

The Double Cover of A_5 and the Flavor Problem

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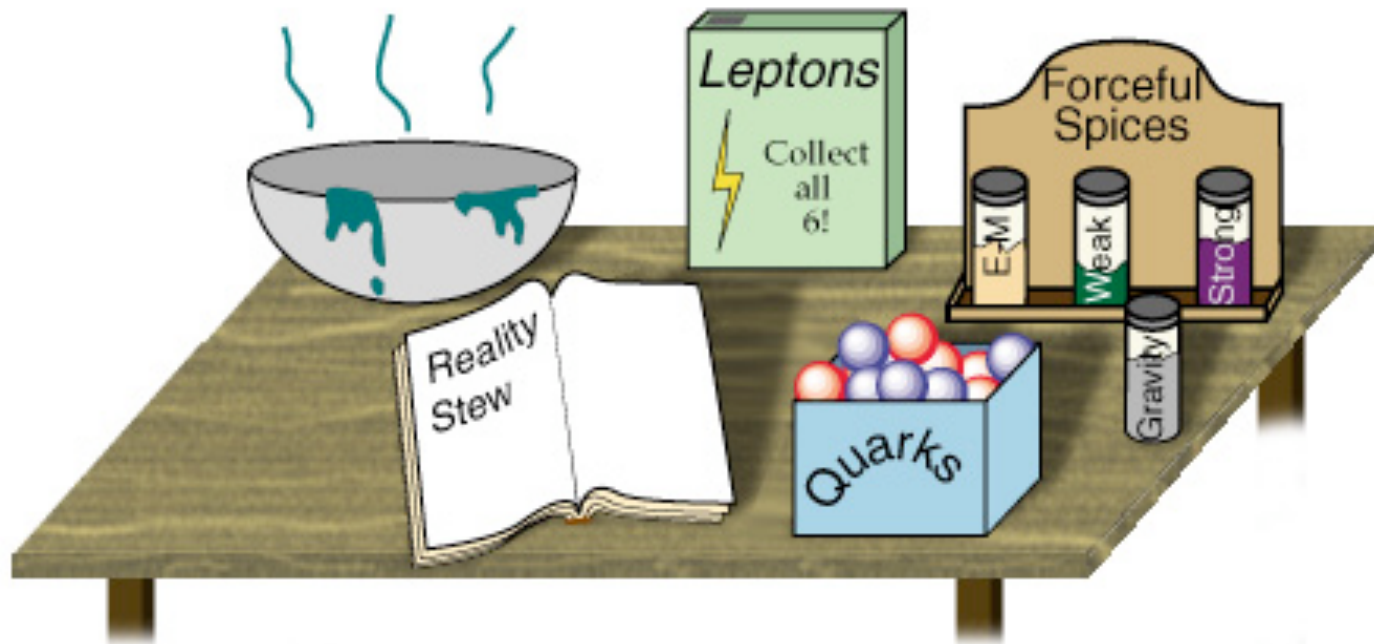
May 10th, 2010

Based on: L. Everett and A. Stuart, 0812.1057, and work in progress

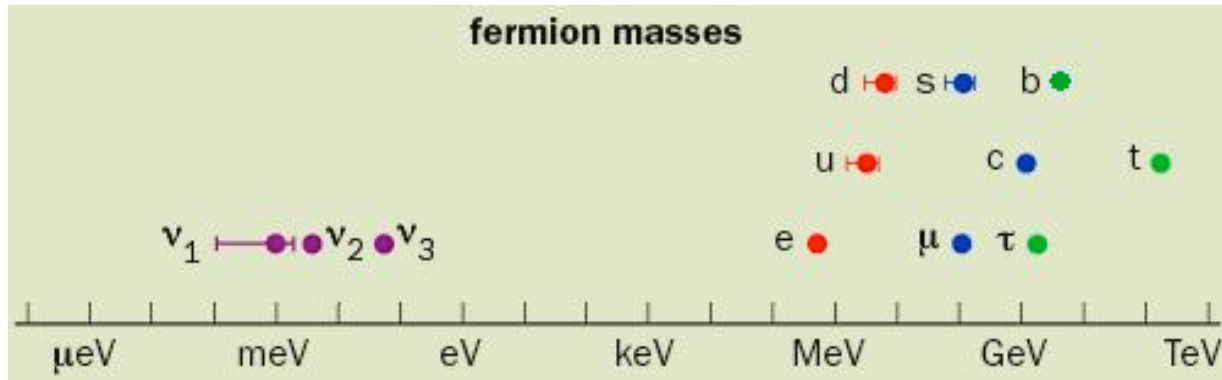
The Standard Model

$$SU(3)_c \times SU(2)_L \times U(1)_Y$$

- Triumph of modern science, but incomplete. Fails to predict measured masses and mixings of fermions (i.e. they are input parameters).
- What exactly do we “taste”?



