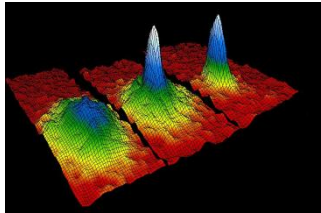


Measurement of Bose-Einstein Correlations in the first LHC-CMS data

Massimo Nespolo
on behalf of
CMS Collaboration

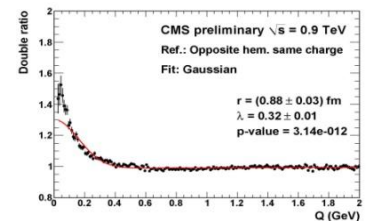
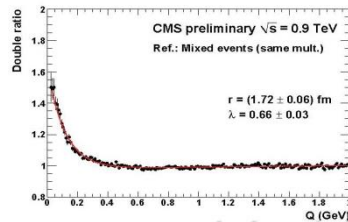
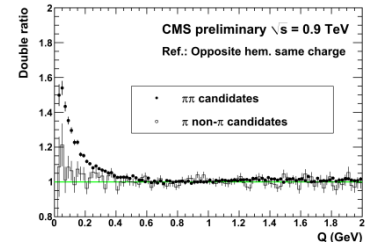
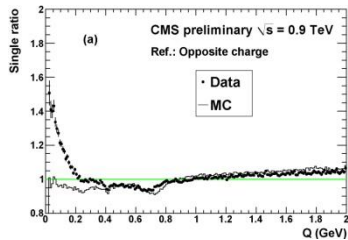
Pheno 2010 (Madison) – May 10-12, 2010

Measurement of Bose-Einstein Correlations in the first LHC-CMS data



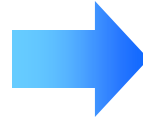
Outline:

- *BEC* in High-Energy Physics:
 - A long story...
- Effect clearly evident in our data:
 - Different reference samples.
 - Particle identification.
- What we measured:
 - Double ratio data/MC.
 - Fit with exp. and gaussian shape.

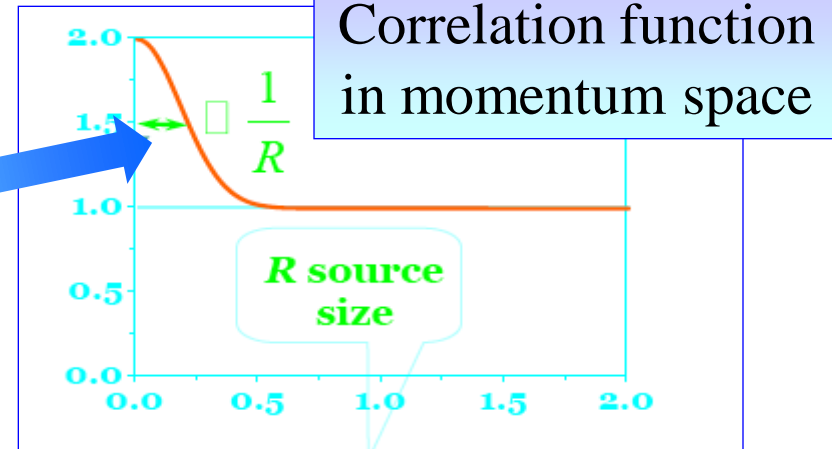
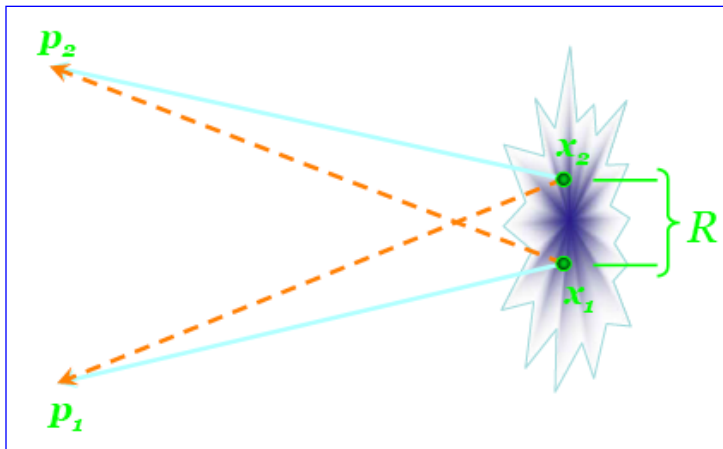


BEC in High Energy Physics

During high-energy collisions, bosons are created at small distance in a “fireball”



Their wave functions overlap, and the Bose-Einstein statistics changes their dynamics



This is essentially the only way to measure the size of a source at the Fermi scale

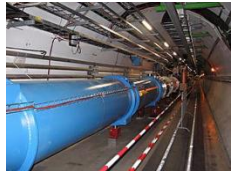
Why measure BEC again?

