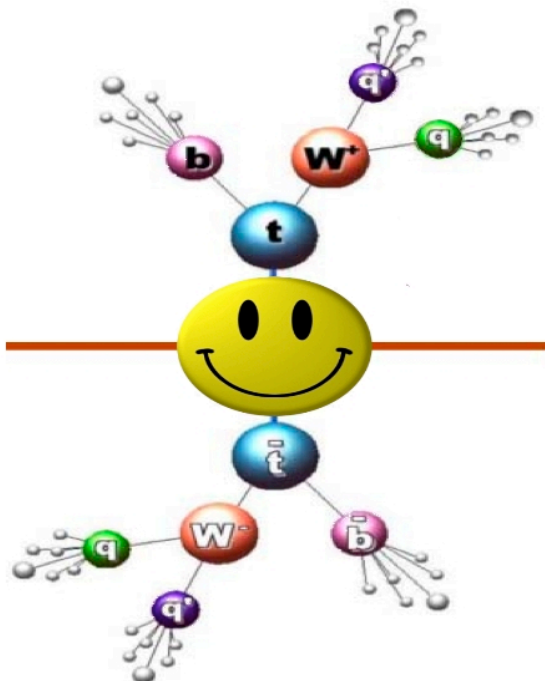




# Resonance search in

$$p\bar{p} \rightarrow X^0 \rightarrow t\bar{t}$$



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# Why and How?



- Goal is to test  $t\bar{t}$  production for possible new sources such as a narrow resonance
  - Top is very heavy, maybe indication of coupling to new physics
  - Top is a young particle
  - Various theoretical models predict it: technicolor, KK gluons
- Search technique:
  - $M_{t\bar{t}}$  spectrum is reconstructed, using FlaME
  - Search for a peak in  $M_{t\bar{t}}$  spectrum
    - Understand SM fluctuation probabilities
    - Calculate UL(Upper Limits)
    - Compare data with our expectations(SM or with new physics)



# Where?



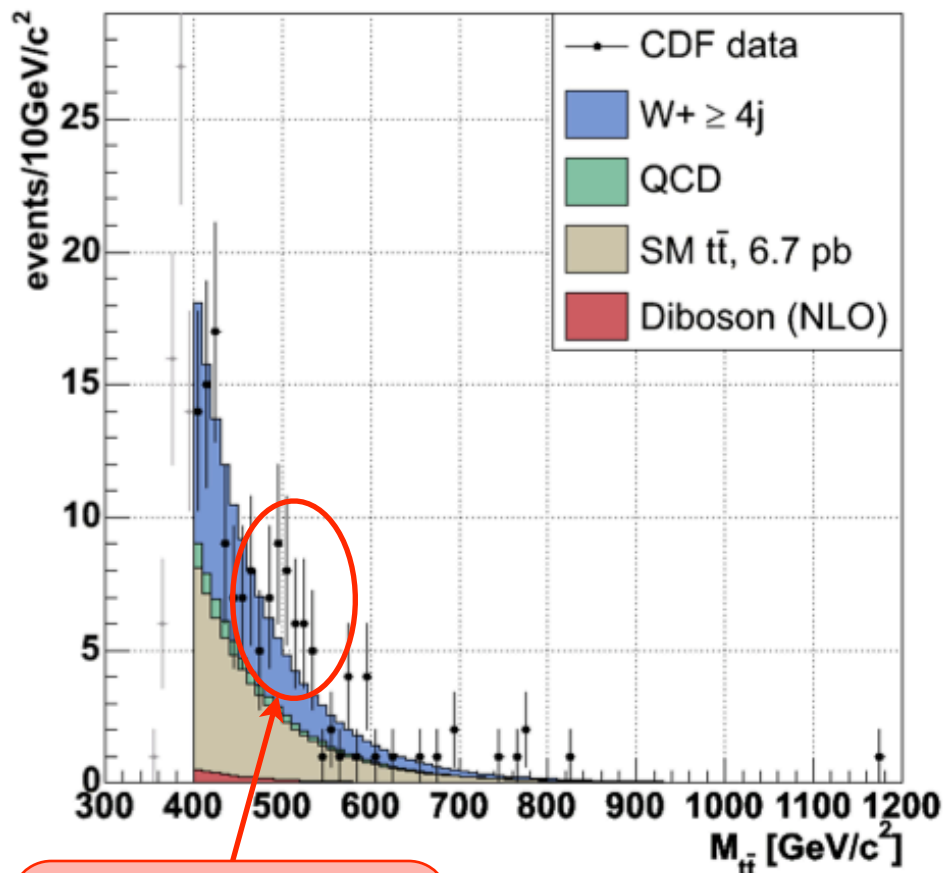
- This is the first  $M_{t\bar{t}}$  analyses in All Hadronic channel
  - Disadvantages
    - Large QCD background
      - » Controlled with good event selection
    - More combinations
  - Advantages
    - Highest branching ratio
      - » Most  $t\bar{t}$  events are here
    - No missing information like neutrino
      - » Better signal templates
  - Future
    - Combined result with lepton+jets channel
      - » Higher sensitivity
    - Cross-check for a possible discovery



# Motivation - previous results

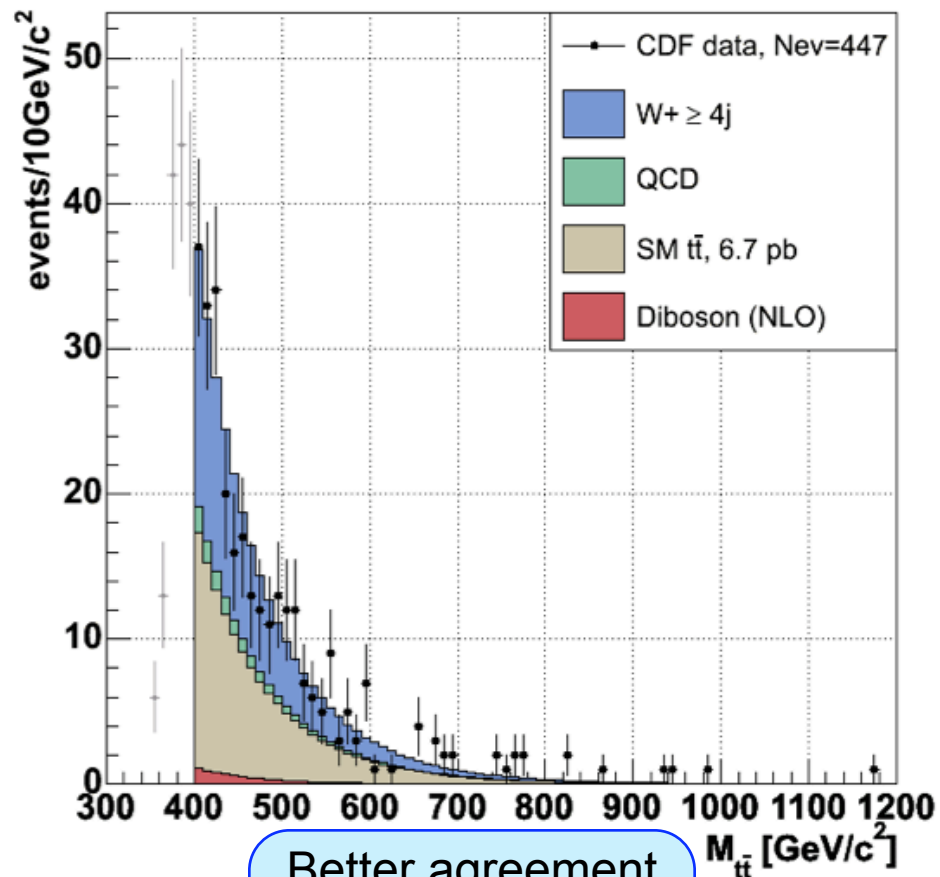


CDF Run 2 preliminary,  $L=319\text{pb}^{-1}$



Excess  $\sim 500\text{GeV}$

CDF Run 2 preliminary,  $L=682\text{pb}^{-1}$



Better agreement with SM :(



# FlaME (Florida Matrix Element)



We calculate the *a priori* probability density for an event to be the result of Standard Model  $t\bar{t}$  production and decay

$$P(j | M_{top}) = \frac{1}{\sigma(M_{top})\epsilon(M_{top})N_{combi}} \sum_{combi} \int dz_a dz_b f(z_a) f(z_b) d\sigma(M_{top}, p) TF(j | p) P_T(p)$$

