Higgs Inflation, Quantum Smearing, and the Tensor to Scalar Ratio

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Motivation

• Non-Supersymmetric Higgs Inflation

• Supersymmetric Higgs (Hybrid) Inflation

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Conclusions

Higgs inflation can be interesting for a number of reasons:

- SM Higgs as inflation? I will not discuss this because of lack of time. In any case the subject is rather controversial.
- ϕ^2 and ϕ^4 inflationary potentials are limiting cases of non-supersymmetric Higgs potential.
- Supersymmetric version of the Higgs potential leads to very different predictions for the tensor to scalar ratio r (measure of primordial gravity waves).
- Last but not least PLANCK, as I will show, will test Higgs inflation models through measurement of the tensor to scalar ratio r.

[Kallosh and Linde, 07; Rehman, Shafi and Wickman, 08]

• Consider the following Higgs Potential:

$$V(\phi) = V_0 \left[1 - \left(\frac{\phi}{M}\right)^2 \right]^2 \quad \longleftarrow \text{(tree level)}$$

Here ϕ is a gauge singlet field.



• WMAP data favors BV inflation.

Limiting Behavior of the Higgs Model

limit	$\phi \longrightarrow M$	$\phi \gg M$	$\phi \ll M$
model	quadratic inflation	quartic inflation	new inflation
V	$\frac{1}{2}m_{\phi}^{2}(\Delta\phi)^{2}$	$\left(rac{V_0}{M^4} ight)\phi^4$	$V_0\left(1-2\left(\frac{\phi}{M}\right)^2\right)$
n_s	$1 - \frac{2}{N_0}$	$1 - \frac{3}{N_0}$	$1 - \frac{2}{N_0} \ln \left(\frac{M}{\sqrt{2}\phi_0} \right)$
r	$4(1-n_s)$	$\frac{16}{3}(1-n_s)$	$8(1-n_s)e^{-N_0(1-n_s)}$

Tree Level Higgs Inflation



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Radiative Corrections in Higgs Inflation

 Consider the following interaction of inflaton φ with some GUT symmetry breaking scalar boson Φ:

$$\mathcal{L}_{int} = \frac{\lambda_{\Phi}^2}{2} \phi^2 \Phi^2$$

Include Radiative Corrections (Quantum Smearing):

$$V = \left(\frac{m^2 M^2}{4}\right) \left[1 - \left(\frac{\phi}{M}\right)^2\right]^2 + A \phi^4 \left[\ln\left(\frac{\phi}{M}\right) - \frac{1}{4}\right] + \frac{A M^4}{4},$$

where
$$V(\phi = 0) \equiv V_0 = \frac{m^2 M^2}{4} + \frac{A M^4}{4}$$
 and $A = \frac{N \lambda_{\Phi}^4}{32 \pi^2}$.

 Note that we can use 'Minkowski space' CW corrections provided the propagating fields have masses >>> H (Hubble constant).

[Rehman and Shafi, 2010]



Note that $r \gtrsim 0.02$ if $n_s \gtrsim 0.96$. Thus, Planck will test Higgs inflation soon!



 $n_s \ge 0.96 \implies A \lesssim 10^{-13.5} \ (\lambda_{\Phi} \lesssim \frac{0.002}{N^{1/4}})$

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Quantum Smearing



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The vacuum energy scale during observable inflation is well below m_P . This implies that the quantum gravity effects are relatively unimportant here.

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Supersymmetric Higgs (Hybrid) Inflation

[Dvali, Shafi, Schaefer; Copeland, Liddle, Lyth, Stewart, Wands '94] [Senoguz, Shafi '04; Linde, Riotto '97]

- Attractive scenario in which inflation can be associated with symmetry breaking $G \longrightarrow H$.
- Simplest inflation model is based on the superpotential

$$W = \kappa S \left(\Phi \,\overline{\Phi} - M^2 \right)$$

S= gauge singlet superfield, $(\Phi\,,\overline{\Phi})$ belong to suitable representation of G

• Need $\Phi, \overline{\Phi}$ pair in order to preserve susy while breaking $G \longrightarrow H$ at scale $M \gg \text{TeV}$, susy breaking scale.