This is *not* a new phenomenon:

1. measured for the first time in HEP by Goldhaber (1960);
2. Since then, many measurements with different detectors and different initial states (e^+e^- , pp , $p\bar{p}$, πN and $\nu_\mu N$).

However...

1. We have a new accelerator (**LHC**), and therefore higher energies (2.36 TeV in December, now 7 TeV).
2. We have a new, hopefully more powerful detector (**CMS**).



This is one of the *first* physics measurements done with real data!

How to measure BEC?

Theoretically, we need to study the ratio between the joint probability of emission of a pair of bosons, and the individual probabilities

$$R = \frac{P(p_1, p_2)}{P(p_1)P(p_2)}$$

Experimentally, we have to produce the distributions of a “proximity” quantity in the data and in a reference sample (*Coulomb corrected*)

$$R = \frac{dN / dQ_{data}}{dN / dQ_{ref}}$$

To measure the proximity between 2 particles, we chose the difference of 4-momentum (assuming all pions):

$$Q = \sqrt{-(p_1 - p_2)^2} = \sqrt{m_{inv}^2 - 4m_\pi^2}$$

To calculate R:

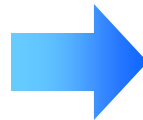
1. Take all (charged) tracks.
2. Construct Q.
3. Repeat for the reference sample.

Reference samples

We used 7 reference samples, mainly taken from literature:

1. *Opposite charge* pairs;
2. *Opposite charge* pairs where one track has its *three-momentum inverted*;
3. *Same-charge* pairs where one track has its *three-momentum inverted* ;
4. *“rotated” pairs*: same charge with one track inverted in the transverse plane;
5. Event mixing 1: every pair has one track from one event, the other from the *following selected event*;
6. Event mixing 2: as above, but events are paired such that they have *similar distribution of $dN_{tracks}/d\eta$* ;
7. Event mixing 3: as above, but events are paired such that they have *similar total invariant mass of charged tracks*;

None of these reference samples is “golden”

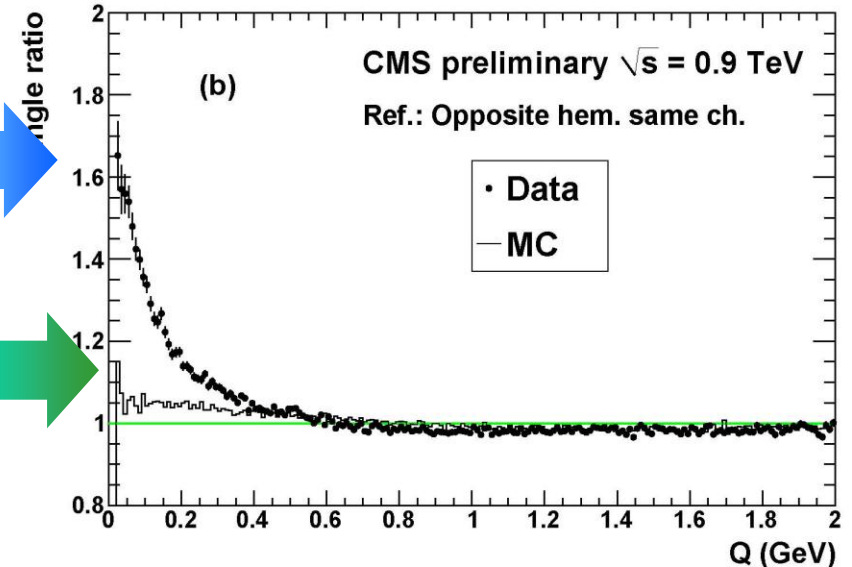
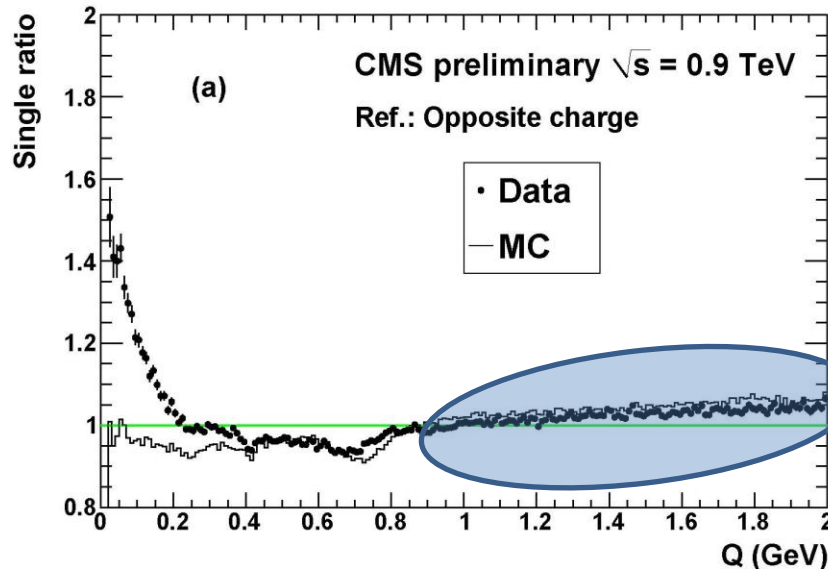


