A Stable Pion

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"No Electroweak Symmetry Breaking"

Matter Pie of Our Universe



Dark Matter 83.2%

From WMAP

WIMP

"WIMP Miracle" relates the DM relic abundance to the DM annihilation cross section.

$$\Omega_{\rm DM} \sim \langle \sigma_{_A} \, v \rangle^{-1} \, \sim \, {m_\chi^2 \over g^4}$$

- Three conditions to have a viable thermal DM candidate.
 - DM is stable: requires an unbroken discrete or continuous symmetry.
 - DM interacts with SM particles: requires DM charged under SM gauge symmetries or via Higgs portal.
 - Not to overclose the universe and not ruled out by the current direct detection experiments.

Discrete and Continuous Symmetries in the SM

- Parity, Charge Conjugation, Time Reversal
- Baryon number and lepton number

Citation: S. Eidelman et al. (Particle Data Group), Phys. Lett. B 592, 1 (2004) (URL: http://pdg.lbl.gov)

$$G(J^{PC}) = 1^{-}(0^{-+})$$

Mass $m = 134.9766 \pm 0.0006$ MeV (S = 1.1) $m_{\pi^{\pm}} - m_{\pi^{0}} = 4.5936 \pm 0.0005$ MeV Mean life $\tau = (8.4 \pm 0.6) \times 10^{-17}$ s (S = 3.0) $c\tau = 25.1$ nm

G-parity in the SM

Selection Rules Imposed by Charge Conjugation.

L. MICHEL

European Council for Nuclear Research (CERN) - Theoretical Study Group (*)

(ricevuto il 24 Gennaio 1953)

Charge Conjugation, a New Quantum Number G, and Selection Rules Concerning a Nucleon-Antinucleon System.

T. D. LEE

Columbia University - New York, N.Y.

C. N. YANG

Institute for Advanced Study - Princeton, N.J.

(ricevuto il 30 Gennaio 1956)

G-parity in the SM

 $G = \mathcal{C} \exp[i\pi I_2]$



• The G-parity in the SM is broken by the chiral weak interaction and nonzero hypercharges of quarks.



Why not to use the G-parity to explain the stability of the dark matter particle?



dark G-parity is unbroken

dark pion is the LGP

New G-parity Odd Pions

- Make "new quarks" vector-like under SU(2) weak gauge interaction.
- Choose zero hypercharge for "new quarks".



General Weak Representations



 $SU(2j+1)_L \times SU(2j+1)_R \times U(1)_B \to SU(2j+1)_V \times U(1)_B$



Extension to Color SU(3)

	$SU(N)_h$	$SU(3)_c$	$SU(2)_W$	$U(1)_Y$
$\psi_{L,R}$	N	8	2	0

only possible for real representation



Phenomenologies

- LGP, a dark matter candidate, is always the neutral part of a weak triplet with hypercharge zero.
- Its mass should be below I-2 TeV to not overclose the universe. Its spin-independent scattering cross section is around $10^{-45} \,\mathrm{cm}^2$.
- Its charged partner is a long-lived particle at colliders.
- The G-even pions behave as SM gauge boson resonances at colliders, while the NLGP pions provide final signatures with SM gauge bosons plus missing transverse energy.

Discussions

- The G-parity is unbroken from all renormalizable operators.
- Two dim-five operators violate G-parity:

 $B_{\mu\nu}\bar{\hat{\psi}}\sigma^{\mu\nu}\hat{\psi}/\Lambda_{\rm UV} \qquad \qquad H^{\dagger}t^{a}H\bar{\hat{\psi}}J^{a}\hat{\psi}/\Lambda_{\rm UV}$

- Introducing a Peccei-Quinn symmetry, the LGP can still have a lifetime longer than the age of our universe.
- Hidden sector baryons are also stable. However, they only contribute a very small fraction of the dark matter density:

$$\frac{\Omega_B}{\Omega_{\Pi}} \sim \frac{\langle \sigma v \rangle_{\Pi}}{\langle \sigma v \rangle_B} \sim 4\pi \, \alpha_2^2 \sim 10^{-2}$$

Conclusions

- An unbroken and nontrivial discrete symmetry emerges in a class of models with vector-like SM gauging of fermions with QCD-like strong dynamics.
- The dark matter particle, LGP, could be a stable pion emerges in this class of models.
- The mechanism leads to testable and distinct dark matter and LHC phenomenologies.