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# Measuring a Light (Dark Matter) Neutralino Mass at the ILD

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## MSSM Neutralino Mixing

- MSSM Neutralino Mass Mixing Matrix in  $(\tilde{B}, \tilde{W}^3,$

$$\mathcal{M}_0 = \begin{pmatrix} M_1 & 0 & -M_z \cos \beta \sin \theta_w \\ 0 & M_2 & M_z \cos \beta \cos \theta_w \\ -M_z \cos \beta \sin \theta_w & M_z \cos \beta \cos \theta_w & 0 \\ M_z \sin \beta \sin \theta_w & -M_z \sin \beta \cos \theta_w & \mu \end{pmatrix}$$

- What do we know about the Mass of  $\tilde{\chi}_1^0$ ?

## Experimental Search at L

- Chargino Search:  $M_{\tilde{\chi}_1^\pm} > 94 \text{ GeV} \Rightarrow |\mu|, M_2 \gtrsim 100 \text{ GeV}$
- Higgs search:  $\tan \beta \gtrsim 1.5$
- Assume:  $M_1 = \frac{5}{3} \tan^2 \theta_w M_2 \Rightarrow M_1 \gtrsim 50 \text{ GeV}$
- Insert into Neutralino Mass Matrix:  $\Rightarrow M_{\tilde{\chi}_1^0} \gtrsim 40 \text{ GeV}$
- **Now drop above assumption on  $M_1, M_2$**

## Massless Neutralino

- Set  $\det(\mathcal{M}_0)=0 \Rightarrow M_1 = \frac{M_2 M_Z^2 \sin(2\beta) s_W^2}{\mu M_2 - M_Z^2 \sin(2\beta) c_W^2}$
- Choose:  $\{M_2, \mu, \tan \beta\} \Rightarrow \exists M_1 : M_{\chi_1^0} = 0$
- Some fine-tuning required

$$M_1 \approx \frac{M_Z^2 \sin(2\beta) s_W^2}{\mu} \approx 2.5 \text{ GeV} \left( \frac{10}{\tan \beta} \right) \left( \frac{150 \text{ GeV}}{\mu} \right)$$

- $\Rightarrow M_{\chi_1^0} = 0$  consistent in MSSM

# $M_{\chi_1^0} = 0$ consistent with all lab

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- Invisible  $Z^0$ -width
- $e^+e^- \rightarrow \chi_1^0\chi_1^0\gamma$
- $e^+e^- \rightarrow \chi_2^0\chi_1^0; \quad \chi_2^0 \rightarrow Z^0\chi_1^0; \quad Z^0 \rightarrow q\bar{q}$
- Precision Observables ( $\delta\Gamma_{\text{inv}}, \delta\Gamma_Z, M_W, \delta a_\mu, \text{EDM}$ )
- Monojets
- Rare Meson Decays
- Supernova Cooling
- Dark Matter: Lee–Weinberg bound:  $M_{\chi_1^0} \gtrsim 6 \text{ GeV}$   
Cowsik–McClellan bound:  $M_{\chi_1^0} \lesssim 0.7 \text{ GeV}$