We used all of them for our analysis

Evidence of the effect

The enhancement at $Q=0$ shows the expected correlations

MC is essentially flat (no BEC in the simulations)



Data are well reproduced by MC

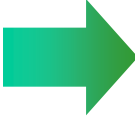
Spectra are *not flat*, in particular for some reference samples

Double ratio data / MC

Every reference sample,
“natural” (opposite-sign) or “artificial”,
is distorted if compared with the real event

Part of the distortion is due to the kinematics or decays,
well described by the MC simulation

To remove the long-range bias,
evident in some of the reference
samples, we used the double
ratio between data and MC


$$\mathcal{R} = \frac{R}{R_{MC}} = \frac{\left(\frac{dN / dQ}{dN / dQ_{ref}} \right)}{\left(\frac{dN / dQ}{dN / dQ_{ref}} \right)_{MC}}$$

Parametrizations

To perform the fit of the double-ratio spectra, we used the function:

$$R(Q) = C[1 + \lambda\Omega(Qr)](1 + \delta Q)$$

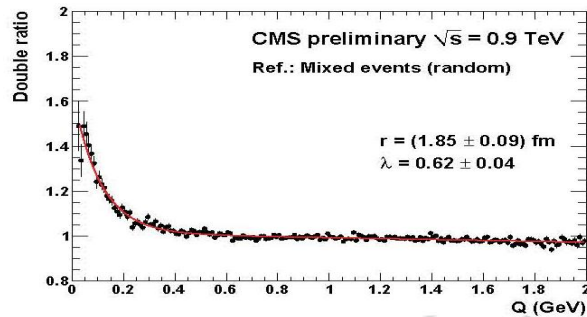
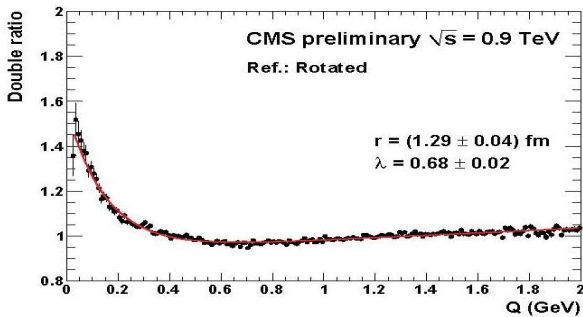
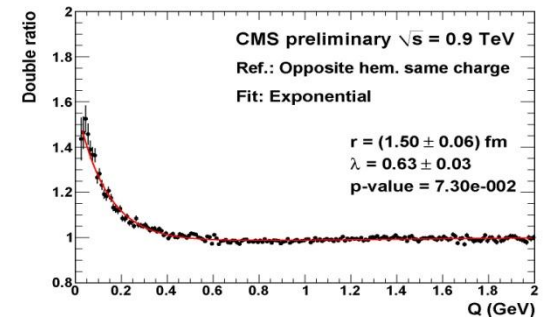
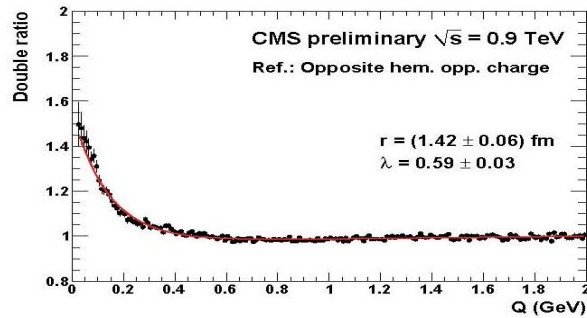
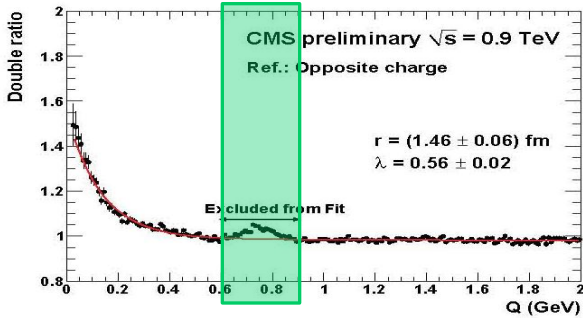
Where λ measures the strength for incoherent boson emission from independent source, δ accounts for long-distance correlations, and C is a normalization factor.

In a static model of particle emission, the $\Omega(Qr)$ function is the *Fourier transform of the emission region*, whose effective size is measured by r .