R-symmetry

$$\Phi \, \overline{\Phi} \to \Phi \, \overline{\Phi}, \ S \to e^{i\alpha} \, S, \ W \to e^{i\alpha} \, W$$

 \Rightarrow W is a unique renormalizable superpotential

- Some examples of gauge groups:
 - $G = U(1)_{B-L}$, (Supersymmetric superconductor)

$$G = SU(5) \times U(1)$$
, ($\Phi = 10$), (Flipped $SU(5)$)

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$$G = 3_c \times 2_L \times 2_R \times 1_{B-L}, \ (\Phi = (1, 1, 2, +1))$$

$$G = 4_c \times 2_L \times 2_R, \ (\Phi = (\overline{4}, 1, 2)),$$

 $G = SO(10), \ (\Phi = 16)$

• Tree Level Potential

$$V_F = \kappa^2 \left(M^2 - |\Phi^2| \right)^2 + 2\kappa^2 |S|^2 |\Phi|^2$$

This is similar to a superconductor with S playing the role of temperature, and $\langle \overline{\Phi} \rangle$, $\langle \Phi \rangle$ are Cooper pairs.

Susy vacua



Take into account radiative corrections (because during inflation $V \neq 0$ and susy is broken by $F_S = -\kappa M^2$)

• Mass splitting in $\Phi - \overline{\Phi}$

$$m_{\pm}^2 = \kappa^2 S^2 \pm \kappa^2 M^2$$
, $m_F^2 = \kappa^2 S^2$

One-loop radiative corrections

$$\Delta V_{1\mathsf{loop}} = \frac{1}{64\pi^2} \mathsf{Str}[\mathcal{M}^4(S)(\ln\frac{\mathcal{M}^2(S)}{Q^2} - \frac{3}{2})]$$

• In the inflationary valley ($\Phi=0)$ with $|S|\gg M$

$$V \approx \kappa^2 M^4 \left(1 + \frac{\kappa^2 \mathcal{N}}{8\pi^2} \ln(|S/M|) \right)$$

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[Dvali, Shafi, Schaefer '94]

Tree Level plus radiative corrections:



$$n_s \approx 1 - \frac{1}{N_0} \approx 0.98$$

 $\delta T/T \propto (M/M_P)^2 \sim 10^{-5} \longrightarrow \text{attractive scenario} (M \sim M_G)$

[Senoguz, Shafi '04]

 Take into account sugra corrections, radiative corrections and soft susy breaking terms:

$$V \simeq \kappa^2 M^4 \left(1 + \left(\frac{M}{m_p}\right)^4 \frac{x^4}{2} + \frac{\kappa^2 \mathcal{N}}{8\pi^2} F(x) + a \left(\frac{m_{3/2} x}{\kappa M}\right) + \left(\frac{M_S x}{\kappa M}\right)^2 \right),$$

where
$$a = 2 |2 - A| \cos[\arg S + \arg(2 - A)], x = |S|/M$$
 and
 $F(x) = \frac{1}{4} \left((x^4 + 1) \ln \frac{(x^4 - 1)}{x^4} + 2x^2 \ln \frac{x^2 + 1}{x^2 - 1} + 2 \ln \frac{\kappa^2 M^2 x^2}{Q^2} - 3 \right).$

Note: No ' η problem' with minimal (canonical) Kähler potential !

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[Rehman, Shafi, Wickman, 2009]



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 $r \lesssim 10^{-4}$ within 2- σ bounds of WMAP data

Conclusions

- One of the most important challenges is to find a "Standard Model of Inflationary Cosmology".
- Radiative corrections are important in the context of precision cosmology.
- Non-supersymmetric GUT inflation models typically predict an 'observable' value for the tensor to scalar ratio $r~(\geq 0.02)$, for $n_s \geq 0.96$)
- Supersymmetric Higgs (hybrid) models by comparison predict 'tiny' values of $r \ (\lesssim 10^{-4})$.
- Results from PLANCK are eagerly awaited!