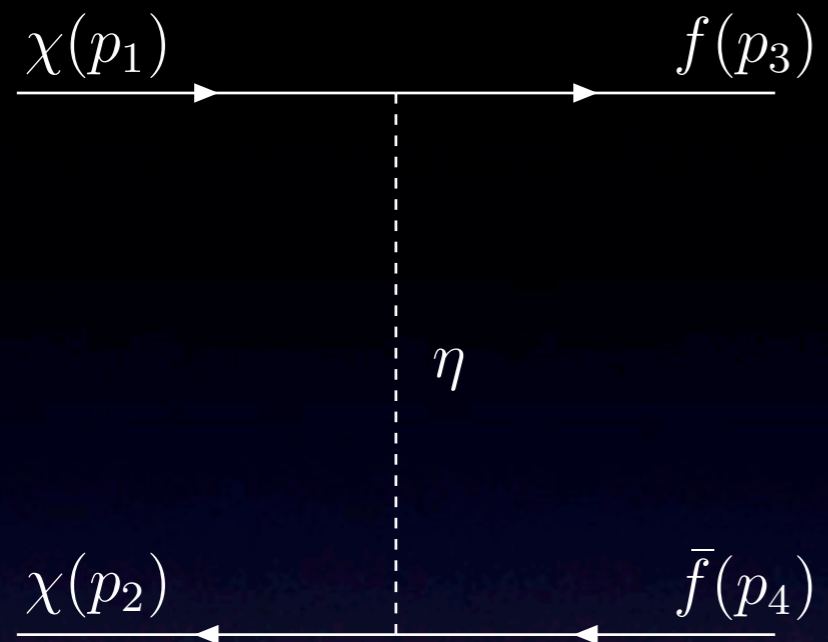
This will introduce a helicity suppression  $m_f/m_\chi$   
in the amplitude<sup>1</sup>

<sup>1</sup> H. Goldberg, *Phys.Rev.Lett.* **50**, 1419, 1983

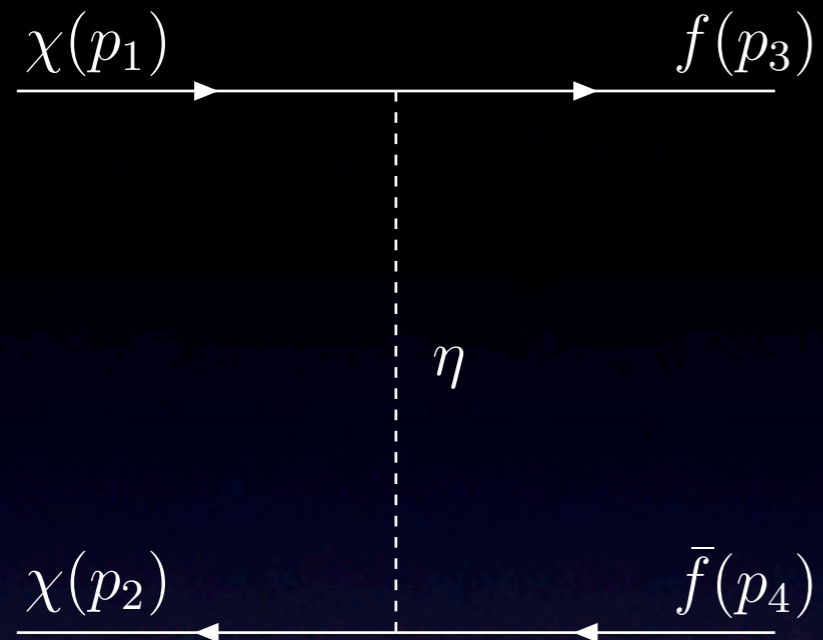


# Fierz Transformations

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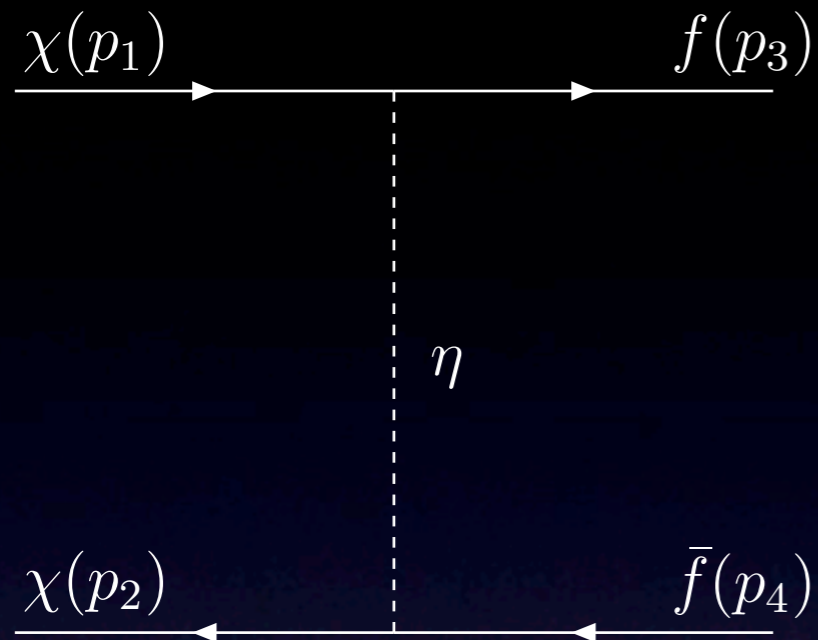
# Fierz Transformations