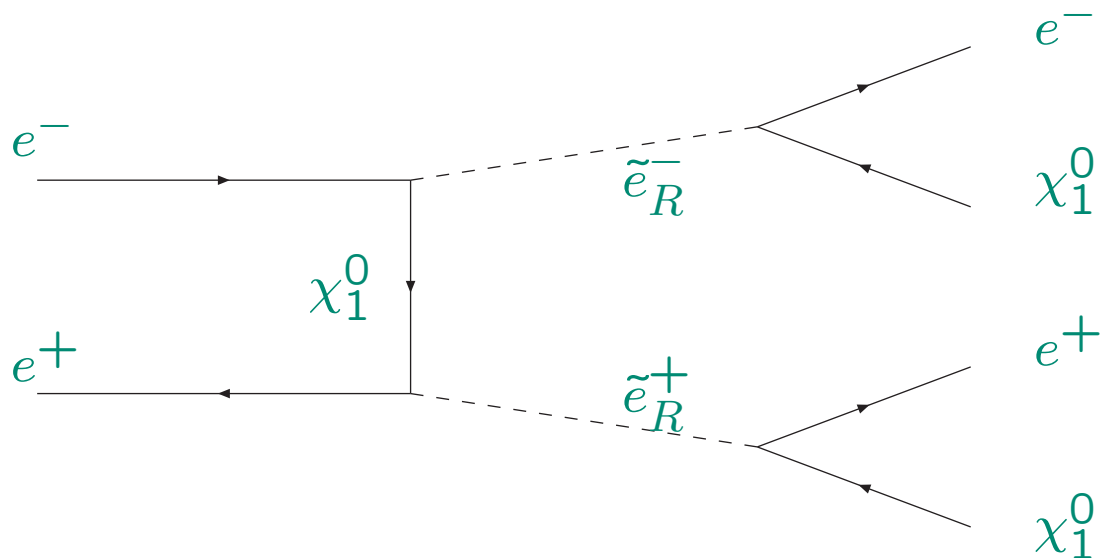
# Publications

- *A Supersymmetric solution to the KARMEN time anomaly*  
D. Choudhury, HD, P. Richardson, Subir Sarkar  
Phys. Rev. D61:095009,2000; e-Print: hep-ph/9911365
- *Supernovae and light neutralinos: SN1987A bounds on super*  
HD, C. Hanhart, U. Langenfeld, D.R. Phillips  
Phys. Rev. D68:055004,2003; e-Print: hep-ph/0304289
- *Discovery potential of radiative neutralino production at the*  
HD, O. Kittel, U. Langenfeld  
Phys. Rev. D74:115010,2006; e-Print: hep-ph/0610020
- *Mass Bounds on a Very Light Neutralino*  
HD, S. Heinemeyer, O. Kittel, U. Langenfeld, A.M. Weber, G  
Eur.Phys.J.C62,2009; e-Print: arXiv:0901.3485
- *Rare Meson Decays to a Light neutralino*  
HD, S. Grab, D. Koschade, M. Krämer, U. Langenfeld, B. O'L  
Phys. Rev. D80:035018,2009

# Measuring a Light Neutralino Mass

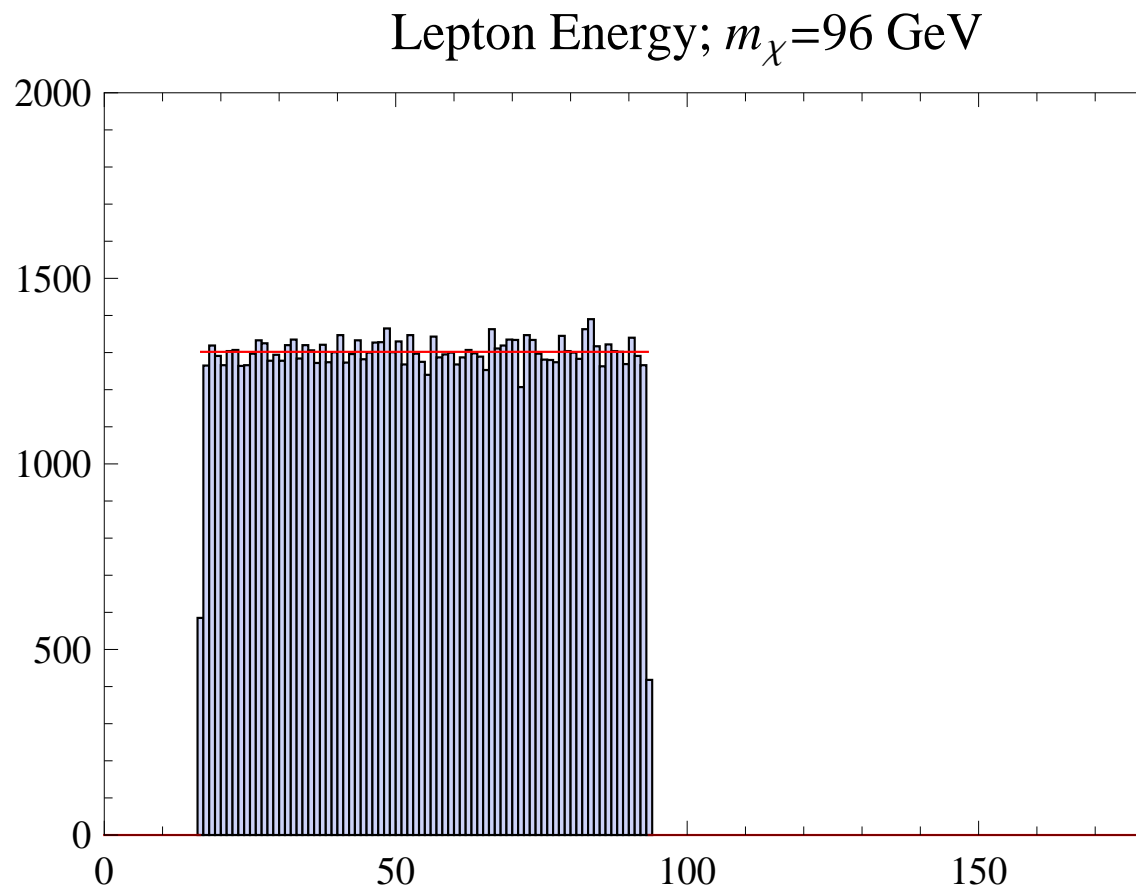
Work in progress with Bonn group: John Conley, HD, Peter W