What We “Taste”



Quark Mixing Angles

Lepton Mixing Angles

$$\mathcal{U}_{\text{MNSP}} = \mathcal{R}_1(\theta_{\oplus})\mathcal{R}_2(\theta_{13}, \delta_{\text{MNSP}})\mathcal{R}_3(\theta_{\odot})\mathcal{P}$$

$$\mathcal{U}_{\text{CKM}} = \mathcal{R}_1(\theta_{23}^{\text{CKM}})\mathcal{R}_2(\theta_{13}^{\text{CKM}}, \delta_{\text{CKM}})\mathcal{R}_3(\theta_{12}^{\text{CKM}})$$

$$\theta_{12}^{\text{CKM}} = 13.0^\circ \pm 0.1^\circ$$

$$\theta_{23}^{\text{CKM}} = 2.4^\circ \pm 0.1^\circ$$

$$\theta_{13}^{\text{CKM}} = 0.2^\circ \pm 0.1^\circ$$

$$\delta_{\text{CKM}} = 60^\circ \pm 14^\circ$$

$$\theta_{\odot} = 33.9^\circ \begin{matrix} +2.4^\circ \\ -2.2^\circ \end{matrix}$$

$$\theta_{\oplus} = 45^\circ \begin{matrix} +10^\circ \\ -10^\circ \end{matrix}$$

$$\theta_{13} < 13^\circ$$

(presently no constraints on CP violation)

Adding Some “Ginger”

(i.e. a discrete flavor symmetry that is spontaneously broken to generate observed masses and mixings)

A_4/T is the most common discrete symmetry used (E. Ma, G. Altarelli et al.,...)

Icosahedral Symmetry

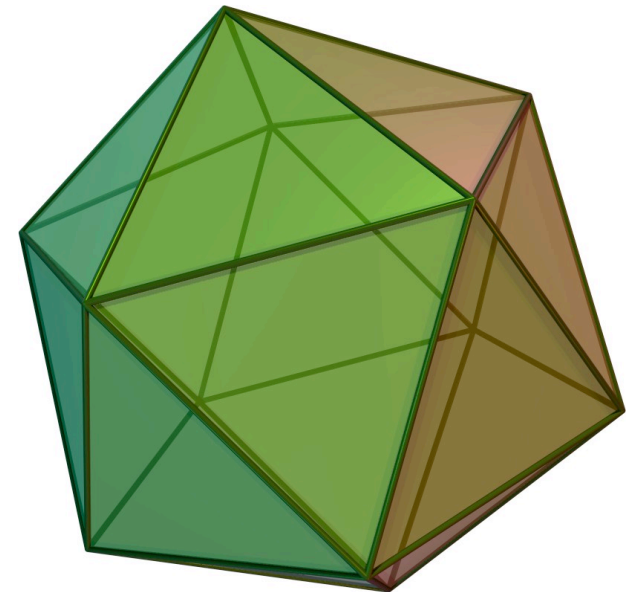
- All over nature
- Provides “natural” setting to look at $\theta_{sol} = \text{ArcTan}\left(\frac{1}{\phi}\right) = 31.7175^\circ$
- Has been applied in this context in **arXiv:0812.1057** [hep-ph]
- Now let’s apply it to the quarks.

What exactly is Icosahedral Symmetry?

The Icosahedral Group, I

- An icosahedron is the Platonic solid that consists of 20 equilateral triangles. $\rightarrow f = 20$
- 20 triangles each have 3 sides $\rightarrow 60$ edges but 2 triangles/edge $\rightarrow 30$ edges $\rightarrow e=30$
- 20 triangles each have 3 vertices $\rightarrow 60$ vertices but 5 vertices/edge $\rightarrow v= 12$
- Are we right? $\chi(g) = 2 - 2g = v - e + f$
- I consists of all rotations that take vertices to vertices i.e. $0, \pi, \frac{2\pi}{3}, \frac{2\pi}{5}, \frac{4\pi}{5}$

$$A_5 \cong I \subseteq SO(3)$$



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Conjugacy Classes of I

Rotation by each angle forms its own conjugacy class.

