

# SUSY-Yukawa Sum Rule at the LHC

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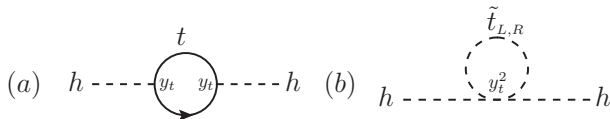


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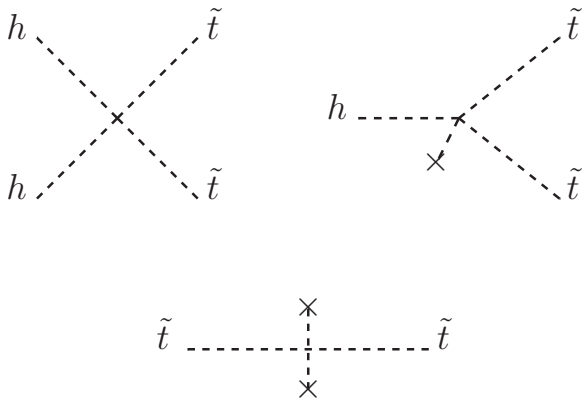
Monday, May 10 2010

- **Hierarchy problem:** In the SM, Higgs mass receives quadratically divergent corrections, most importantly from the top quark
- In **SUSY**, top contribution cancelled by stop



- This relies on both **particle content** and **coupling relations**. **We want to test the coupling relations.**

# How to probe the Quartic Higgs Coupling?



**Look at diagonal sfermion mass terms!**

# SUSY-Yukawa Sum Rule

Look at stop/sbottom  $LL$  mass terms at tree level:

$$M_{\tilde{t}_L \tilde{t}_L}^2 = M_L^2 + \hat{m}_t^2 + g_{uL} \hat{m}_Z^2 \cos 2\beta = m_{t1}^2 c_t^2 + m_{t2}^2 s_t^2 \quad (1)$$

$$M_{\tilde{b}_L \tilde{b}_L}^2 = M_L^2 + \hat{m}_b^2 + g_{bL} \hat{m}_Z^2 \cos 2\beta = m_{b1}^2 c_b^2 + m_{b2}^2 s_b^2 \quad (2)$$

Soft masses Higgs Quartic Coupling D-term contributions measurable

(1) – (2) eliminates the soft mass:

$$\hat{m}_t^2 - \hat{m}_b^2 = m_{t1}^2 c_t^2 + m_{t2}^2 s_t^2 - m_{b1}^2 c_b^2 - m_{b2}^2 s_b^2 - \hat{m}_Z^2 \cos^2 \theta_w \cos 2\beta$$

We call this the **SUSY-Yukawa Sum Rule**: It has its origins in the same coupling relations that cancel higgs mass corrections.

**We want to test this sum rule at a collider!**

# Defining an observable to test the Sum Rule

- Assume **SUSY-like particle content** ( $\tilde{t}_L, \tilde{b}_L$ ),  $\tilde{t}_R, \tilde{b}_R$  but **not the SUSY coupling relations**.
- Before EWSB,  $M_{\tilde{t}}^2 = \begin{pmatrix} M_L^2 & \\ & M_t^2 \end{pmatrix}$ ,  $M_{\tilde{b}}^2 = \begin{pmatrix} M_L^2 & \\ & M_b^2 \end{pmatrix}$
- After EWSB, can parameterize quartic higgs coupling 'model-independently':

$$(M_{\tilde{t}}^2)_{11} \rightarrow M_L^2 + v^2 Y_{11}^t \quad , \quad (M_{\tilde{b}}^2)_{11} \rightarrow M_L^2 + v^2 Y_{11}^b$$

Define a new observable to probe the quartic higgs coupling:

$$\begin{aligned} \Upsilon &\equiv \frac{1}{v^2} \left( m_{t1}^2 c_t^2 + m_{t2}^2 s_t^2 - m_{b1}^2 c_b^2 - m_{b2}^2 s_b^2 \right) \\ &= Y_{11}^t - Y_{11}^b \quad \text{at tree level} \end{aligned}$$

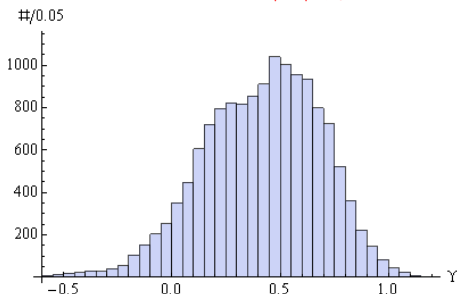
# SUSY prediction for $\Upsilon$ & Radiative Corrections