$$\mathcal{M}_t \propto \bar{v}(p_2) P_L v(p_4) \bar{u}(p_3) P_R u(p_1)$$



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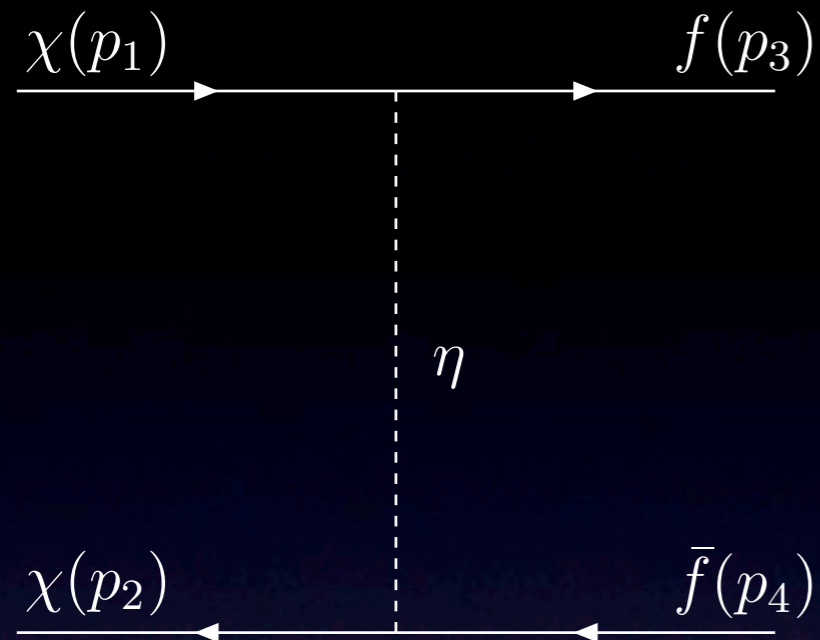


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Using the Fierz transformation

