

# HIGGS LOOK-ALIKES AT THE LHC

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**Joseph Lykken**  
**Fermilab**

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# OUTLINE

- **Higgs look-alikes: nonstandard spin and CP**
- **Higgs look-alikes: nonstandard electroweak but custodially symmetric**
- **Higgs look-alikes: nonstandard everything**

## References:

### Higgs look-alikes:

- Alvaro De Rujula, J.L., Maurizio Pierini, Chris Rogan, Maria Spiropulu, arXiv:1001.5300
- Y. Gao, A. Gritsan, Z. Guo, K. Melnikov, M. Schulz, N. V. Tran, arXiv:1001.3396
- Ian Low and J.L., arXiv:1005.0872

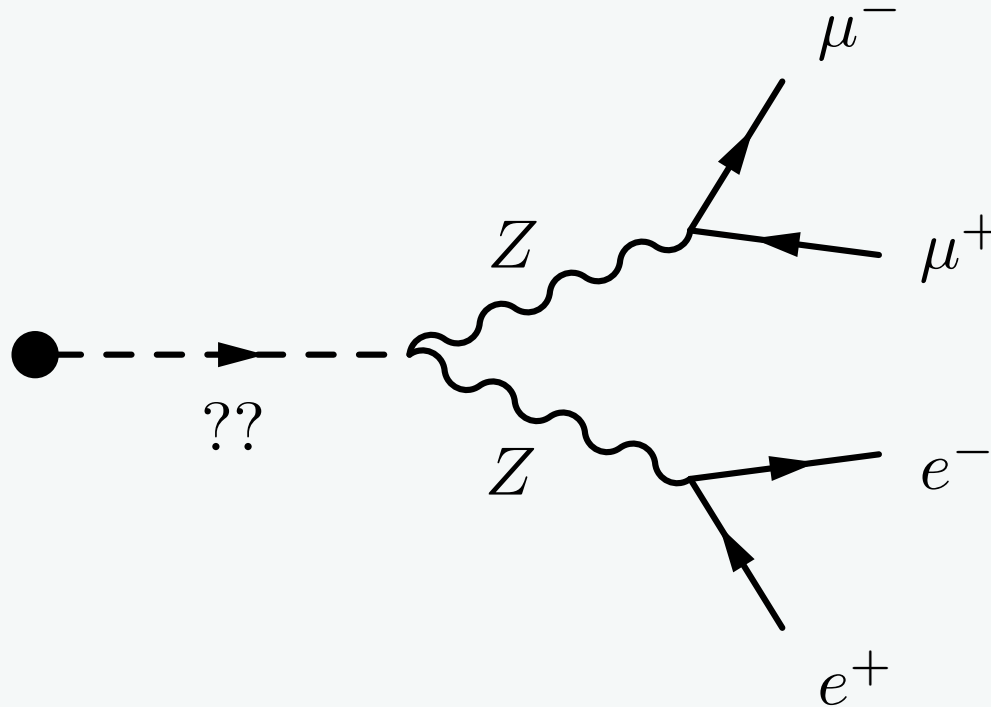
see also R. Lafaye, T. Plehn, M. Rauch, D. Zerwas, M. Duehrssen, arXiv:0904.3866  
and Stephen Godfrey and Ken Moats, arXiv.1003.3033

### MSSM Higgs with general dimension 6 operators:

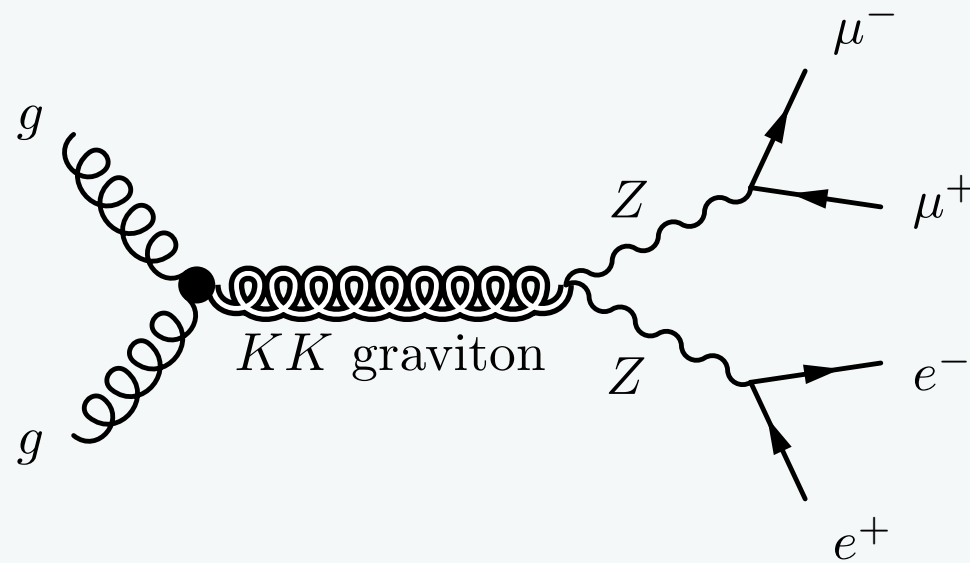
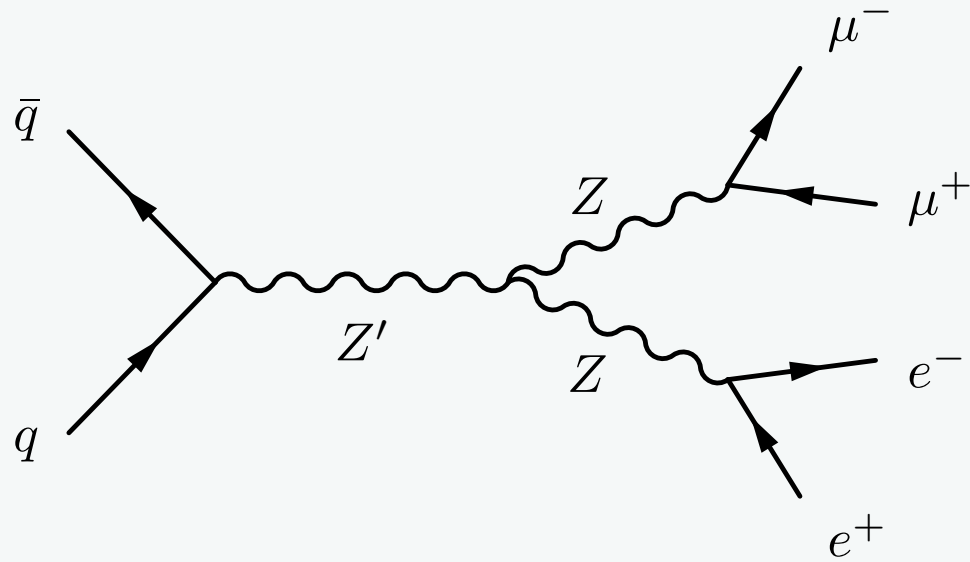
- Marcela Carena, K.C. Kong, Eduardo Ponton, Jose Zarita, arXiv:0909.5434, and TBA

## Higgs look-alikes

- Suppose your favorite LHC experiment sees a resonant signal
- How do we determine that this is the neutral CP-even spin 0 component of a  $(2_L, 2_R)$  of  $SU(2)_L \times SU(2)_R$  predicted by the Standard Model, or a look-alike?



## We need to exclude other spins, etc



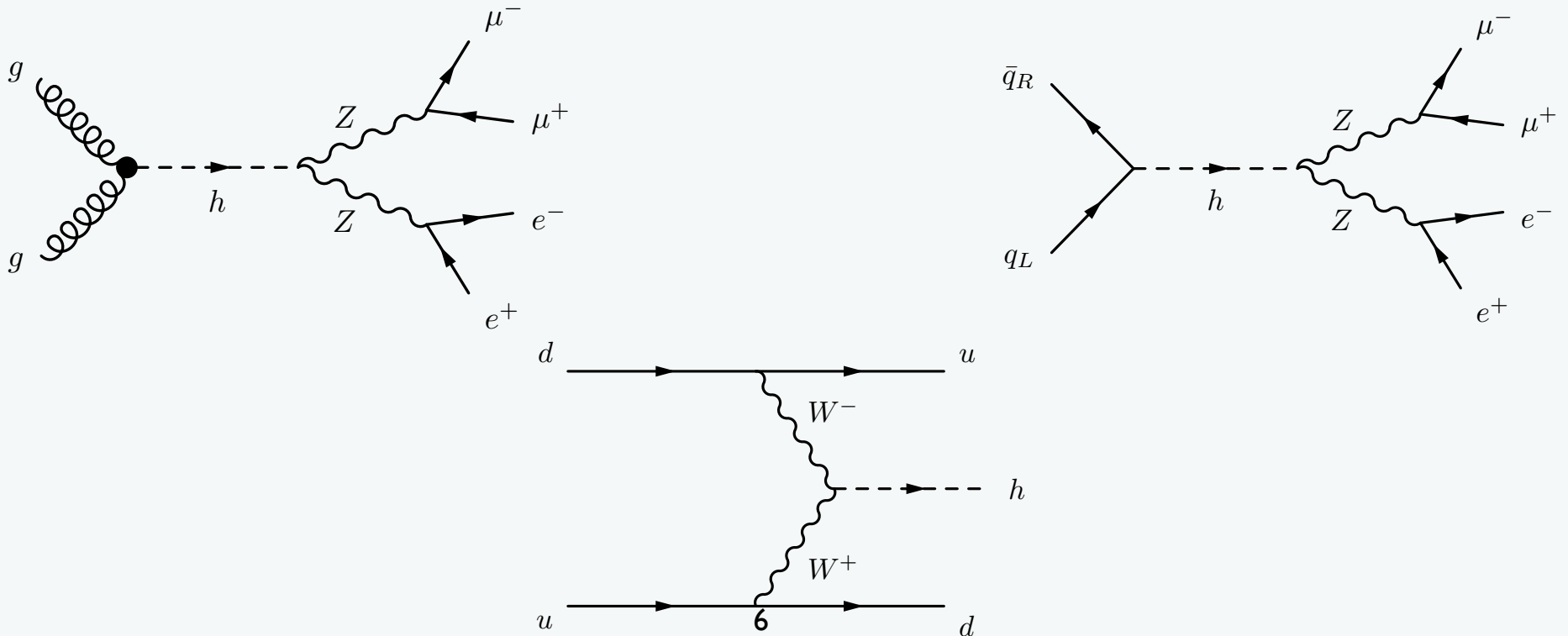
## The post-discovery LHC Higgs challenge

- You have ~100 signal events that could be a Higgs
- How many Higgs look-alike candidates can you eliminate *at or around the time of discovery?*
- A simpler question: How many Higgs look-alike candidates can you eliminate at or around the time of discovery *by looking at distributions and correlations in the 4 lepton final state?*