## Tree-Level Prediction for $\Upsilon$ from SUSY-Yukawa Sum Rule

$$\begin{aligned}\Upsilon_{\text{SUSY}}^{\text{tree}} &= \frac{1}{\sqrt{2}} \left( \hat{m}_t^2 - \hat{m}_b^2 + m_Z^2 \cos^2 \theta_W \cos 2\beta \right) \\ &= \begin{cases} 0.39 & \text{for } \tan \beta = 1 \\ \mathbf{0.28} & \text{for } \tan \beta \rightarrow \infty \text{ (converges quickly for } \tan \beta \gtrsim 5) \end{cases}\end{aligned}$$

- In a generic theory, only 'requirement' is  $|\Upsilon| \lesssim 16\pi$ .
- Radiative Corrections wash out SUSY tree-level prediction for  $\Upsilon$ .
- Worst case scenario (SuSpect)  $\rightarrow$
- Can narrow predicted range with more measurements (see later).

**TeV-scale SUSY:  $|\Upsilon| \lesssim 1$ .**



# Prospects at the LHC

- Fully measuring  $\Upsilon$  requires lepton collider.
- Can make some progress at LHC in favorable regions of MSSM parameter space.  $\Rightarrow$  Could then use  $\Upsilon$  to constrain stop/sbottom parameters.
- Demonstrate feasibility of partial  $\Upsilon$ -measurement with a particular Benchmark Point:

## Parameters:

$\tan \beta$	$M_1$	$M_2$	$M_3$	$\mu$	$M_A$	$M_{Q3L}$	$M_{tR}$	$A_t$
10	100	450	450	400	600	310.6	778.1	392.6

## Spectrum: (GeV)

$m_{t1}$	$m_{t2}$	$s_t$	$m_{b1}$	$m_{b2}$	$s_b$	$m_{\tilde{g}}$	$m_{\tilde{\chi}_1^0}$
371	800	-0.095	341	1000	-0.011	525	98

# Measuring part of $\Upsilon$

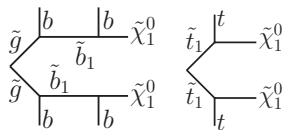
- **Small mixing angles** and **light  $\tilde{t}_1, \tilde{b}_1$**   $\implies$  rewrite

$$\Upsilon = \underbrace{\frac{1}{v^2} (m_{\tilde{t}1}^2 - m_{\tilde{b}1}^2)}_{\Upsilon'} + \underbrace{\frac{s_t^2}{v^2} (m_{\tilde{t}2}^2 - m_{\tilde{t}1}^2)}_{\Delta\Upsilon_t} - \underbrace{\frac{s_b^2}{v^2} (m_{\tilde{b}2}^2 - m_{\tilde{b}1}^2)}_{\Delta\Upsilon_b}$$

- Most of  $\Upsilon = 0.423$  comes from  $\Upsilon' = 0.350$ .  $\Delta\Upsilon_t \lesssim O(0.1)$  can be estimated <sup>1</sup>.  $\Delta\Upsilon_b$  can't be measured at LHC.

- **We will measure  $\Upsilon'$**

- Need to determine  $m_{\tilde{t}1}, m_{\tilde{b}1}$
- Analyse gluino & stop pair production & decay
- Extract kinematic- and  $M_{T2}$ -edges to **get all the masses  $\implies \Upsilon'$**

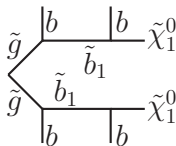


<sup>1</sup>MP, Weiler 2008



# (I) Gluino Pair Production

- Analyze the process<sup>2</sup>  
 $\tilde{g}\tilde{g} \rightarrow 2\tilde{b}_1 + 2b \rightarrow 4b + 2\tilde{\chi}_1^0$ .
- $\sigma_{\tilde{g}\tilde{g}} \approx 11.6 \text{ pb @ } \sqrt{s} = 14 \text{ TeV}$ .

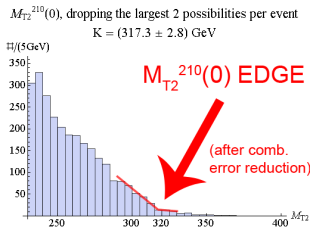


- Impose basic  $p_T$ , MET-cuts and require **4  $b$ -tags**.
- Use  $\mathcal{L} = 10 \text{ fb}^{-1}$ . After cuts we are left with 4800 signal events.
- No SUSY-BG. SM-BG suppressed by  $b$ -tag requirement.
- Even with parton-level pure signal, full mass extraction is challenging!