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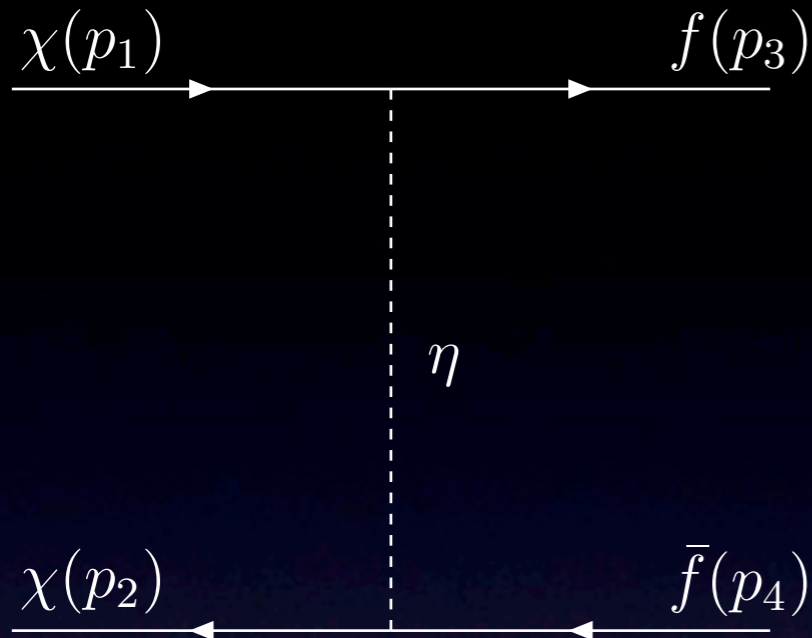
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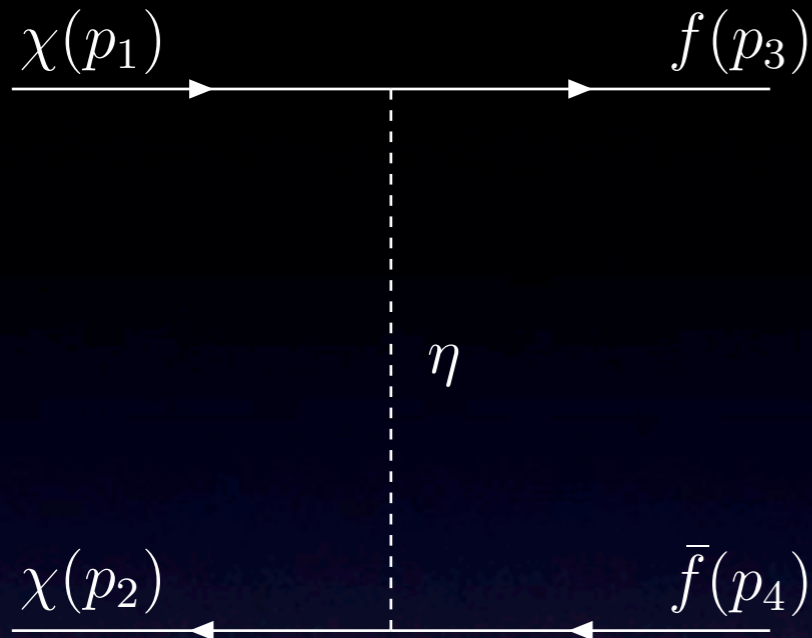
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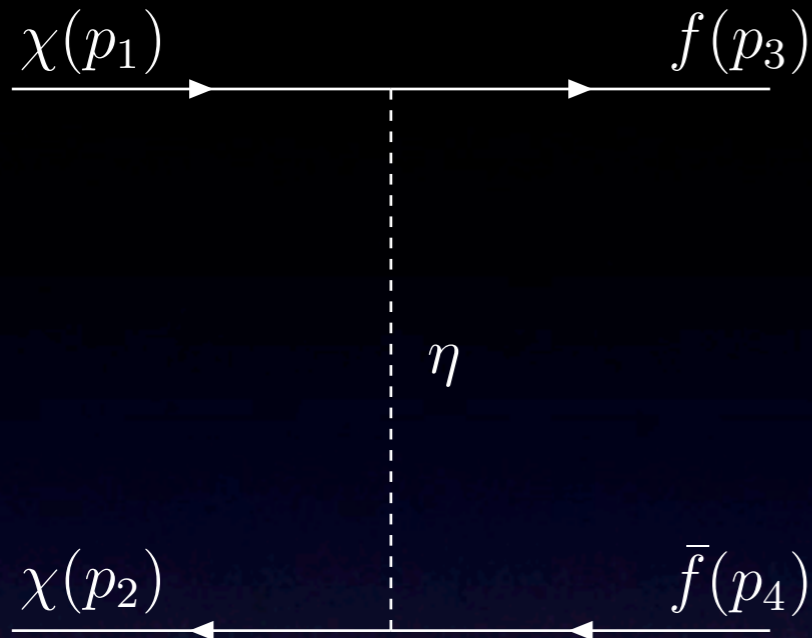
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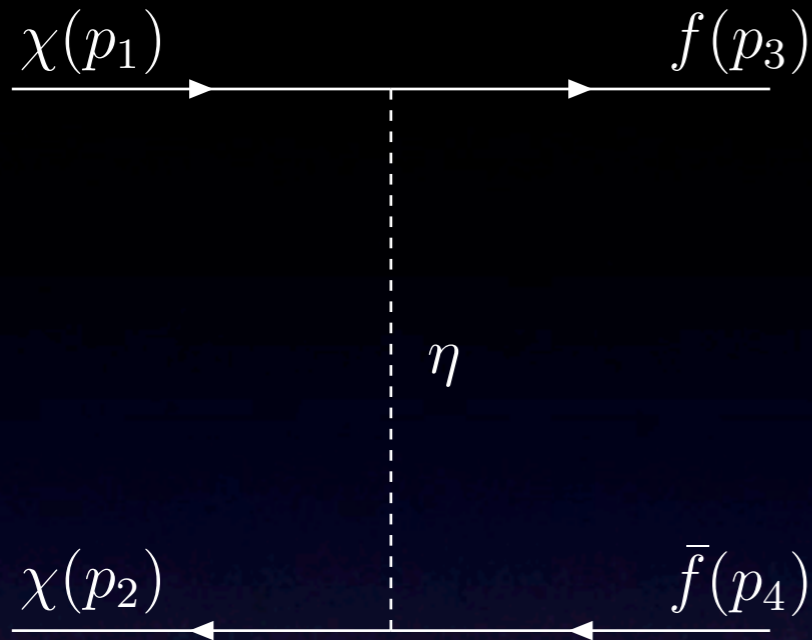