Schoenflies Notation: C_n^k is a rotation by $\frac{2\pi k}{n}$
in front = # of elements in class

So for the icosahedral group we have:

$$e, 12C_5, 12C_5^2, 20C_3, 15C_2$$

Note:

$$1 + 12 + 12 + 15 + 20 = 60 = 1^2 + 3^2 + 3^2 + 4^2 + 5^2 \quad \text{Two triplets.}$$

The Double Cover of I, I'

- To each element $g \in I$, associate another element $gR \ni R^2 = e$
- Therefore, $|I'| = 2|I| = 120$
- The characters (traces) of the new elements are related to the old by: $\chi(gR) = -\chi(g)$
- Furthermore, each conjugacy class gets a partner: $C_n^k R$
- One notable exception: $15C_2$
- We get 4 more (spinorial) irreps: $2^2 + 2^2 + 4^2 + 6^2 = 60$
- What about its character table?

$$I' \subseteq SU(2)$$

I' Character Table

\mathcal{I}'	1	3	3'	4	5	2	2'	4'	6
e	1	3	3	4	5	2	2	4	6
$12C_5$	1	ϕ	$1 - \phi$	-1	0	ϕ	$1 - \phi$	1	-1
$12C_5^2$	1	$1 - \phi$	ϕ	-1	0	$\phi - 1$	$-\phi$	-1	1
$20C_3$	1	0	0	1	-1	1	1	-1	0
$30C_2$	1	-1	-1	0	1	0	0	0	0
R	1	3	3	4	5	-2	-2	-4	-6
$12C_5R$	1	ϕ	$1 - \phi$	-1	0	$-\phi$	$\phi - 1$	-1	1
$12C_5^2R$	1	$1 - \phi$	ϕ	-1	0	$1 - \phi$	ϕ	1	-1
$20C_3R$	1	0	0	1	-1	-1	-1	1	0

...So now what?

Notable Kronecker Products of I' .

Use Character Table to easily obtain Kronecker Products (known)

$$3' \otimes 3' = 1 \oplus 3' \oplus 5 \quad 3 \otimes 3 = 1 \oplus 3 \oplus 5 \quad 3 \otimes 3' = 4 \oplus 5$$

$$2 \otimes 2 = 1 \oplus 3 \quad 2 \otimes 2' = 4 \quad 2' \otimes 2' = 1 \oplus 3'$$

$$4 \otimes 4 = 1 \oplus 3 \oplus 3' \oplus 4 \oplus 5$$

$$5 \otimes 5 = 1 \oplus 3 \oplus 3' \oplus 4 \oplus 4 \oplus 5 \oplus 5$$

These Kronecker Products will allow us to construct a simple Lagrangian (superpotential) that is invariant under the discrete symmetry that generates our observed quark masses and mixings.

**All of this is abstract. We need actual representations.
Invariants not available in literature.... Work to be done.**

Tensor Product Decomposition (I)

$$3 \otimes 3$$

$$3 = \begin{pmatrix} a_1 \\ a_2 \\ a_3 \end{pmatrix}$$

$$3 \otimes 3'$$

$$1 = a_1 b_1 + a_2 b_2 + a_3 b_3$$

$$3 = (a_3 b_2 - b_2 a_3, a_1 b_3 - a_3 b_1, a_2 b_1 - a_1 b_2)^T$$

$$4 = \begin{pmatrix} \frac{1}{\phi} a_3 b_2 - \phi a_1 b_3 \\ \phi a_3 b_1 + \frac{1}{\phi} a_2 b_3 \\ -\frac{1}{\phi} a_1 b_1 + \phi a_2 b_2 \\ a_2 b_1 - a_1 b_2 + a_3 b_3 \end{pmatrix}$$

$$5 = \begin{pmatrix} a_2 b_2 - a_1 b_1 \\ a_2 b_1 + a_1 b_2 \\ a_3 b_2 + a_2 b_3 \\ a_1 b_3 + a_3 b_1 \\ -\frac{1}{\sqrt{3}}(a_1 b_1 + a_2 b_2 - 2a_3 b_3) \end{pmatrix}$$

$$5 = \begin{pmatrix} \frac{1}{2}(\phi^2 a_2 b_1 + \frac{1}{\phi^2} a_1 b_2 - \sqrt{5} a_3 b_3) \\ -(\phi a_1 b_1 + \frac{1}{\phi} a_2 b_2) \\ \frac{1}{\phi} a_3 b_1 - \phi a_2 b_3 \\ \phi a_3 b_2 + \frac{1}{\phi} a_1 b_3 \\ \frac{\sqrt{3}}{2}(\frac{1}{\phi} a_2 b_1 + \phi a_1 b_2 + a_3 b_3) \end{pmatrix}$$

