

Precision predictions for Beyond the Standard Model processes.

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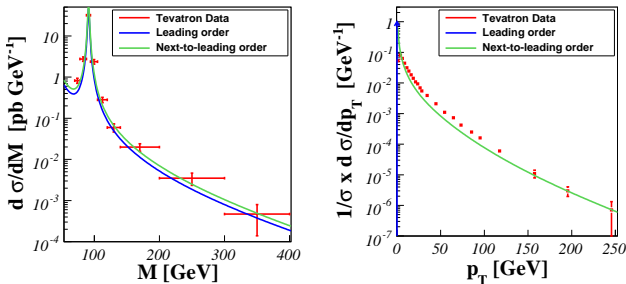
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Outline

- 1 Motivation for precision calculations
- 2 Resummation formalism
- 3 Numerical results, including uncertainties, for supersymmetry and Z'
- 4 Summary - conclusions

Need for precision: Drell-Yan process at the Tevatron

- Confrontation between theory and Tevatron data [D ϕ collaboration (2005, 2008)].



- * LO calculation: **Disagreement between theory and experiment.**
- * NLO invariant-mass distribution: **good agreement.**
- * NLO p_T -distribution:
 - ◇ **Very good agreement in the large- p_T region.**
 - ◇ **Underestimation in the intermediate- p_T region.**
 - ◇ **Divergence in the small- p_T region.**

- How to improve NLO predictions [in particular for the small- p_T region]?

Investigation of the next-to-leading order contributions

- **Partonic invariant-mass and transverse-momentum distributions at $\mathcal{O}(\alpha_s)$,**

$$\frac{d\hat{\sigma}}{dM} = \hat{\sigma}^{(0)}(M) \delta(1-z) + \alpha_s \hat{\sigma}^{(1)}(M, z) + \mathcal{O}(\alpha_s^2),$$

$$\frac{d^2\hat{\sigma}}{dM dp_T} = \hat{\sigma}^{(0)}(M) \delta(p_T) \delta(1-z) + \alpha_s \hat{\sigma}^{(1)}(M, z, p_T) + \mathcal{O}(\alpha_s^2),$$

where $z = M^2/s$.

- $\hat{\sigma}^{(1)}$ contains different pieces.

- * **Real gluon emission** diagrams.

$$iM \approx g_s T^a \left[\frac{\epsilon \cdot k_2}{k_2^0 \mathbf{k}_g^0 (1 + \cos \theta)} - \frac{k_1 \cdot \epsilon}{k_1^0 \mathbf{k}_g^0 (1 - \cos \theta)} \right] iM^{\text{Born}}$$

- * **Virtual loop** contributions.

$$iM \approx (i g_s^2) \int dk_g \frac{k_1 \cdot k_2}{k_g^2 (k_1^0 \mathbf{k}_g^0 (1 - \cos \theta)) (k_2^0 \mathbf{k}_g^0 (1 + \cos \theta))} iM^{\text{Born}}$$

Soft and collinear radiation diverges and factorizes.

The problem of the soft and collinear radiation

- **Sum of the two contributions.**

$$\hat{\sigma}^{(1)} = \hat{\sigma}^{(1,\text{loop})} + \hat{\sigma}^{(1,\text{real})}$$

- * **Cancellation** of the poles.
- * **Infrared behaviour:** logarithmic terms in the distributions,

$$\alpha_s \left(\frac{\ln(1-z)}{1-z} \right)_+ \quad \text{and} \quad \frac{\alpha_s}{p_T^2} \ln \frac{M^2}{p_T^2}$$

- * **Problems at $z \lesssim 1$ or small p_T .**

The fixed-order theory is unreliable in these kinematical regions.

Improvements

Improvements of the next-to-leading order calculation.

- Matching with a resummation calculation.
 - * **Correct treatment** of the soft and collinear radiation.
 - * **Perturbative method.**
 - * **Parton-level calculation.**
 - * Three formalisms:
 - **p_T -resummation**: universally resums $\frac{\alpha_s}{p_T} \ln \frac{M^2}{p_T}$.
 - **Threshold resummation**: universally resums $\left(\frac{\ln(1-z)}{1-z} \right)_+$.
 - **Joint resummation**: universally resums both logs.

