Lepton and Quark Flavor Violation in a Minimal $SO(10) \times A_4 SUSY GUT$

Abdelhamid Albaid

Under the supervision of

Prof. K. S. Babu

Oklahoma State University

Outline



- A brief description of $SO(10) \times A_4$ SUSY GUT Model
- Lepton and quark flavor mixings



- A-terms mass insertion
- The sfermion mass insertion



Universality hypothesis at the Planck scale.



Motivation



To investigate the flavor mixing of quarks and charged leptons in the minimal $SO(10) \times A_4$ SUSY GUT.



Consequently, to obtain the constraints on the soft terms which might gives us a hint about how supersymmetry is broken.

A brief description of $SO(10) \times A_4$ SUSY GUT Model

Features of SO(10):

$$\Leftrightarrow (Q_i, L_i, u_i^c, d_i^c, e_i^c) + v^c \to 16$$

 $\checkmark M_{u} \propto M_{d} \rightarrow V_{CKM} = diag \ \{1, 1, 1\}$

See-Saw Mechanism is automatic

Features of A4:

A4 is the smallest discrete group that has a three dimensional irreducible representation [E. Ma, G. Rajasekaran, 2001]

4 A4 flavor symmetry very easily gives tri-bi-maximal mixing matrix [P.F Harrison et al, 2002]

		16 _i	$16_1, 1\overline{6}$	16	2 , 1 6 ₂	16 ₃ , 1 6	<u>5</u> ₃	1_i^c
A4		3	1		1	1		3
$Z_2 \times Z_4 \times Z_4$	Z ₂ +	, + , +	+, -, +	- –,	+,+	+ , + , -	-	+ , + , +
SO(10)	10 _i	1 0 '		$1 \ 0 ''_{i}$	1 0 '''		1 _i
A4		3	1		1	1		3
$Z_2 \times Z_4 \times Z_4$	Z ₂ +	- , <i>i</i> , +	+, -i, +	+ +	, <i>i</i> , –	+, -i, -i	- +,	-i, +
SO(10)	10 _{<i>H</i>}	45 _{<i>H</i>}	16 _{<i>H</i>}	$1\overline{6}_{H}$	1_{Hi}	$1'_{Hi}$	$1''_{Hi}$	$1_{Hi}^{\prime\prime\prime}$
A4	1	1	1	1	3	3	3	3
$Z_2 \times Z_4 \times Z_2$	-, +, -	+,-,-	+,- <i>i</i> , +	+,- <i>i</i> ,+	+, -, +	-, +, +	+, +, -	+,- <i>i</i> ,+

A brief description of $SO(10) \times A_4 \, \text{SUSY GUT}$ Model



A brief description of SO(10) $\times A_4$ SUSY GUT Model

$$M_{F}^{o} = C_{F} \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & Q_{F} s \\ 0 & Q_{F^{c}} s & (Q_{F^{c}} + Q_{F}) c \end{bmatrix}$$

Where,
$$Q_F = 2I_{3R} + \frac{6}{5}\delta\left(\frac{Y}{2}\right)$$
,

$$\begin{array}{ccc} & & & & \\ & & & \\ & & & \\ & \swarrow & & \\ & & & \\ & \swarrow & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & \\ & & & \\ & &$$

A brief description of $SO(10) \times A_{_4}\,\text{SUSY}$ GUT Model

By including the couplings with the vector 10-plet fermions, the model leads to the doubly lopsided structure.

$$M_{L} \approx M_{D}^{T} \approx m_{d}^{0} \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ \delta_{2} & \delta_{1} & 1 \end{bmatrix}$$

A brief description of $SO(10) \times A_{_4}\,\text{SUSY}$ GUT Model

51.5 - A.S.	Model predictions	Experiment	Pull
$m_e(m_e)$	0.511×10^{-3}	0.511×10^{-3}	
$m_{\mu}(m_{\mu})$	105.6×10^{-3}	105.6×10^{-3}	
$m_{\tau}(m_{\tau})$	1.776	1.776	
\overline{m}_{ud}	4.32×10^{-3}	$(3.85 \pm 0.52) \times 10^{-3}$	0.9
$m_c(m_c)$	1.4	$1.27^{+0.07}_{-0.11}$	1.85
$m_t(m_t)$	172.5	171.3 ± 2.3	0.52
$\frac{m_s}{\overline{m}_{s-d}}$	25.36	27.3 ± 1.5	1.29
$m_s(2Gev)$	109.6×10^{-3}	$105^{+25}_{-35} \times 10^{-3}$	0.184
$m_b(m_b)$	4.31	$4.2^{+0.17}_{-0.07}$	0.58
V_{us}	0.2264	0.2255 ± 0.0019	0.473
V_{cb}	39.2×10^{-3}	$(41.2\pm1.1)\times10^{-3}$	1.82
V_{ub}	4.00×10^{-3}	$(3.93 \pm 0.36) \times 10^{-3}$	0.194
η	0.3569	$0.349^{+0.015}_{-0.017}$	0.526
$\sin \theta_{12}^{sol}$	0.551	0.566 ± 0.018	0.83
$\sin \theta_{23}^{atm}$	0.776	0.707 ± 0.108	0.63
$\sin \theta_{13}$	0.154	< 0.22	-