- Consider the process:



- Measure  $e^\pm$  energies  $\implies$  determine  $M_{\chi_1^0}$
- Earlier work by Uli Martyn (DESY):  $M_{\chi_1^0} \gtrsim 95 \text{ GeV}$
- How well does this work for light neutralinos?

# Electron Energy Distribution



- Electron energy distribution is flat, with sharp cut
- Here chosen: SPS1a:  $M_{\tilde{e}_R} = 143$  GeV and  $\sqrt{s} = 4$



# Simple Kinematics

- $\theta_0$ : angle between  $\vec{p}(e)$  in slepton rest-frame and

$$E_e = \frac{\sqrt{s}}{4} \left( 1 - \frac{M_{\chi_1^0}^2}{M_{\tilde{e}}^2} \right) (1 + \beta_{\tilde{e}} \cos \theta_0), \quad \beta_{\tilde{e}}$$

- Max/Min Electron energy:  $E_{\pm}$  for  $\cos \theta_0 = \pm 1$
- Solve for the SUSY Masses

$$M_{\tilde{e}} = \sqrt{s} \frac{\sqrt{E_+ E_-}}{E_+ + E_-}, \quad M_{\chi_1^0} = M_{\tilde{e}} \sqrt{1 - \frac{E_+}{\sqrt{s}}}$$

- Measure  $E_+$  and  $E_-$ : thus determine  $M_{\chi_1^0}$

# Neutralino Mass Sensitivity

$$E_{\pm} = \frac{\sqrt{s}}{4} \left( 1 - \frac{M_{\chi_1^0}^2}{M_{\tilde{e}}^2} \right) (1 \pm \beta_{\tilde{e}}), \quad \beta_{\tilde{e}} = \sqrt{1 - \frac{M_{\chi_1^0}^2}{M_{\tilde{e}}^2}}$$

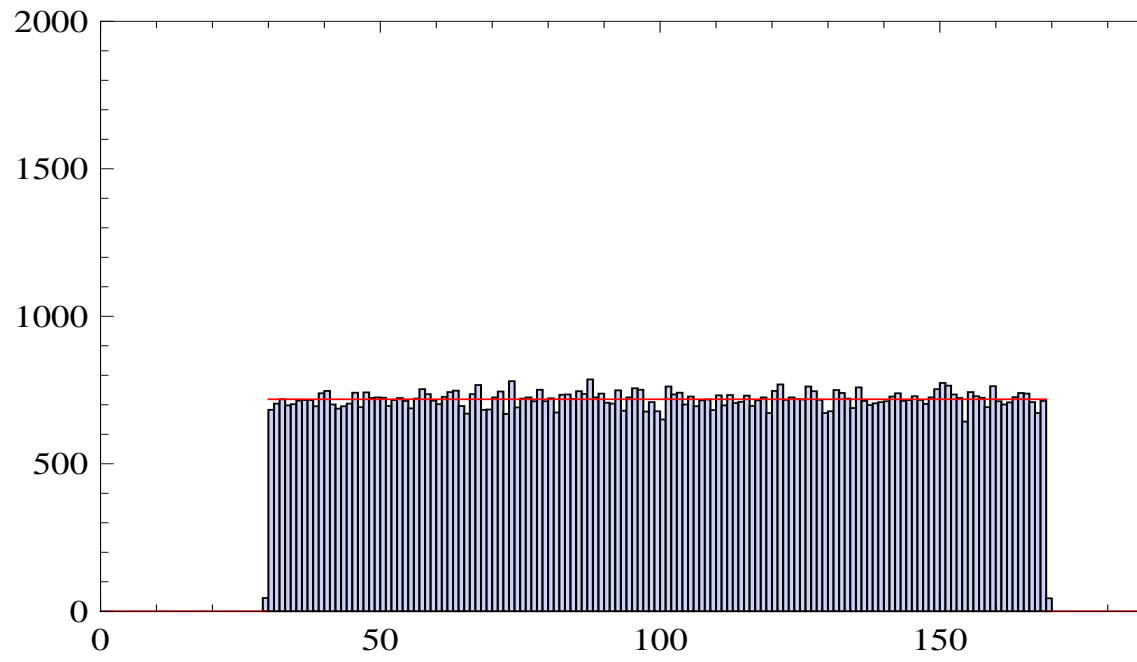
- Detailed ILC study by Uli Martyn for heavy neutralinos