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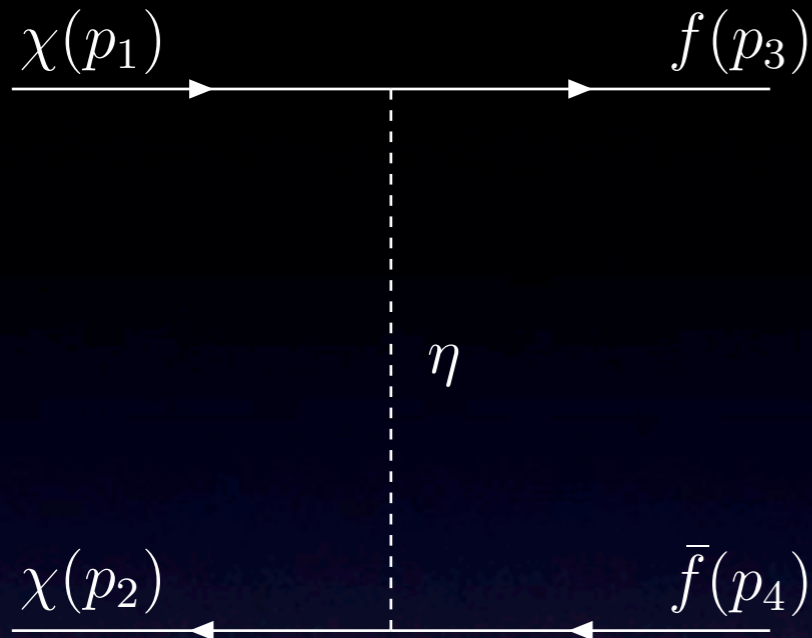
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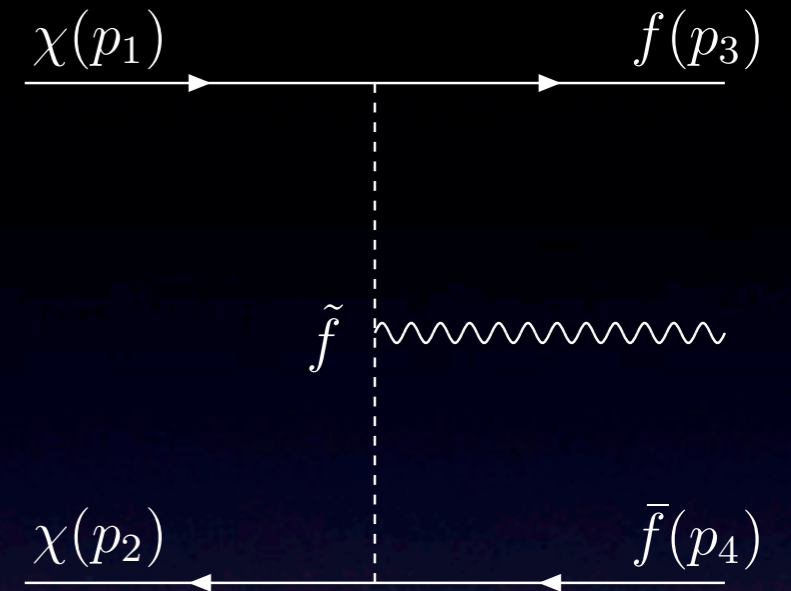
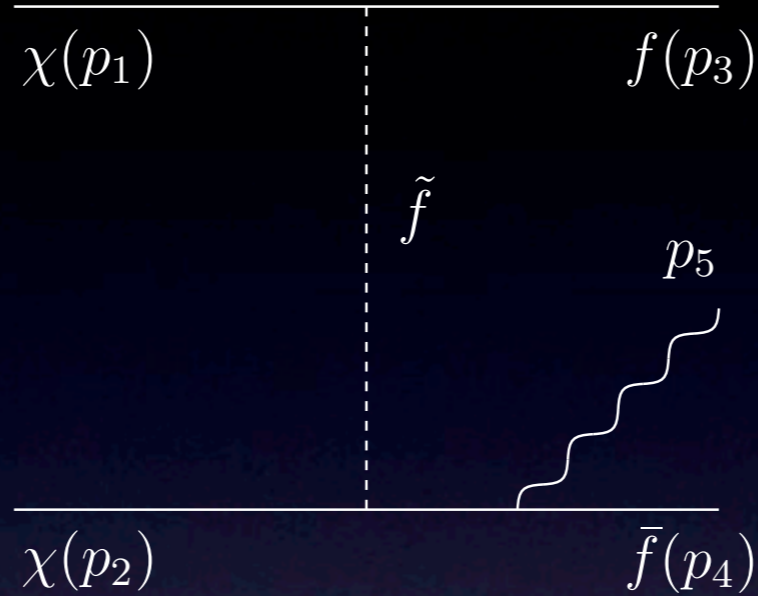
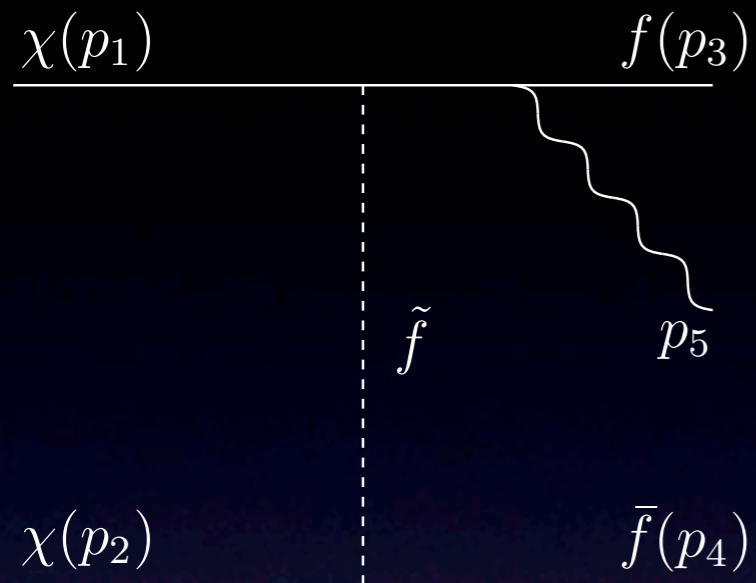
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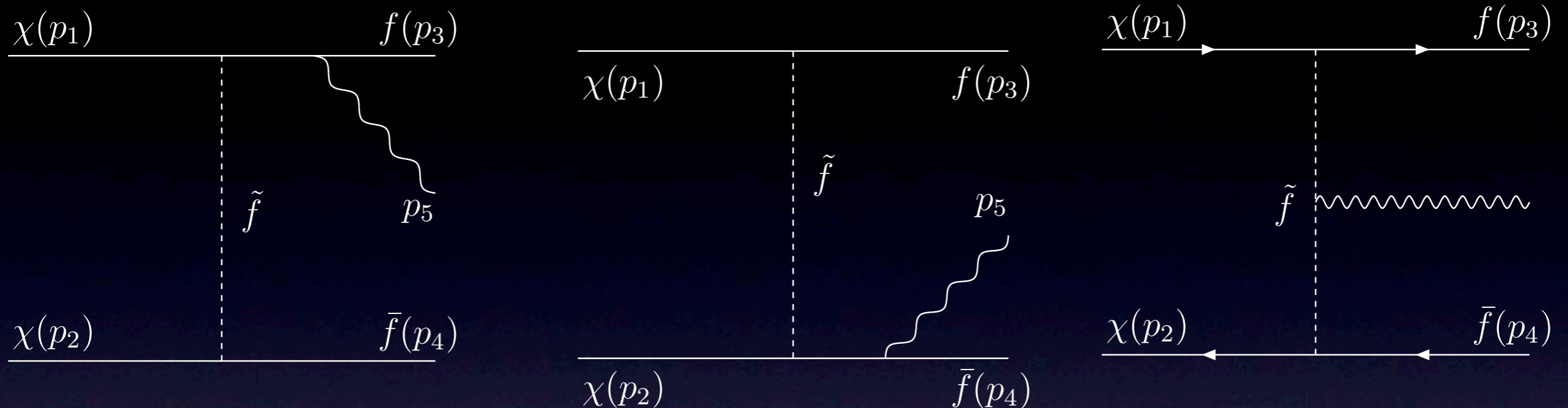
# Photon Bremsstrahlung