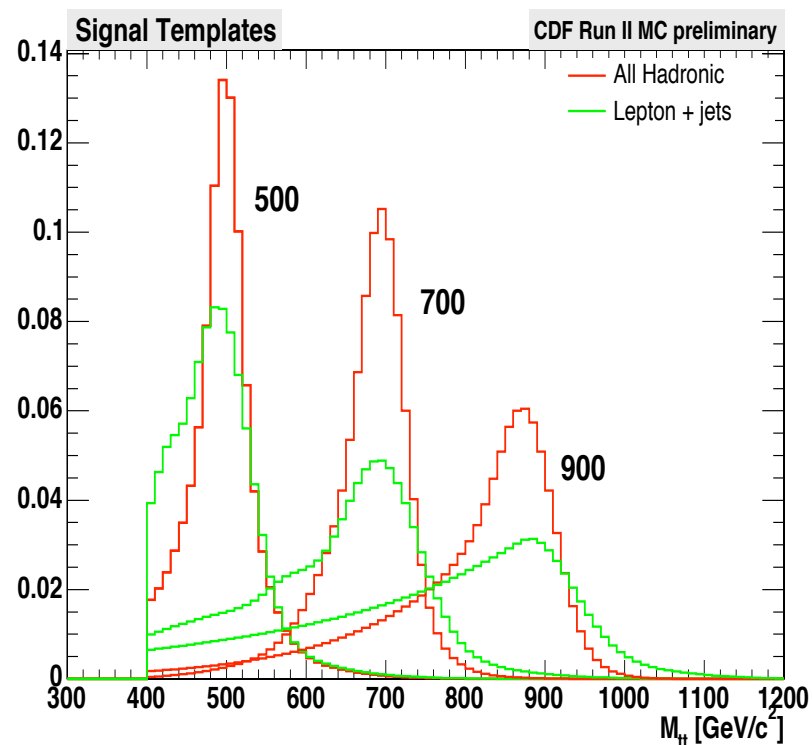
To calculate the  $M_{t\bar{t}}$  probability density, we modify the integral above:

$$\rho(x | j) = \frac{1}{\sigma(M_{top})\epsilon(M_{top})N_{combi}} \sum_{combi} \int dz_a dz_b f(z_a) f(z_b) d\sigma(M_{top}, p) TF(j | p) P_T(p) \delta(x - M_{t\bar{t}}(p))$$

As  $M_{t\bar{t}}$  estimator we use average of this distribution:

$$M_{t\bar{t}} = \langle \rho(x | j) \rangle$$

- Signal samples:
  - Pythia generated narrow resonant  $t\bar{t}$  samples with masses 450, 500 ... 900 GeV
- Background Samples:
  - SM  $t\bar{t}$  MC sample
  - QCD
    - Data driven





# Trigger & Prerequisites



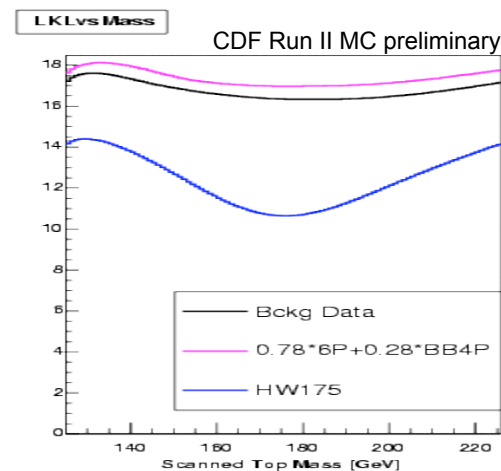
- Multijet Trigger
  - L1:  $\geq 1$  tower with  $E_T \geq 10$  GeV
  - L2:  $\geq 4$  clusters with  $E_T^{cl} \geq 15$  GeV,  $\Sigma E_T \geq 125$  GeV
  - L3:  $N_{jet} \geq 4$ , with  $E_T^{jet} \geq 10$  GeV
    - $\sigma \approx 14$  nb,  $\sim 85\%$  all hadronic efficiency
- Prerequisites
  - Good run list
  - Vertex:  $|z| < 60$ cm &  $|z - z_p| < 5$ cm
  - Missing Et Significance:  $< 3$  (GeV)<sup>1/2</sup>
  - Tight lepton veto
  - 6,7 tight jets -  $E_T^{jet} \geq 15$ GeV,  $|\eta| < 2.0$
- After prerequisites we have  $t\bar{t}/\text{QCD} \sim 1/1000!$



# Neural Net Idea



- Neural net event selection:
  - Uses a Root class TMultiLayerPerceptron
  - 11 inputs, 2 hidden layers with 20/10 nodes and 1 output
- SumEt - total transverse energy
- SumEt3 - sub-leading transverse energy  $\sum E_T - E_{T1} - E_{T2}$
- C - centrality:  $\sum E_T / \sqrt{\hat{s}}$
- A - aplanarity:  $3/2 * (\text{smallest eigenvalue})$  of  $M_{ab} = \sum_j P_a^j P_b^j / \sum_j |\vec{P}^j|$
- $E_N^*$  - geom average of transverse energy of the N-(2 leading jets)
- $E_{T1}^*$  - transverse energy of the leading jet
- $M_{2j}^{\min}$  - the minimum dijet mass
- $M_{2j}^{\max}$  - the maximum dijet mass
- $M_{3j}^{\min}$  - the minimum trijet mass
- $M_{3j}^{\max}$  - the maximum trijet mass
- FlAME variable,  $\sum -\text{Log}(P(M_{\text{top}}=155, 160 \dots 195 \text{ GeV}))$







# QCD background



- We build tag matrix from events from 4,5 jet events.
- Each element in the matrix defined as:

$$P^+(N_{v12}, E_{tjet}, N_{tr,jet}) \equiv \frac{N_{j,tag}(N_{v12}, E_{tjet}, N_{tr,jet})}{N_{j,tgbl}(N_{v12}, E_{tjet}, N_{tr,jet})}$$

- The probability to single/double tag an event:

$$\sum_i \left[ P_i^+ \cdot \prod_{j \neq i} (1 - P_j^+) \right] \quad \sum_{i,j \neq i} \left[ P_i^+ \cdot P_j^+ \cdot \prod_{k \neq i,j} (1 - P_k^+) \right]$$