$$\frac{\Delta M_{\chi_1^0}}{M_{\chi_1^0}} < 0.2\%$$

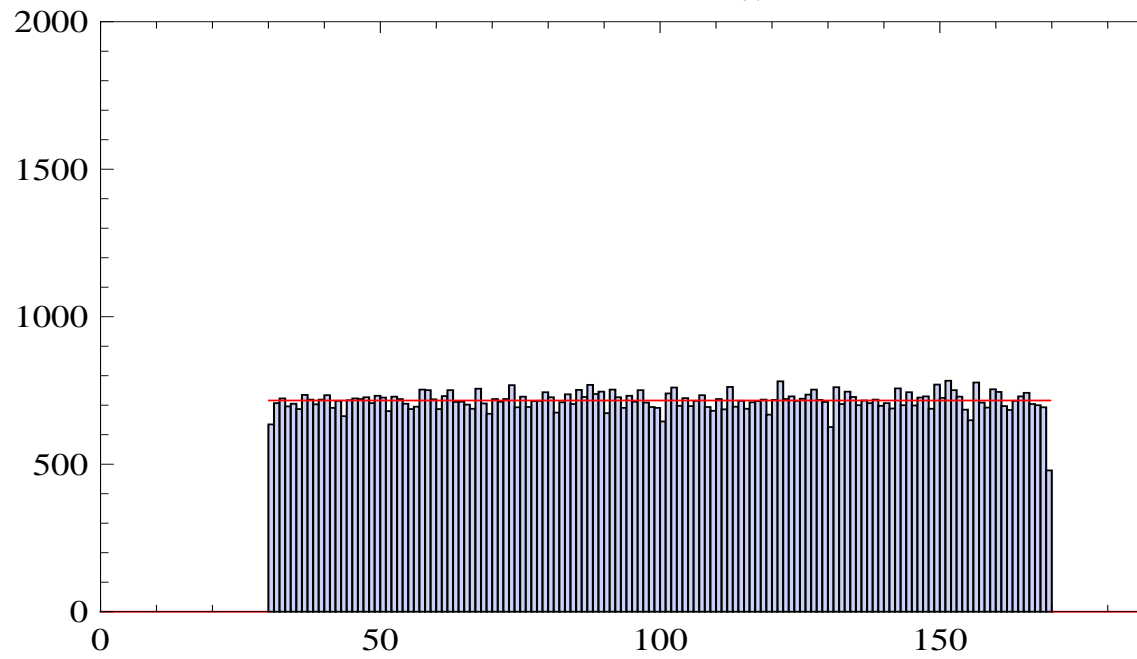
- $E_{\pm}$  depends on  $\frac{M_{\chi_1^0}^2}{M_{\tilde{e}}^2} \implies$  Expect less sensitivity

$$M_{\chi_1^0} \ll M_{\tilde{e}} \quad \text{or as} \quad M_{\chi_1^0} \longrightarrow 0$$

Lepton Energy;  $m_\chi=10$  GeV



Lepton Energy;  $m_\chi=5$  GeV



## Work in Progress: Simple Sim

- Do not consider full detector simulation. Instead s
- Consider  $\tilde{e}_R^- \tilde{e}_R^+$  and  $\tilde{\mu}_R^- \tilde{\mu}_R^+$ -Production
- $\tilde{e}_R^- \tilde{e}_R^+$  dominant
- $\sqrt{s} = 500$  GeV
- Beam polarisations  $(\mathcal{P}_{e^-}, \mathcal{P}_{e^+}) = (+80\%, -60\%)$
- Include Beam Strahlung:  $\sqrt{s} \longrightarrow \sqrt{s'} < \sqrt{s}$

This smears out the  $E_{\pm}$ -edges

## Further Details of Simula

- Approximate detector resolutions

$$\Delta \left( \frac{1}{p_T} \right) = 1 \cdot 10^{-4} \text{ GeV}^{-1} \quad (\text{tra})$$

$$\frac{\Delta E}{E} = \frac{0.166}{\sqrt{E/\text{GeV}}} \oplus 0.011 \quad (\text{EC})$$

- Smear electron energy according to minimum of t
- For muons always choose momentum resolution
- This further smoothes out the edges
- We then fit the edges using basically the convolution and an upward or downward step function.