Phys. Rev. D 80, 093002 (2009) - Published November 12, 2009

- \Rightarrow Introduction
 - \checkmark Sources of quark and lepton flavor violation
 - The sfermion mass insertions $\delta_{LL,RR}^{ij}$
 - The chirality-flipping mass insertion (A-terms) $\delta_{LR,RL}^{ij}$

Both sources arise from the mixing of the light states with the heavy ones.

- \checkmark There are two transformations
 - Transformation needed to block diagonalize the femion mass matrix
 - Transformation needed to diagonalize the 3 by 3 light femion mass matrix

□ Introduction

In order to block diagonalize the fermion mass matrix, we define the light and heavy states as follows:

$$\begin{split} & L_{2} = N_{1}(\psi_{2} - \boldsymbol{s}_{\theta}\boldsymbol{T}_{2}\chi_{2}) \\ & L_{3} = \frac{N_{1}}{R_{q}}(-\psi_{3} + \frac{c_{\theta}s_{\theta}T_{2}^{2}}{N_{1}^{2}}\psi_{2} + \boldsymbol{T}_{1}\chi_{1} - \boldsymbol{Q}\boldsymbol{T}_{3}\boldsymbol{T}_{1}\chi_{3} + \frac{c_{\theta}\boldsymbol{T}_{2}}{N_{1}^{2}}\chi_{2}) \\ & H_{1} = \frac{1}{K_{q}}(\boldsymbol{Q}T_{3}\chi_{1} + \chi_{3}) \\ & H_{2} = \frac{1}{N_{2}}(\chi_{2} + T_{2}(s_{\theta}\psi_{2} + c_{\theta}\psi_{3})) \\ & H_{3} = \frac{T_{1}K_{q}}{N_{2}R_{q}}(-N_{1}^{2}\psi_{3} + c_{\theta}s_{\theta}T_{2}^{2}\psi_{2} + c_{\theta}T_{2}\chi_{2}) + \frac{N_{2}}{K_{q}R_{q}}(\boldsymbol{Q}T_{3}\chi_{3} - \chi_{1}) \end{split}$$
Where, $\begin{aligned} & W_{1} = \sqrt{1 + s_{\theta}^{2}T_{2}^{2}}, \quad N_{2} = \sqrt{1 + T_{2}^{2}} \\ & R_{q} = \sqrt{1 + T_{2}^{2} + T_{1}^{2}(1 + s_{\theta}^{2}T_{2}^{2})(1 + \boldsymbol{Q}^{2}T_{3}^{2})}, \\ & K_{q} = \sqrt{1 + Q^{2}T_{3}^{2}}, \quad T_{1} = \frac{\langle S \rangle}{M_{1}}, \quad T_{2} = \frac{\langle E \rangle}{M_{2}}, \quad T_{3} = \frac{\langle 45_{H} \rangle}{M_{3}}. \end{aligned}$

$$s_{\theta} \approx 0.04, s_{\theta}T_2 \ll 1$$

The first and second generations are almost degenerate

$$\frac{c_{\theta}T_2}{N_1^2}, Q T_3 T_1 \approx 1 \qquad \Longrightarrow$$

The third generation splits and

Causes flavor violation

 \implies A-terms mass insertion

First Transformation gives

$$\widetilde{M}_{RL}^2 = \widetilde{a}_1 M_{1f} + \widetilde{a}_2 M_{2f}$$

Second diagonalization $V_R^*M_fV_L = M_f^{diag}$

$$V_{R}^{*} \widetilde{M}_{RL}^{2} V_{L} = \widetilde{a}_{1} M_{f}^{diag.} + \widetilde{a}_{21} V_{R}^{*} M_{2f} V_{L}. \quad \text{Where,} \quad \widetilde{a}_{21} = \widetilde{a}_{2} - \widetilde{a}_{1}$$

$$\delta_{12}^{l} = (0.075, 0.01) \text{GeV} \frac{\tilde{a}_{21}}{\tilde{m}^{2}} \le 10^{-5}, \ \delta_{12}^{d} = (0.34, 0.69) \text{GeV} \frac{\tilde{a}_{31}}{\tilde{m}^{2}} \le 9 \times 10^{-5}$$

$$\delta_{13}^{l} = (0.035, 0.07) \text{GeV} \frac{\tilde{a}_{13}}{\tilde{m}^{2}} \le 0.04, \\ \delta_{13}^{d} = (0.079, 0.35) \text{GeV} \frac{\tilde{a}_{31}}{\tilde{m}^{2}} \le 1.7 \times 10^{-2}$$

$$\delta_{23}^{l} = (0.056, 0.88) \text{GeV} \frac{\tilde{a}_{13}}{\tilde{m}^{2}} \le 0.03, \\ \delta_{23}^{d} = (0.055, 0.8) \text{GeV} \frac{\tilde{a}_{31}}{\tilde{m}^{2}} \le (6, 4.5) \times 10^{-3}$$