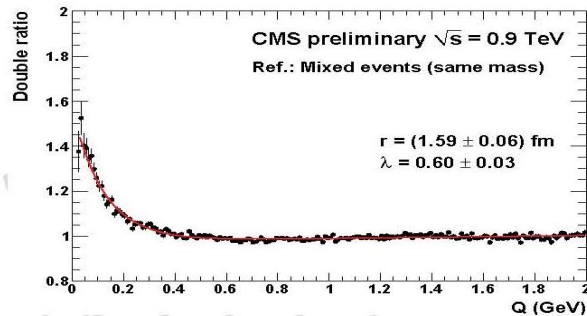
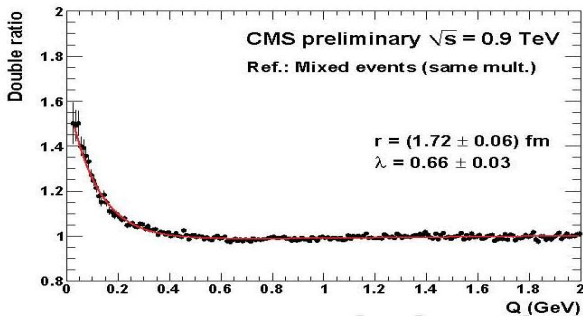
We chose two parametrizations:

1. $\Omega(Qr) = \exp(-Qr)$ – **Exponential**, our default.
2. $\Omega(Qr) = \exp(-Q^2r^2)$ – **Gaussian**, widely used.

Fits for data @ 900 GeV



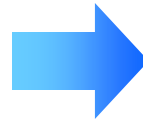
Results are between
1.29 and 1.85 for r ,
0.56 and 0.68 for λ



For the opposite
charge, we removed
the region with
sizeable contribution
from $\rho \rightarrow \pi^+ \pi^-$

Additional check

In our data, many particles have small momentum ($p < 2$ GeV)

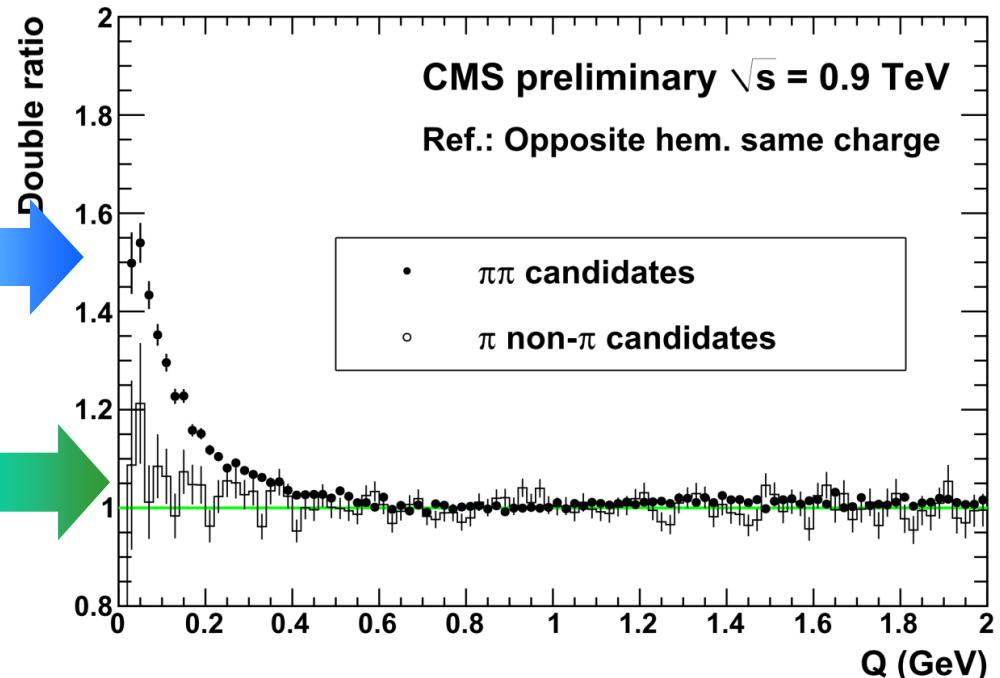


The energy loss in the tracker allows to select π and non- π

Signal is *really due* to BEC (enhanced for $\pi\pi$ pairs)



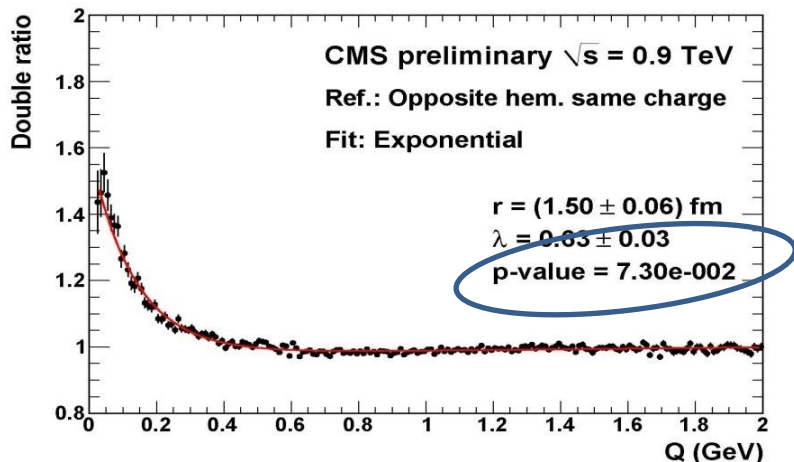
The plot is almost flat for particles of different type