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It has been known for some time that radiative corrections to the dark matter annihilation process can be enormous when compared to the lowest order rate<sup>1,2,3,4</sup>.

<sup>1</sup>L. Bergstrom, Phys.Lett. B **225**, 372 (1989)

<sup>2</sup>R. Flores, K.A. Olive, and S. Rudaz, Phys.Lett. B **232**, 377 (1989)

<sup>3</sup>T. Bringmann, L. Bergstrom, and J. Edsjo, JHEP **0801** 049 (2008)

<sup>4</sup>V. Barger, Y. Gao, W.-Y. Keung, and D. Marfatia, Phys.Rev.D **80**, 063537 (2009)

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These studies were done assuming no helicity suppression. We will now investigate these effects in conjunction with the suppressed processes.

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We will examine the leptophilic model<sup>1</sup>

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For Majorana dark matter the cross-section in the massless lepton limit is

$$\sigma v = \frac{f^4(r^2 - 2r^3 + 2r^4)v^2}{24m_\chi^2\pi} \quad \text{where} \quad r = \left( \frac{m_\chi^2}{m_\eta^2 + m_\chi^2} \right).$$

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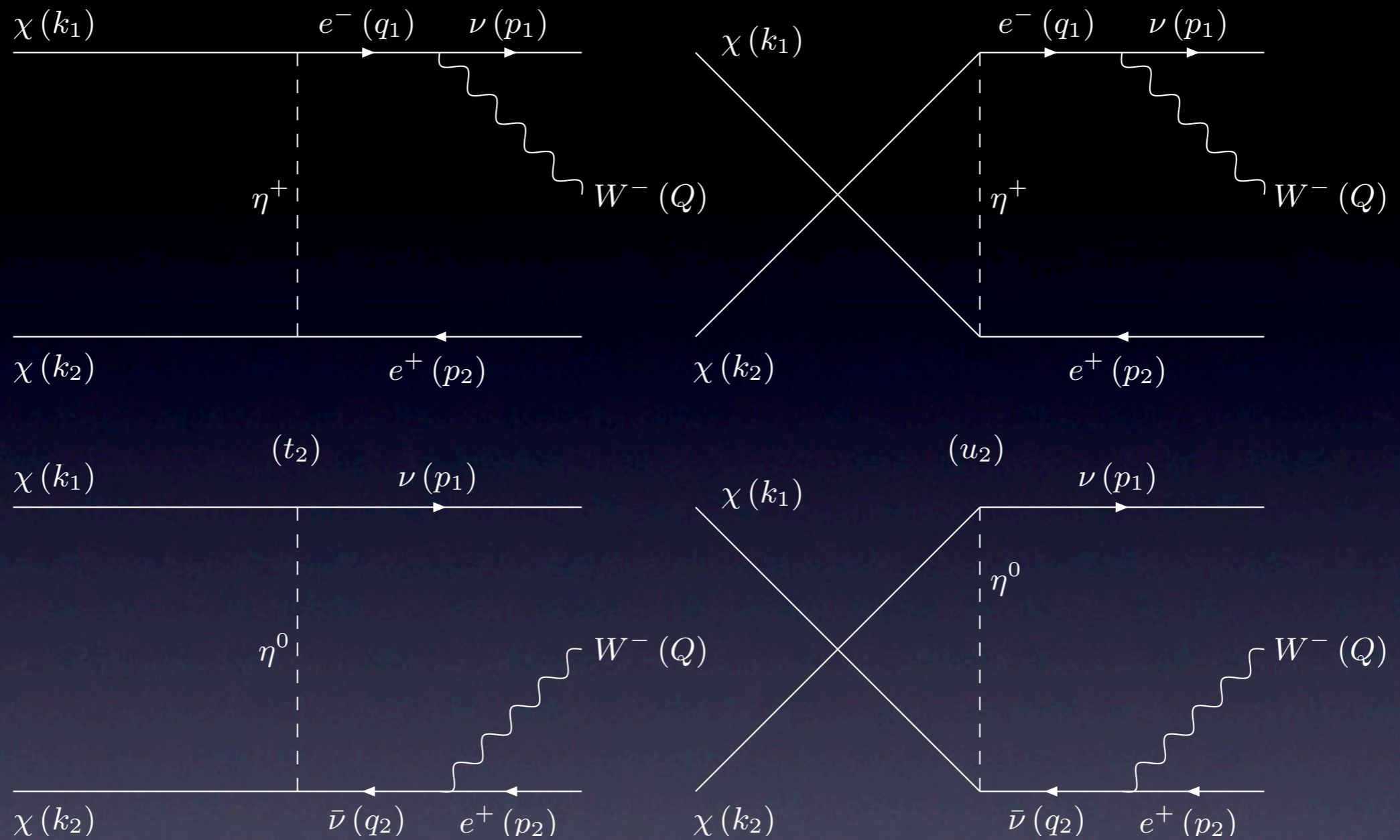
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We see that the expected velocity suppression arises