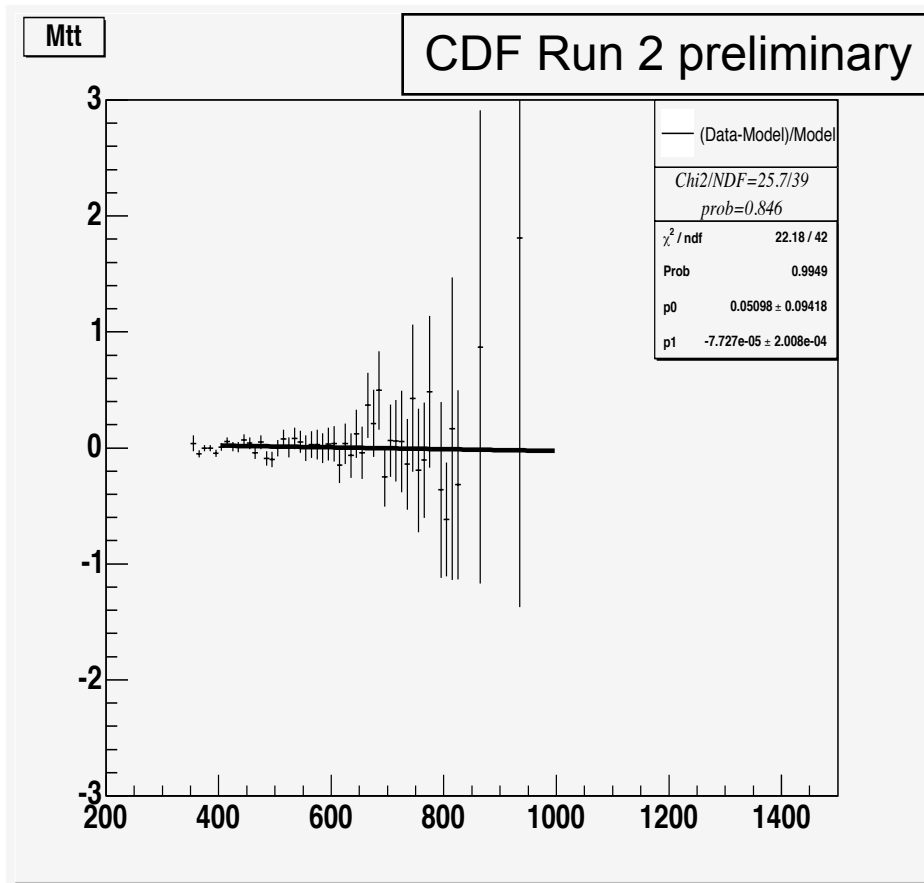
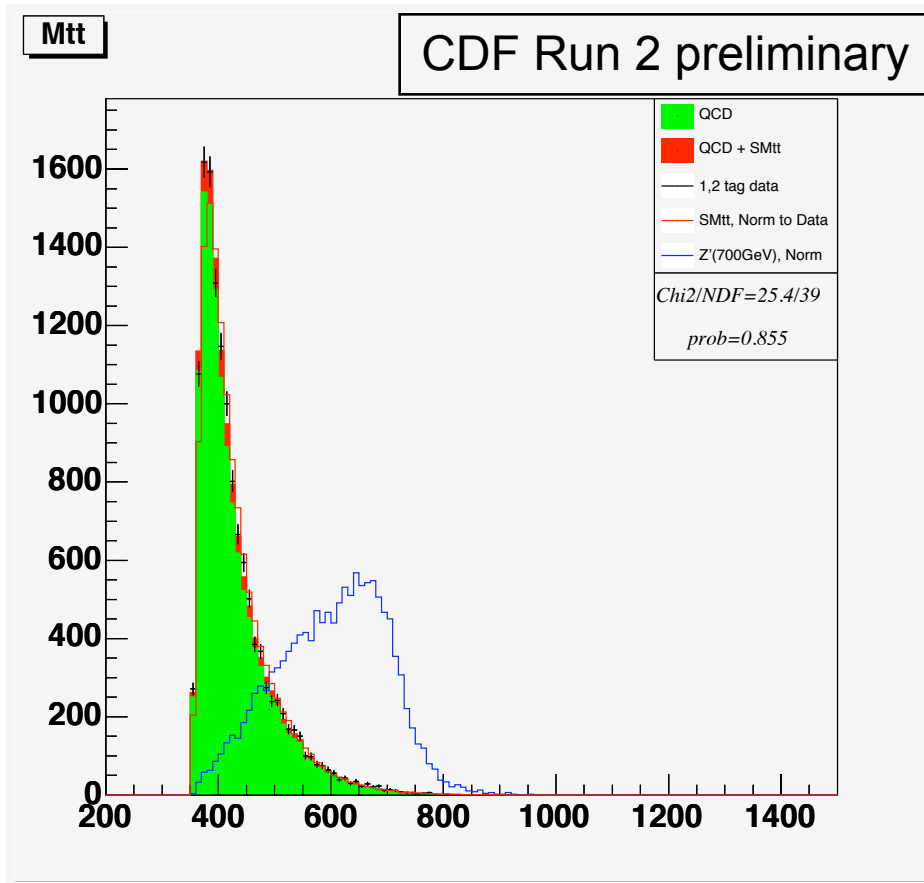
- We weight each event in pre-tagged data sample to get the prediction for 1, 2 tagged events
- Finally, we define several control region and test our modeling with observation
- For all control regions we get a very good agreement
- Biggest impact on final result comes from possible signal contamination, using this procedure



# Crosscheck with data



$0.75 < \text{NNetOut} < 0.93$ .





# Limit Setting Methodology



- Template event weighting
  - $N_{X0}$ : based on assumed cross-section and acceptance
  - $N_{tt}$ : based on theoretical cross-section and acceptance
  - $N_{QCD}$ : Balance from data

$$N_{cdf}^{tot} = \int L dt \cdot (\sigma_{X0} A_{X0} + \sigma_{t\bar{t}} A_{t\bar{t}}) + N_{QCD}$$

- Likelihood

- $N_{X0}$ ,  $N_{tt}$ ,  $N_{QCD}$  are used to compute the expected number of events in mass bin “i”:

$$\mu(i) = N_{X0} T_{X0}(i) + N_{tt} T_{tt}(i) + N_{QCD} T_{QCD}(i)$$

- Given the observed number of events  $n(i)$  and expected  $\mu(i)$  in bin “i”, the likelihood is equal to:

$$L(\sigma_{X0}, \vec{v} | \vec{n}) = \prod e^{-\mu_i} \frac{\mu_i^{n_i}}{n_i!}$$



# Posterior density function



- Acceptance uncertainties accounting

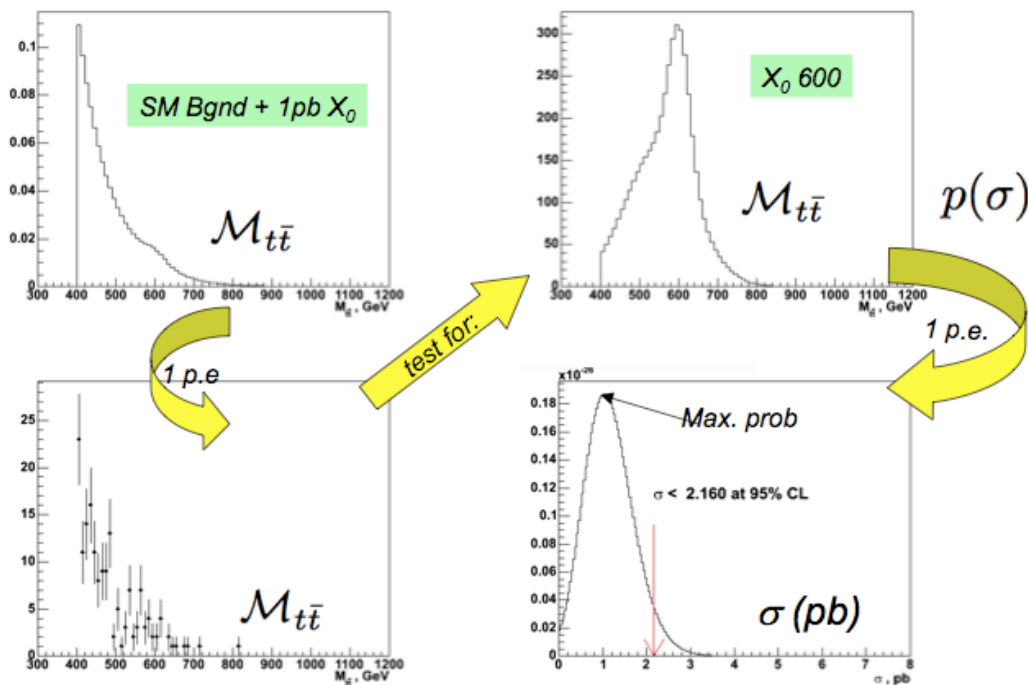
$$p(\sigma_{X_0}, \vec{n}) = \int d\vec{v} \cdot L(\sigma_{X_0}, \vec{v} | \vec{n}) \cdot \pi(\sigma, \vec{v})$$

- We integrate over the nuisance parameters, uncertainties for:

- Signal acceptance
- Background acceptance
- Background cross-section

- Given  $p(\sigma|\vec{n})$  we define:

- $\sigma_{X_0}$  - max of PDF
- 95% confidence level upper limit (UL)  $\frac{1}{Area} \int_0^{UL} p(\sigma | \vec{n}) d\sigma = 0.95$
- Values are calculated as median after 1000 PE's