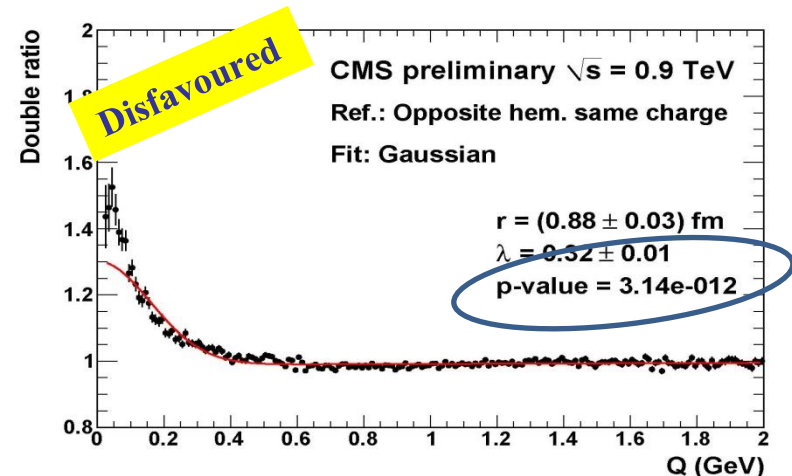
Different parametrizations

We tested: exponential, Gaussian, Levy,
and those described by Kozlov and Biyajima

$$\Omega(Qr) = \exp(-Qr) - \textit{Exponential}$$



$$\Omega(Qr) = \exp(-Q^2 r^2) - \textit{Gaussian}$$



All functions with the Gaussian form have to be ruled out:
the others give equally good results

Combined reference sample

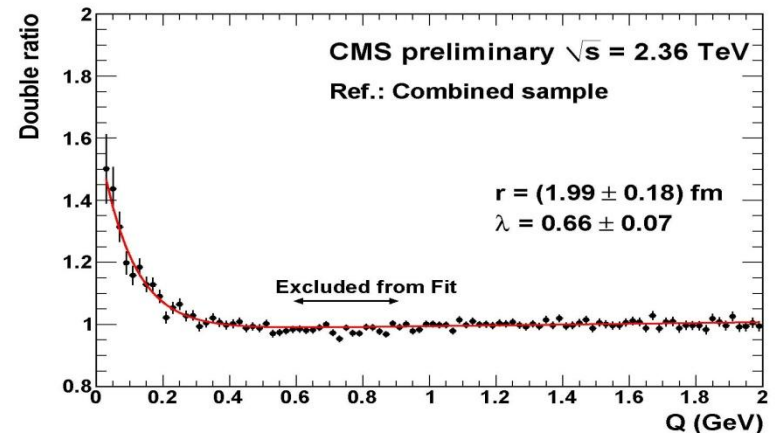
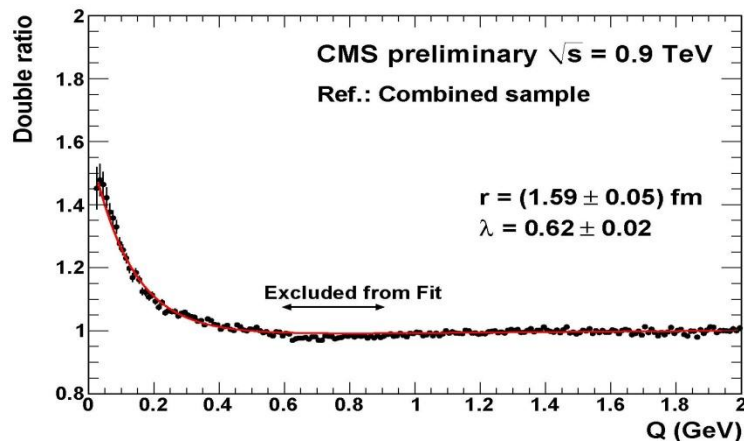
We consider it valuable to provide a single value, together with a conservative estimate of the systematic error

Since the samples are correlated, we cannot simply calculate the average of r and λ



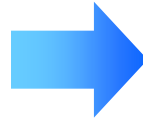
$$\mathcal{R}^{avg} = \frac{dN / dQ}{dN / dQ_{MC}} \frac{\sum_{i=1}^m dN / dQ_{MC}^i}{\sum_{i=1}^m dN / dQ^i}$$

Then we performed an exp. fit for both sets of data:



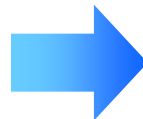
Systematic uncertainties

No reference sample is perfect,
and none can be discarded



r.m.s. of fit results
($\pm 7\%$ for λ , $\pm 12\%$ for r)

Gamow factor to correct
Coulomb effects ($\pm 15\%$)



Propagate agreement margin
($\pm 2.8\%$ for λ , $\pm 0.8\%$ for r)

By comparing MC (*with BEC*) at the reconstruction and generation level, we noticed that variations are within errors.