# Fit Functions

$$f_{-}(E) = \begin{cases} \frac{1}{2} \left[ \operatorname{erf}\left(\frac{E - \hat{E}_{-}}{\sqrt{2}\sigma_{1}^{-}}\right) + 1 \right] & : E < \hat{E}_{-} \\ \frac{1}{2} \left[ \operatorname{erf}\left(\frac{E - \hat{E}_{-}}{\sqrt{2}\sigma_{2}^{-}}\right) + 1 \right] & : E > \hat{E}_{-} \end{cases}$$

$$f_{+}(E) = \begin{cases} \frac{1}{2} \left[ \operatorname{erfc}\left(\frac{E - \hat{E}_{-}}{\sqrt{2}\sigma_{1}^{-}}\right) \right] & : E < \hat{E}_{-} \\ \frac{1}{2} \left[ \operatorname{erfc}\left(\frac{E - \hat{E}_{-}}{\sqrt{2}\sigma_{2}^{-}}\right) \right] & : E > \hat{E}_{-} \end{cases}$$

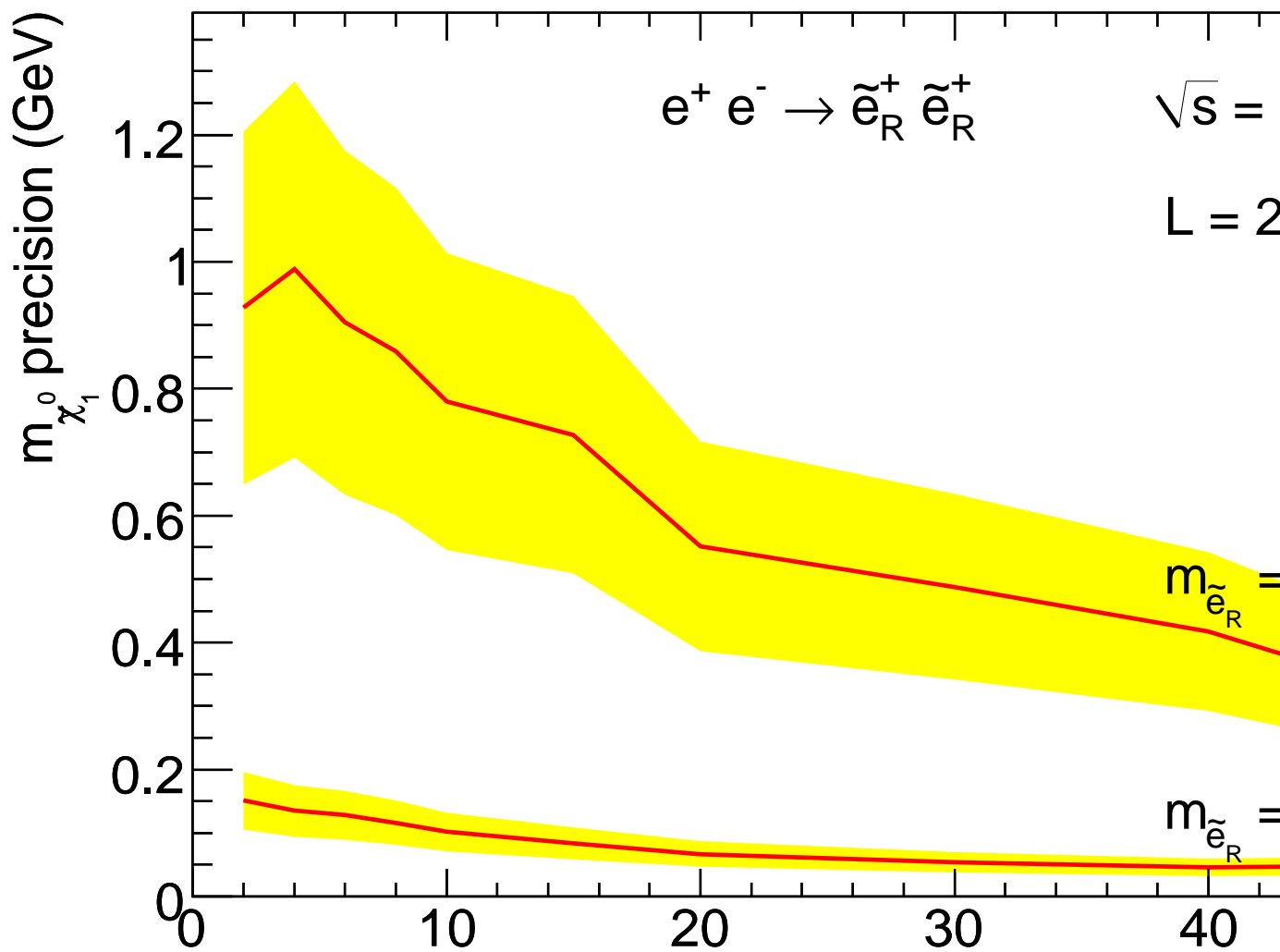
$$\operatorname{erf}(x) = \frac{2}{\sqrt{\pi}} \int_0^x e^{-t^2} dt$$