# Factorizing the problem

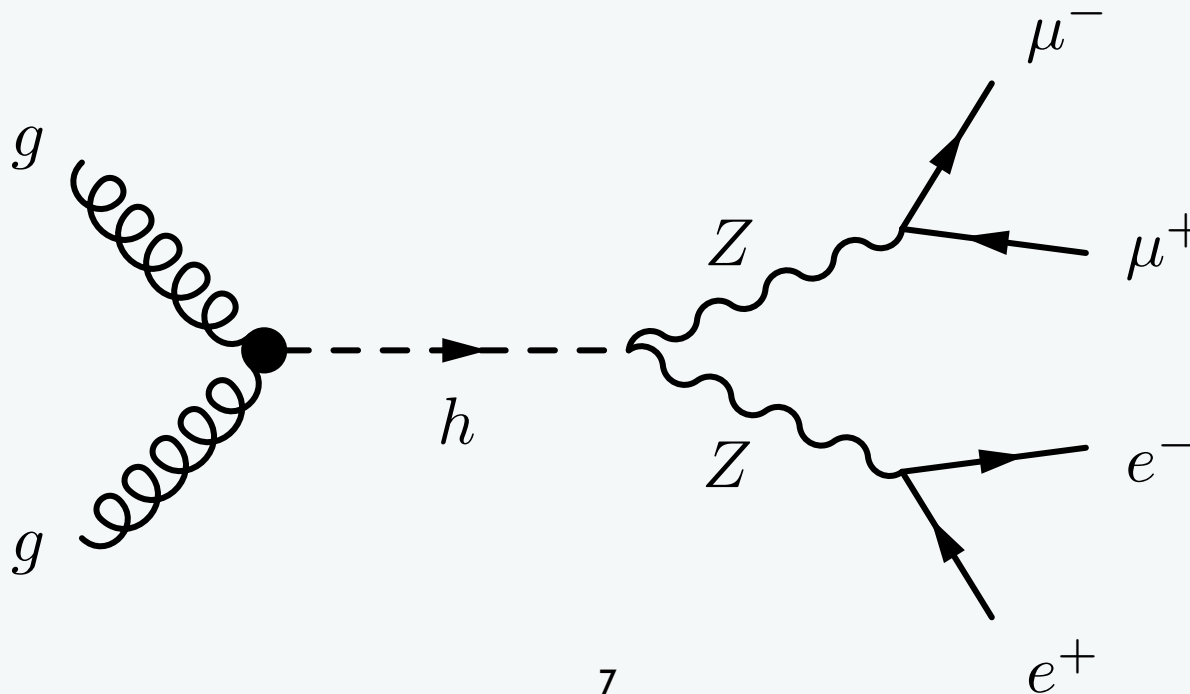
## ✓ Distributions and correlations in the 4 lepton final state

- Production (gluon fusion, VBF, ...)
- Correlations with signals (or lack of signals) in other channels



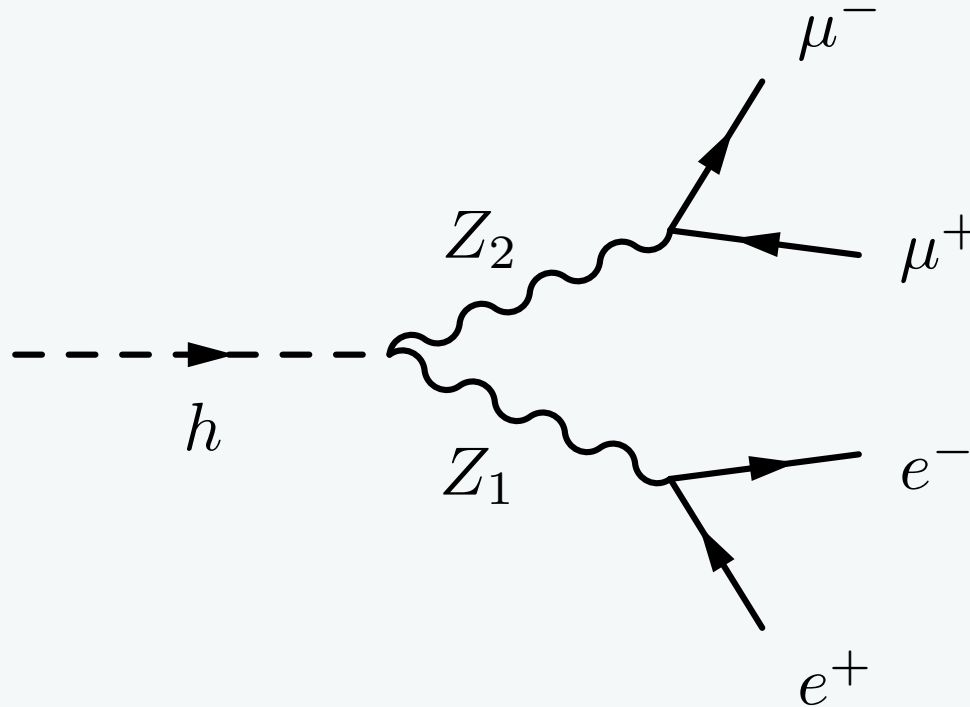
# The golden Higgs channel at the LHC

- The leptonic decay  $h \rightarrow ZZ \rightarrow 4\ell$  has a small branching fraction but provides a (relatively) clean and fully-reconstructable final state
- The Z bosons don't have to be on shell
- Relevant for SM Higgs mass above about  $\sim 130$  GeV



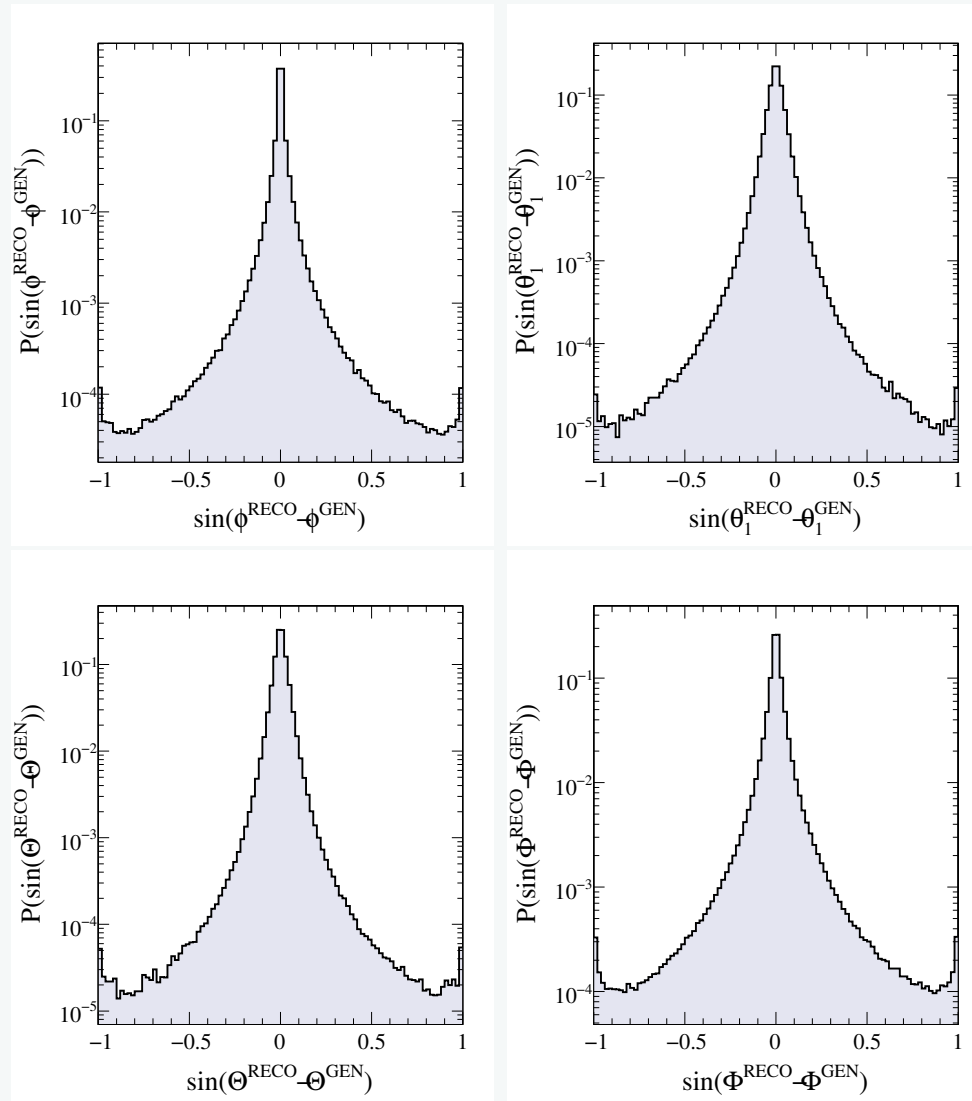
# The 12 observables of the fully reconstructed event

- Treating the leptons as massless, there are 12 momentum components per event measured in the final state
- Each event is fully reconstructable in terms of these 12 observables





# ATLAS and CMS can measure the 4-lepton final state with exquisite precision



So you can choose any basis you want for your 12 observables without losing experimental realism

## The 12 observables of the fully reconstructed event

- To get from the lab frame to the Higgs rest frame, I need to specify a boost and the direction of the boost, which is given by two angles:

$$\gamma_h, \theta_h, \phi_h$$

- I need to specify the reconstructed Higgs mass

$$M_h$$

- In the Higgs rest frame, by convention, take the positive z-axis to be along the direction of motion of  $Z_2$ , then use two angles to specify the direction of one of the incoming partons (note 2-fold ambiguity)

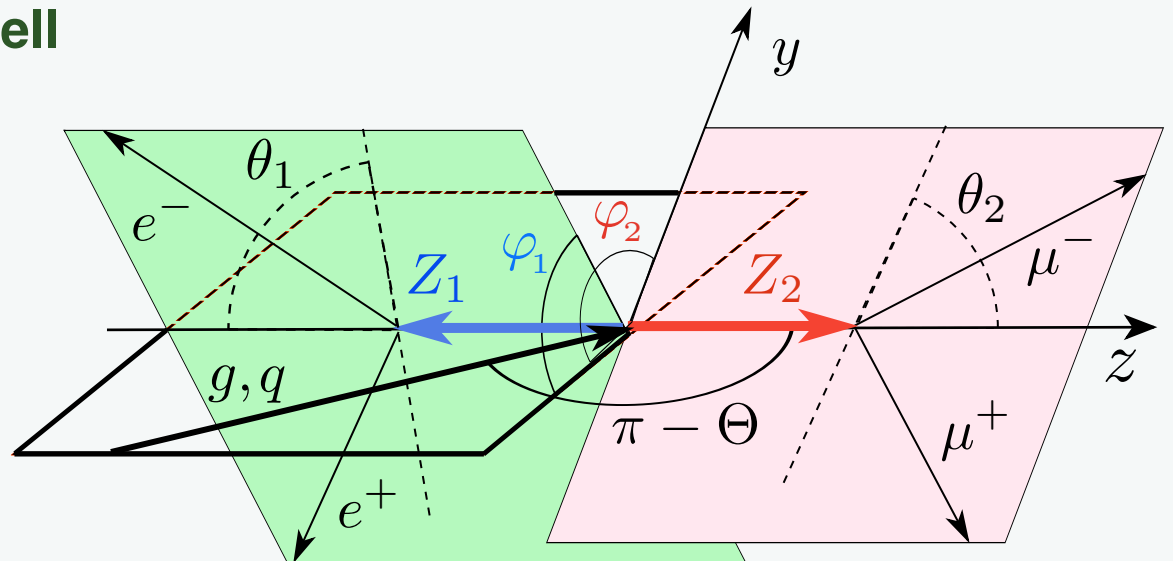
$$\Theta, \Phi$$

- Z decay involves another pair of angles measured in the Z rest frame, with the polar angle measured wrt the z-axis defined above. We also need the two boosts from the Higgs rest frame to the Z rest frames,  $\gamma_1, \gamma_2$ , which is equivalent to specifying the (possibly off-shell) Z masses:

$$m_1, \theta_1, \phi_1, \quad m_2, \theta_2, \phi_2$$

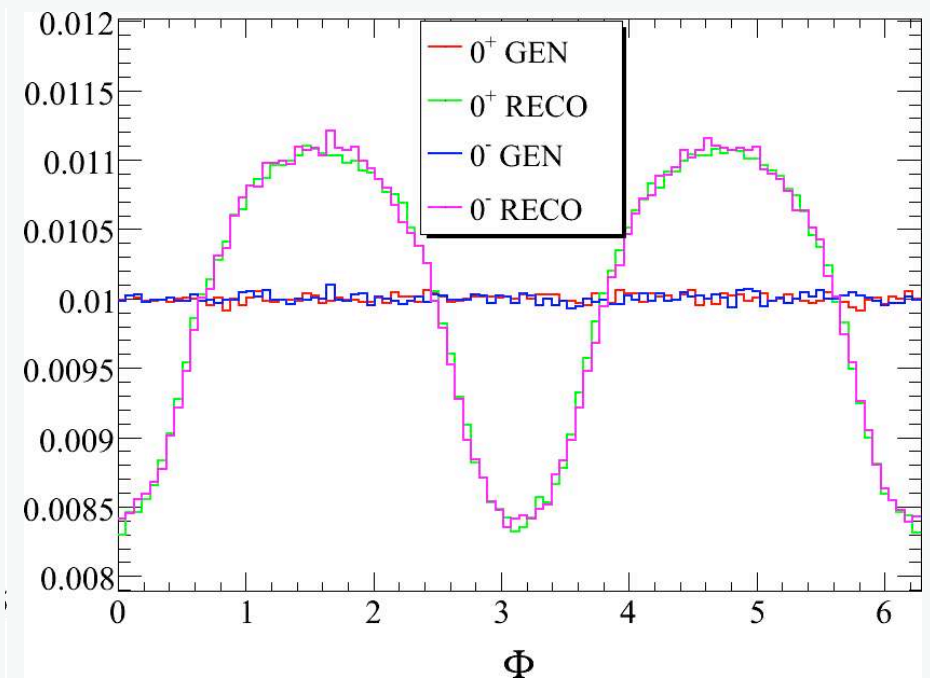
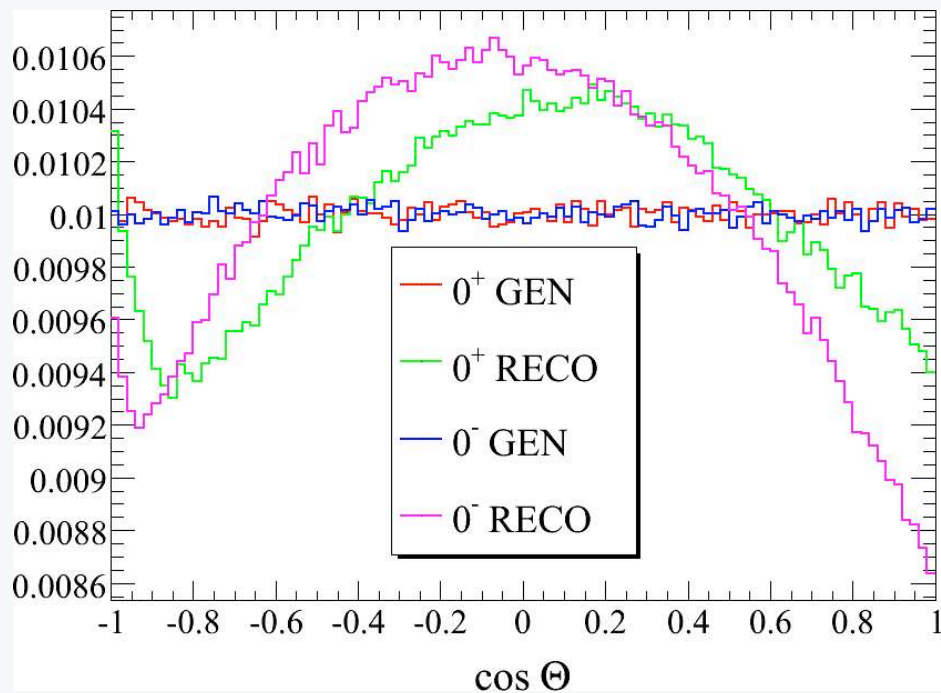
## 8 angles!

- In the spirit of factorization, we will (for now) ignore the two production angles  $\theta_h, \phi_h$
- If the resonance is a spin 0 particle, the signal distribution will be isotropic (i.e. flat) in the  $h \rightarrow ZZ$  angles  $\Theta, \Phi$
- Twenty-year-old common wisdom says that therefore we should ignore these angles as well
- Is this reasonable?

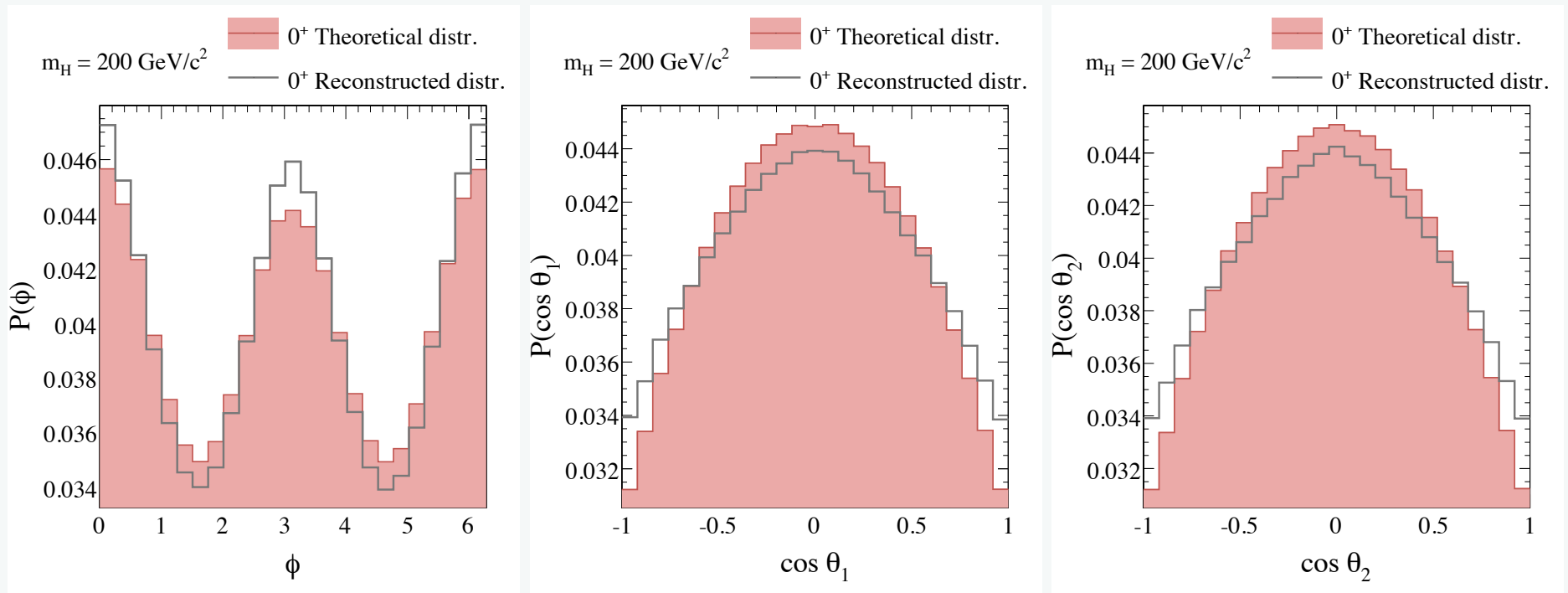


No!

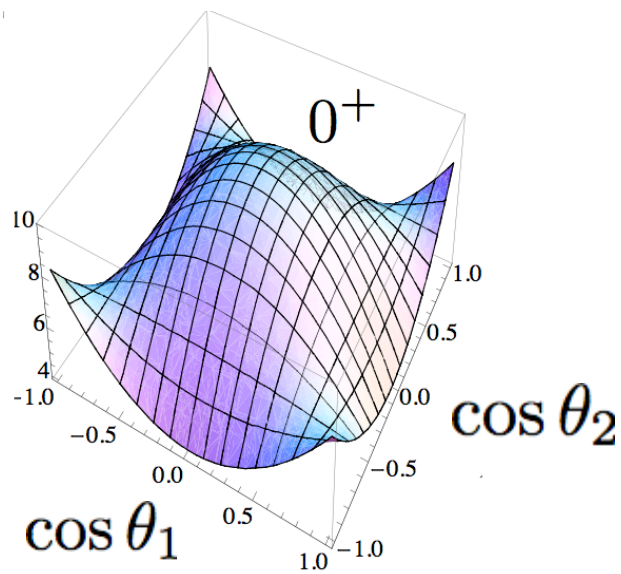
- If we want to test that the Higgs is a Higgs, and not a higher spin look-alike, then we should use the  $h \rightarrow ZZ$  angles  $\Theta$ ,  $\Phi$  as discriminators
- Furthermore, even for the spin 0 case, it is NOT TRUE that the distributions are flat in these angles, after we take into account realistic detector effects:



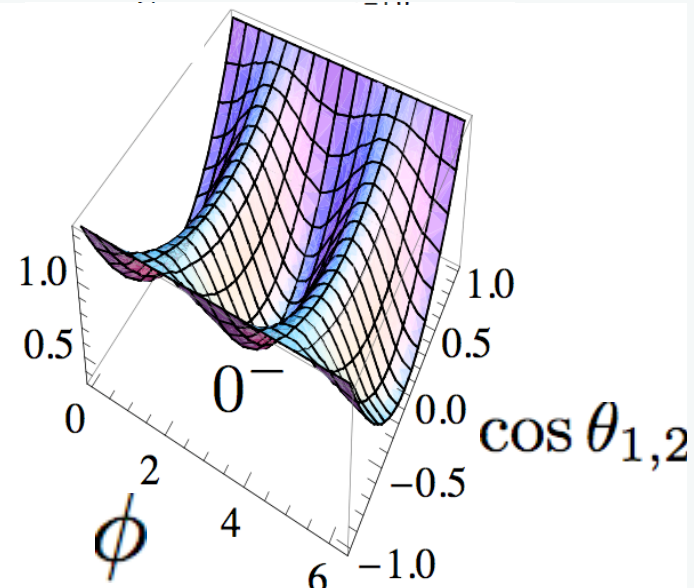
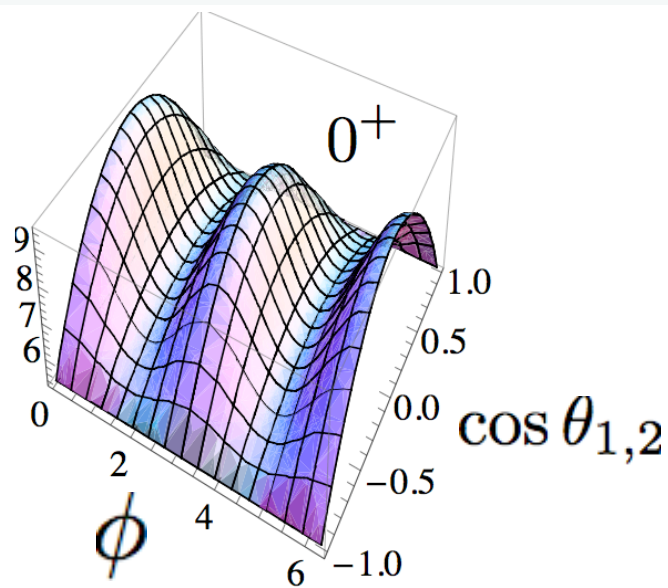
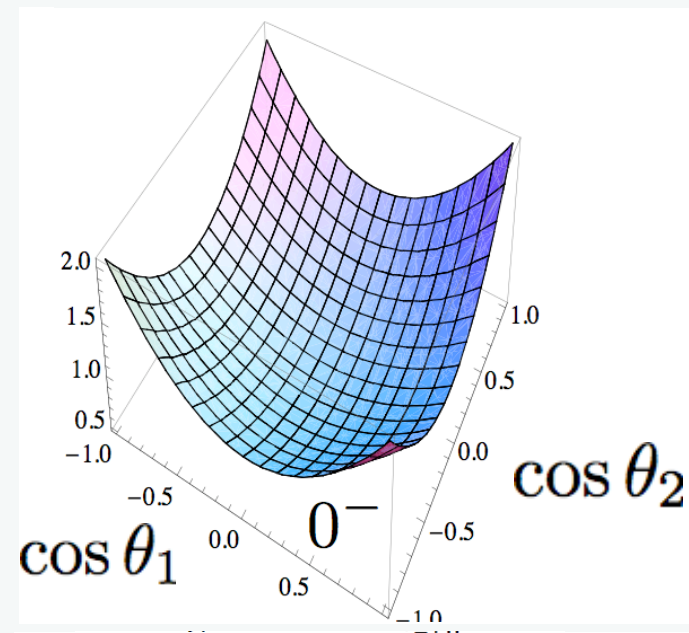
- For the three nontrivial  $ZZ \rightarrow \mu^- \mu^+ e^- e^+$  decay angles  $\theta_1, \theta_2, \phi$ , detector effects flatten the polar distributions slightly and sharpen the azimuthal distribution slightly
- Higher order corrections have similar (computable) effects