Data @ 900 GeV:

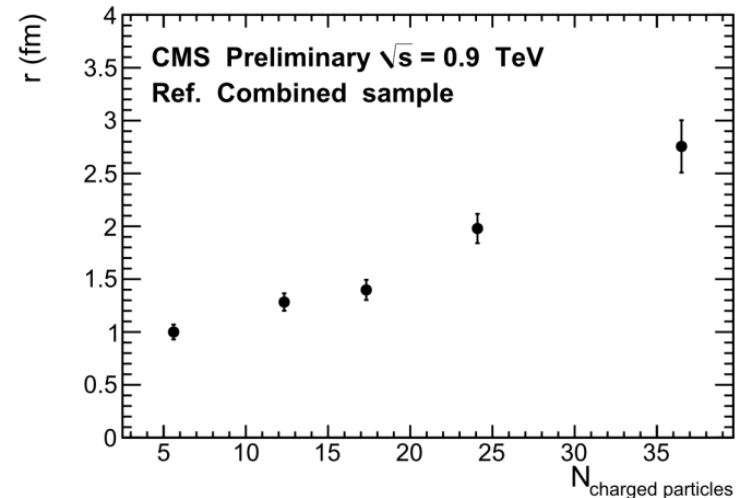
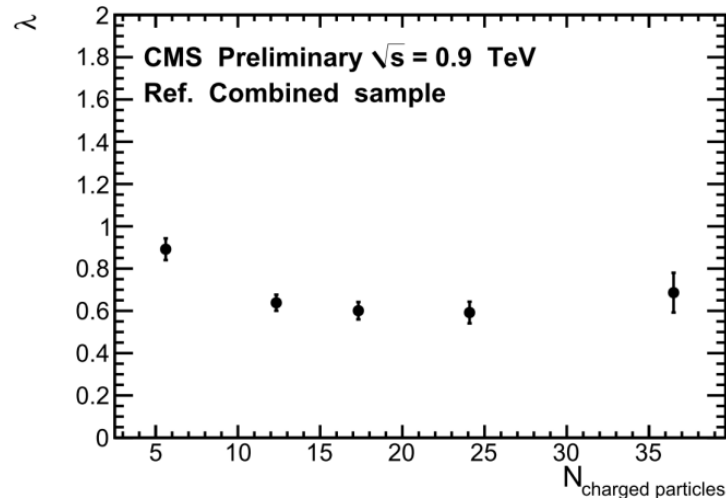
$$r = 1.59 \pm 0.05_{stat.} \pm 0.19_{syst.} \text{ fm}; \lambda = 0.625 \pm 0.021_{stat.} \pm 0.046_{syst.}$$

Data @ 2.36 TeV:

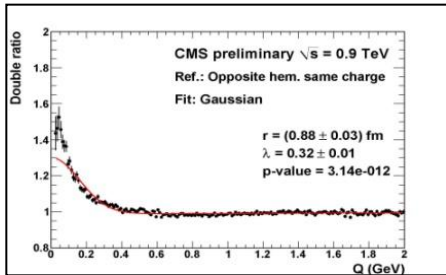
$$r = 1.99 \pm 0.18_{stat.} \pm 0.24_{syst.} \text{ fm}; \lambda = 0.663 \pm 0.073_{stat.} \pm 0.048_{syst.}$$

Dependence on event topology

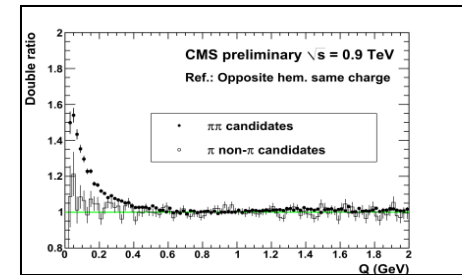
Significant dependence of the r parameter with the charged-particle multiplicity in the event for *all reference samples*



Our results confirm what was noticed with previous experiments in a wide range of energies and initial states



In short...