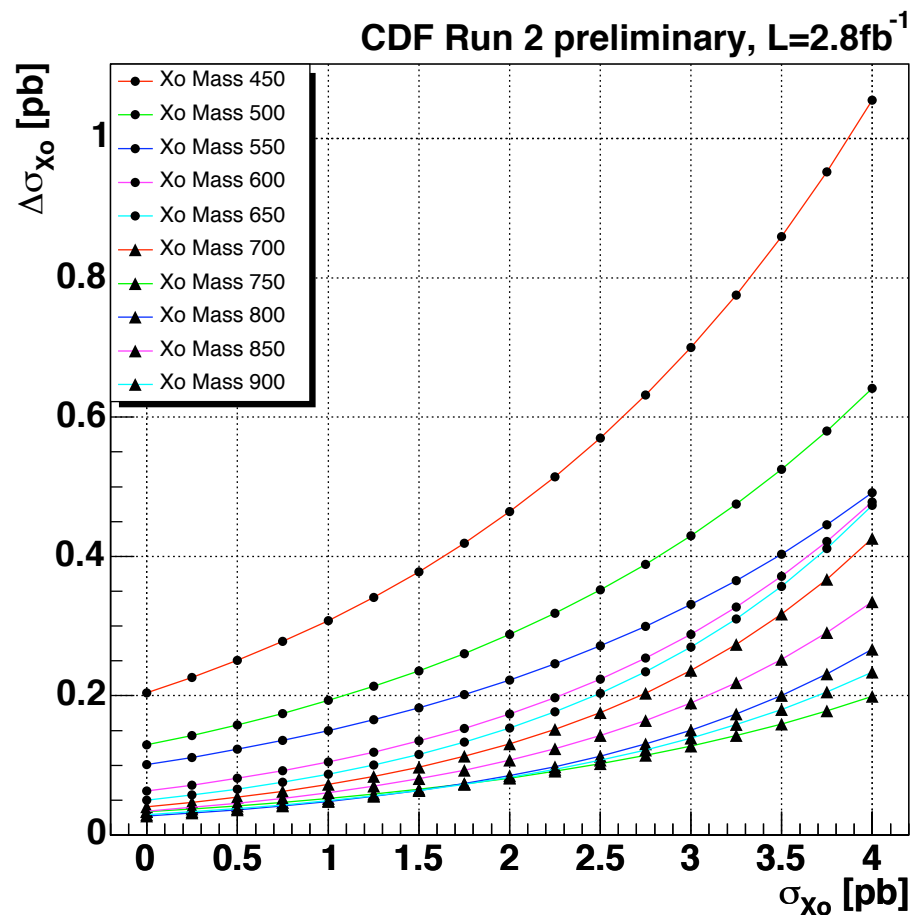




# Systematics



- To consider systematics, which both affect shape and acceptances, we:
- Consider the shift on cross-section by:
- Running PE from shifted templates and fit them with nominal ones
- We considered systematics due to JES, ISR/FSR. PDF found to be negligible

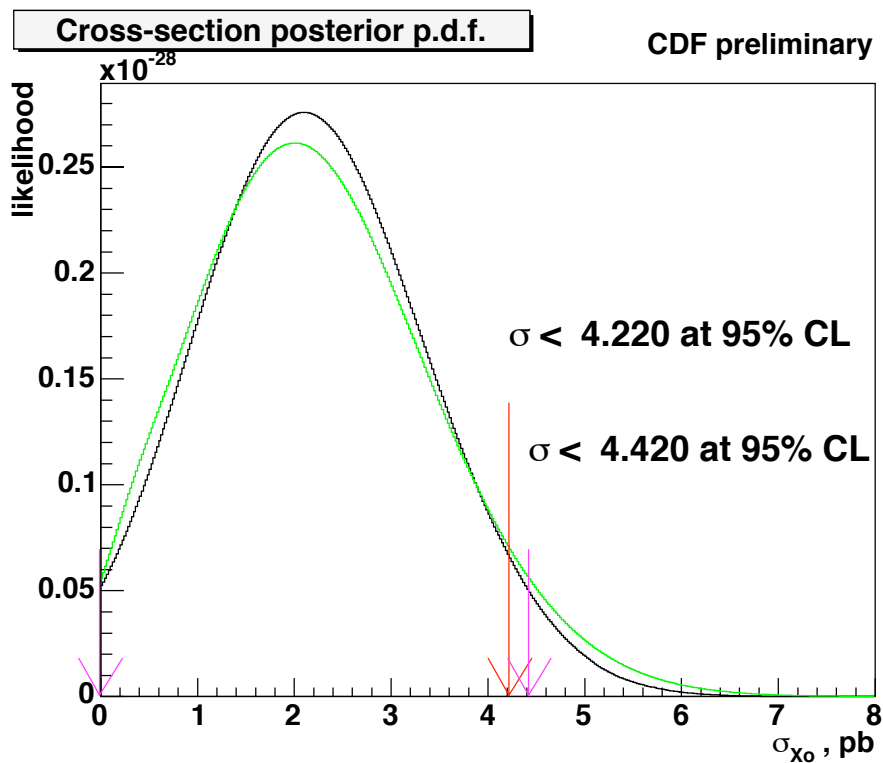




# Applying systematics

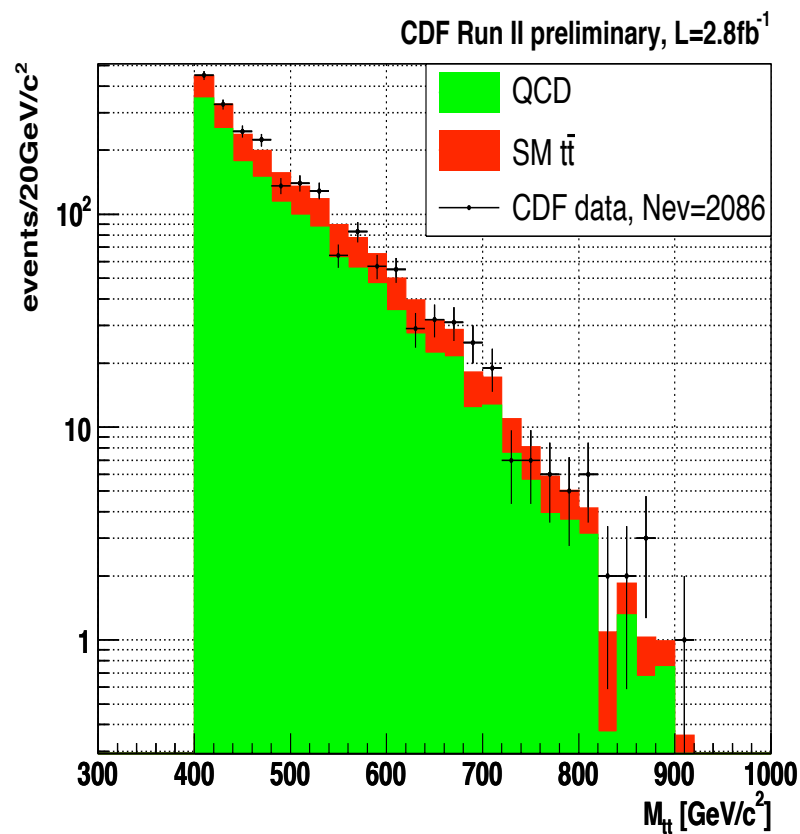
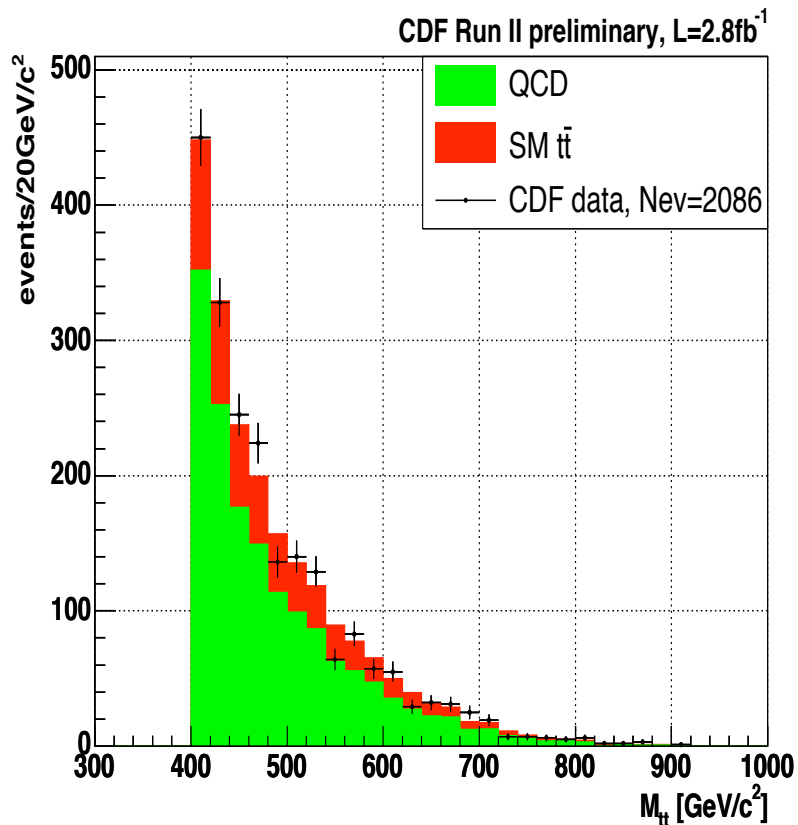


$$PDF_{SYS}(\sigma_{X_0}) = \int_0^{\infty} \frac{1}{\delta\sigma_{X_0} \sqrt{2\pi}} \exp\left(-\frac{1}{2} \left(\frac{\sigma_{X_0} - \sigma'}{\delta\sigma_{X_0}}\right)^2\right) PDF(\sigma') \cdot d\sigma'$$