# Correlations are important

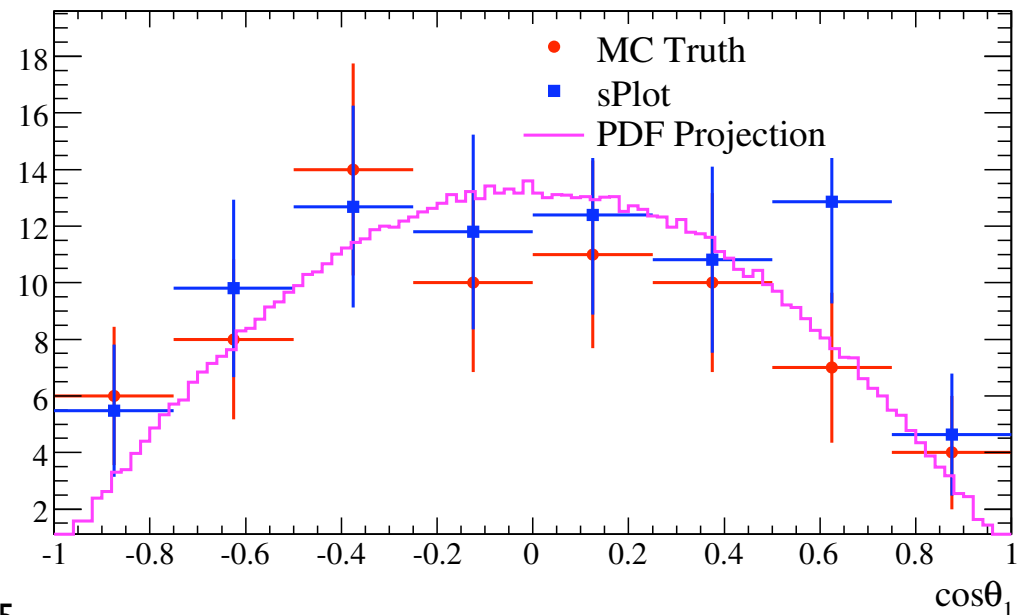
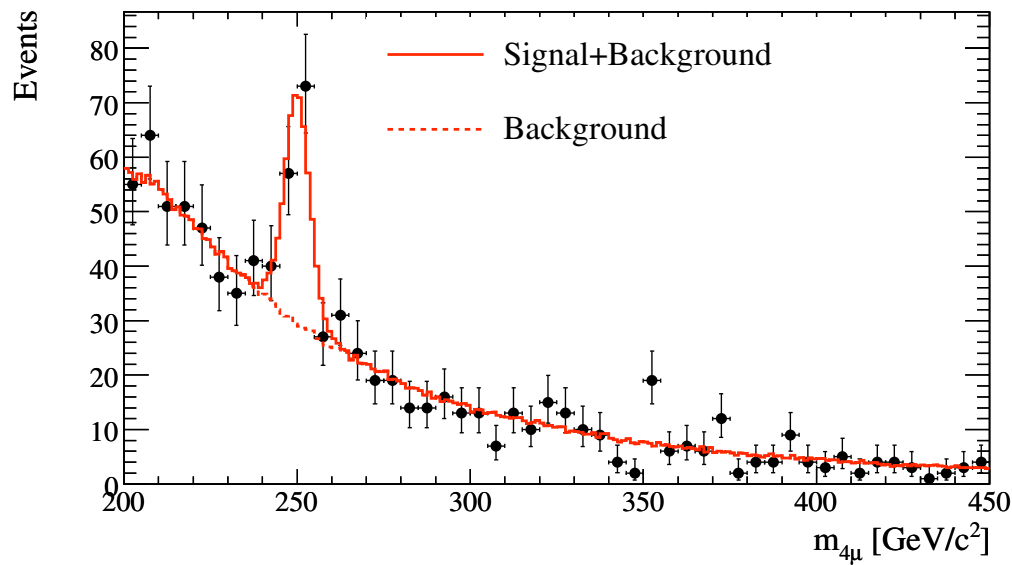


$m_H = 200 \text{ GeV}$



# What about the backgrounds?

- Although there is a fairly large irreducible background from ZZ production, this can be “subtracted” using a fit+weighting scheme called sPlots



# General couplings of Higgs Look-alikes to ZZ

- Allow couplings up to dimension 6
- Allow spin 0, 1, 2, and all possible C and P
- Note includes derivative couplings as would occur e.g. from expanding the form factor of a composite spin 0

$$\mathbf{L}_{\mu\nu}^0 = \mathbf{X} \mathbf{g}_{\mu\nu} - (\mathbf{Y} + \mathbf{iZ}) \frac{\mathbf{p}_\mu^h \mathbf{p}_\nu^h}{\mathbf{M}_Z^2} + (\mathbf{P} + \mathbf{iQ}) \epsilon_{\mu\nu\rho\sigma} \frac{\mathbf{p}_1^\rho \mathbf{p}_2^\sigma}{\mathbf{M}_Z^2}$$

$$\mathbf{L}_1^{\mu\nu\rho} = \mathbf{X} (\mathbf{g}^{\mu\nu} \mathbf{p}_1^\rho + \mathbf{g}^{\mu\rho} \mathbf{p}_2^\nu) + (\mathbf{P} + \mathbf{iQ}) \epsilon_{\rho\sigma}^{\mu\nu} (\mathbf{p}_1^\sigma - \mathbf{p}_2^\sigma)$$

$$\begin{aligned} \mathbf{L}_2^{\mu\nu\rho\sigma} = & \mathbf{M}_h^2 \mathbf{X}_0 \mathbf{g}^{\mu\rho} \mathbf{g}^{\nu\sigma} + (\mathbf{X}_1 + \mathbf{iY}_1) (\mathbf{p}_1^\nu \mathbf{p}_2^\rho \mathbf{g}^{\sigma\mu} + \mathbf{p}_2^\mu \mathbf{p}_1^\rho \mathbf{g}^{\sigma\nu}) \\ & + (\mathbf{X}_2 + \mathbf{iY}_2) \mathbf{g}^{\mu\nu} \mathbf{p}_1^\rho \mathbf{p}_2^\sigma + (\mathbf{P} + \mathbf{iQ}) \epsilon_{\alpha}^{\rho\mu\nu} (\mathbf{p}_1^\alpha \mathbf{p}_2^\sigma - \mathbf{p}_2^\alpha \mathbf{p}_1^\sigma) \end{aligned}$$



# fully-differential decay widths (tree level)

- SM Higgs

$$\frac{d\Gamma[0^+]}{dc_1 dc_2 d\phi} \propto m_1^2 m_2^2 m_H^4 \left[ 1 + c_1^2 c_2^2 + (\gamma_b^2 + c^2) s_1^2 s_2^2 \right. \\ \left. + 2\gamma_a c s_1 s_2 c_1 c_2 + 2\eta^2 (c_1 c_2 + \gamma_a c s_1 s_2) \right]. \quad (14)$$

- pure 1-

$$4m_1^2 m_2^2 X^2 \gamma_b^2 \left[ g_1 S^2 s_1^2 s_2^2 (2\ell_0^2 m_d^4 - \ell^2 m_H^2 [m_1^2 \cos(2\varphi_1) + m_2^2 \cos(2\varphi_2)]) \right. \\ \left. + g_1 \ell^2 m_H^2 (1 + C^2) [2m_2^2 s_1^2 + 2m_1^2 s_2^2 - (m_1^2 + m_2^2) s_1^2 s_2^2] + 4\ell \ell_0 g_1 m_H m_d^2 C S [m_1 c_1 s_1 s_2^2 \sin \varphi_1 - m_2 c_2 s_2 s_1^2 \sin \varphi_2] \right. \\ \left. - 2\ell^2 m_H^2 m_1 m_2 s_1 s_2 ((1 + C^2)(g_1 c_1 c_2 - g_{\sigma\sigma}) \cos(\varphi_1 - \varphi_2) + S^2 (g_1 c_1 c_2 + g_{\sigma\sigma}) \cos(\varphi_1 + \varphi_2)) \right].$$

- pure 1+

$$P^2 \left[ \ell^2 g_1 m_H^2 S^2 s_1^2 s_2^2 [M_2^4 m_1^2 \cos(2\varphi_1) + M_1^4 m_2^2 \cos(2\varphi_2)] \right. \\ \left. + 8\ell_0^2 m_1^2 m_2^2 m_d^4 S^2 [g_1 (c_1^2 + c_2^2 + s_1^2 s_2^2 \sin(\varphi_1 - \varphi_2)^2) + 2g_{\sigma\sigma} c_1 c_2] \right. \\ \left. + (1 + C^2) \ell^2 g_1 m_H^2 [2M_1^4 m_2^2 s_1^2 + 2M_2^4 m_1^2 s_2^2 - (M_2^4 m_1^2 + M_1^4 m_2^2) s_1^2 s_2^2] \right. \\ \left. - 8\ell \ell_0 m_H m_d^2 m_1 m_2 C S [M_2^2 m_1 s_2 (g_1 c_2 s_1^2 \sin \varphi_1 \cos(\varphi_1 - \varphi_2) + c_1 (g_1 c_1 c_2 + g_{\sigma\sigma}) \sin \varphi_2) \right. \\ \left. - M_1^2 m_2 s_1 (g_1 c_1 s_2^2 \sin \varphi_2 \cos(\varphi_1 - \varphi_2) + c_2 (g_1 c_1 c_2 + g_{\sigma\sigma}) \sin \varphi_1) \right. \\ \left. + 2\ell^2 m_H^2 M_1^2 M_2^2 m_1 m_2 s_1 s_2 [(1 + C^2)(g_1 c_1 c_2 - g_{\sigma\sigma}) \cos(\varphi_1 - \varphi_2) - S^2 (g_1 c_1 c_2 + g_{\sigma\sigma}) \cos(\varphi_1 + \varphi_2)] \right].$$

