## $Z_{B-L}^{\prime}$ phenomenology at LHC and

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## Introduction

- In this work we study the phenomenology of two models with $S U(2)_{L} \otimes U(1)_{1} \otimes U(1)_{2}$ gauge symmetry for the colliders LHC and ILC. We will explore reactions like:
$\checkmark p+p \rightarrow \mu^{+} \mu^{-}+X$
$\checkmark \mathrm{e}^{+}+\mathrm{e}^{-} \rightarrow \mathrm{f} \overline{\mathrm{f}}$
- In order to perform these studies we will consider important observables for LHC as: total cross sections, number of events, forward-backward asymmetry, rapidity and transverse momentum distribution related to the final states. For ILC we will consider some asymmetry distributions.


## The Models

- B-L Secluded: $\operatorname{SU}(2)_{L} \otimes U(1)_{Y} \otimes U(1)_{Z}$
- B-L Flipped: $S U(2)_{L} \otimes U(1)_{Y} \otimes U(1)_{B-L}$


## Charge Operator

$$
\begin{array}{c|c}
\text { B-L Secluded } & \text { B-L Flipped } \\
\hline \mathrm{Q} / \mathrm{e}=\mathrm{I}_{3}+\mathrm{Y} / 2 & \mathrm{Q} / \mathrm{e}=\mathrm{I}_{3}+1 / 2\left[\mathrm{Y}^{\prime}+(\mathrm{B}-\mathrm{L})\right]
\end{array}
$$

- These models have two massive neutral vector bosons that will be denoted as $Z_{1}$ and $Z_{2}$ and their weak neutral currents will be parameterized as:

$$
L^{N C}=-\frac{g}{2 c_{w}} \sum \bar{\psi}_{i} \gamma_{\mu}\left[\left(g_{V}^{i}-g_{A}^{i} \gamma_{5}\right) Z_{1}^{\mu}+\left(f_{V}^{i}-f_{A}^{i} \gamma_{5}\right) Z_{2}^{\mu}\right] \psi_{i}
$$

## B-L Secluded Model

- In this model the masses of the neutral gauge bosons arise from the following terms in the covariant derivatives $\quad z_{\varphi}=2$

$$
g^{2} \frac{v^{2}}{8}\left(W_{3}^{\mu}-t_{W} B_{Y}^{\mu}-z_{H} t_{Z} B_{Z}^{\mu}\right)^{2}+\frac{u^{2}}{8}\left(z_{\varphi} g_{Z} B_{Z}^{\mu}\right)^{2} \quad t_{Z}=\frac{g_{Z}}{g} \quad t_{W}=\frac{g_{Y}}{g}
$$

- In the basis $W_{3}, B_{Y}$ and $B_{Z}$ the mass square matrix for the three electrically neutral gauge bosons is:

$$
M_{\text {neutral }}^{2}=\frac{g^{2} u^{2}}{4}\left(\begin{array}{ccc}
\bar{v}^{2} & -t_{W} \bar{v}^{2} & -2 t_{Z} z_{H} \bar{v}^{2} \\
-t_{W} \bar{v}^{2} & t_{W}^{2} \bar{v}^{2} & 2 t_{W} t_{Z} z_{H} \bar{v}^{2} \\
-2 t_{Z} z_{H} \bar{v}^{2} & 2 t_{W} t_{Z} z_{H} \bar{v}^{2} & 4 t_{Z}^{2}\left(1+z_{H}^{2} \bar{v}^{2}\right)
\end{array}\right) \quad g_{Z}>0
$$

## B-L Flipped Model

- The masses of the neutral gauge bosons arise from the following terms in the covariant derivatives:

$$
\frac{v^{2}}{8}\left(g W_{3}^{\mu}-g^{\prime} B_{Y^{\prime}}^{\mu}-g_{B-L} B_{B-L}^{\mu}\right)^{2}+\frac{u^{2}}{8}\left(g^{\prime} Y_{\phi}^{\prime} B_{Y^{\prime}}^{\mu}-g_{B-L} Y_{\phi}^{\prime} B_{B-L}^{\mu}\right)^{2} \quad Y_{\phi}^{\prime}=-2
$$

- The mass square matrix for the three electrically neutral gauge bosons in the basis $W_{3}, B_{Y^{\prime}}, B_{\mathrm{B}-\mathrm{L}}$ is:

$$
M_{\text {neutral }}^{2}=g^{2} u^{2}\left(\begin{array}{ccc}
\bar{v}^{2} / 4 & t^{\prime} \bar{v}^{2} / 4 & 0 \\
t^{\prime} \bar{v}^{2} / 4 & t^{\prime 2}\left(1+\bar{v}^{2} / 4\right) & -t^{\prime} t_{B-L} \\
0 & -t^{\prime} t_{B-L} & t_{B-L}^{2}
\end{array}\right)
$$

$$
\begin{gathered}
t^{\prime}=\frac{g^{\prime}}{g} \\
t_{B-L}=\frac{g_{B-L}}{g}
\end{gathered}
$$