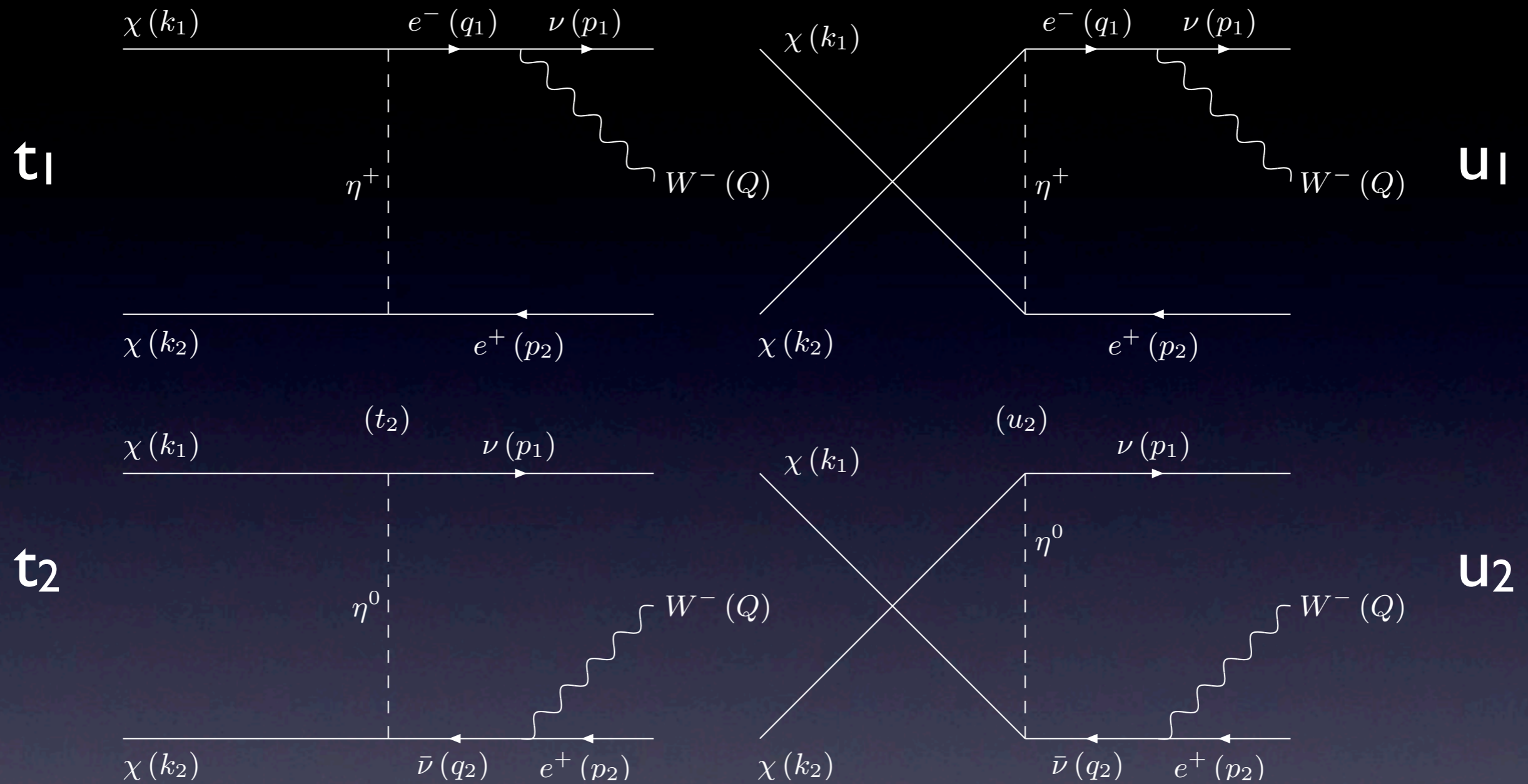
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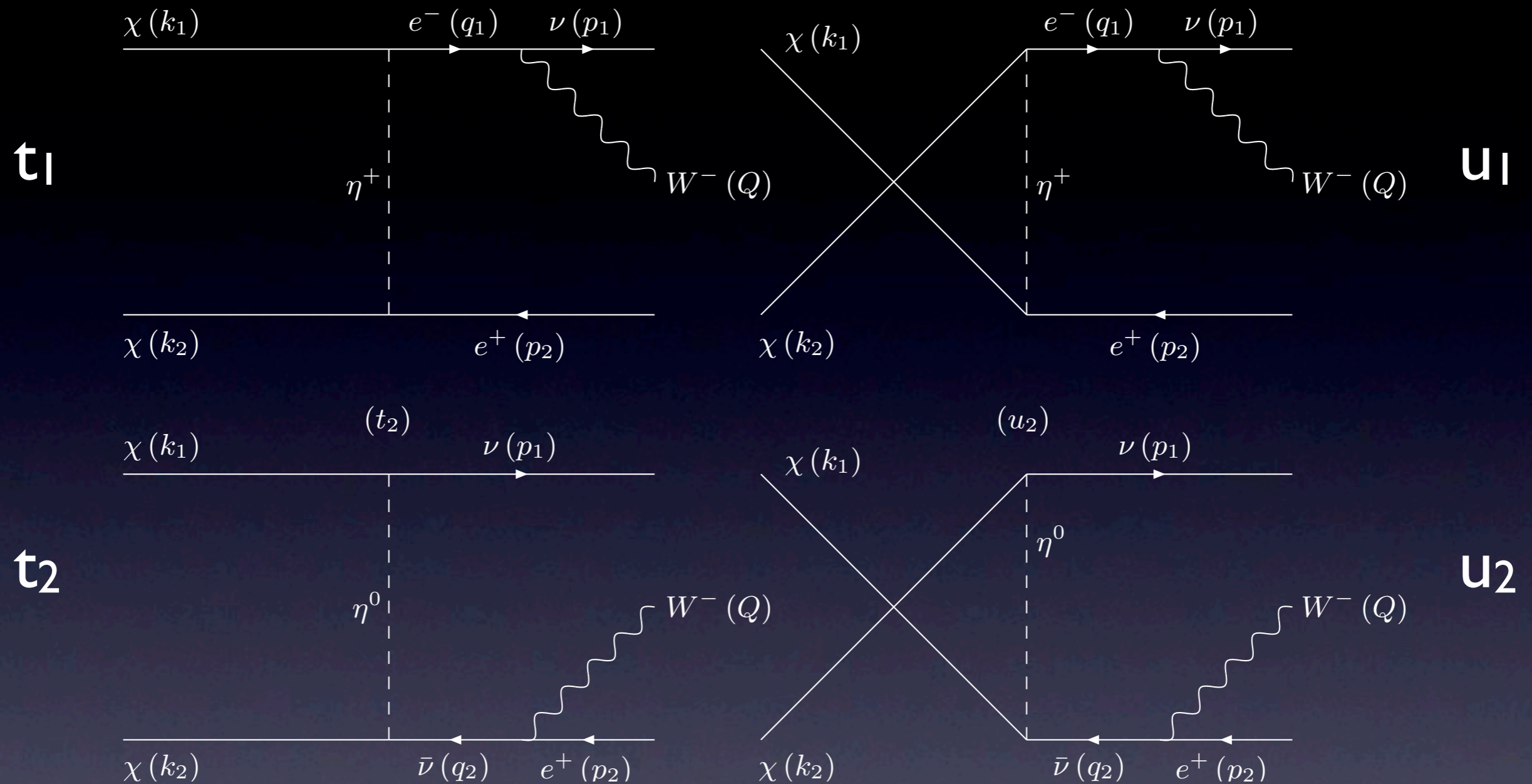
# Lifting the Suppression with W or Z emission



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Contributions to  $\chi\chi \rightarrow e^+\nu W^-$



# Applying the Fierz Relations

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$$\mathcal{M}_{t_1} = \frac{igf^2}{\sqrt{2}q_1^2} \frac{1}{t_1 - m_\eta^2} \left( \bar{v}(k_2) P_L v(p_2) \right) \left( \bar{u}(p_1) \gamma^\mu P_L \not{q}_1 u(k_1) \right) \epsilon_\mu^Q$$

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 &= \frac{igf^2}{\sqrt{2}q_1^2} \frac{1}{t_1 - m_\eta^2} \epsilon_\mu^Q \frac{1}{4} \left[ \left( \bar{v}(k_2) u(k_1) \right) \left( \bar{u}(p_1) P_L \gamma^\mu P_L \not{q}_1 v(p_2) \right) \right. \\
 &\quad \left. + \left( \bar{v}(k_2) \gamma_5 u(k_1) \right) \left( \bar{u}(p_1) P_L \gamma_5 \gamma^\mu P_L \not{q}_1 v(p_2) \right) \right. \\
 &\quad \left. - \left( \bar{v}(k_2) \gamma_5 \gamma_\alpha u(k_1) \right) \left( \bar{u}(p_1) \gamma^\alpha \gamma^\mu P_L \not{q}_1 v(p_2) \right) \right]
 \end{aligned}$$