## Hypothesis testing with likelihood ratios

- Statistical approach - Neyman-Pearson hypothesis test:
  - Each 'experiment' corresponds to some number of observed signal events,  $N$
  - Each event,  $i$ , corresponds to a set of observables,  $\vec{X}_i$
  - Construct likelihoods for the two different hypotheses based on the multidimensional PDF's  $P_{0+}$  and  $P_{0-}$

$$\mathcal{L}_{0+} = \prod_i^N P_{0+}(\vec{X}_i) \quad \mathcal{L}_{0-} = \prod_i^N P_{0-}(\vec{X}_i)$$

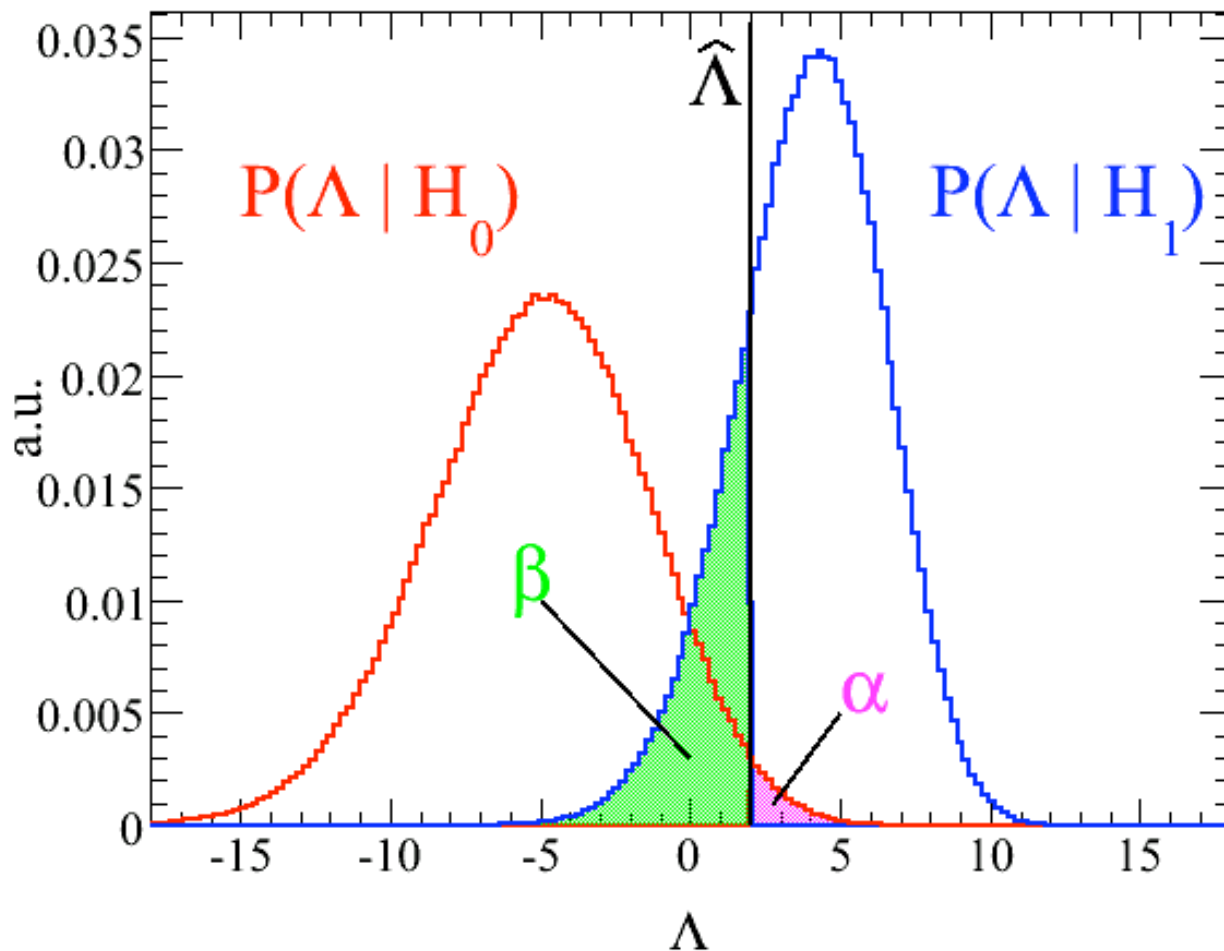
- Construct a test-statistic based on those likelihoods

$$\Lambda = \log(\mathcal{L}_{0+} / \mathcal{L}_{0-})$$

## Hypothesis testing with likelihood ratios

$$H_0 = 0^- \quad H_1 = 0^+ \quad \Lambda = \log(\mathcal{L}_{0+}/\mathcal{L}_{0-})$$

Neyman-Pearson (NP) simple hypothesis test



Risk of the 1st type:

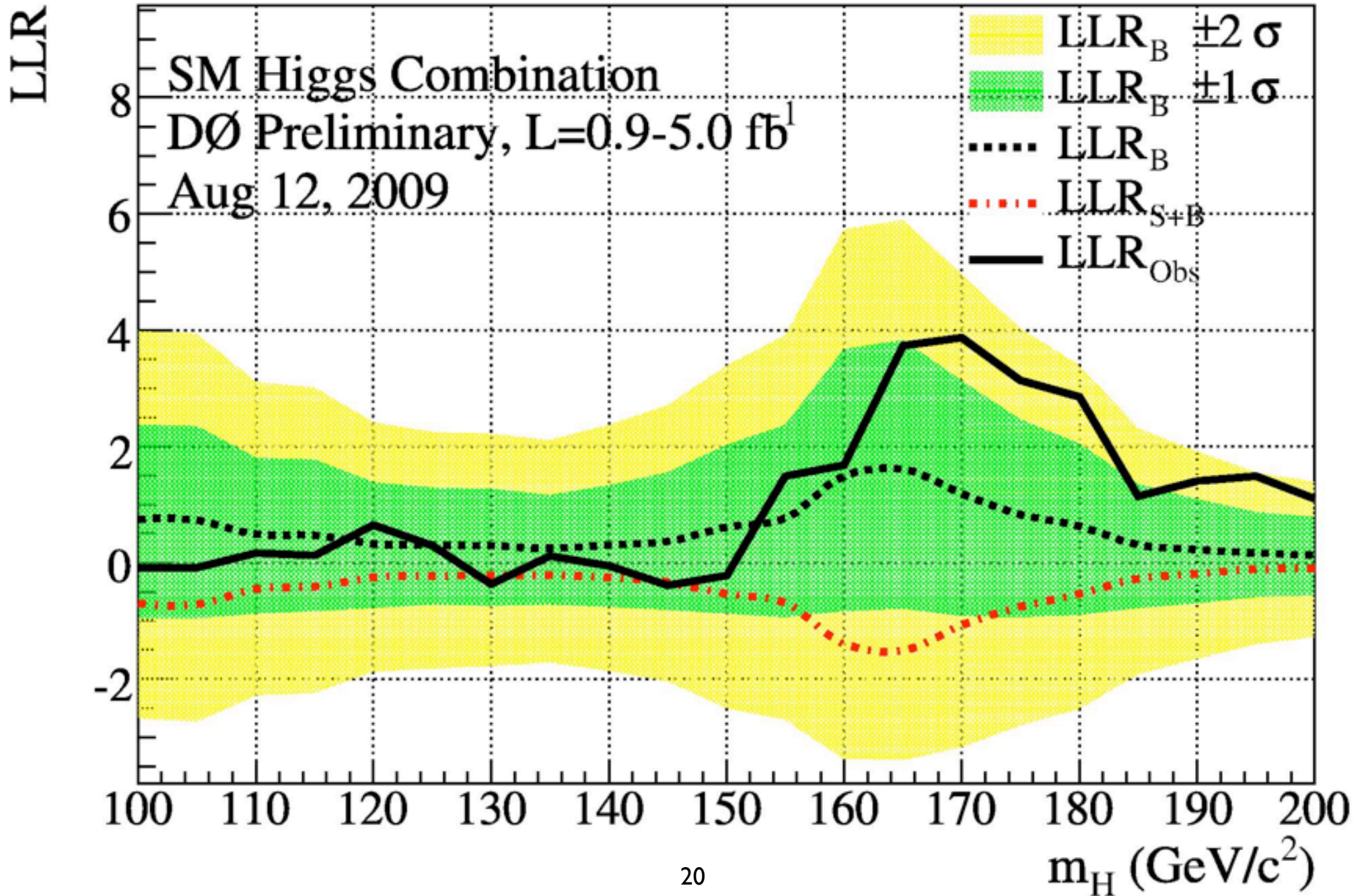
$$\alpha = \int_{\hat{\Lambda}}^{\infty} P(\Lambda | H_0) d\Lambda$$

Risk of the 2nd type:

$$\int_{-\infty}^{\hat{\Lambda}} P(\Lambda | H_1) d\Lambda = \beta$$

Power of the test:  $1 - \beta$

# Example of hypothesis testing: Higgs or no Higgs?



# Example: $0^+$ vs. $0^-$

- Consider the case when we are trying to distinguish between  $0^+$  vs.  $0^-$  resonances:

$$\gamma_a = \frac{1}{2m_1 m_2} [m_H^2 - m_1^2 - m_2^2]$$

$$\cos \theta_i = c_i, \sin \varphi = s$$

$$\eta \equiv \frac{2c_v v_a}{(c_v^2 + c_a^2)} \approx 0.15$$

**The standard Higgs,  $J^{PC} = 0^{++}$**

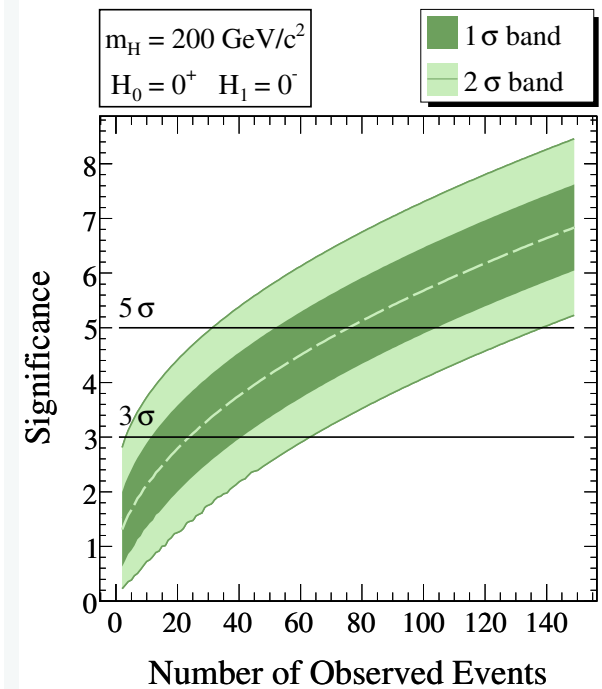
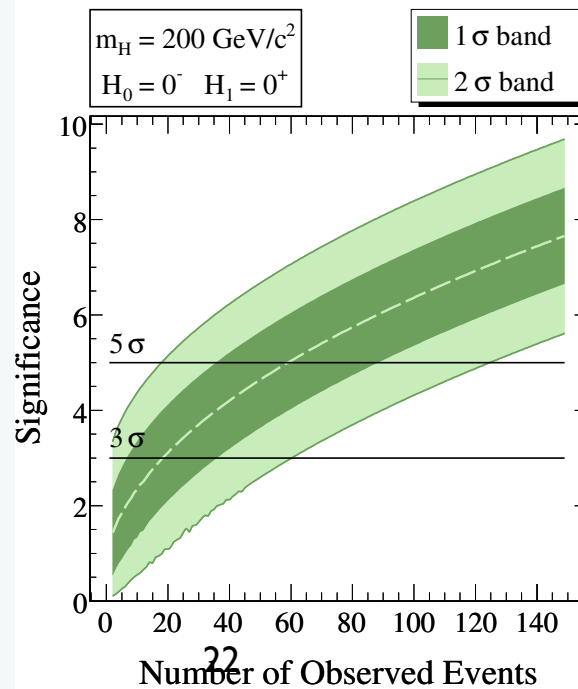
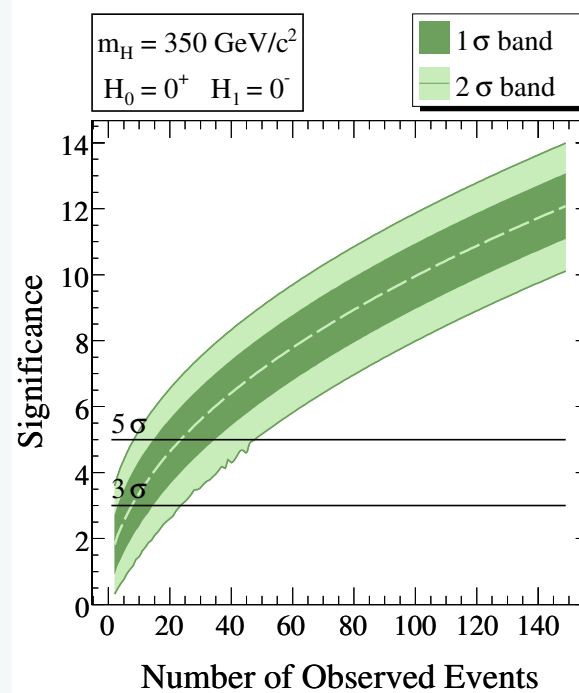
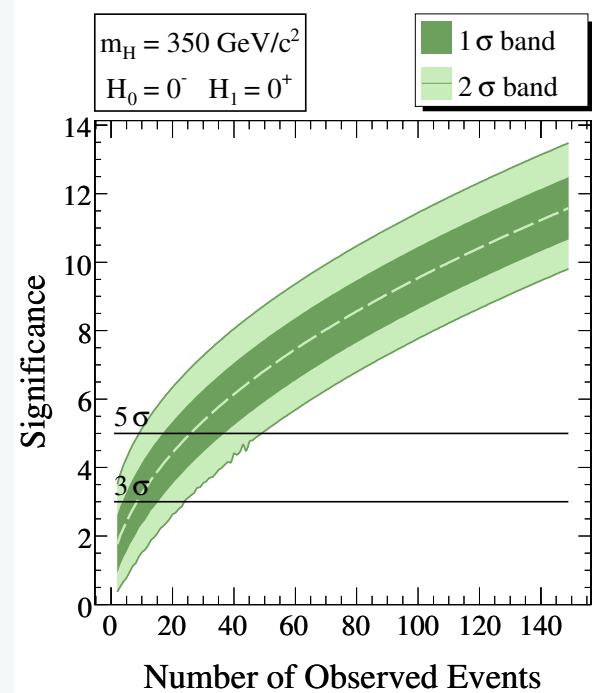
$$|\mathcal{M}[0^+]|^2 \equiv \frac{d\Gamma[0^+]}{dc_1 dc_2 d\varphi} \propto m_1^2 m_2^2 \{ 2(c_1 c_2 + c s_1 s_2 \gamma_a) \eta^2 + s_1^2 s_2^2 \gamma_a^2 + \frac{1}{2} [(2c^2 - 1) s_1^2 s_2^2 + (c_1^2 + 1)(c_2^2 + 1)] + 2c c_1 c_2 s_1 s_2 \gamma_a \}$$

**VS**

**A pure pseudoscalar,  $J^{PC} = 0^{-+}$**

$$|\mathcal{M}[0^-]|^2 \equiv \frac{d\Gamma[0^-]}{dc_1 dc_2 d\varphi} \propto m_1^4 m_2^4 \gamma_b^2 (c_1^2 c_2^2 + 2\eta^2 c_1 c_2 - c^2 s_1^2 s_2^2 + 1)$$

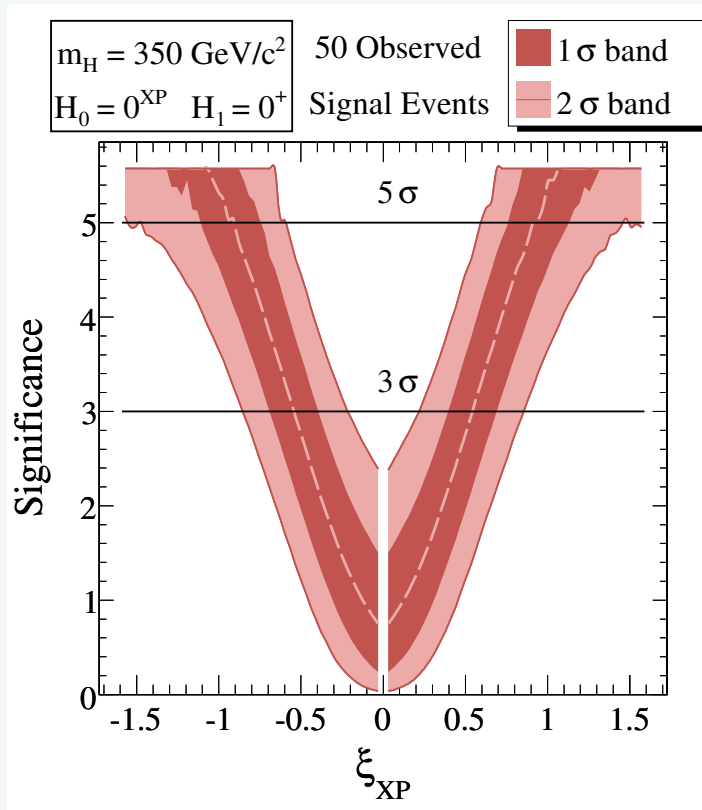
# Simple hypothesis test results, $0+$ versus $0-$



## 0+ versus a little bit of mixed CP

$$\mathcal{L}_{\mu\alpha} \propto \cos(\xi_{XP}) g_{\mu\alpha} + \sin(\xi_{XP}) \epsilon_{\mu\alpha} p_1 p_2 / M_Z^2$$

how small an admixture can I exclude when in fact it is an SM Higgs?



how large does the admixture have to be before I will be able to exclude the SM?

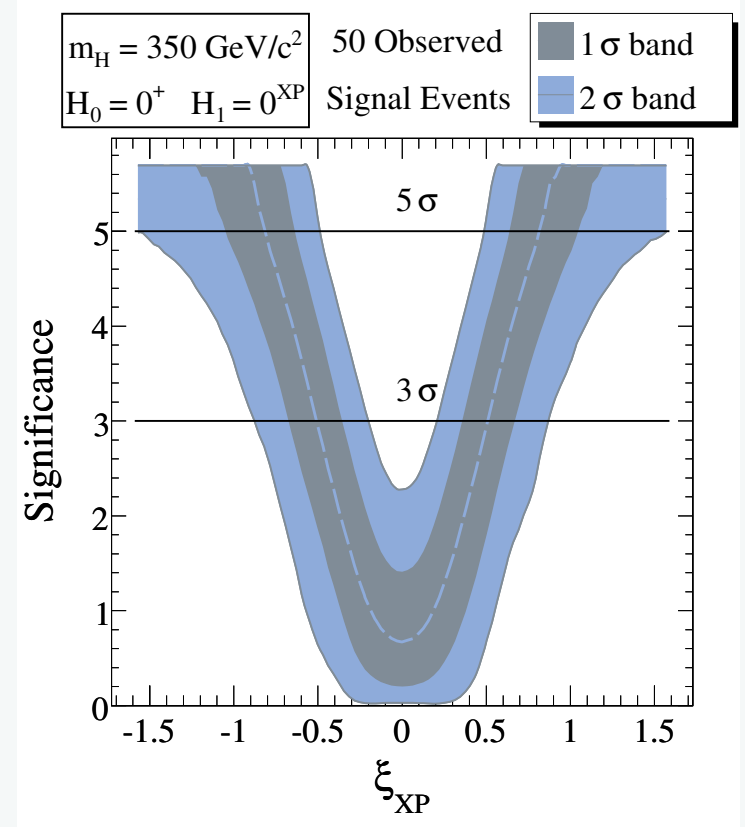


FIG. 37: Significance for excluding values of  $\xi_{XP}$  in the  $CP$ -violating  $J=0$  hypothesis in favor of the  $0^+$  one, assumed to be correct, for  $m_H=350 \text{ GeV}/c^2$  and  $N_S=50$ . The dashed line corresponds to the median of the significance. The 1 and 2  $\sigma$  bands correspond to 68% and 95% confidence intervals centered on the median value.

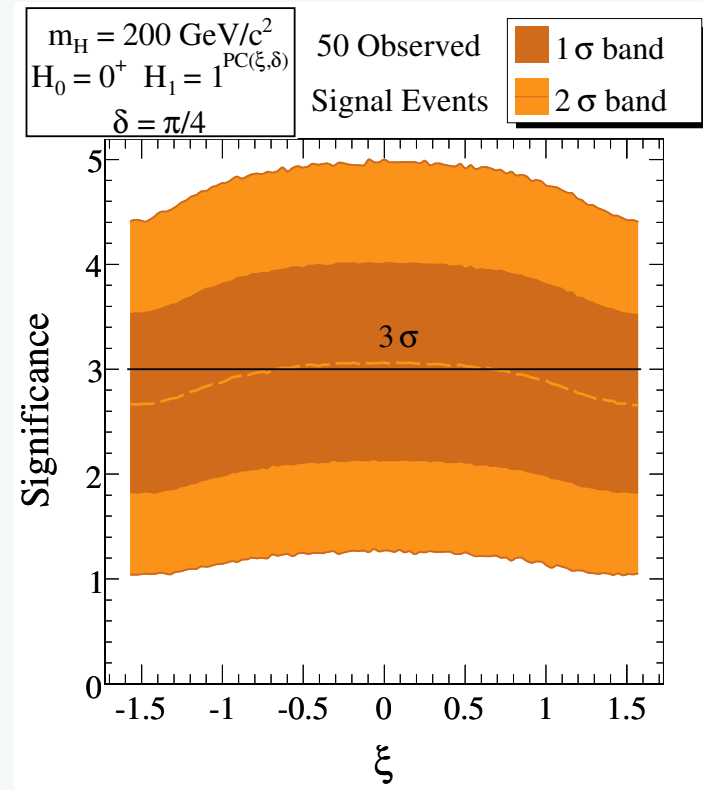
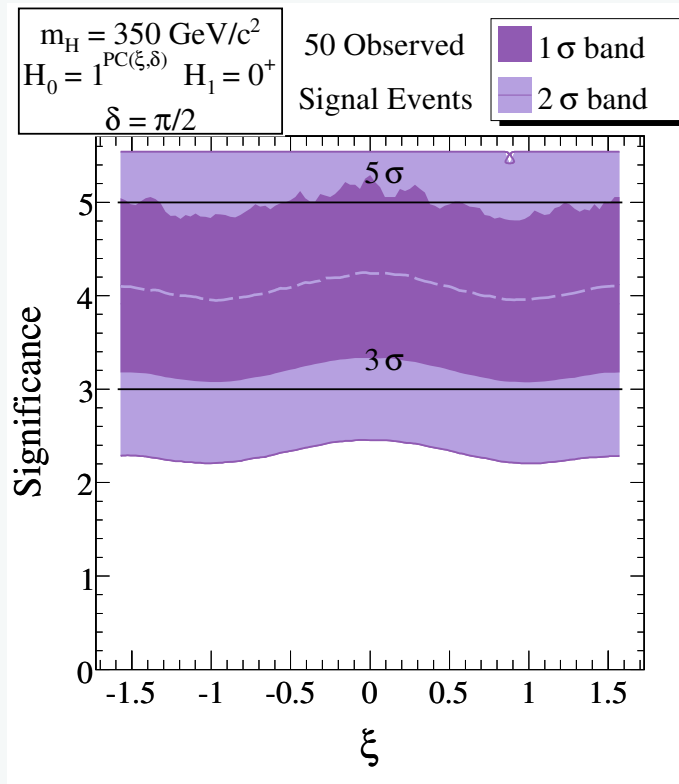
FIG. 38: The significance for excluding a pure  $0^+$  in favor of a  $CP$ -violating  $HZZ$  coupling ( $\xi_{XP} \neq 0$ ), assuming the latter to be correct, with  $\xi_{XP}$  given by its  $x$ -axis values. Example for  $N_S=50$ ,  $m_H=350 \text{ GeV}/c^2$ . Dashed line and bands as in Fig. 37.