<sup>2</sup>MadGraph/Madevent & BRIDGE

# Edge Extraction & Mass Measurement

- To measure masses at hadron colliders with **invisible massive particles in the final state**, we go **Edge Hunting!**
- Distributions of  $M_{T2}$ -**subsystem-variables**<sup>3</sup> and  $M_{bb}$  show **edges** which tell us mass combinations.
- Big Problem: **Combinatorial Error** (especially for  $M_{T2}$ 's).



We are able to successfully measure  $M_{bb}$ ,  $M_{T2}^{210}(0)$  and  $M_{T2}^{220}(0)$  edges

⇒

mass	th.	68 % c.l.
$m_{b1}$	341	(316, 356)
$m_{\tilde{g}}$	525	(508, 552)
$m_{\tilde{\chi}_1^0}$	98	(45*, 115)

<sup>3</sup>Barr, Lester, Stephens, 2003; Cho, Choi, Kim, Park 2008; Burns, Kang, Matchev, Park 2009

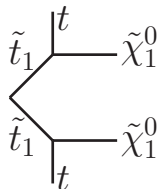
\* LEP bound

## (II) Stop Pair Production

- Analyze the process  $\tilde{t}_1 \tilde{t}_1^* \rightarrow t\bar{t} + 2\tilde{\chi}_1^0$ .

- $\sigma_{\tilde{t}_1 \tilde{t}_1^*} \approx 2 \text{ pb} @ \sqrt{s} = 14 \text{ TeV}$ .

- Impose standard cuts & use hadronic tops<sup>4</sup>.



- Use  $\mathcal{L} = 100 \text{ fb}^{-1}$ . After cuts: 1481 signal and 105 BG events.

- Easy to extract  $M_{T2}^{\text{max}}$  edge  $\implies$  Gives  $m_{t\tilde{t}_1}(m_{\tilde{\chi}_1^0})$

- Combine with (I)  $\implies$

	th.	68 % c.l.
$m_{t\tilde{t}_1}$	371	(356, 414)

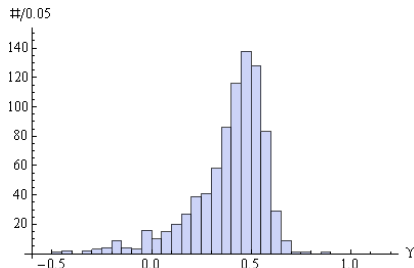
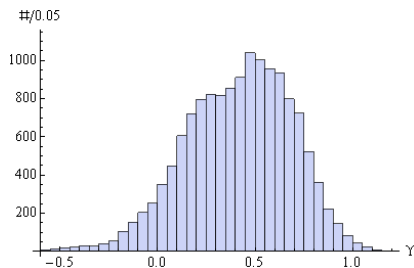
<sup>4</sup>Meade, Reece 2006

# $\Upsilon'$ Measurement and SUSY-prediction for $\Upsilon$

Putting all these measurements together, we get

	th.	meas.
$\Upsilon'$	0.350	$0.525^{+0.20}_{-0.15}$
$\Upsilon$	0.423	—

The measurements of the  $\tilde{b}_1, \tilde{t}_1, \tilde{g}, \tilde{\chi}_1^0$  masses also allow us to make the SUSY-prediction for  $\Upsilon$  more precise:



# Summary & Conclusions

- Confirmation of **SUSY-Yukawa Sum Rule**

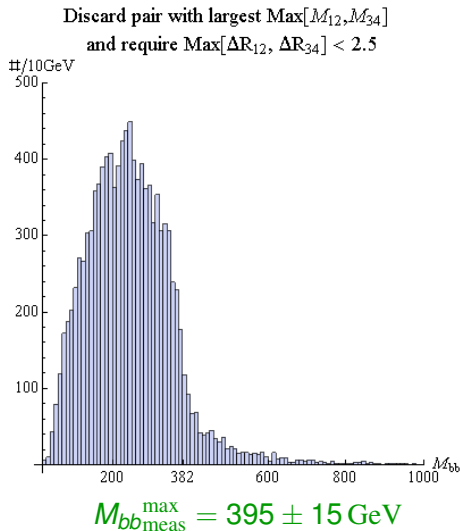
$$\hat{m}_t^2 - \hat{m}_b^2 = m_{t_1}^2 c_t^2 + m_{t_2}^2 s_t^2 - m_{b_1}^2 c_b^2 - m_{b_2}^2 s_b^2 - \hat{m}_Z^2 \cos^2 \theta_w \cos 2\beta$$

would be strong support for TeV-scale SUSY as the solution for hierarchy problem.