[Catani, de Florian, Grazzini (2001); Bozzi, Catani, de Florian, Grazzini (2006); Sterman (1987); Catani, Trentadue (1989,1991); Bozzi, BenjF, Klasen (2008)]
- Matching with a parton shower algorithm.
 - * **Approximation of the resummation calculation.**
 - * **Suitable for a proper description of the collision.**

The resummed component

- **Based on factorization properties.**

- * **Holds in non-physical conjugate spaces.**
- * Mellin N -space (N conjugate to M^2/S_h).
- * Impact parameter b (conjugate to p_T).

$$d\sigma^{(\text{res})}(N, b) = \sum_{a,b} f_{a/h_1}(N+1) f_{b/h_2}(N+1) \mathcal{W}_{ab}(N, b),$$

$$\mathcal{W}_{ab}(N, b) = \mathcal{H}_{ab}(N) \exp\{\mathcal{G}(N, b)\}.$$

- The \mathcal{H} -coefficient:

- * Contains **real and virtual collinear radiation, hard contributions.**
- * **Can be computed perturbatively as series in α_s ,** from fixed-order results.
- * Is process-dependent.

- The Sudakov form factor \mathcal{G} :

- * Contains **the soft-collinear radiation** (the logarithmic terms).
- * Resummed to all orders in α_s .
- * **Can be computed perturbatively as series in $\alpha_s \log$.**
- * **Is process-independent (universal).**

The finite component - matching to the fixed order

- **Fixed-order calculations.**
 - * **Reliable far from the critical kinematical regions.**
 - * **Spoiled in the critical regions.**
- **Resummation.**
 - * **Needed in the critical regions.**
 - * **Not justified far from the critical regions.**
- **Intermediate kinematical regions:** both should contribute.

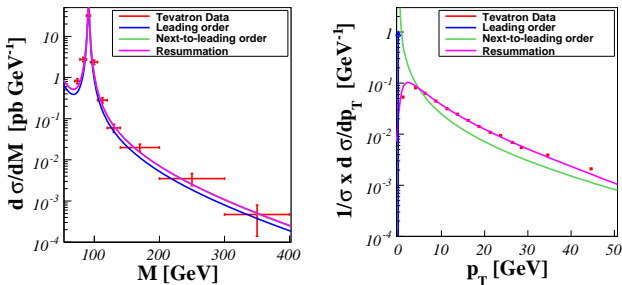
Information from both fixed order and resummation is required.
⇒ consistent matching procedure.

- **Matching procedure:**
 - * Addition of both resummation and fixed-order results.
 - * Subtracting the **expansion** in α_s of the resummed result.
 - * No double-counting of the logarithms.

$$d\sigma = d\sigma^{(\text{F.O.})} + d\sigma^{(\text{res})} - d\sigma^{(\text{exp})}.$$

Resummation vs. Tevatron data

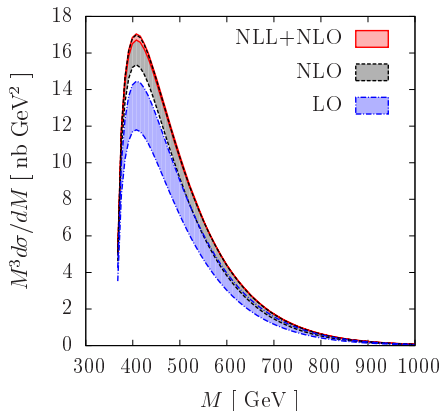
- Confrontation between theory and Tevatron data [DØ collaboration (2005, 2008)].



- * Invariant-mass distribution: **good agreement**.
(no change with respect to next-to-leading order).
- * p_T -distribution: **good agreement**.
(big improvement with respect to next-to-leading order).

Chargino-neutralino associated production: mass-spectrum

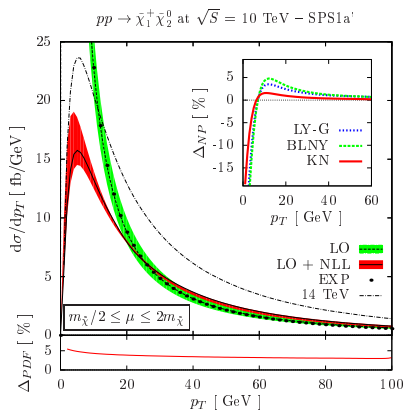
$$p\bar{p} \rightarrow \tilde{\chi}_1^+ \tilde{\chi}_2^0 \text{ at } \sqrt{S} = 1.96 \text{ TeV}$$



[Debove, BenjF, Klasen (in prep.)]

- **Scenario:** ≈ 180 GeV gauginos.
- **Invariant-mass spectrum**
 - * **NLO:** 20-25% increase.
 - * **Resummation:** moderate increase.
- **Scale dependence** ($M/2 \leq \mu_R = \mu_F \leq 2M$).
 - * **LO:** very large dependence.
 - * **NLO:** large dependence.
 - * **Resummation:** reduced dependence.
(higher order terms in the Sudakov).