## 0+ versus any possible spin 1 look-alike

$$\mathcal{L}^{\rho\mu\alpha} \propto \cos \xi (g^{\rho\mu} p_1^\alpha + g^{\rho\alpha} p_2^\mu) + e^{i\delta} \sin \xi \epsilon^{\rho\mu\alpha} (p_1 - p_2)$$

how well do I exclude arbitrary spin 1 when in fact I have a SM Higgs?

how well do I exclude an SM Higgs when in fact I have some arbitrary spin 1?



for SM Higgs masses (145, 200, 350) GeV we can exclude the general spin 1 hypothesis at 5 sigma with (60, 200, 85) signal events



# discriminating Higgs look-alikes at the moment of discovery

- number of signal events required for (median) 3 sigma discrimination:

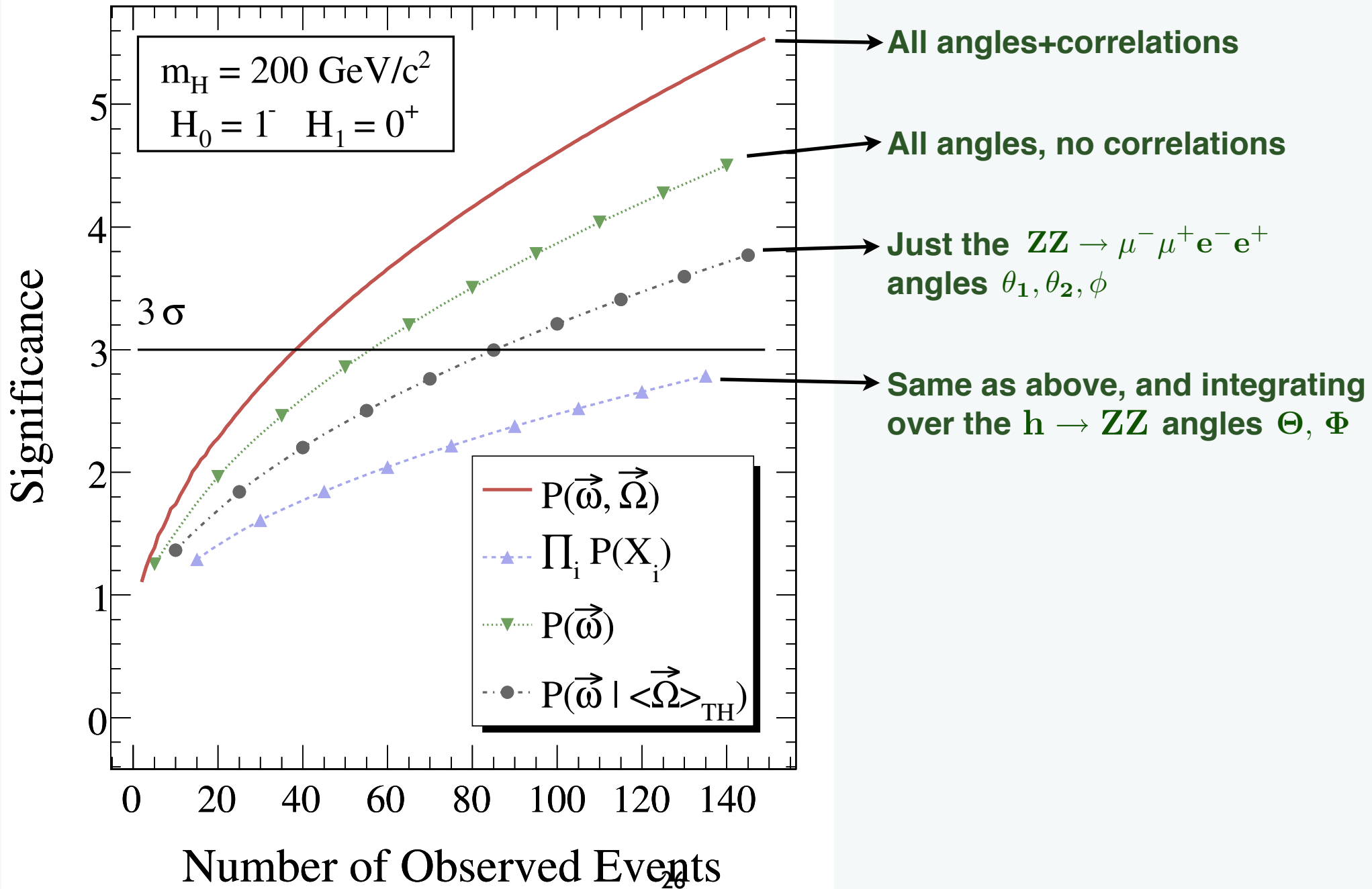
$\mathbb{H}_0 \downarrow \mathbb{H}_1 \Rightarrow$	$0^+$	$0^-$	$1^-$	$1^+$
$0^+$	–	17	12	16
$0^-$	14	–	11	17
$1^-$	11	11	–	35
$1^+$	17	18	34	–

TABLE I: Minimum number of observed events such that the median significance for rejecting  $\mathbb{H}_0$  in favor of the hypothesis  $\mathbb{H}_1$  (assuming  $\mathbb{H}_1$  is right) exceeds  $3\sigma$  with  $m_H=145 \text{ GeV}/c^2$ .

$\mathbb{H}_0 \downarrow \mathbb{H}_1 \Rightarrow$	$0^+$	$0^-$	$1^-$	$1^+$	$2^+$
$0^+$	–	24	45	62	86
$0^-$	19	–	19	19	38
$1^-$	40	18	–	90	48
$1^+$	56	19	85	–	66
$2^+$	86	45	54	70	–

TABLE III: Minimum number of observed events such that the median significance for rejecting  $\mathbb{H}_0$  in favor of the hypothesis  $\mathbb{H}_1$  (assuming  $\mathbb{H}_1$  is right) exceeds  $3\sigma$  with  $m_H=200 \text{ GeV}/c^2$ .

# The importance of using all the information



## The importance of using all the information

- **Should also put the production information back in (how to handle the uncertainties?)**
- **Should also use NLO to compute the likelihoods (do these calculations exist for all the possible look-alikes?)**
- **Should also include the other possible decay channels!**

**we have to measure other decay modes to determine the electroweak properties of the putative Higgs**

## Higgs electroweak look-alikes

- OK so you discovered a neutral resonance and used the first 20 events in the ZZ golden mode to exclude higher spins, large CP admixtures, etc.
- But is this particle the SM Higgs of electroweak symmetry breaking?
- Can we pin down the electroweak properties of the neutral resonance by measuring its branching fractions into electroweak vector bosons?

$$h \rightarrow W^+ W^-, ZZ, \gamma\gamma, Z\gamma$$

- what look-alikes should we worry about?
- do we need to measure all four branching fractions?

## Higgs electroweak look-alikes

$$\mathbf{h} \rightarrow \mathbf{W}^+ \mathbf{W}^-, \mathbf{ZZ}, \gamma\gamma, \mathbf{Z}\gamma$$

- Can do a general analysis making one additional assumption: the look-alike electroweak sector still respects custodial symmetry
- Thus the only look-alikes we have to worry about transform like some  $(\mathbf{N}_L, \mathbf{N}_R)$  under the global  $\mathbf{SU}(2)_L \times \mathbf{SU}(2)_R$  of which custodial  $\mathbf{SU}(2)_C$  is the diagonal remnant after EWSB

**Go to Ian Low's talk this afternoon!!**

## what look-alikes should we worry about?

$$\mathbf{h} \rightarrow \mathbf{W}^+ \mathbf{W}^-, \mathbf{ZZ}, \gamma\gamma, \mathbf{Z}\gamma$$

- $(\mathbf{1}_L, \mathbf{1}_R)$  an electroweak singlet with dimension 5 couplings to  $\mathbf{VV}$
- $(\mathbf{2}_L, \mathbf{2}_R)$  the SM case
- $(\mathbf{3}_L, \mathbf{3}_R)$  the custodial symmetry preserving combination of a real and a complex  $\mathbf{SU}(2)_L$  triplet
- $(\mathbf{4}_L, \mathbf{4}_R)$  some weird thing nobody bothers to talk about

In the last three cases we have dimension 4 couplings to  $\mathbf{WW}$  and  $\mathbf{ZZ}$

## do we need to measure all four branching fractions?

$$\mathbf{h} \rightarrow \mathbf{W}^+ \mathbf{W}^-, \mathbf{ZZ}, \gamma\gamma, \mathbf{Z}\gamma$$

Yes

$m_S$ (GeV)	$Br(\gamma\gamma/WW)$	$Br(ZZ/WW)$	$Br(Z\gamma/WW)$
115	$2.7 \times 10^{-2}$ ( $2.7 \times 10^{-2}$ )	$5.1 \times 10^{-2}$ (0.11)	39 ( $9.0 \times 10^{-3}$ )
120	$1.7 \times 10^{-2}$ ( $1.7 \times 10^{-2}$ )	$5.7 \times 10^{-2}$ (0.11)	35 ( $8.2 \times 10^{-3}$ )
130	$7.8 \times 10^{-3}$ ( $7.8 \times 10^{-3}$ )	$6.7 \times 10^{-2}$ (0.13)	26 ( $6.7 \times 10^{-3}$ )
140	$4.0 \times 10^{-3}$ ( $4.0 \times 10^{-3}$ )	$7.1 \times 10^{-2}$ (0.14)	18 ( $5.1 \times 10^{-3}$ )
150	$2.0 \times 10^{-3}$ ( $2.0 \times 10^{-3}$ )	$6.4 \times 10^{-2}$ (0.12)	10 ( $3.5 \times 10^{-3}$ )
170	$1.6 \times 10^{-4}$ ( $1.6 \times 10^{-4}$ )	$1.4 \times 10^{-2}$ ( $2.3 \times 10^{-2}$ )	0.81 ( $4.1 \times 10^{-4}$ )

TABLE II: Ratios of branching fractions for an electroweak singlet scalar when  $Br(\gamma\gamma/WW)$  is tuned to the SM value. The value in the parenthesis is for the corresponding SM prediction.