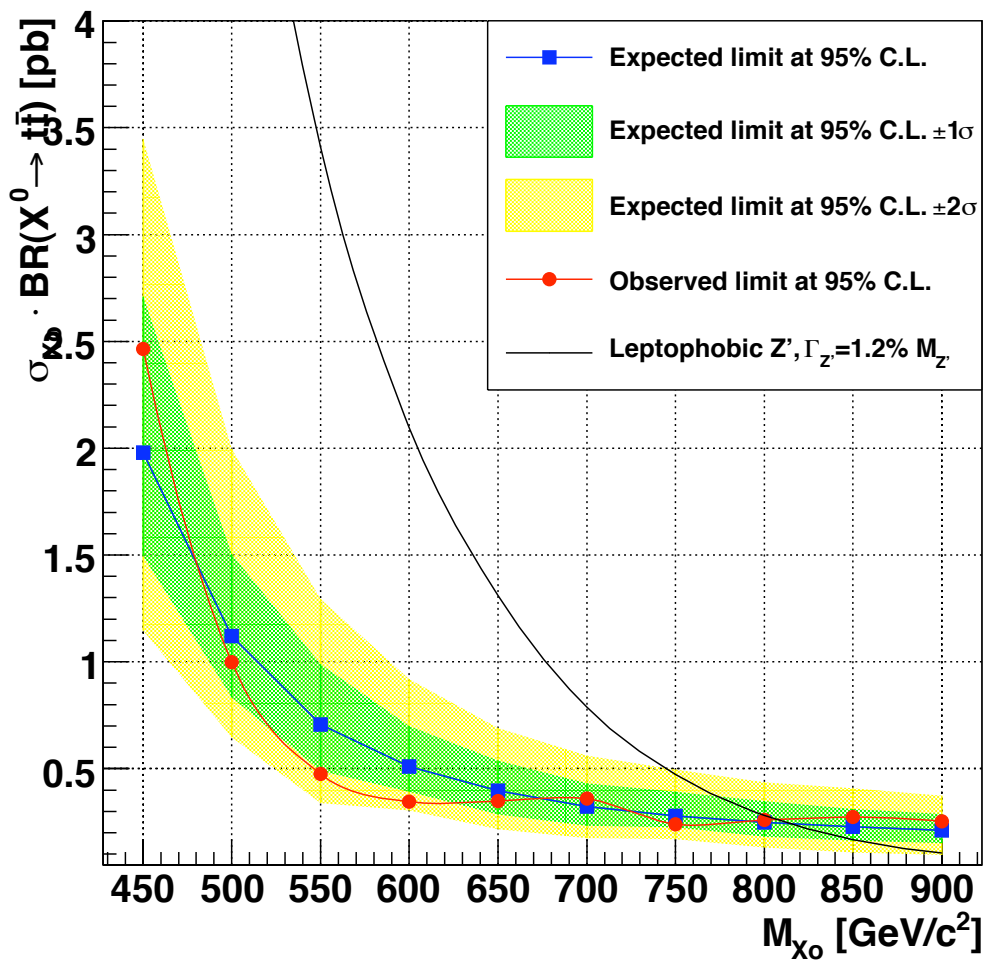


# Data/BG prediction





# Upper Limits







# Conclusions & Plans



- First search for narrow  $t\bar{t}b\bar{a}$  resonance in all jets final state
  - No excess found in 2.8/fb of CDF data
  - We set observed upper limit on leptophobic  $Z'$  mass up to 805 GeV
  - Used various tools to test SM for very small % "contaminations of new physics"
- Analysis has been reviewed and no unresolved issues were found
- Plan
  - PRD publication
  - Updated result with improvements in lepton plus jets



# Backup slides



# Signal contamination



- Signal contribution to QCD shape will be treated as following:

From Equation 4 we have, number of events in bin “i”:

$$\mu = \sigma_s A_s T_s + \sigma_{tt} A_{tt} T_{tt} + N_{QCD}^{pure} T_{QCD}^{pure}$$

$$N_{QCD}^{pure} T_{QCD}^{pure} = N_{QCD}^{cont} T_{QCD}^{cont} - \sigma_s A_s^{cont} T_s^{cont} - \sigma_{tt} A_{tt}^{cont} T_{tt}^{cont}$$

Comparing signal templates of predicted and observed values we can assume:

$$T_s = T_s^{cont}$$

So, finally we get:

$$\mu = \sigma_s (A_s - A_s^{cont}) T_s + \sigma_{tt} A_{tt} T_{tt} + N_{QCD}^{cont} T_{QCD}^{cont} - \sigma_{tt} A_{tt}^{cont} T_{tt}^{cont}$$

- In the end, it decreases signal acceptances by the values we get from TRM, which is about 1-1.5%
- It will obviously result in the worse sensitivity.



# Simplifications



- To calculate that probability we need to compute 28 integrals:
  - $P_t$  and  $P_z$  of incoming partons
  - 4-momenta of 6 final partons
- To reduce CPU time, we made some assumptions:
  - $P_t$  of the incoming partons is 0. -2 integrals
  - All quarks except top are massless. -8 integrals
  - Partons and jets have the same direction. -12 integrals
  - $W$ 's and top's are on shell. -4 integrals
- Only 2 integrals in total. We'll do more in the future.



# Improvements



- Some of the improvements made:
  - Added events with 7 jets, considering last 3 jets in  $E_t$  as extra jets from radiation
    - results in better signal acceptance
  - Used refined binning for transfer functions, both in  $E_t$  in  $E_t$ 
    - results in better signal templates
- Both improvements should result in better sensitivity



# Details



Jet-parton assignments

Differential xsection

$P_t$  of  $t\bar{t}$  system

$$P(j | M_{top}) = \frac{1}{\sigma(M_{top}) \epsilon(M_{top}) N_{combi}} \sum_{combi} \int dz_a dz_b f(z_a) f(z_b) d\sigma(M_{top}, p) TF(j | p) P_T(p)$$

Normalization factor

Integration over PDF's

Transfer functions



# $d\sigma$ calculation



$$d\sigma(M_{top}, p) = \frac{|\mathcal{M}(M_{top}, p)|^2}{4E_a E_b |v_a - v_b|} (2\pi)^4 \delta^4(p_a + p_b - \sum_{i=1}^6 p_i) \prod_{i=1}^6 \frac{d^3 p_i}{(2\pi)^3 2E_i}$$

$$\mathcal{M}(M_{top}, p) \approx \frac{\bar{v}(p_{\bar{q}}) \gamma^\mu u(p_q)}{(p_q + p_{\bar{q}})^2} \cdot \frac{\bar{u}(p_u) \gamma^\alpha (1 - \gamma^5) v(p_{\bar{d}})}{P_{W^+}^2 - m_W^2 + im_W \Gamma_W} \cdot \frac{\bar{u}(p_d) \gamma^\sigma (1 - \gamma^5) v(p_{\bar{u}})}{P_{W^-}^2 - m_W^2 + im_W \Gamma_W} \cdot \bar{u}(p_b) \gamma_\alpha (1 - \gamma^5) \frac{\not{p}_t + m_t}{p_t^2 - m_t^2 + im_t \Gamma_t} \gamma_\mu \frac{\not{p}_{\bar{t}} + m_t}{p_{\bar{t}}^2 - m_t^2 + im_t \Gamma_t} \gamma_\sigma (1 - \gamma^5) v(p_{\bar{b}})$$

- We use uubar --> 6 exact tree level ME
- Spin-correlations are included
- We compute the amplitudes directly using explicit Dirac matrices and spinors

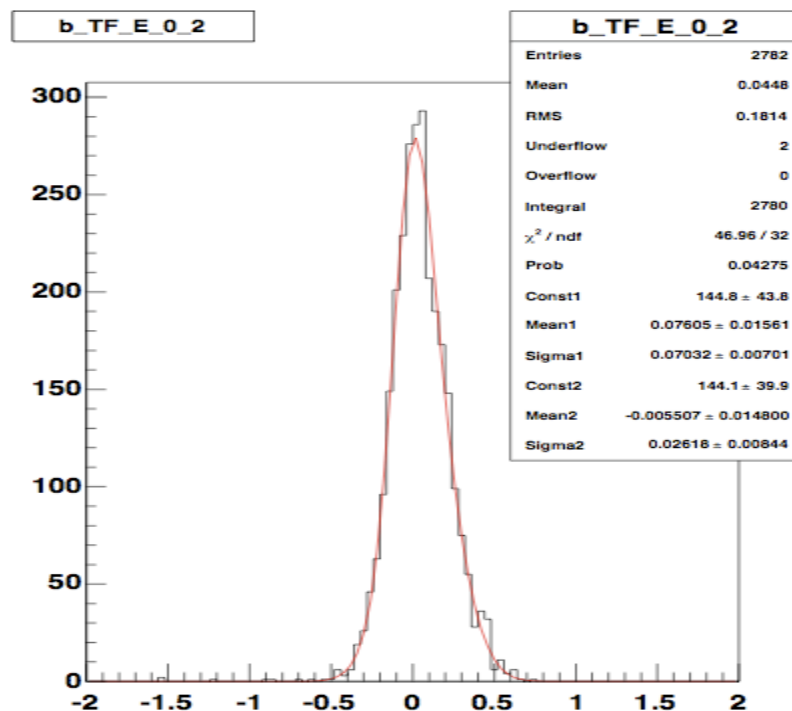


# Transfer functions



- From MC calculate the probability density function  $TF(E_j|E_p)$ 
  - $\xi = 1 - E_{jet} / E_{parton}$
- Use different TF's for different regions in  $\eta$ , energy, quark types