 $\tilde{a}_{21} \approx 100 \,\text{GeV}$, and $\tilde{m} \approx 500 \,\text{GeV}$ The Correct FCNC Suppression

 \Rightarrow The sfermion mass insertion

A4 discrete symmetry is broken at GUT scale. From top Yukawa coupling

 $1.5 \le T_2, T_3 \le 2.4$ $T_1 \approx 1$

The quadratic mass soft terms are give by:

$$L_{\rm soft} = \overline{m}^2 \psi_i^* \psi_i + \overline{M}_{\alpha}^2 \chi_{\alpha}^* \chi_{\alpha}$$

In terms of the new orthogonal eigenstates, The soft mass matrix for light state is modified as follows:

Where,

$$\delta = \frac{\overline{m}^2 + \overline{M}_1^2 T_1^2 + \overline{M}_2^2 T_2^2 + \overline{M}_3^2 T_1^2 T_3^2 Q^2}{R_q^2} - \overline{m}^2, \quad \alpha = \frac{(\overline{m}^2 - \overline{M}_2^2) s_\theta T_2^2}{R_q}$$

The sfermion mass insertion

$$\begin{bmatrix} T_1 = 1, & T_2 = 2, \\ T_3 = 2.3, s_{\theta} = 0.042 \\ M_1 = M_2 = M_3 \end{bmatrix} \longrightarrow \begin{bmatrix} \delta = 0.87(\overline{M}^2 - \overline{m}^2) \\ \alpha = -0.059(\overline{M}^2 - \overline{m}^2) \end{bmatrix}$$

Second diagonalization $V_R^* M_f V_L = M_f^{diag}$

 $(\delta_{13}^{l})_{LL} = 0.027 \ \sigma \le 0.15, \\ \delta_{12}^{d} = (2.3 \times 10^{-4}, 0.13) \ \sigma \le (1.4, 90) 10^{-4}, \\ (\delta_{23}^{l})_{LL} = 0.36 \ \sigma \le 0.12, \\ \delta_{13}^{d} = (0.013, 0.077) \ \sigma \le (0.09, 0.07), \\ (\delta_{12}^{l})_{RR} = 7 \times 10^{-4} \ \sigma \le 0.09, \\ \delta_{23}^{d} = (0.06, 0.36) \ \sigma \le (0.16, 0.22).$

$(\delta_{12}^{l})_{LL} = 0.051 \sigma \le 6 \times 10^{-4}$

The above stringent upper bounds comes from $\Gamma(\mu \rightarrow e\gamma)$

Where,
$$\sigma = \frac{\tilde{m}^2 - \tilde{M}^2}{\tilde{m}^2} \approx 0.01$$
 So, it is required high degree of degeneracy to suppress $\Gamma(\mu \to e\gamma)$

Note: The above numerical upper bound is obtained from arXive:hep-ph/0702144v2

 \Rightarrow Universality hypothesis at the Planck scale.

Can the universality hypothesis of masses at the Planck scale reproduce this high degeneracy?

At Planck scale $\tilde{M}_1^2 = \tilde{M}_2^2 = \tilde{M}_3^2 = \tilde{m}^2$

RGE
$$\widetilde{M}_2^2 = \widetilde{M}_3^2, \quad \frac{\widetilde{m}^2 - \widetilde{M}_1^2}{\widetilde{m}^2} \approx 1, \quad \frac{\widetilde{m}^2 - \widetilde{M}_2^2}{\widetilde{m}^2} \approx 0.07$$

Gauge Mediating Supersymetry Breaking (GMSB):

FV appears only if there is mixing between the messenger fields and ordinary fields :



then, the FCNC constraints are satisfied (Dine et al., 1996)

Conclusion

- $\overset{\wedge}{\boxtimes}$
- The supersymmetric model of SO(10)XA4 with minimum Higgs representation makes the unified gauge coupling perturbative to the Planck scale, It successfully describes the fermion masses, CKM and neutrino mixings.



The sources of the flavor mixings arise from the order one mixing of the third light generation and the heavy states. Consequently, the mass of the third generation splits.



FCNC of $\mu \rightarrow e\gamma$ is only suppressed if we make the soft mass degeneracy arrangement of the light and heavy state. However, this arrangement is not consistent with universality hypothesis of PMSB.



Therefore, the model adopts the GMSB. It contains the fields that serve as SUSY breaking messengers.