$$\operatorname{erfc}(x) = 1 - \operatorname{erf}(x)$$

- $\sigma_1 \neq \sigma_2$  because of asymmetric beam strahlung

## Results

- Reproduced detailed detector study by Uli Martyn and  $\sqrt{s} = 400 \text{ GeV}$  to within  $\pm 30\%$
- Used this as a systematic error for our analysis of



- $\sqrt{s} = 500$  GeV

- Integrated Luminosity:  $\int \mathcal{L} dt = 250$  fb $^{-1}$

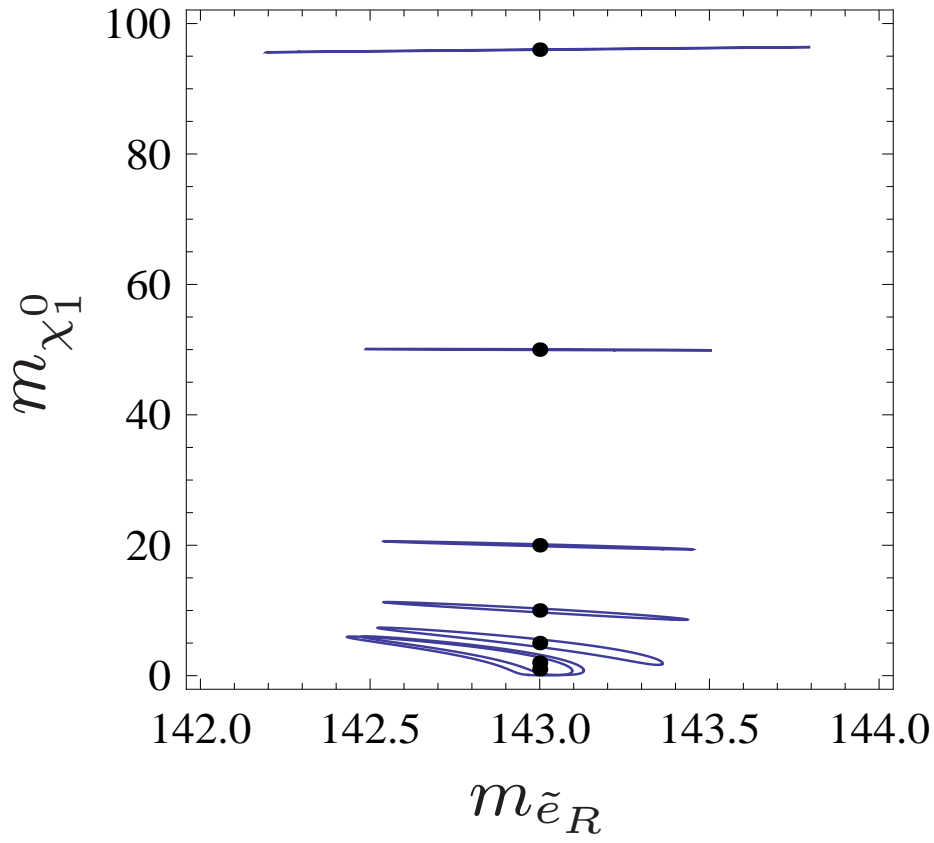


## Summary & Conclusion

- A massless neutralino is consistent with all data
- For  $M_{\tilde{e}_R} = 100$  GeV can measure  $M_{\chi_1^0}$  down to less
- For  $M_{\tilde{e}_R} = 200$  GeV can measure  $M_{\chi_1^0}$  down to ab
- Implications for Dark Matter



95% CL contours



95% CL contours