Tensor Product Decomposition (II)

$$2 \otimes 2 = 1 \oplus 3$$

$$2' \otimes 2' = 1 \oplus 3'$$

$$2 \otimes 2' = 4$$

$$1 = a_2 b_1 - a_1 b_2$$

$$1 = a_2 b_1 - a_1 b_2$$

$$3 = \begin{pmatrix} -ia_1 b_1 + ia_2 b_2 \\ a_1 b_1 + a_2 b_2 \\ ia_2 b_1 + ia_1 b_2 \end{pmatrix}$$

$$3' = \begin{pmatrix} -ia_1 b_1 - ia_2 b_2 \\ -a_1 b_1 + a_2 b_2 \\ a_2 b_1 + a_1 b_2 \end{pmatrix}$$

$$4 = \begin{pmatrix} -a_2 b_1 - a_1 b_2 \\ ia_2 b_1 - ia_1 b_2 \\ -ia_1 b_1 - ia_2 b_2 \\ -a_1 b_1 + a_2 b_2 \end{pmatrix}$$

Now we have the explicit invariants... Onto the best part...

How to Build an I' Flavor Model

Seek scenario with diagonal CKM matrix (at leading order) and a hierarchical mass pattern. This “eliminates” embedding in a 3 and 3', and tells us to embed in 2's.

Our approach: Let 3rd generation quarks be singlets, and embed the 1st and 2nd generation quarks in doublets.

(Standard approach in A'_4/T' models: A. Aranda, et al., F. Feruglio et al., ...)

$$\begin{array}{l} U_{L,R} \rightarrow 2 \\ D_{L,R} \rightarrow 2 \end{array} \quad 2 \otimes 2 = 1 \oplus 3$$

Forbid tree-level mass term (of 1st and 2nd generations) with an additional Z_3 symmetry

Introduce flavon fields to generate mass at higher level through spontaneous symmetry breaking

Let's Start "Cooking"

Field	\bar{t}_L	t_R	\bar{b}_L	b_R	\bar{U}_L	U_R	\bar{D}_L	D_R	H_u	H_d	ρ_u	ρ_d	ξ_u	ξ_d	ρ_u^0	ρ_d^0	ξ_u^0	ξ_d^0
I'	1	1	1	1	2	2	2	2	1	1	1	1	3	3	1	1	3	3
Z_3	0	0	0	0	1	1	2	2	0	0	1	2	1	2	1	2	1	2
$U(1)_R$	1	1	1	1	1	1	1	1	0	0	0	0	0	0	2	2	2	2

With these charge assignments, our effective mass terms to 'leading order' become:

$$w_q = \bar{t}_L t_R + \bar{b}_L b_R + \bar{U}_L U_R \frac{\rho_u}{\Lambda} + \bar{U}_L U_R \frac{\xi_u}{\Lambda} + \bar{D}_L D_R \frac{\rho_d}{\Lambda} + \bar{D}_L D_R \frac{\xi_d}{\Lambda}$$

Breaking the Flavor

$$\langle \xi_d \rangle = \begin{pmatrix} 0 \\ 0 \\ v_d \end{pmatrix} \quad \langle \rho_d \rangle = \rho_d \quad \langle \xi_u \rangle = \begin{pmatrix} 0 \\ 0 \\ v_u \end{pmatrix} \quad \langle \rho_u \rangle = \rho_u$$

Spontaneously Breaking the flavor symmetry, as above, yields the following mass matrices:

$$m_{u,d}^2 \approx \begin{pmatrix} (v_{u,d} - i\rho_{u,d}^2)(v_{u,d}^* + i\rho_{u,d}^{2*}) & 0 & 0 \\ 0 & (v_{u,d} + i\rho_{u,d}^2)(v_{u,d}^* - i\rho_{u,d}^{2*}) & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

Freedom to assume singlets are real and triplets are purely imaginary

$$m_{u,d} = |v_{u,d} - \rho_{u,d}| \quad m_{c,s} = |v_{u,d} + \rho_{u,d}|$$

$$U_{CKM} = \mathbf{1}_{3 \times 3}$$

Outlook and Conclusion

- Summary:
 - The Double Icosahedral Group provides a rich setting for investigating the flavor puzzle.
 - Our work provides a toolbox for using Icosahedral Symmetry as a flavor symmetry group.
- Where to next?
 - Flavon Potential and generate Cabibbo Angle.
 - Apply in Grand Unified Theory contexts and/or look into higher dimensional irreducible representation.
 - Write Papers.