- Full measurement will have to wait for **Lepton Collider**.
- Can make **significant progress at LHC** in some regions of parameter space.
- We developed **new techniques for reducing  $M_{T2}$ -combinatorial background**, allowing us to measure  $\tilde{t}_1, \tilde{b}_1, \tilde{g}, \tilde{\chi}_1^0$  masses at our benchmark point.

# Glino Pair Production: Kinematic Edge

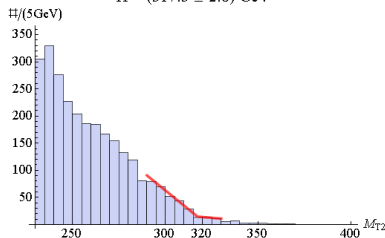
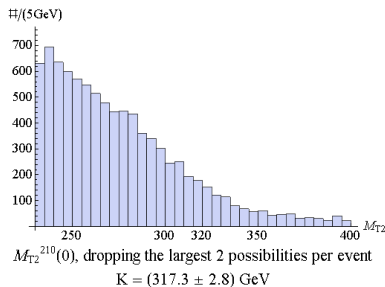
- $M_{bb}^{\max} = \sqrt{\frac{(m_{\tilde{g}}^2 - m_{b_1}^2)(m_{b_1}^2 - m_{\tilde{\chi}_1^0}^2)}{m_{b_1}^2}}$
- With known decay chain assignments get  $(M_{b_1 b_2}, M_{b_3 b_4})$  for each event, plot  $M_{bb}$ -distribution  $\Rightarrow$  edge at 382 GeV.
- Main problem: **Combinatorial Background!**
- Can reduce CB with  $\Delta R$  cuts and dropping largest  $M_{bb}$ 's per event.



# Glino Pair Production: $M_{T2}$ -subsystem Edges

- The distributions of  $M_{T2}$  subsystem variables<sup>5</sup> also have edges we can measure. Look at  $M_{T2}^{210}(0)$ .
- **Combinatorial Background** is more dangerous.
  - To calculate  $M_{T2}^{210}$ , have to divide  $4b$  into an **upstream** and **downstream** pair: **6 possibilities**.
  - The  $M_{T2}$ -distribution for **wrong** pairings is **more featured** than  $M_{bb}$ .
- One way to reduce CB:  
**Drop largest 2  $M_{T2}^{210}$ 's per event** →

$M_{T2}^{210}(0)$  without combinatorial error reduction



<sup>5</sup>Barr, Lester, Stephens, 2003; Cho, Choi, Kim, Park 2008; Burns, Kang, Matchev, Park 2009

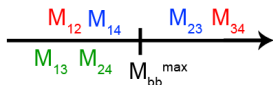
# Glino Pair Production: $M_{T2}$ -subsystem Edges

Another way to reduce CB:

Use Kinematic Edge Measurement!

Possible  $M_{bb}$  pairs:  $(M_{12}, M_{34})$ ,  $(M_{13}, M_{24})$ ,  $(M_{14}, M_{23})$

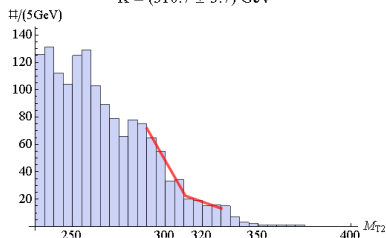
For  $\sim 30\%$  of events, situation like:



Can deduce correct decay chain assignment!

$M_{T2}^{210}(0)$  with known decay chain assignment

$K = (310.7 \pm 3.7) \text{ GeV}$



For edge measurement, require two methods to agree!

edge	th.	measurement
$M_{bb}$	382	$395 \pm 15$
$M_{T2}^{210}(0)$	321	$314 \pm 13 \text{ GeV}$
$M_{T2}^{220}(0)$	507	$492 \pm 14 \text{ GeV}$

$\Rightarrow$

mass	th.	68 % c.l.
$m_{b1}$	341	(316, 356)
$m_{\tilde{g}}$	525	(508, 552)
$m_{\tilde{\chi}_1^0}$	98	(45, 115)

(Imposed  $m_{\tilde{\chi}_1^0} > 45 \text{ GeV}$  bound from LEP measurement of invisible  $Z$  decay width.)