## what happens if you have violations of the custodial symmetry in the “Higgs” sector?

- **This is more complicated...**
- **But there is a general analysis for the case of the MSSM**
- **Here you have to consider all operators up to dimension 6**

**See the slides from Jose Zurita's talk yesterday!!**



## what happens if you have violations of the custodial symmetry in the “Higgs” sector?

- There are generically violations of custodial symmetry at both tree level and loop level (and perhaps cancellations)
- There are also generic custodial-preserving effects that, e.g. raise the Higgs mass

Custodially violating (tree level) :

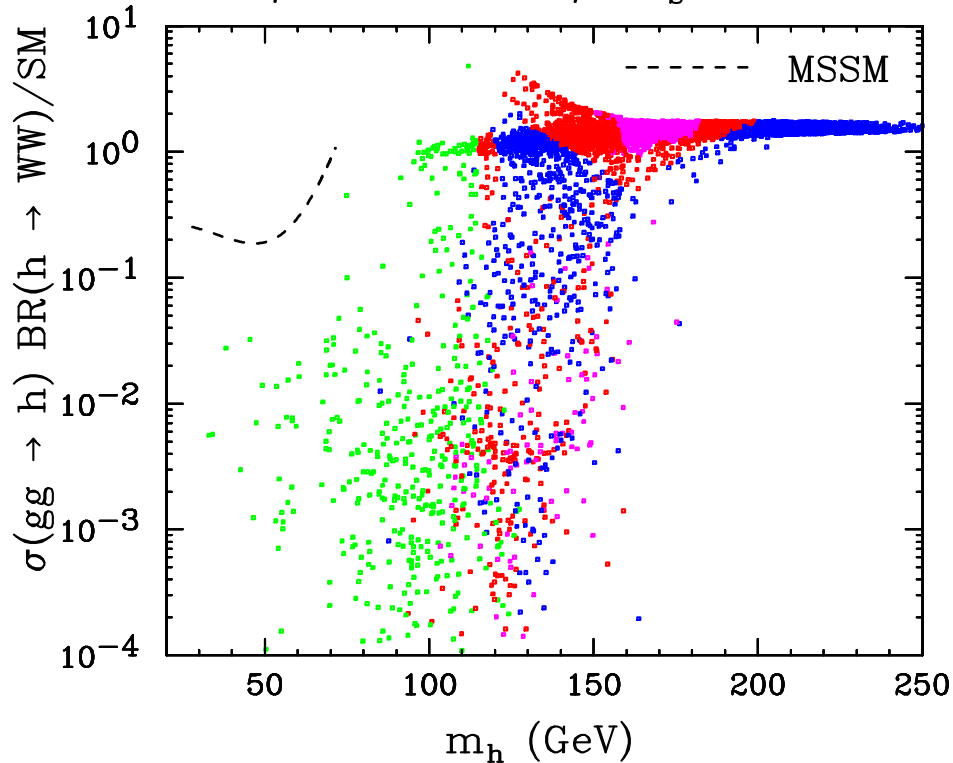
$$\Delta K^{\text{CV}} = \frac{c_1}{2|M|^2} (H_d^\dagger e^{2V} H_d)^2 + \frac{c_2}{2|M|^2} (H_u^\dagger e^{2V} H_u)^2 + \frac{c_3}{|M|^2} (H_u^\dagger e^{2V} H_u)(H_d^\dagger e^{2V} H_d)$$

Custodially preserving (tree level) :

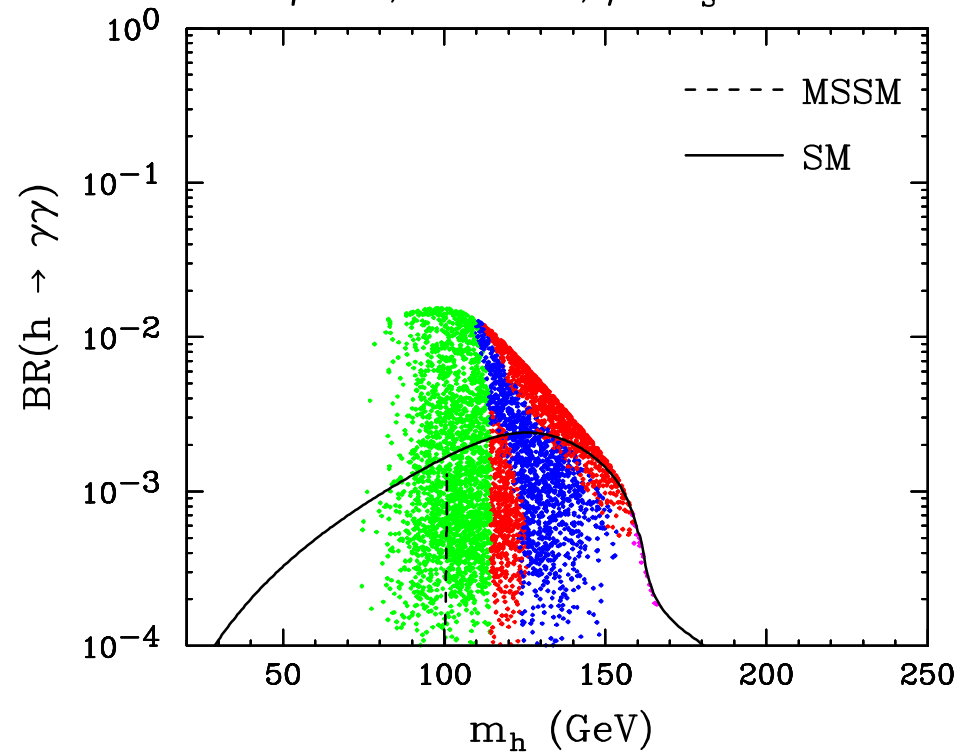
$$\Delta K^{\text{Cust}} = \frac{c_4}{|M|^2} |H_u H_d|^2 + \left[ \frac{c_6}{|M|^2} H_d^\dagger e^{2V} H_d + \frac{c_7}{|M|^2} H_u^\dagger e^{2V} H_u \right] (H_u H_d) + \text{h.c.}$$

# large differences in the branching fractions to VV, compared to the MSSM or SM

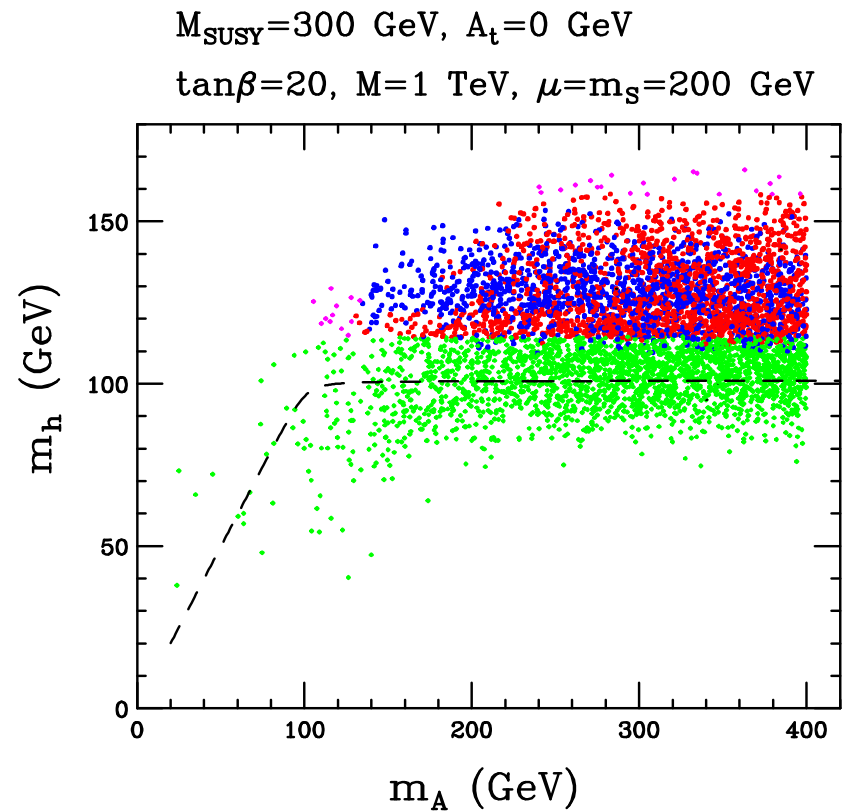
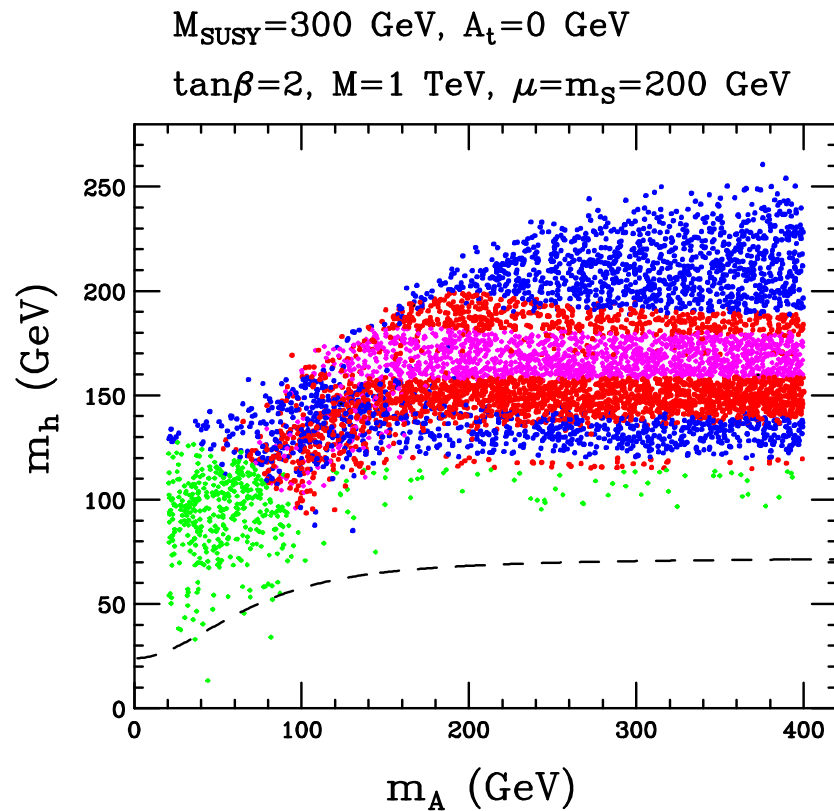
$M_{\text{SUSY}}=300 \text{ GeV}$ ,  $A_t=0 \text{ GeV}$   
 $\tan\beta=2$ ,  $M=1 \text{ TeV}$ ,  $\mu=m_s=200 \text{ GeV}$



$M_{\text{SUSY}}=300 \text{ GeV}$ ,  $A_t=0 \text{ GeV}$   
 $\tan\beta=20$ ,  $M=1 \text{ TeV}$ ,  $\mu=m_s=200 \text{ GeV}$



## also changes the bounds from LEP and Tevatron Higgs searches



- |   |   |   |
|---|---|---|
| <span style="color: green;">■</span> Excluded by LEP        | <span style="color: red;">■</span> Tevatron upgrade | → $10 \text{ fb}^{-1} + 50\% \text{ efficiency in } b\bar{b}, WW$ |
| <span style="color: magenta;">■</span> Excluded by Tevatron | <span style="color: blue;">■</span> Allowed         |   |

## Conclusion

- **The LHC will (we hope) discover Higgs-like resonances**
- **We have powerful tools to figure out the identity of what we find**
- **Most of this does not require 1 ab<sup>-1</sup> or an ILC, but it will require (i) more work to get ready, (ii) multi-channel searches, (iii) luck**