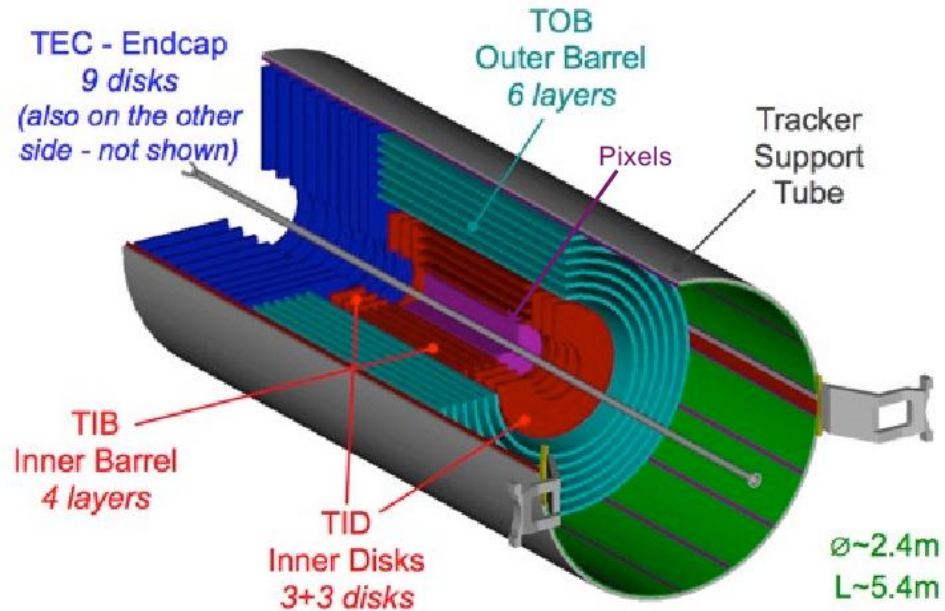


- We immediately saw the effect (as expected):
 - Clearly visible also in the single ratio.
 - Checked with particle identification.
- Measurement and systematic uncertainties:
 - We used the double ratio data/MC.
 - Fit (exp. and Gaussian) with many reference samples.
- We tried to give a single number:
 - Combined reference sample.
 - Dependence from the (charged) track multiplicity.

Backup slides

The CMS Silicon Tracker

- **Pixel:**
 - $\sim 1 \text{ m}^2$ of Si sensors;
 - 66M channels, 1440 modules;
 - 3 barrel layers ($R=4, 7, 11 \text{ cm}$), 2 endcap disks;
- **Strips:**
 - $\sim 198 \text{ m}^2$ of Si sensors;
 - 9.6M channels, 15148 modules;
 - 10 barrel layers, 9+3 endcap wheels per side;
 - $|\eta| < 2.5$.



- **From simulation studies**
 - Tracking efficiency $> 99\%$ (μ), $> 90\%$ (hadrons)
 - Resolution: $\Delta p/p \sim 1\text{-}2\%$ (@100 GeV, $|\eta| < 1.6$)

Additional selections

A track was used if:

1. $p_T > 200$ MeV (to cross all 3 layers of *pixel* detectors);
2. $|\eta| < 2.4$;
3. $N_{\text{dof}} > 5$ and $\chi^2/N_{\text{dof}} < 5.0$;
4. $|d_{xy}| < 0.15$ cm and $R_{\text{innermost}} < 20$ cm.

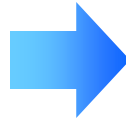
$$0.02 \text{ GeV} < Q < 2.0 \text{ GeV}$$

Avoid not well-separated
or duplicated tracks

Allows to check a good
matching data – ref. sample

Gamow factors

The Coulomb interaction modifies the Q-value distribution of same-charge and different-charge pairs



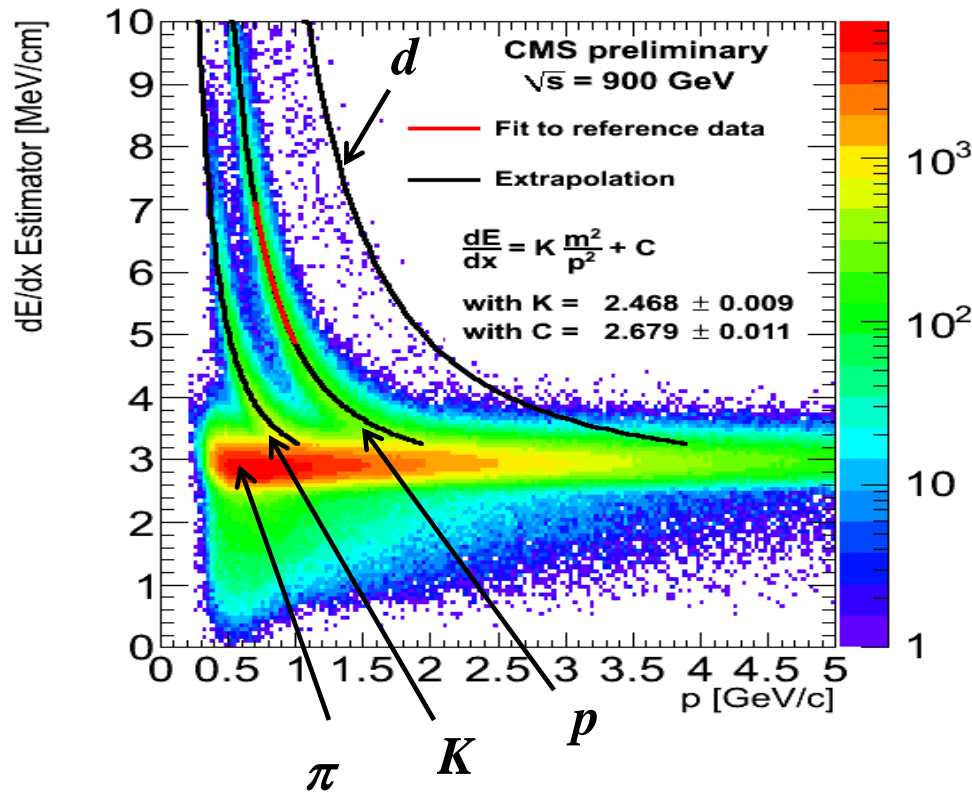
We need to correct this effect, and this is usually done by applying the Gamow factors W_S and W_D