Example of b-quark  
Transfer Function for  
 $1.3 \leq |\eta| \leq 2$







# Prereqs effects



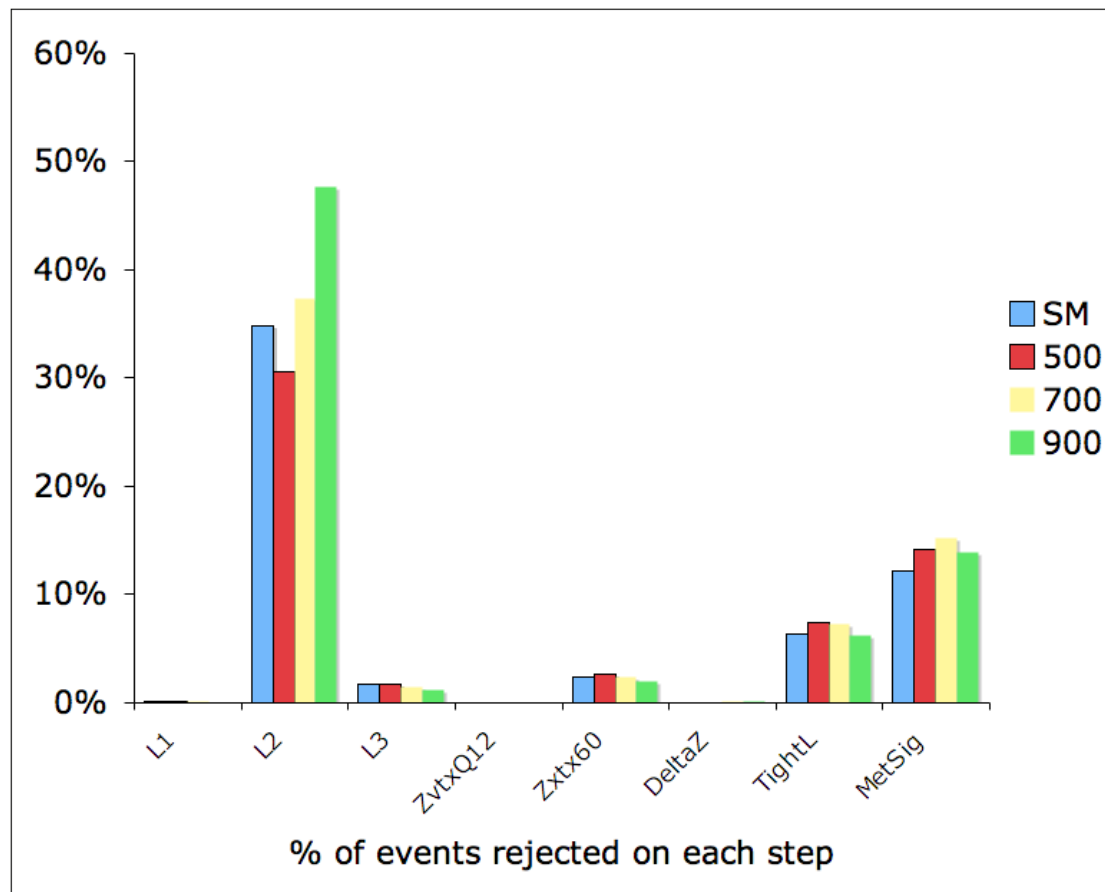
- Total efficiencies:  $\varepsilon_{tt}=42\%$ ,  $\varepsilon_{500}=43\%$ ,  $\varepsilon_{700}=36\%$ ,  $\varepsilon_{900}=28\%$
- Where do we lose events for high masses?

More interesting: Why?

Most of the events are lost on L2, which requires at least 4 clusters

For higher resonance masses, decay products are boosted more=> higher chance to merge in one cluster

See backup slides for details

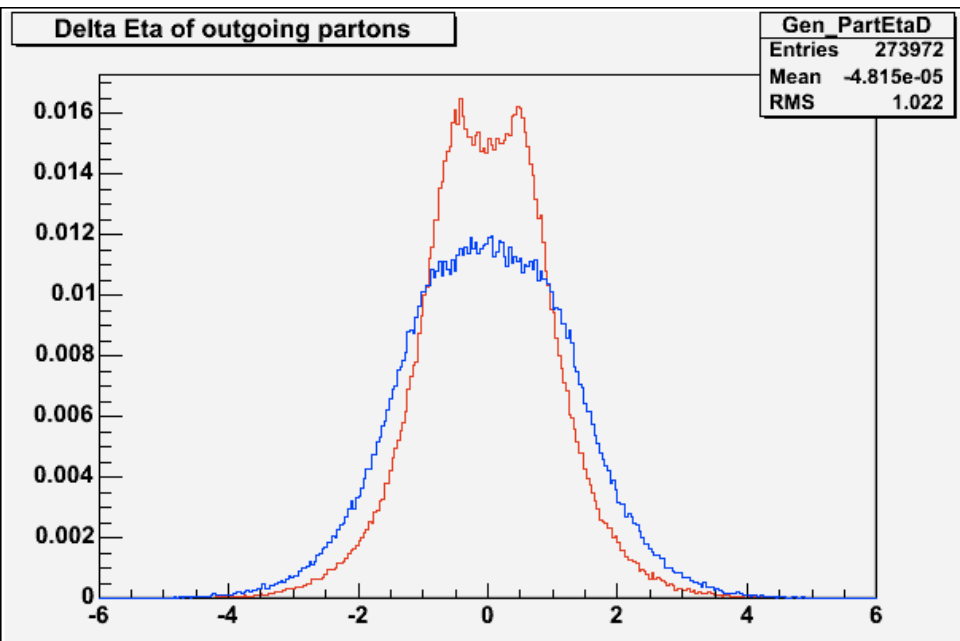
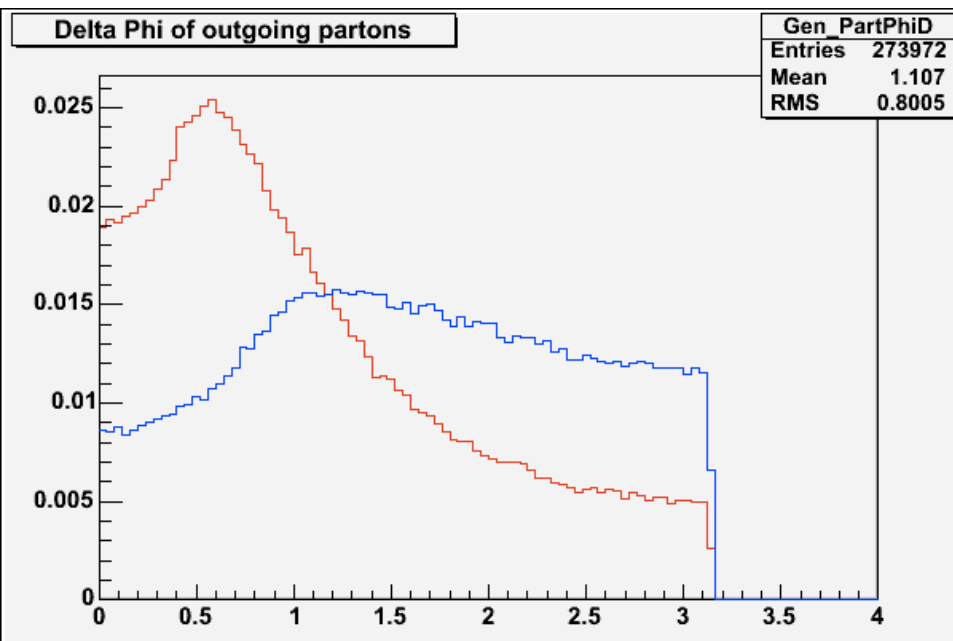