## Imputs Chosen for Both Models

2 scenarios: First: $M_{Z}=1000 \mathrm{GeV}$; Second $M_{Z},=1500 \mathrm{GeV}$

## B-L Flipped

- $g^{\prime}=0.44$
- $g_{\mathrm{B}-\mathrm{L}}=0.6132$
- $u=1324.4$ / 1987
- $\Gamma_{Z}=26.37 \mathrm{GeV} / 38.87 \mathrm{GeV}$



## B-L Secluded

- $g_{z}=0.2$
- $z_{q}=1 / 3$
- $u=5000 / 7500$
- $\mathrm{Z}_{\mathrm{H}}=0 \longrightarrow f_{\mathrm{A}}$ 's vanish
- The vectorial couplings of Z' to fermions will be given by:

$$
f_{V}^{v}=f_{V}^{l}=-3 f_{V}^{u}=-3 f_{V}^{d}=t_{z} c_{W}
$$

- $\Gamma_{Z}=9.55 \mathrm{GeV} / 10.48 \mathrm{GeV}$


## Neutral Coupling Constants $\mathrm{f}_{\mathrm{V}, \mathrm{A}}$ and Decay Widths for Both Models

| $\mathrm{M}_{\mathrm{Z}}=1000 / 1500$ | B-L Flipped |  | B-L Secluded |  |
| :---: | :---: | :---: | :---: | :---: |
| GeV | $f_{\mathrm{V}}$ | $f_{\mathrm{A}}$ | $f_{\mathrm{V}}$ | $f_{\mathrm{A}}$ |
| neutrinos | $0.8412 / 0.8420$ | $-0.1739 /-0.1732$ | $0.2690 / 0.2690$ | $0 / 0$ |
| leptons | $0.4977 / 0.4977$ | $0.1739 /-0.1732$ | $0.2690 / 0.2690$ | $0 / 0$ |
| u-quarks | $-0.0510 /-0.0511$ | $-0.1739 /-0.1732$ | $-0.0897 /-0.0897$ | $0 / 0$ |
| d-quarks | $-0.3949 /-0.3955$ | $0.1739 /-0.1732$ | $-0.0897 /-0.0897$ | $0 / 0$ |


| $\mathrm{M}_{\mathrm{z}}=1000 / 1500$ <br> GeV | B-L Flipped | B-L Secluded |
| :---: | :---: | :---: |
| $Z^{\prime} \rightarrow \sum_{i} \bar{v}_{i} v_{i}$ | $36 \% / 36 \%$ | $23.5 \% / 23.6 \%$ |
| $Z^{\prime} \rightarrow \sum_{i} \bar{l}_{i} l_{i}$ | $18.6 \% / 18.6 \%$ | $45.1 \% / 45.5 \%$ |
| $Z^{\prime} \rightarrow \sum_{i} \bar{q}_{i} q_{i}$ | $42.4 \% / 42.6 \%$ | $31.4 \% / 30.9 \%$ |
| $Z^{\prime} \rightarrow W^{+} W^{-}$ | $3 \% / 2.8 \%$ | $0 \% / 0 \%$ |

## Observables of Z' at Colliders

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## LHC

 <br> $\checkmark$ Total cross sections; <br> $\checkmark$ Forward-backward Asymmetry; <br> $\checkmark$ Rapidity Distributions; <br> $\checkmark$ Transverse moment Distributions; <br> $\checkmark$ Lepton angular Distribution.}

ILC
$\checkmark$ Forward-backward Asymmetry;
$\checkmark$ Left-right Asymmetry;
$\checkmark$ Polarization Asymmetry;
$\checkmark$ Mixed Asymmetries.

Drell-Yan Channel: $\mathrm{p}+\mathrm{p} \rightarrow \mu^{+}+\mu^{-}+\mathrm{X}$

$$
\mathrm{e}^{+}+\mathrm{e}^{-} \rightarrow \mu^{+}+\mu^{-}
$$

## Results - LHC






## Results - LHC






## Results - LHC



## Results - LHC





## Results - ILC ( $\mathrm{M}_{\text {z2 }}=1000 \mathrm{GeV}$ )






## Conclusions

- The B-L Secluded Model is leptofilic, its cross section near the $Z_{2}$-peak, is larger for leptons if compared to quarks;
- In both models, $Z_{2}$ decays preferentially to leptons compared to the SM;
- The $Z_{2}$ widths are very different in each model and are larger in the flipped model;
- According to the chosen parameters, the Flipped model has better chances to be disentangled from the background of the standard model, due to the nature of $Z_{2}$ couplings to fermions;


## References

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