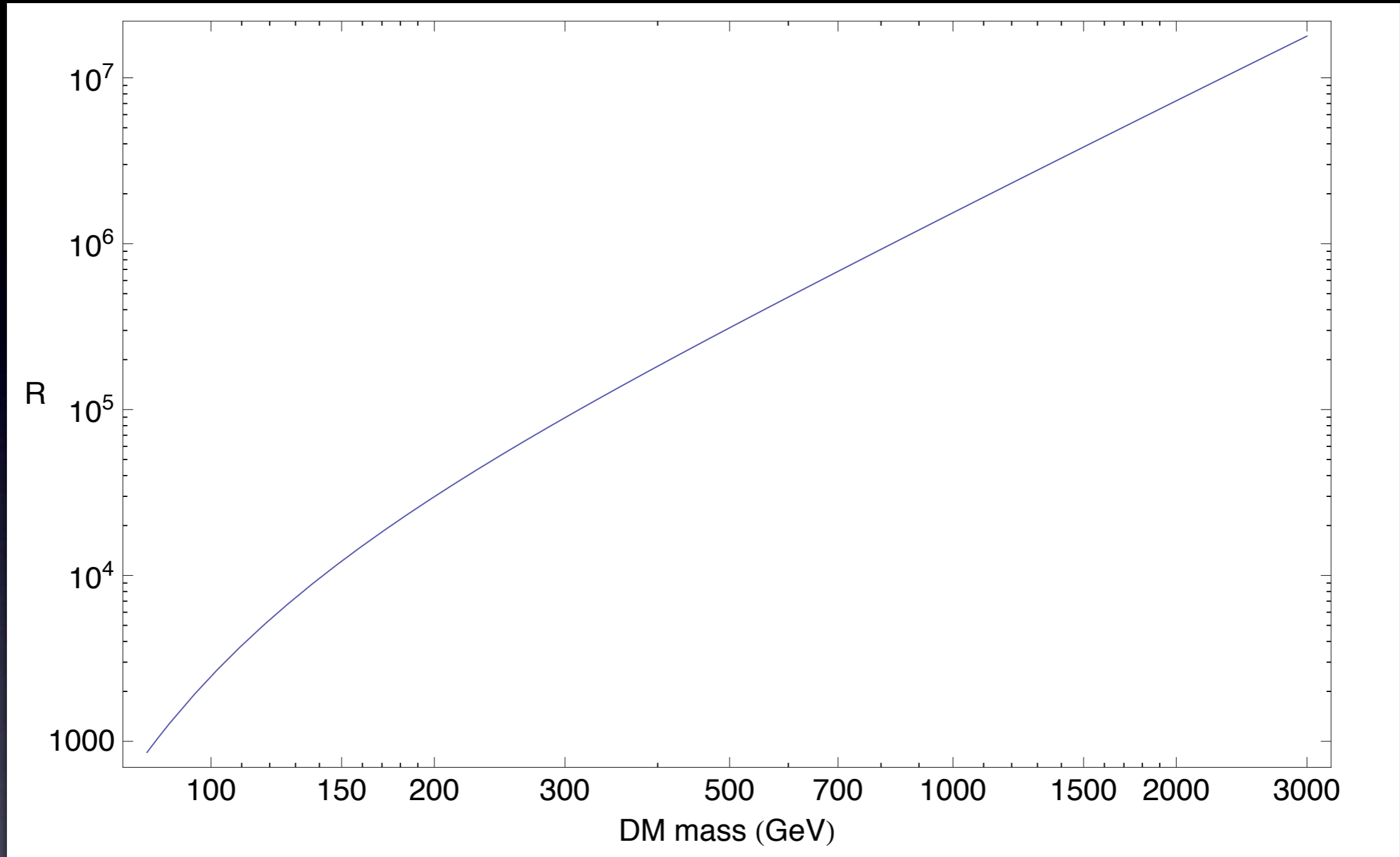
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 \end{aligned}$$