# Support for L2 issue

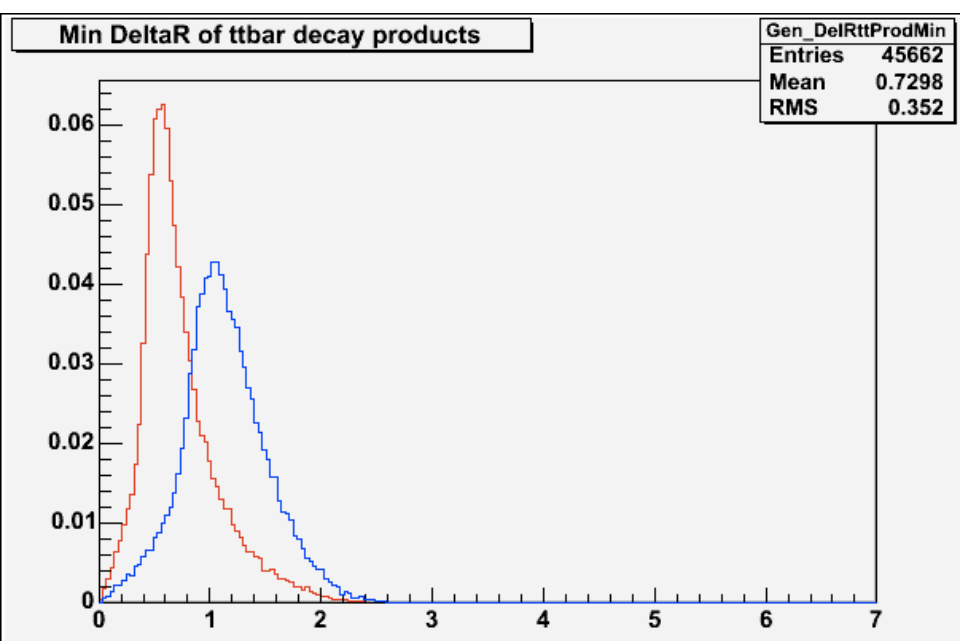
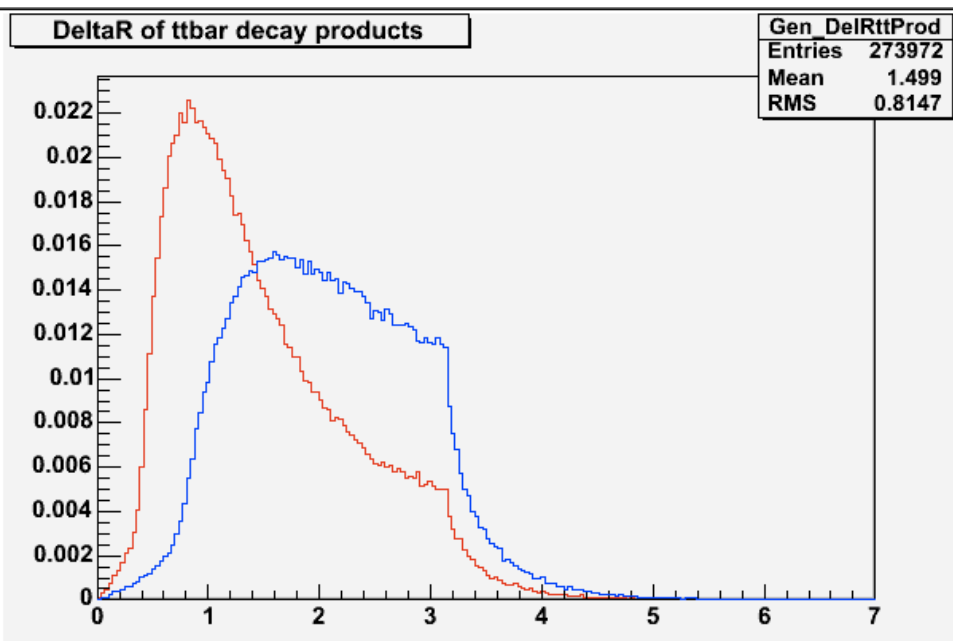


Blue 500 GeV resonance  
Red 900 GeV resonance.



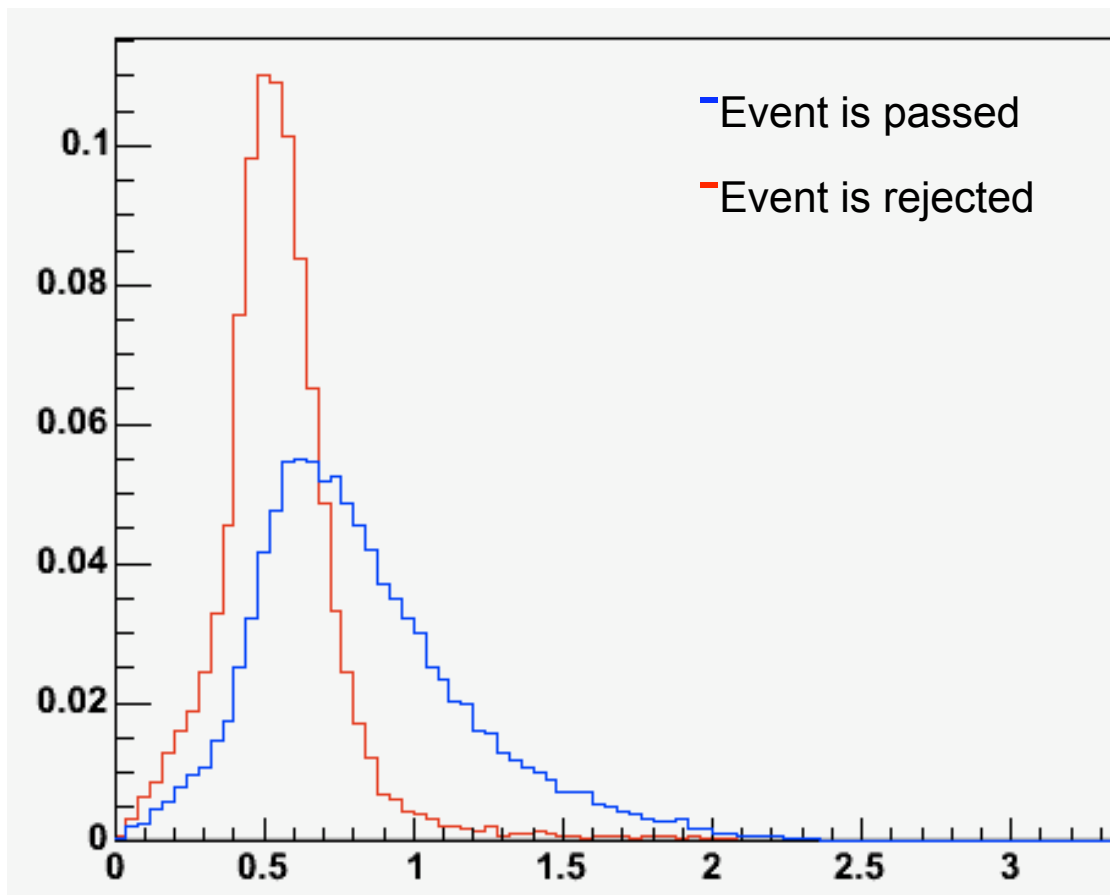


# L2 continued





# L2 final

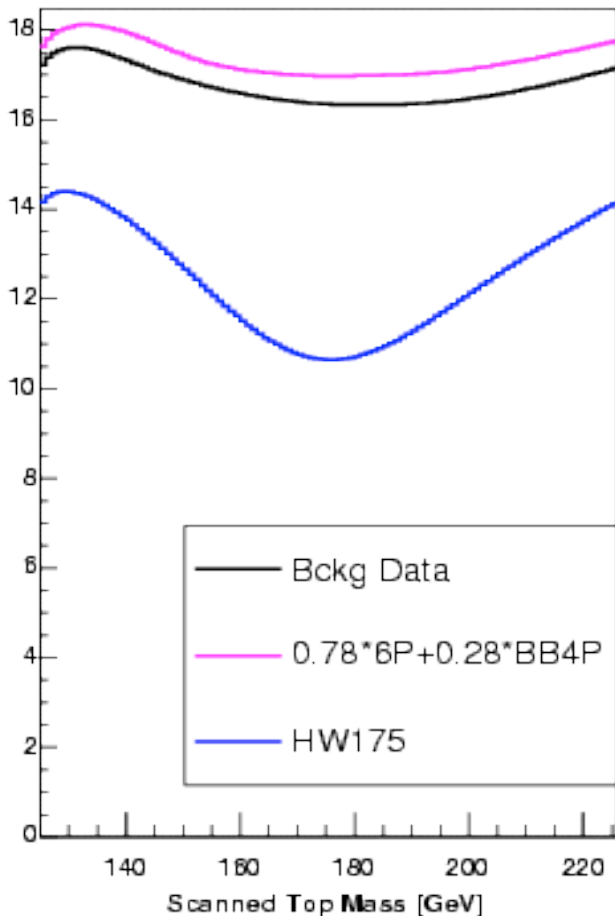




# FlaME Variable



LKLvs Mass



FlaME gives the probability of an event to come from SM  $t\bar{t}$ . Let's take advantage of it!