$$W_S(\eta) = \frac{e^{2\pi\eta} - 1}{2\pi\eta} \quad W_D(\eta) = \frac{1 - e^{-2\pi\eta}}{2\pi\eta} \quad \text{with } \eta = \frac{\alpha_{em} m_\pi}{Q}$$

We checked this correction on the ratio data/MC with opposite-charge tracks

Particle Identification

Charged particles with low momentum can be identified using the energy loss in the silicon due to ionization



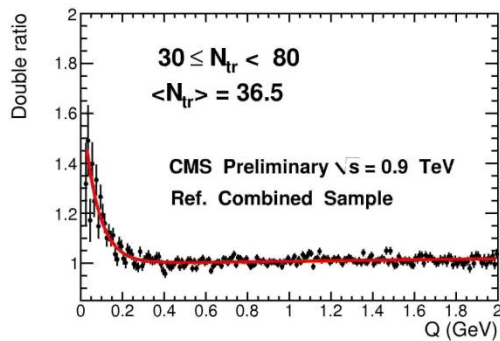
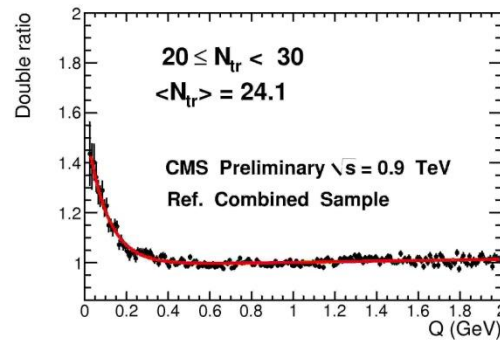
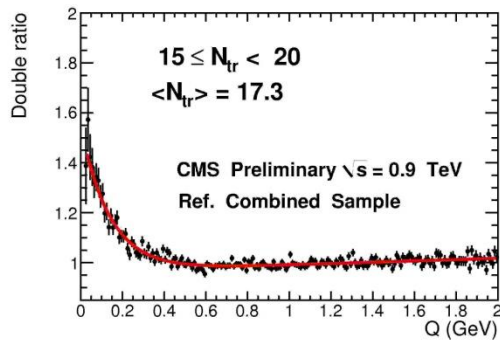
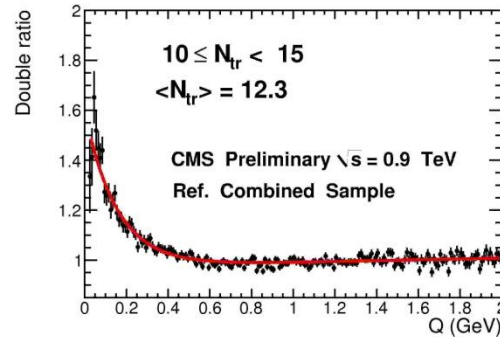
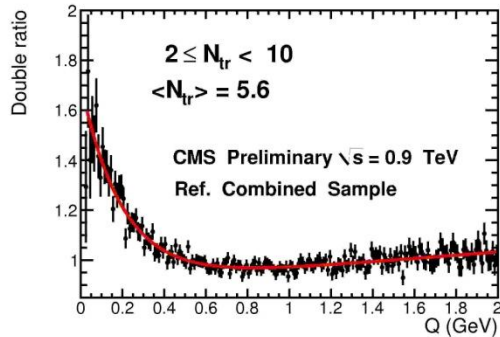
We considered a particle a π if:

- $dE/dx < 3.6 \text{ MeV}$;
- **Or if:**
 - $dE/dx > 3.6 \text{ MeV}$ *and*
 - $M < M_K - 200 \text{ MeV}$.

A non- π has, at the same time:

- $dE/dx > 4.15 \text{ MeV}$ *and*
- $M > M_K - 200 \text{ MeV}$.

Dependence on event topology



Exponential fits of
combined samples for
different values of the track
multiplicity



Shape which clearly
changes with the charged-
tracks multiplicity