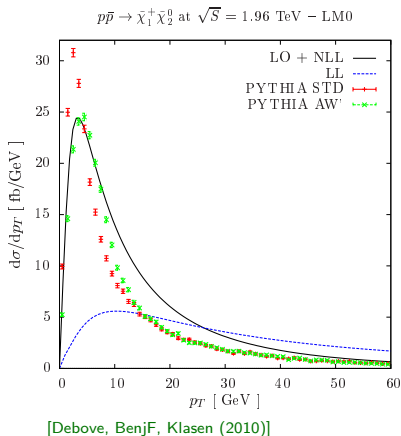
Chargino-neutralino associated production: p_T -spectrum



[Debove, BenjF, Klasen (2010)]

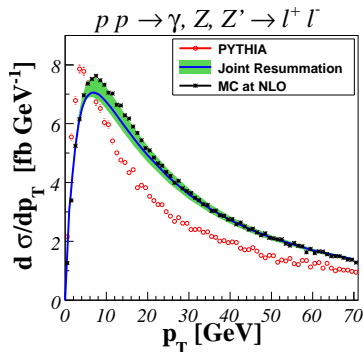
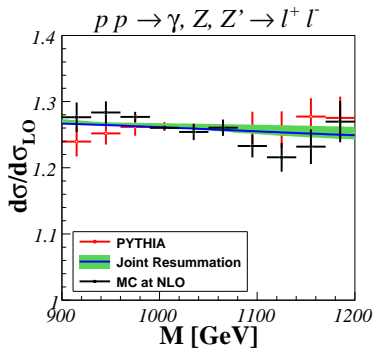
- **Scenario:** ≈ 180 GeV gauginos.
- **Matching effects**
 - * **Small p_T :** expansion \approx fixed-order.
 - * **Large p_T :** expansion \approx resummation.
 - * **Intermediate p_T :** enhancement.
- **Scale dependence** ($M/2 \leq \mu_R = \mu_F \leq 2M$).
 - * **Reduction** of the uncertainties.
 - * Less than **5%** for $p_T > 5$ GeV.
- **Parton densities dependence** (44 CTEQ sets).
 - * **4-5%** uncertainties for all p_T .
 - * Similar to **weak boson** production.
- **Non perturbative effects at low p_T .**
 - * Intrinsic p_T of the partons.
 - * **Under control** for $p_T > 5$ GeV.
- **Uncertainties under control for $p_T > 5$ GeV.**

Comparison: PYTHIA and p_T -resummation



- **Scenario** ≈ 110 GeV gauginos.
- **PYTHIA predictions.**
 - * Used for SUSY **experimental analyses.**
 - * **Leading log** Sudakov form factor.
 - * **Two tunes.**
 - ◇ CDF-AW.
 - ◇ Our tune AW'.
- **Two set of resummed predictions.**
 - * **Leading logarithmic** approximation.
 - * **Next-to-leading logarithmic** results.
- **PYTHIA results.**
 - * **Improves** the LL picture.
 - * **Intrinsic p_T** helps to reproduce NLL.
 - * **Underestimation** for intermediate p_T .
 - * **Direct impact for experimental analyses.**

Comparison: PYTHIA, MC@NLO and joint resummation



[BenjF, Klasen, Ledroit, Li, Morel (2008)]

- 1 TeV Z' ; PYTHIA (LO/LL $_+$), MC@NLO (NLO/LL), resummation (NLO/NLL).
- **Mass-spectrum normalized to leading order:**
 - * PYTHIA (*power shower*): mass-spectrum multiplied by a K -factor of 1.26.
 - * **Good agreement between MC@NLO and resummation.**
- **Transverse-momentum distribution:**
 - * **PYTHIA spectrum much too soft, peak not well predicted.**
 - * **Good agreement between MC@NLO and resummation.**

Summary - conclusions

● Soft and collinear radiation:

- * Large logarithmic corrections in p_T - and invariant-mass spectra.
- * **Need for resummation (or parton showers).**

● p_T , threshold and joint resummations have been implemented.

- * Reliable perturbative results.
- * Correct quantification of the soft-collinear radiation.
- * **Important effects**, even far from the critical regions.
- * **Uncertainties from scales and parton densities under good control.**
- * **Reduced dependence on non-perturbative effects.**

● Comparison with Monte Carlo generators

- * **Significant shortcomings in normalization and shapes for PYTHIA.**
- * **MC@NLO reaches (almost) the same precision level as resummation.**
BUT: easier implementation in the analysis chains of any experiment.