Here we plot  $-\log(P)$  vs top mass for various samples. As you see there is a difference between  $t\bar{t}$  and QCD

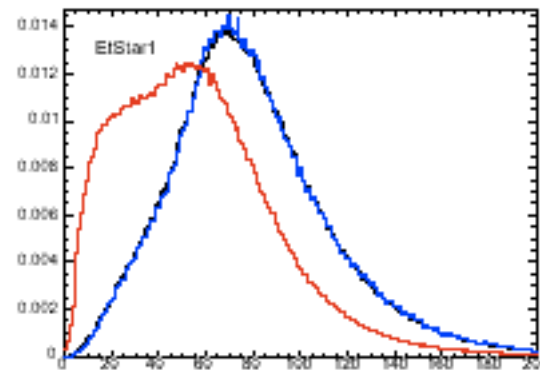
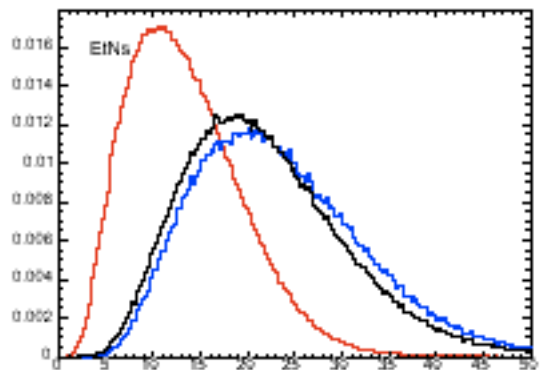
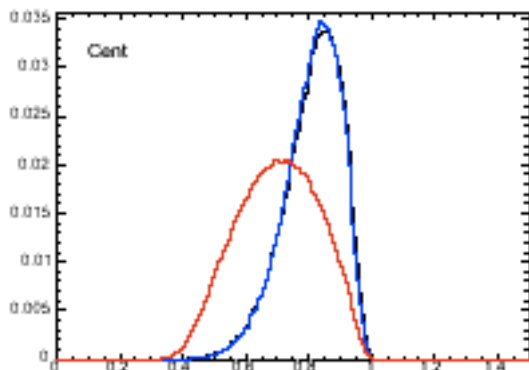
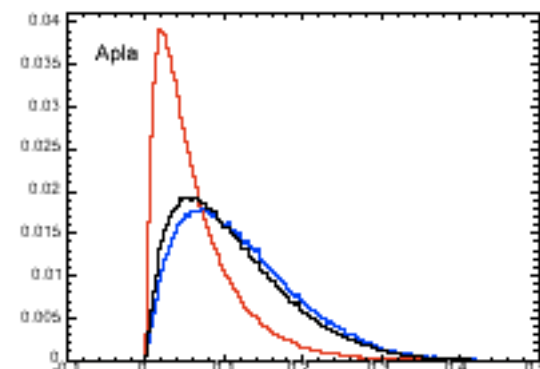
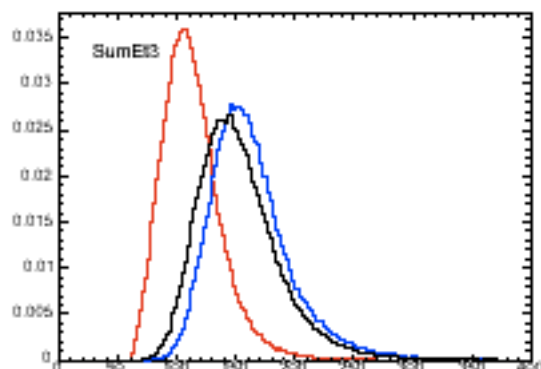
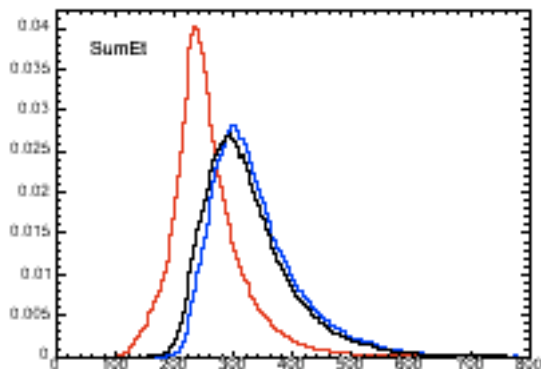
Lets calculate  $-\log(P)$  for 9 mass points: 155,160...195GeV. Decided to use their sum



# ...and their distributions

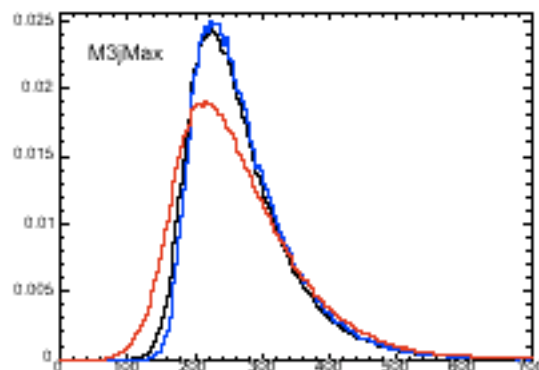
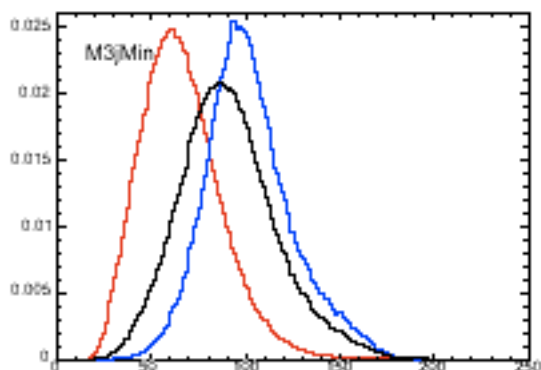
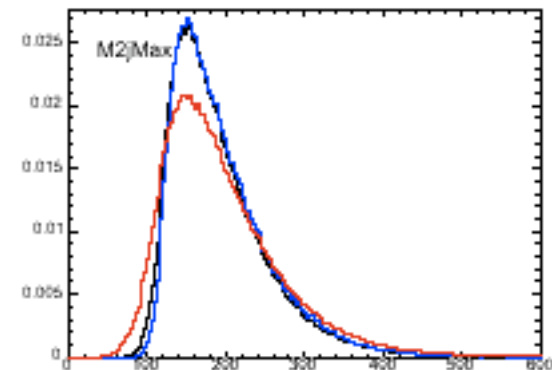
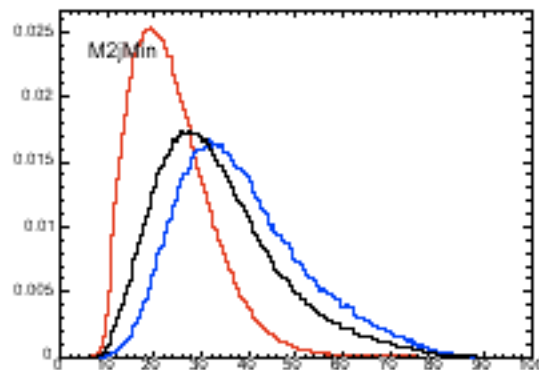
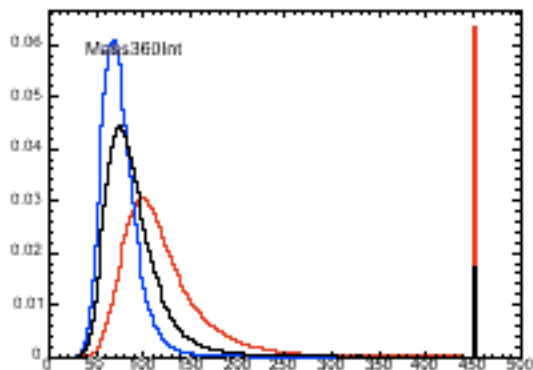


Red lines correspond to data.  
Black lines correspond to SMtt  
Blue lines correspond to SMtt matched only





# ...and their distributions





# Plug&Play



Black with FlaME, Red without FlaME, green kin. ev. sel.