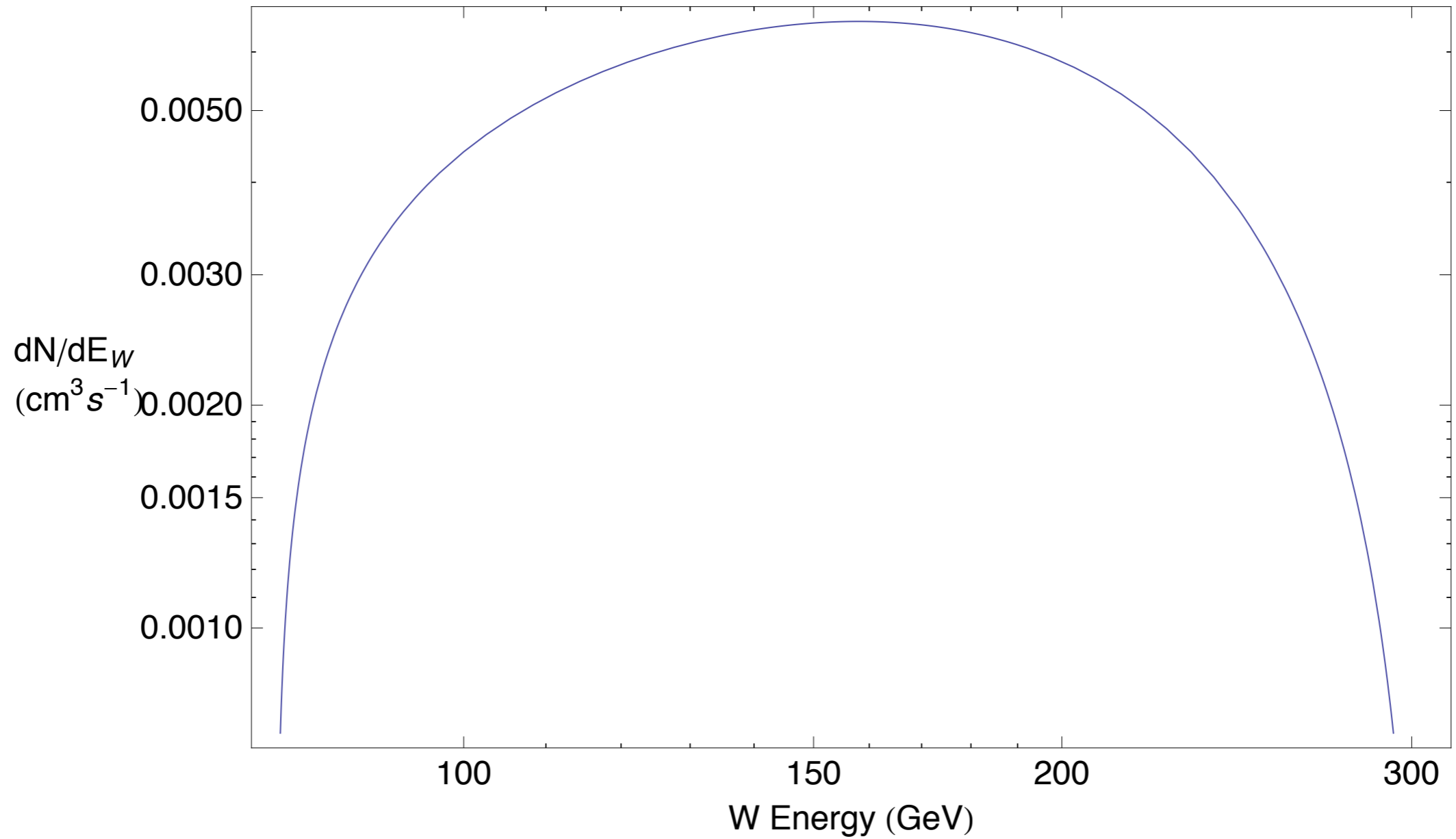
# Results



$\sigma v(\chi\chi \rightarrow e^+ \nu W^-) / \sigma v(\chi\chi \rightarrow e^+ e^-)$  vs.  $m_\chi$  (GeV)  $m_\eta = 10^{10}$  GeV

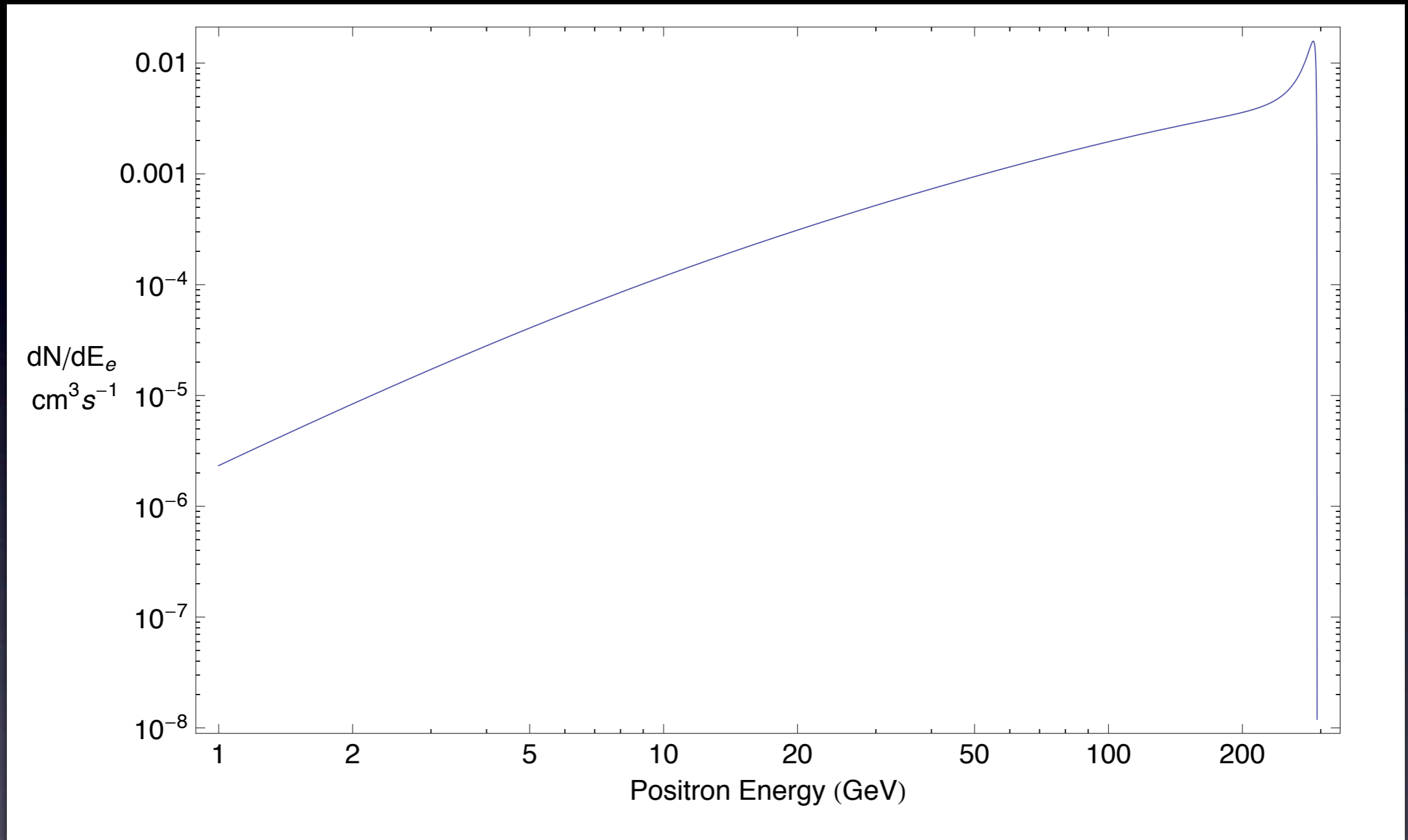


# Spectra



$$M_\chi = 300 \text{ GeV}, M_\eta = 10^{10} \text{ GeV}$$

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We have recapitulated under what circumstances these suppressions may arise and how they may be circumvented

We have demonstrated in an example model that electroweak bremsstrahlung may have dramatic effects on suppressed processes leading to possibly strong constraints on various models